# Intramolecular Rearrangement in $\mathrm{IF}_{7}$ and $\mathrm{XeF}_{6}$ 

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No single molecular geometry has proved capable of accounting for the electron-diffraction patterns reported for $\mathrm{IF}_{7}{ }^{1,2}$ and $\mathrm{XeF}_{6}{ }^{3,4}$ We present here for consideration an explanation for the reported data that involves intramolecular rearrangement between several molecular geometries. We feel that the agreement between the synthetic radial distribution curves obtained with our model and those reported is so good that the model should be given serious consideration.
$\mathrm{XeF}_{6}$ can be considered seven-co-ordinated if the valence electron lone pair manifests steric activity as predicted by the Gillespie-Nyholm ${ }^{5,6}$ theory. Therefore it is reasonable to seek for a common structural behaviour for both $\mathrm{XeF}_{8}$ and $\mathrm{IF}_{7}$. The calculations of Claxton and Benson ${ }^{7}$ show that in seven-co-ordination there is very little difference in energy between $D_{5 k}$ and $C_{3 v}$ configurations. In Figure 1 a transformation between $D_{5 h}$ (or its subgroup $C_{2 v}$ ) and $C_{3 v}$ along a path of $C_{s}$ configurations is illustrated. This transformation path is likely to be the path of minimum activation energy. By two successive transformations, either

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D_{5 h}\left(\text { or } C_{2 v}\right) \rightarrow C_{3 v} \rightarrow D_{5 h}\left(\text { or } C_{2 v}\right)
$$

or

$$
C_{3 v} \rightarrow D_{5 h}\left(\text { or } C_{2 v}\right) \rightarrow C_{3 v},
$$

a molecular rearrangement may be effected. This model uses $D_{5 \hbar} \longleftrightarrow C_{3 v}$ for $\mathrm{IF}_{7}$ and $C_{2 v} \longleftrightarrow C_{3 v}$ for $\mathrm{XeF}_{6}$.


Figure 1. Transformation of configurations. The seventh ligand represented by an open circle containing two dots is either a fuovine atom in $\mathrm{IF}_{7}$ or an electron lone pair in $\mathrm{XeF}_{6}$. There are 7! ways in which the ligands can be numbered. If the numbering is permuted after two or more transformations then a rearrangement has occurred. For example, at $\mathrm{C}_{3 v}$, if ligand 7 reorients along the symmetry plane between ligands 2 and 5 then on transformation to $\mathrm{D}_{5 n}\left(\right.$ or $\left.\mathrm{C}_{2 v}\right)$ the axial ligands will be 1 and 4 instead of 5 and 6.

To test the model synthetic radial distribution curves were calculated for comparison with the published curves. ${ }^{1-3}$ The latter were derived from experimental scattering curves by a Fourier inversion process. No single configuration such as
$D_{5 h}$ or $C_{3 v}$ taken by itself was able to provide curves akin to the reported distributions. Weighted averages were formed of several configurations including $D_{5 h}$ (or $C_{2 v}$ ) and $C_{3 v}$ and $C_{s}$ configurations at points removed one quarter, one half, and three quarters from the terminal configurations. To define the ligand co-ordinates the simplest possible assumptions were made including, for $\mathrm{IF}_{7}$, a constant I-F distance ${ }^{1}$ of $1.823 \AA$ and a minimum F-F distance of $2.145 \AA$ at both $D_{5 h}$ and $C_{3 v}$, for $\mathrm{XeF}_{8}$, a constant $\mathrm{Xe}-\mathrm{F}$ distance ${ }^{3,8}$ of $1 \cdot 90 \AA$ and a minimum F-F distance ${ }^{3,8}$ of $2 \cdot 60 \AA$ at $C_{3 v}$. Furthermore, we have assumed that the xenon nonbonding valence electron pair dipole moment has a value equal to $\sim 0.44$ of that of an Xe-F bond. With this value, all configurations chosen to represent the intramolecular rearrangement of the molecule possessed negligible dipole moments. This condition satisfied a requirement imposed by the findings of Falconer et al. ${ }^{9}$ that the $\mathrm{XeF}_{6}$ dipole moment is too small to be revealed by electric deflection experiments. The minimum F-F distance for $\mathrm{XeF}_{6}$ at $C_{2 v}$ was $2.48 \AA$. To compute the synthetic curves each interatomic distance was represented by a lambda peak with a constant base width and a height of $Z_{i} Z_{j} / r_{i j}$, where $Z_{i}$ and $Z_{j}$ are the atomic numbers of the atoms and $r_{i j}$ is the distance between them. The value of the base width for the lambda peaks was extrapolated from the main peak of a reported curve, $0.56 \AA$ for $\mathrm{IF}_{7}{ }^{1,2}$ and $0.5 \AA$ for $\mathrm{XeF}_{6} .{ }^{3}$

