

## HgH<sub>4</sub> and HgH<sub>6</sub>: further candidates for high-valent mercury compounds

Pekka Pyykkö, Michal Straka and Michael Patzschke

Department of Chemistry, University of Helsinki, P.O.B. 55 (A. I. Virtasen aukio 1), FIN-00014 Helsinki, Finland. E-mail: Pekka.Pyykko@helsinki.fi

Received (in Cambridge, UK) 15th April 2002, Accepted 11th June 2002

First published as an Advance Article on the web 16th July 2002

Mercury tetrahydride ( $D_{4h}$ ) is calculated to have similar bond lengths and vibrational frequencies as the already known HgH<sub>2</sub> and to lie energetically 200 kJ mol<sup>-1</sup> above HgH<sub>2</sub> + H<sub>2</sub>, in a local well, about 40 kJ mol<sup>-1</sup> below a transition state.

Each new oxidation state of an element is now a chemical achievement and a larger achievement at that, if a new atomic shell is formally oxidized or reduced for that particular element. Granted the existence of Au(III), which in fact is the most stable oxidation state of gold, the possible existence of Hg(IV), partially oxidizing the 5d shell, has been expected for some time. Jørgensen mentioned both Hg(III) and Hg(IV) as future possibilities.<sup>1</sup> Kaupp *et al.*<sup>2,3</sup> indeed calculated the formation reaction HgF<sub>2</sub> + F<sub>2</sub> → HgF<sub>4</sub> ( $D_{4h}$ ) to be exothermic. Although the prediction was confirmed by later calculation,<sup>4</sup> no successful experiments have been reported.

How about hydrogen? The inner bonding of H<sub>2</sub> is much stronger than that of F<sub>2</sub>; the  $D_0$  are 432.07 and 154.57 kJ mol<sup>-1</sup>, respectively.<sup>5</sup> While this would disfavour the mercury tetrahydride over the mercury tetrafluoride, the small size of the hydrogen might actually lead to a better hybridization with the mercury 5d shell. The relativistic destabilisation of the Hg 5d shell is expected to make Hg the best candidate in Group 12. Furthermore hydrogen as reagent has experimental advantages over fluorine. We have in mind either matrix spectroscopic or high-pressure experiments.

The known mercury hydride species comprise the diatomic HgH and HgH<sup>+</sup> as well as HgH<sub>2</sub> that has been studied both theoretically<sup>6,7</sup> and using matrix spectroscopy.<sup>8</sup> The latter authors also present evidence for HHgHgH. No earlier studies on HgH<sub>4</sub> were found. It should be noted that it is valence isoelectronic with the known square planar PtH<sub>4</sub><sup>2-</sup> anion, present *e.g.* in K<sub>2</sub>PtD<sub>4</sub><sup>9</sup> (Pt–D 163 pm).

We here report *ab initio* pseudo-potential calculations on HgH<sub>4</sub> and related systems. For technical calibration we add HgH<sup>+</sup> and HgH<sub>2</sub>. Also, noting that the 18-valence electron rule would favour HgH<sub>6</sub>, we have included it.

We used the hybrid density functional B3LYP,<sup>10,11</sup> MP2, and CCSD(T) methods implemented in Gaussian 98 software package<sup>12</sup> together with the Stuttgart small-core relativistic pseudopotential and corresponding basis set for mercury,<sup>13</sup> augmented by 2f functions with  $\alpha = 1.7$  and  $\alpha = 0.6$ ,<sup>7</sup> and by one g function with  $\alpha = 1.0$ . The TZVP basis set for hydrogen<sup>14</sup> was used.

The calculated Hg–H distances are listed in Table 1 and vibrational frequencies in Table 2. Somewhat surprisingly, the Hg–H bond lengths actually decrease from HgH<sub>2</sub> to HgH<sub>4</sub> and HgH<sub>6</sub>. The symmetric stretch frequencies are nearly constant along the same series and so are the antisymmetric vibrations ( $\Sigma_u$ ,  $E_u$  and  $T_{1u}$ , respectively). The energies of formation (compared with Hg(g) +  $n$ H<sub>2</sub>(g)) are included in Table 1. They increase from the (CCSD(T)) 116 kJ mol<sup>-1</sup> for the observed HgH<sub>2</sub> to 340 and 575.3 kJ mol<sup>-1</sup> for HgH<sub>4</sub> and HgH<sub>6</sub>, respectively. Noting that  $3D_0(\text{H}_2)$  is 1296 kJ mol<sup>-1</sup> even the hexahydride is far below the atoms. Note that H atoms in rare-gas matrices are easy to create.

The planar  $C_{2v}$  transition state for dissociation of HgH<sub>4</sub>

(B3LYP) is shown in Fig. 1. It lies 39 kJ mol<sup>-1</sup> above the  $D_{4h}$  minimum at B3LYP level (including the zero-point energy correction), or 34 kJ mol<sup>-1</sup> at the MP2 level. The dissociation leads to the in-plane loss of one H<sub>2</sub> molecule. HgH<sub>6</sub> also has a  $C_{2v}$  transition state, leading to H<sub>2</sub> + HgH<sub>4</sub>. It lies 30 kJ mol<sup>-1</sup> above the hexahydride at both B3LYP and MP2 levels.

