## Predictions of Structures and Stabilities for Cl<sub>2</sub>F<sub>2</sub>, Cl<sub>3</sub>F, and Cl<sub>4</sub>

## **Colin J. Marsden**

*Department of Inorganic Chemistry, University of Melbourne, Parkville, Victoria 3052, Australia* 

Extensive ab *initio* calculations have predicted structures and energies for the T-shaped molecules Cl<sub>2</sub>F<sub>2</sub>, Cl<sub>3</sub>F, and Cl<sub>4</sub>; each is thermodynamically stable to atomisation, but not to dissociation into the appropriate mixture of CIF and  $Cl<sub>2</sub>$ .

Although  $CIF_3$  has been well known for many years, there are only fragmentary reports of any derivatives. Photolysis of a mixture of  $Cl_2$  and  $F_2$  in an inert matrix produces new i.r. bands, one of which has been assigned to  $Cl_2F_2$ .<sup>1,2</sup> Interaction of Cl<sub>2</sub> and IrF<sub>6</sub> leads to an unstable volatile substance, of stoicheiometry  $Cl_3F$ .<sup>3</sup> To help characterise these systems, I have performed an extensive series of *ab initio* calculations, and report stable local minima in the potential energy surfaces of  $Cl_2F_2$  (2 isomers),  $Cl_3F$  (2 isomers), and  $Cl_4$ .

Geometries were optimised by gradient techniques, using the programs Gaussian  $80<sup>4</sup>$  and Gaussian  $82<sup>5</sup>$  and a mediumsized 'double-zeta plus polarisation' basis **A;** (10,6,1)/[6,3,1]

**Table 1.** Calculated energies for  $F_2$ , ClF, Cl<sub>2</sub>, ClF<sub>3</sub>, Cl<sub>2</sub>F<sub>2</sub>, Cl<sub>3</sub>F, and Cl<sub>4</sub>.

			Atomisation <sup>b</sup>			Dissociation <sup>b,c</sup>		
	<b>SCFa</b>	$\triangle MP2^a$	<b>SCF</b>	MP2	Expt. <sup>d</sup>	<b>SCF</b>	MP <sub>2</sub>	Expt. <sup>d</sup>
F <sub>2</sub>	$-198.7326$	$-0.3198$	$-164$	143	155			
CIF	$-558.8470$	$-0.2811$	24	243	253			
Cl <sub>2</sub>	$-918.9047$	$-0.2697$	63	197	239			
CIF <sub>3</sub>	$-757.5854$	$-0.6382$	$-124$	500	523	15	113	113
$Cl_2F_2$ (A) <sup>c</sup>	$-1117.6269$	$-0.6154$	$-128$	416		$-176$	$-71$	
(B)	$-1117.6073$	$-0.6187$	$-180$	373		$-228$	$-114$	
Cl <sub>3</sub> F (A)e	$-1477.6516$	$-0.5977$	$-176$	300		$-263$	$-140$	
ΈB)	$-1477.6379$	$-0.5959$	$-212$	260		$-299$	$-180$	
$Cl_4$	$-1837.6871$	$-0.5724$	$-196$	193		$-321$	$-200$	

<sup>a</sup> In a.u.; 1 a.u.  $\sim$ 2.6256  $\times$  10<sup>3</sup> kJ mol<sup>-1</sup>. **b** In kJ mol<sup>-1</sup>. *c* To appropriate mixture of F<sub>2</sub>, CIF, Cl<sub>2</sub>. *d* From refs. 7, 15, corrected for zero-point vibrational effects. *e* Isomers A and **B** identified in Figure 1.



**Figure 1.** Scaled drawings of calculated structures for CIF,  $CI_2$ , CIF<sub>3</sub>,  $Cl_2F_2$ , FCI<sub>3</sub>, and CI<sub>4</sub>. Bond lengths in Å, angles in degrees, converged to  $0.001$  Å or  $0.1^{\circ}$ . Experimental structural information in parentheses, where available. Net atomic charges, in units of e, shown by atoms.

on C1, **(8,4,1)/[4,2,1]** on F.6 To demonstrate the level of reliability in structural predictions achievable with this basis for interhalogens, I also report calculated geometries for ClF, C12, and ClF3 in Figure **1.** All five structural parameters are within 1.7% of their experimental value,<sup>7,8</sup> with an average discrepancy of only 1.0%.

