The Anisotropies of the C-C, C-H, and C=C Bonds

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The shielding effect of a distant bond (possessing axial symmetry) upon a resonating nucleus may be calculated from the McConnell equation¹ (Term I).

We have modified this equation to include a correction, Term II (as below), for the situations where the resonating nucleus is near to the bond

Term I

Term II

$$\sigma = \left[\frac{\Delta\chi}{3R^3} \left(1 - 3\cos^2\theta\right)\right] + \frac{S^2}{R^5} \left[-\frac{(\chi_{\rm L} + 2\chi_{\rm T})}{2} + 5\left(\chi_{\rm L}\cos^2\theta + \chi_{\rm T}\sin^2\theta\right) - \frac{35}{6}\left(\chi_{\rm L}\cos^4\theta + \chi_{\rm T}\sin^4\theta\right)\right] (1)$$

(*i.e.*, ~ 3 Å or less.) In (1) S is half the length of the induced dipole: for the C–C and C–H bonds we use S = 0.6 Å and S = 0.25 Å respectively.

In calculating the shift of a distant nucleus, produced by the introduction of a substituent into a molecule, both the substituted and unsubstituted molecules must be considered since the resultant shift

$$\Delta \sigma = {}_{1}^{n} \Sigma \sigma_{\text{substituted}} - {}_{1}^{n} \Sigma \sigma_{\text{unsubstituted}}$$

where ${}_{1}^{n}\Sigma$ is the sum of the screening effects of all *n* bonds in the molecules. It follows that the net chemical shift is caused only by those bonds which differ in the substituted and unsubstituted molecules. Hence the chemical shift for any proton in proceeding from an unsubstituted to a substituted molecule is the algebraic sum of the screening constants

$$egin{array}{lll} \Delta\chi^{ ext{C-C}} = + \ 13.98 imes 10^{-30}; \ \chi_{ ext{L}}^{ ext{C-C}} = + \ 3.15 imes 10^{-30}; \ \chi_{ ext{T}}^{ ext{C-C}} = - \ 10.83 imes 10^{-30}; \end{array}$$

for all the bonds displaced together with all the bonds introduced.

We have deduced values for the C-H and C-C anisotropies by applying this concept and Equation (1) to cyclohexane where the chemical shift difference ($\Delta \sigma_{a-e}$) between the axial and equatorial protons has been accurately measured and for which the relationship (2) has been derived.² For

$$0.114\Delta\chi^{\text{C-C}} - 0.094 \Delta\chi^{\text{C-H}} = 0.46 \times 10^{-6}$$
 (2)

any pair of geminal protons several of the screening bonds responsible for the differential *axial equatorial* shift are less than 3 Å distant. An initial estimate of the correction factor [Term II in (1)] was made using the susceptibilities of Davies.³ From our measurements and calculations $\Sigma\Delta\sigma_{a-e}$ (Term II) = 0.093 p.p.m.

Since $\Sigma\Delta\sigma_{a-e}$ (Term I) + $\Sigma\Delta\sigma_{a-e}$ (Term II) = 0.46×10^{-6}

$$\Sigma\Delta\sigma_{\mathrm{a-e}}~(\mathrm{Term}~\mathrm{I}) = 0.367 imes 10^{-6}$$

Hence Equation (2) becomes

$$0.114\Delta\chi^{\text{C-C}} - 0.094\Delta\chi^{\text{C-H}} = 0.367 \times 10^{-6}$$
 (3)

From magnetic-birefringence measurements upon ethane it has been deduced⁴ that

$$\Delta \chi^{\text{C-C}} - 2\Delta \chi^{\text{C-H}} = -4.8 \times 10^{-6}$$
 (4)