The uppermost curve, labelled $\mathrm{O}_{h}$ in Figure 2, illustrates the trivial case of the $O_{h}$ synthetic radial distribution curve for $\mathrm{XeF}_{6}$. Six Xe-F distances form the main peak at $1.90 \AA$, twelve $F-F$ distances form a large peak at $2.69 \AA$, and three $F-F$ distances form a small peak at $3 \cdot 80 \AA$. In this simple example the three peaks are widely separated and show the discontinuous lambda shape (the main peak is truncated to fit into the Figure). In the $C_{3 v}$ curve there are three $\mathrm{F}-\mathrm{F}$ peaks with only the one at $3.79 \AA$ fully resolved. In the $C_{2 v}$ curve there are five F-F peaks with only the one at $3 \cdot 21 \AA$ fully resolved. The curve for a $C_{s}$ configuration would contain nine F-F peaks. When three intermediate $C_{s}$ configurations are combined with $C_{2 v}$ and $C_{3 v}$ the resultant curve contains $35 \mathrm{~F}-\mathrm{F}$ peaks. By assigning various weights to the different configurations the resultant curve was modified by successive approximations to resemble the reported curve. ${ }^{3}$ The model curve in Figure 2 is based on uniform weights of $1: 1: 1: 1: 1$ for configurations from $C_{2 v}$ to $C_{3 v}$. In addition to the completely resolved Xe-F peak at $1.90 \AA$ it has a very broad peak at $2 \cdot 63 \AA$ with an extraordinary shoulder extending from about $2 \cdot 95 \AA$ to a minimum
at $3 \cdot 50 \AA$ followed by a small peak at $3 \cdot 87 \AA$. These features are in excellent agreement with those reported by Bartell et al. ${ }^{3}$


Figure 2. Synthetic radial distribution curves for $\mathrm{XeF}_{6}$. All $\mathrm{Xe}-\mathrm{F}$ and $\mathrm{F}-\mathrm{F}$ distances are represented by lambda peaks with $0.5 A$ base width. Each of the curves labelled $\mathrm{O}_{h}, C_{3 v}$, and $\mathrm{C}_{2 v}$ is based on a single configuration of the designated symmetry. The Model curve is based on weighting of $1: 1: 1: 1: 1$ of configurations from $\mathrm{C}_{2}$ to $\mathrm{C}_{3 \mathrm{v}}$.

The uppermost curve, labelled $C_{3 v}$, in Figure 3 illustrates a synthetic radial distribution for a single configuration of $\mathrm{IF}_{7}$, seven I-F distances form the main peak (truncated) at $1.825 \AA$. There are five F-F peaks, two of which overlap the base of the main peak and three of which form a separate overlapping group. In the $D_{5 h}$ curve there are four F-F peaks, two of which form a separate overlapping group. The curve for a $C_{s}$ configuration would contain $13 \mathrm{~F}-\mathrm{F}$ peaks. When three intermediate $C_{s}$ configurations are combined with $D_{5 h}$ and $C_{3 v}$ the resultant curve contains $48 \mathrm{~F}-\mathrm{F}$ peaks. Again, by assigning various weights to the different configurations the resultant curve was modified to resemble the reported curves. ${ }^{1,2}$ The model curve in Figure 3 is based on weights of $5: 4: 3: 2: 1$ for
configurations from $D_{5 h}$ to $C_{3 v} \cdot \dagger$ The main peak at $1.825 \AA$ is followed by a barely resolved minor maximum at $2 \cdot 17 \AA$, then a very broad maximum


