## **Triplet Excitons in Fremy's Salt**

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Summary Evidence for a thermally accessible triplet state in Fremy's salt is described, including measurements of the zero-field splitting, the temperature dependence of the magnetic susceptibility, and the singlet-triplet energy gap.

SEVERAL authors have inferred the diamagnetic nature of the yellow-orange salt  $K_4[(SO_3)_2NO]_2$  (Fremy's salt),<sup>1-4</sup> and attributed this to dimerization of monomeric nitrosodisulphonate ions in the salt. The most recent of these reports provides evidence for the existence of two distinct crystalline forms of Fremy's salt at room temperature.<sup>4</sup> Our observations are restricted to the yellow-orange modification of Fremy's salt (1) described therein.

We have noted that the colour of (1) changes from yellow-orange  $(20^{\circ})$  to brown-orange  $(100^{\circ})$ . We find that the colour change is reversible, and presumably arises from a change in the electronic configuration of (1). An excited state has been detected by e.s.r. and static magnetic susceptibility measurements.

At low temperature, powdered samples of (1) show a weak, complex spectrum near g = 2, which is attributed

changes are reversible. The e.s.r. spectrum of a powdered sample of (1) at 363 K is shown in Figure 1. The three pairs of absorptions (X, Y,

thousands of Gauss appear and grow in intensity. These



FIGURE 1. X-band first derivative e.s.r. spectrum of a powdered sample of (1) at 363 K.

and Z) centred about g = ca. 2 are ascribed to dimeric nitrosodisulphonate ions with S = 1, orientated with the three distinct principal axes of the zero-field splitting tensor lying, respectively, along the external magnetic field.<sup>5</sup> Assuming an isotropic g tensor with g = 2.0023, the



FIGURE 2. Temperature dependence of the absorption area A (in arbitrary units) of the half-field transition of a powdered sample of (1).

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zero-field splitting parameters D and E are  $\pm 0.0746$  $(\pm 0.0005)$  and  $\pm 0.0044$   $(\pm 0.0002)$  cm<sup>-1</sup>, respectively, and are temperature independent (290-390 K) within the experimental error. The zero-field splitting parameter  $(D^2 + 3E^2)^{1/2}$  is  $\pm 0.0776$  cm<sup>-1</sup> as calculated<sup>5,6</sup> from the position of  $H_{\min}$  (Figure 1), v = 9.315 GHz. The residual monomer obscures the observation of any double quantum transition in the g = 2 region.

We have measured the temperature dependence of the absorption area A of the half-field transition, and find that it increases exponentially with temperature. Here A is proportional to the static magnetic susceptibility, and is taken to be of the form<sup>6,7</sup>  $CT^{-1}[3 + \exp(\hat{E}_{j}/kT)]^{-1}$ . The relation assumes that the singlet-triplet energy gap  $E_j$  is independent of temperature and that interaction between neighbours is negligible. With the approximation [3+ $\exp(E_j/kT)$ ]  $\simeq \exp(E_j/kT)$ , the data plotted in Figure 2 gave a value of  $0.27 \pm 0.01$  eV for  $E_i$ .

Moreover, the exponential population growth of the paramagnetic excited state with temperature was observed by static magnetic susceptibility measurements (Faraday method) for two independently prepared powdered samples of (1), as shown in Figure 3. For the above approximated relation,  $E_{\mathbf{a}}$  is  $0.27 \pm 0.01$  and  $0.23 \pm 0.01$  eV for these data. The observed molar diamagnetic correction of (1) is  $-194 \pm 4 \times 10^{-6}$  c.g.s. units.

To our knowledge there are no previous examples of inorganic salts with a thermally accessible triplet state.



FIGURE 3. Temperature dependence of the corrected molar susceptibility for two independently prepared powdered samples of (1).

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