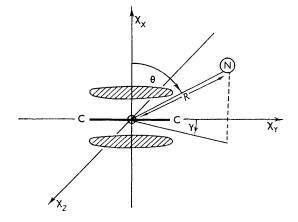
Thus, solving (3) and (4) simultaneously $\Delta \chi^{C-C} =$ + 8.86 × 10⁻⁶ = + 14.77 × 10⁻³⁰ and $\Delta \chi^{C-H} =$ + 6.84 × 10⁻⁶ = + 11.40 × 10⁻³⁰. Using the relationships⁵

 $\frac{1}{3}(\chi_L+2\chi_T)=$ - $6{\cdot}2$ \times 10^{-30} cm.3 molecule^-1 for the C–C bond

and $\frac{1}{3}(\chi_L+2\chi_T)=-6\cdot41\times10^{-30}\,\text{cm.}^3\,\text{molecule}^{-1}$ for the C–H bond,

magnetic susceptibilities corresponding to these anisotropies were then calculated for the C–C and C–H bonds. These susceptibilities were then used to compute the correction Term II. Successive iterations to constant values gave (5).

$$\begin{aligned} \Delta \chi^{\text{C-H}} &= + 11 \cdot 00 \times 10^{-30} \\ \chi_{\text{L}}^{\text{C-H}} &= + 0.92 \times 10^{-30} \\ \chi_{\text{T}}^{\text{C-H}} &= - 10.08 \times 10^{-30} \end{aligned} \tag{5}$$



All values are in cm.³ molecule⁻¹. The C-H anisotropy is thus not zero^{6,7} or negligible⁸ but is of a value similar to that suggested $(+11\cdot33 \times 10^{-30} \text{ cm}^3 \text{ molecule}^{-1})$ by Pople.⁷ For $\Delta \chi^{\text{C-C}}$ our value compares satisfactorily with Pople's estimate⁹ of $+18\cdot33 \times 10^{-30}$ cm.³ molecule⁻¹ and that of $+16\cdot7 \times 10^{-30}$ cm.³ molecule⁻¹ based¹⁰ on ¹³C-H coupling constants.

This treatment has been extended to the C=Cbond, which does not possess cylindrical or conical symmetry. Hence three susceptibilities (i.e., two anisotropies) are necessary to describe the anisotropy in the magnetic susceptibility. It is mathematically convenient to redefine the angles involved as in the Figure. The relevant equation for non-axially symmetrical groups is then (7).

The susceptibilities corresponding to (8) are: $\chi_{\rm x} = +3.68 \times 10^{-30}, \ \chi_{\rm y} = -9.07 \times 10^{-30}, \ {\rm and}$ $\chi_z = -10.59 \times 10^{-30}$ cm³. molecule⁻¹. The anisotropies (5) and (8) were used to calculate the chemical shifts of the C-18 and C-19 methyl groups of thirteen representative steroidal examples with double bonds in various positions. In a plot of calculated/observed results the line of best fit is

$$= \frac{1}{3R^3} \left[\Delta \chi_1 (1 - 3 \cos^2 \theta) + \Delta \chi_2 (1 - 3 \sin^2 \theta \cdot \sin^2 \gamma) \right]$$
(7)

where $\Delta \chi_1 = \chi_Y - \chi_X$, $\Delta \chi_2 = \chi_Z - \chi_X$; the X-axis coincides with the bond axis and the Y-axis is the direction orthogonal to the nodal plane of the π orbitals.

σ

Using the data for twenty-one steroidal examples twenty-one equations were set up and solved by the method of least squares. Value (5) for $\Delta \chi^{\text{C-C}}$ and $\Delta \chi^{\text{C-H}}$ gave (8).

approximately that of the theoretical line: with only two exceptions all calculated values of the methyl resonances are within \pm 3 c./sec. of the observed values.

Calculations were carried out on the Atlas computer of the University of London.

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$$\Delta \chi_1^{C=C} = -12.75 \times 10^{-30} \text{ and } \Delta \chi_2^{C=C} = -14.27 \times 10^{-30} \text{ cm}^3. \text{ molecule}^{-1}$$
(8)

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