The bonding of HgH<sub>4</sub> is delightfully simple. The B3LYP molecular orbitals are shown in Fig. 2. Letting the Hg–H bonds

**Table 1** The Hg–H distances and zero-point energy corrected energies of formation for systems considered

System	Method	Hg–H/pm	$\Delta E/\text{kJ mol}^{-1}$
HgH <sub>6</sub>	B3LYP	163.6	607.1
	MP2	161.9	598.2
	CCSD(T)	163.5	575.3
HgH <sub>4</sub>	B3LYP	164.1	312.1
	MP2	161.8	348.7
	CCSD(T)	163.5	339.7
HgH <sub>2</sub>	B3LYP	165.3	104.4
	MP2	162.5	129.0
	CCSD(T)	164.5	115.7
HgH <sup>+</sup>	CCSD(T) <sup>a</sup>	163.9	100.6
	B3LYP	161.0	
	MP2	157.9	
	CCSD(T)	160.7	
	CCSD(T) <sup>a</sup>	159.6	
	Exptl. <sup>b</sup>	159.4	

<sup>a</sup> Ref. 7. <sup>b</sup> Ref. 15.

**Table 2** Harmonic vibrational frequencies with intensities in parentheses; frequencies in cm<sup>-1</sup>, intensities in km mol<sup>-1</sup>

Species	Symmetry	B3LYP	MP2	Exptl.
HgH <sub>6</sub>	$E_g$	2090	2165	
	$A_{1g}$	2063	2129	
	$T_{1u}$	1984 (266)	2080 (187)	
	$T_{1u}$	810 (56)	789 (69)	
	$T_{2g}$	755	759	
	$T_{2u}$	695	688	
	$B_{1g}$	2085	2194	
HgH <sub>4</sub>	$A_{1g}$	2060	2154	
	$E_u$	1959 (356)	2058 (310)	
	$B_{2u}$	884	909	
	$A_{2u}$	839 (10)	842 (7)	
	$B_{2g}$	759	772	
	$E_u$	764 (40)	757 (58)	
	$B_{1g}$	1475		
HgD <sub>4</sub>	$A_{1g}$	1457		
	$E_u$	1393 (178)		
	$B_{2u}$	625		
	$A_{2u}$	599 (5)		
	$B_{2g}$	537		
	$E_u$	543 (21)		
	$\Sigma_g$	2033	2164	
HgH <sub>2</sub>	$\Sigma_u$	1922 (442)	2023 (425)	1896 <sup>a</sup>
	$\Pi_u$	799 (32)	821 (31)	773 <sup>a</sup>
	$\Sigma_g$	1987 (137)	2092 (158)	2028 <sup>b</sup>

<sup>a</sup> Ref. 15. <sup>b</sup> Ref. 5.

lie in the  $x$  and  $y$  directions, the lowest bonding MO is the  $b_{1g}$ ,  $5d_{x^2-y^2} + 1s$ . The next one, HOMO-4 is the  $a_{1g}$   $5d_{3z^2-r^2} + 1s$ . Its

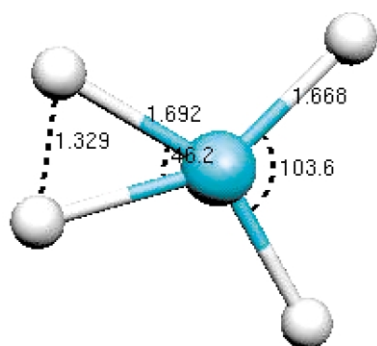


Fig. 1 The geometry of the transition state for  $\text{HgH}_4$  dissociation.

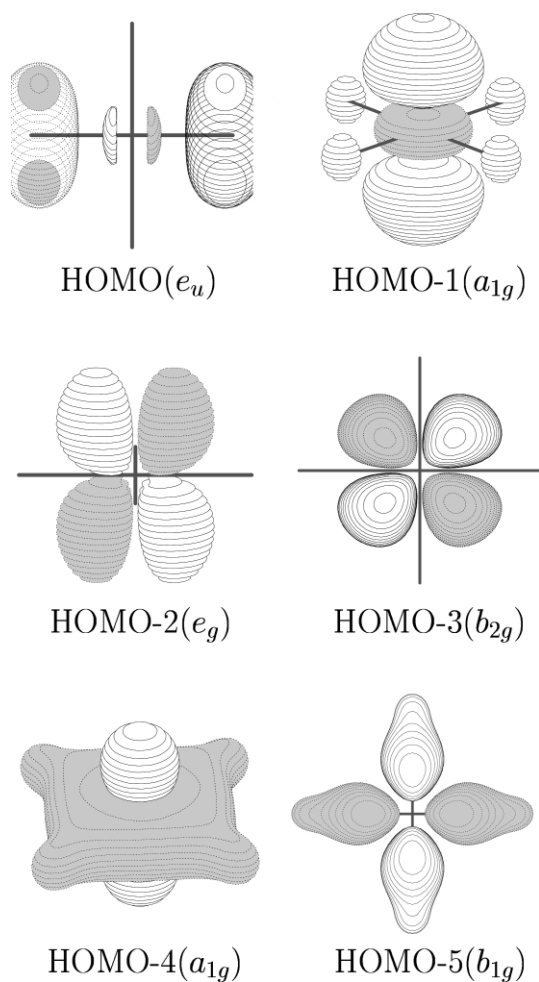


Fig. 2 The molecular orbitals of  $\text{HgH}_4$ .