Predicted geometries of  $Cl_2F_2$ ,  $Cl_3F$ , and  $Cl_4$  are displayed in Figure **1.** Each was identified as a true local minimum, by the eigenvalue criterion of the analytically calculated secondderivative matrix. (However, the analogous species  $F_4$  was shown not to have a similar local minimum.) All structures are planar, and can be described as trigonal-bipyramids, with two equatorial lone pairs. Isomerism is possible for  $Cl_2F_2$  and  $CI<sub>3</sub>F$ ; isomers A have axial fluorines only, and isomers B have an equatorial fluorine. Consistent structural trends are apparent in this series of chlorine fluorides, as in phosphoranes which are valence-shell isoelectronic.9 As the number of fluorine atoms decreases, bonds lengthen at an increasing rate, with axial bonds more severely affected. Angles involving bonds to fluorine are smaller than to chlorine, consistent with valence shell electron pair repulsion (VSEPR) principles,lo though the prediction of angles greater than **90"** in  $\overline{Cl}_3F$  (isomer A) and  $\overline{Cl}_4$  is a little surprising.

Final SCF calculations were performed using a fairly large basis B; **(11,7,2)/[6,4,2]** on C1, **(9,5,1)/[5,3,1]** on F.6 Since the Hartree-Fock method gives a poor account of the energies of interhalogens,  $11-13$  the influence of electron correlation was estimated using second-order Møller-Plesset perturbation theory,l4 at the SCF geometries, and basis B, but with only a single roughly optimised polarisation function on C1, and a double-zeta s,p contraction for F. Core orbitals and their virtual counterparts were frozen. Although this procedure **is**  less than exact, more rigorous treatments were not feasible for the largest systems studied, and I thought it important to use a uniform level of theory throughout. In Table **1,** I report absolute and relative energies for the various species, and compare with experimental data where possible.<sup>7,15</sup> Binding energies calculated at the MP2 level are of useful accuracy in all cases where experimental results are available. Although this accuracy may be partly fortuitous,  $12,13$  one expects errors such as superposition to be comparable in a series of closely related molecules, so one may reasonably suppose that comparable accuracy will be achieved throughout.

Fluorine preferentially occupies axial sites in  $Cl_2F_2$  and C13F, as **is** generally found in trigonal-bipyramidal molecules.<sup>9,16</sup> Each of Cl<sub>2</sub>F<sub>2</sub>, Cl<sub>3</sub>F, and Cl<sub>4</sub> is found to be stable towards atomisation, but the results show that all three compounds are likely to be thermodynamically *unstable* with respect to the appropriate mixture of ClF and  $Cl<sub>2</sub>$ ; the least unstable species, isomer A of  $Cl_2F_2$ , is predicted to lie 71 kJ mol<sup>-1</sup> higher than 2ClF. It is most unlikely that the errors in the calculated dissociation energies are as large as **70** kJ  $mol<sup>-1</sup>$ . However, kinetic stability for any or all of these species is not excluded, because I have not investigated whether the energy wells are deep enough to support at least one vibrational level.

The electronic structure of  $CIF<sub>3</sub>$  has been discussed in detail,<sup>12,17</sup> and the derivatives  $Cl_2F_2$ ,  $Cl_3F$ , and  $Cl_4$  are qualitatively similar. Although the central chlorine atom may appear in a similar environment in each compound, stability towards atomisation decreases steadily as the number of fluorine substituents decreases. This trend may be understood by considering the net atomic charges,<sup>18</sup> displayed in Figure 1. Since terminal F atoms bear a greater negative charge than terminal C1, the central chlorine becomes progressively less positive in the series  $CIF_3 \rightarrow Cl_4$ . Simultaneously, its d orbital population decreases from **0.39** e in ClF3, through **0.36** e in

 $Cl_2F_2$ , 0.31 e in  $Cl_3F$ , to only 0.23 e in  $Cl_4$ . The quantitative importance of d orbitals in  $CIF_3$  has already been established.<sup>12,17</sup> We now see that decreasing thermodynamic stability in the series  $CIF_3 \rightarrow Cl_4$  is associated with, and perhaps caused by, decreasing d orbital involvement of the central chlorine, as are the trends in bond lengths already mentioned.