Figure 3. Synthetic radial distribution curves for $\mathrm{IF}_{7}$. All I-F and $\mathrm{F}-\mathrm{F}$ distances are represented by lambda peaks with 0.56 A base width. Each of the curves labelled $\mathrm{C}_{30}$ and $\mathrm{D}_{5 \hbar}$ is based on a single configuration of the designated symmetry. The Model curve is based on weighting of 5:4:3:2:1 of configurations from $\mathrm{D}_{5 h}$ to $\mathrm{C}_{3 v}$.
at $2.52 \AA$ that does not fall to a minimum until $3 \cdot 15 \AA$, and finally another broad maximum at $3 \cdot 50 \AA$. Again these features are in good agreement
with the reported curves with one qualification. The reported curves are quite sensitive to the theoretical allowance that is made for the phase shift in the electron scattering atomic form factors for atoms of high $Z$ and low $Z$, and to the angular range of scattering data that is included in the Fourier inversion. The minor maximum at $2 \cdot 17 \AA$ in the synthetic curve which has a height about $20 \%$ less than the $2.52 \AA$ maximum is a highly variable feature in the reported curves. It appears as a shoulder on the main peak of the two curves reported by LaVilla and Bauer ${ }^{1}$ with an inflection point at the same height as the $2.52 \AA$ maximum in one curve, and with about $20 \%$ less height in the other curve. In the curve of Thompson and Bartell ${ }^{2}$ it appears as a barely resolved maximum with about $20 \%$ greater height than the next maximum.
From the viewpoint of its effect on a diffraction experiment the model can be interpreted either in terms of continuous ligand motions or in terms of a collection of molecules which have a weighted distribution over a series of configurations. In principle, microwave spectroscopy should reveal the transitions between states of different configuration but these states may be numerous or the energy separations small in which cases the detection of discrete peaks will be difficult. All available data suggests that there can be very little directional quality in the bonding in these molecules and our model is consistent with this interpretation. The uniform weighting found for the $\mathrm{XeF}_{6}$ configurations suggests that the reorientation of the valence lone pair is very ready. The steric activity of the valence lone pair appears to be significantly less than that of the seventh fluorine ligand in $\mathrm{IF}_{7}$. The effective repulsive forces between ligands in $\mathrm{IF}_{7}$ may be quite different from those in $\mathrm{XeF}_{6}$.
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[^0]${ }^{1}$ R. E. LaVilla and S. H. Bauer, J. Chem. Phys., 1960, 33, 182.
${ }^{2}$ H. B. Thompson and L. S. Bartell, Trans. Amer. Cryst. Assoc., 1966, 2, 190.
${ }^{3}$ L. S. Bartell, R. H. Gavin, and M. B. Thompson, J. Chem. Phys., 1965, 43, 2547.
${ }^{4}$ K. Hedberg, S. H. Peterson, and R. R. Ryan, J. Chem. Phys., 1966, 44, 1726.
${ }^{5}$ R. J. Gillespie and R. S. Nyholm, Quart. Rev., 1957, 11, 339.
${ }^{6}$ R. J. Gillespie, J. Chem. Educ., 1963, 40, 295.
7 T. A. Claxton and G. C. Benson, Canad. J. Chem., 1966, 44, 157.
${ }^{8}$ N. Bartlett, F. Einstein, D. F. Stewart, and J. Trotter, Chem. Comm., 1966, 550.
${ }^{\circ}$ W. E. Falconer, A. Büchler, J. L. Stauffer, and W. Klemperer, J. Chem. Phys., 1968, 48, 312.


[^0]:    $\dagger$ These ratios gave a satisfactory fit but with more precise data there could be a considerable deviation from integer ratios. The only significance to be attached to these ratios is that the $D_{5 h}$ configuration predominates.