Table 3 The natural orbital occupancies and charges

System	H (1s)	Hg (5d)	Hg (6s)	Hg (6p)	$Q(\text{H})$	$Q(\text{Hg})$
$\text{HgH}_6$	1.20	9.59	1.00	0.17	-0.21	1.24
$\text{HgH}_4$	1.28	9.72	1.03	0.10	-0.29	1.14
$\text{HgH}_2$	1.39	9.89	1.26	0.04	-0.40	0.81
$\text{HgH}^+$	0.94	9.93	1.12	0.01	0.05	0.95

weakly antibonding counterpart is the HOMO-1. Note that the two lowest, strongly bonding orbitals lie below the pure-5d core orbitals. Finally, the bonding  $e_u$   $6p + 1s$  HOMO and non-bonding, purely d orbitals HOMO-(2,3), complete the picture. The four bonding molecular orbitals  $2 \times e_u$ ,  $a_{1g}$  and  $b_{1g}$  essentially result in four single Hg-H bonds for  $\text{HgH}_4$ . The Natural Bond Orbital analysis is shown in Table 3 and gives an idea of the 5d hole and 6p occupation.

We conclude, that the  $\text{Hg}(\text{iv})$  compound  $\text{HgH}_4$  and even the  $\text{Hg}(\text{vi})$  compound  $\text{HgH}_6$  are possible synthetic objects for matrix spectroscopic experiments. Deuteration may be necessary to slow down decay by tunneling.

This study is supported by The Academy of Finland and the European Research Training Network 'Molecular Properties and Molecular Materials' (MOLPROP), contract No. HPRN-2000-00013. We thank Dr Eran Sterer for an inspiring discussion.

## Notes and references

- C. K. Jørgensen, *Z. Anorg. Allg. Chem.*, 1986, **540**–541, 91.
- M. Kaupp, M. Dolg, H. Stoll and H. G. von Schnering, *Inorg. Chem.*, 1994, **33**, 2122.
- M. Kaupp and H. G. von Schnering, *Angew. Chem., Int. Ed.*, 1993, **32**, 861.
- W. Liu, R. Franke and M. Dolg, *Chem. Phys. Lett.*, 1999, **302**, 231.
- K. P. Huber and G. Herzberg, *Constants of Diatomic Molecules*, Van Nostrand, New York, 1979.
- P. Pyykkö, *J. Chem. Soc., Faraday Trans. 2*, 1979, **75**, 1256.
- T. M. Greene, W. Brown, L. Andrews, A. J. Downs, G. V. Chertihin, N. Runeberg and P. Pyykkö, *J. Phys. Chem.*, 1995, **99**, 7925.
- N. Legay-Sommaire and F. Legay, *J. Phys. Chem.*, 1995, **99**, 16945.
- W. Bronger, *Angew. Chem., Int. Ed.*, 1991, **30**, 759.
- A. D. Becke, *J. Chem. Phys.*, 1993, **98**, 5648.
- C. Lee, W. Yang and R. G. Parr, *Phys. Rev. B*, 1988, **37**, 785.
- Gaussian 98, Revision A.7, M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, V. G. Zakrzewski, J. A. Montgomery, R. E. Stratmann, J. C. Burant, S. Dapprich, J. M. Millam, A. D. Daniels, K. N. Kudin, M. C. Strain, O. Farkas, J. Tomasi, V. Barone, M. Cossi, R. Cammi, B. Mennucci, C. Pomelli, C. Adamo, S. Clifford, J. Ochterski, G. A. Petersson, P. Y. Ayala, Q. Cui, K. Morokuma, D. K. Malick, A. D. Rabuck, K. Raghavachari, J. B. Foresman, J. Cioslowski, J. V. Ortiz, B. B. Stefanov, G. Liu, A. Liashenko, P. Piskorz, I. Komaromi, R. Gomperts, R. L. Martin, D. J. Fox, T. Keith, M. A. Al-Laham, C. Y. Peng, A. Nanayakkara, C. Gonzalez, M. Challacombe, P. M. W. Gill, B. G. Johnson, W. Chen, M. W. Wong, J. L. Andres, M. Head-Gordon, E. S. Replogle and J. A. Pople, Gaussian, Inc., Pittsburgh PA, 1998.
- D. Andrae, U. Häussermann, M. Dolg, H. Stoll and H. Preuss, *Theor. Chim. Acta*, 1990, **77**, 123.
- A. Schäfer, H. Horn and R. Ahlrichs, *J. Chem. Phys.*, 1992, **97**, 2571.
- N. Legay-Sommaire and F. Legay, *Chem. Phys. Lett.*, 1993, **207**, 123.