Experimental characterisation of  $Cl_2F_2^{1,2}$  and  $Cl_3F^3$  rests on i.r. absorptions at **6361,2** and 625 cm-1,3 which presumably are C1-F stretching modes. Unfortunately there is no direct relationship between bond lengths and stretching frequencies in polyatomic molecules. It is therefore premature to enquire how these observed frequencies compare with properties predicted for  $Cl_2F_2$  and  $Cl_3F$ .

It is a pleasure to acknowledge extremely rewarding discussions with Professor N. Bartlett.

*Received, 21st December 1984; Com. 1787* 

## **References**

- 1 M. R. Clarke, W. H. Fletcher, G. Mamantov, E. J. Vasini, and D. G. Vickroy, *Znorg. Nucl. Chem. Lett.,* 1972, **8,** 611.
- 2 E. **S.** Prochaska, L. Andrews, N. R. Smyrl, and G. Mamantov, *Inorg. Chem.,* 1978, **17,** 970.
- 3 L. Graham, Ph.D. thesis, University of California, Berkeley, 1978; N. Bartlett, personal communication, 1983.
- 4 **J. S.** Binkley, R. A. Whiteside, R. Krishnan, R. Seeger, D. **J.**  DeFrees, H. B. Schlegel, **S.** Topiol, L. R. Kahn, and J. A. Pople, *Quantum Chemistry Program Exchange,* 1981, **13,** 406.
- *5* J. **S.** Binkley, M. J. Frisch, D. J. DeFrees, K. Raghavachari, R. A. Whiteside, H. B. Schlegel, E. **M.** Fluder, and J. A. Pople, Chemistry Department, Carnegie-Mellon University, Pittsburgh, U.S.A., 1983.
- 6 **S.** Huzinaga, Technical Report, Chemistry Department, University of Alberta, Canada, 1971.
- 7 **K.** P. Huber and G. Herzberg, 'Constants of Diatomic Molecules,' Van Nostrand-Reinhold, New York, 1979.
- **8** D. F. Smith, *J. Chem. Phys.,* 1953, **21,** 609.
- 9 R. R. Holmes, 'Spectroscopy and Structures of Pentacoordinated Phosphorus Compounds,' ACS Monograph 175, 1980.
- 10 R. J. Gillespie, 'Molecular Geometry,' Van-Nostrand-Reinhold, New York, 1972.
- 11 B. Jonsson, **B.** 0. Roos, P. R. Taylor, and P. E. M. Siegbahn, *J. Chem. Phys.,* 1981, **74,** 4566 and references therein.
- 12 L. G. M. Pettersson, P. E. M. Siegbahn, and 0. Gropen, *Mol. Phys.,* 1983,48, 871.
- 13 M. R. A. Blomberg and P. E. M. Siegbahn, *Chem. Phys. Lett.,*  1981, **81,** 4.
- 14 J. A. Pople, J. S. Binkley, and R. Seeger, Int. J. Quantum Chem., *Symp.,* 1976, No. **10,** 1.
- 15 JANAF Thermochemical Tables, Air Force Contract 04(611)- 7554.
- 16 E. L. Muetterties, W. Mahler, **K.** J. Packer, and R. Schmutzler, *Znorg. Chem.,* 1964, *3,* 1298.
- 17 0. Gropen and E. Wislgff-Nilssen, J. *Mol. Struct.,* 1982,88,243.
- 18 R. *S.* Mulliken, *1. Chem. Phys.,* 1955, **23,** 1833.