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A new endofullerene, Cu@C<sub>60</sub>, was synthesized and isolated. **Ring-current induced diamagnetic interior and paramagnetic** exterior of  $C_{60}$  have been quantitatively described using the Cu **probe.**

Endofullerenes are of great interest due to their doping-element dependent structural varieties. Applications, such as to field effect transistor (FET), MRI agent and quantum computer, have been proposed, based on different endofullerenes.1–3 As for material research of endofullerenes, much progress has been made over the past decade. However, the elements that can be entrapped inside fullerenes are still only those within the groups of the left (Group 1–4) and the right (Groups 15 and 18) sides of the periodic table.3–5 Whether or not other elements that occupy the majority and the middle portion of the periodic table can be encapsulated inside fullerenes and what kinds of electronic structures they may possess remains an unsolved mystery due partially to difficulties in their preparation. In this work, we report the first synthesis, isolation and characterization of an unconventional endofullerene,  $Cu@C<sub>60</sub>$ . The endohedral nature and electronic structure of Cu have been clarified. In addition, ring-current induced magnetism for both interior and exterior of  $C_{60}$  cage has been obtained using the Cu probe.

The experimental setup for producing  $Cu@C_{60}$  was based on capacitor-type plasma that was similar to that used to produce fullerene polymers and  $N@C_{60}.^{6-8}$  A sheet of copper foil was used to cover the cathode electrode, and subsequently copper ions were sputtered from the copper foil during plasma processing.  $C_{60}$  (100 mg, purity > 99.95%) was evaporated through resist heating, and collided with the copper ions in the plasma cathode-sheath region. High-purity nitrogen (99.99995%) was used as the buffer gas, and was kept at 10 Pa during plasma processing. The rf-plasma was operated at 13.56 MHz and 50 W. After evaporation of  $C_{60}$  in the rf-plasma, the deposit was collected and dissolved in  $CS_2$  *via* ultrasonication. The soluble portion was filtered through a  $0.2 \mu m$ PTFE syringe filter and used for analysis and isolation. ESR *g*value measurements were taken using precision magnetic-field and frequency counters.

The ESR spectrum of the soluble portion was closely dependent on the ESR operation parameters. For example, at a very low microwave power ( $\sim 10 \mu$ W) and very narrow modulation width (0.002 mT), only three sharp lines with a hyperfine constant of 0.567 mT could be observed, and they originated from  $N@C_{60}$ . When the microwave power was increased to 1.0 mW and the modulation width was set at 0.2 mT, besides the peaks from  $N@C_{60}$ , new ESR signals with four major peaks appeared, as is shown in Fig. 1. These ESR signals (without those from  $N@C_{60}$ ) were also observed when Ar was used as the buffer gas, but not observed when aluminum foil (instead of copper foil) was used to cover the rf-plasma cathode. These additional ESR signals apparently originate from copper.

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The origin of these ESR signals is further clarified by an analysis of the spectrum. First, the four major ESR peaks with equal integrated area correspond to the nuclear spin quantum number of  $Cu (I = 3/2)$ . Second, when the integrated signal of the peak around 334 mT is deconvoluted into two components, the ratio of the two components is 2.246, which is in excellent agreement with the ratio of the natural abundance of  ${}^{63}Cu$  and  ${}^{65}Cu$  (2.247). Since Cu<sup>1+</sup>  $(3d<sup>10</sup>)$  is ESR silent while Cu<sup>0</sup> is unlikely due to its very different ESR saturation behavior from  $N@C_{60}$  as well as HPLC elution behavior that will be described below, Cu can only be in its oxidation state of  $2+(3d^9)$ , after transferring two electrons to the  $C_{60}$  cage. Each ESR peak height and shape of  $Cu^{2+}$  is the combined result of the two Cu isotopes and the nuclear-spin quantum-number (m*I*) -dependent line broadening.

Because the ESR spectrum was obtained from a solution sample at room temperature, anisotropic *g*-tensors of Cu ions that are usually observed in the solid phase are averaged out. After deconvoluting each ESR peak, we obtain isotropic hyperfine coupling (hfc) constants of 7.757 mT and 8.277 mT for  $63Cu$  and  $65Cu$ , respectively. The ratio of the two hfc constants (0.937) is in an excellent agreement with the ratio of theoretical gyromagneticratio  $(y_n({}^{63}Cu)/y_n({}^{65}Cu) = 0.935)$  of the two Cu<sup>2+</sup> ions.<sup>9</sup>

Evidence for the endohedral nature of the copper ions is provided by both ESR and TOF-mass spectra. The ESR peaks from Cu<sup>2+</sup> ions are much more clearly isolated from each other compared to the ESR peaks of Cu ions doped in other inorganic matrices or ligands.10–11 We have also investigated various materials containing Cu2+, such as Cu-doped zeolite and Cu amine complexes. The spectra of none of these materials show well-isolated quadruple peaks attributable to Cu ions because of the strong spin–dipolar interaction among  $Cu^{2+}$  ions. The Cu-ion signals reported here might be the first x-band cw-ESR spectrum with completely resolved hyperfine lines originating from m<sub>*I*</sub>-dependent magnetic interaction. This improved resolution can only be explained by reduced spin–dipolar interaction among  $Cu^{2+}$  ions after being encaged inside  $C_{60}$  cages.

Evidence strongly indicating the endohedral nature of Cu2+ comes from the *g*-values of the Cu<sup>2+</sup> ion signals when compared to those of other Cu<sup>2+</sup> complexes and exohedral CuC<sub>60</sub>.<sup>12</sup> The *g* values of 63Cu and 65Cu in Fig. 1 are 2.0483 and 2.0477, respectively. The  $g$ -values for other Cu<sup>2+</sup> complexes in solution that possess no extra molecular magnetism, such as diethanoamine–Cu2+ (traced at 2.1210), are generally in range of 2.10–2.13.13 The *g*-value shift becomes more evident when compared to that of the exohedral  $CuC<sub>60</sub>$  complex. ESR spectrum of exohedral  $CuC<sub>60</sub>$  shows a single broad peak centered at a *g*-value of 2.2016. These *g*-value data



Fig. 1 ESR spectrum of vacuum-sealed  $Cu@C_{60}$  in CS<sub>2</sub>. The quartet splittings of 63Cu and 65Cu are indexed. Over the ESR spectrum two views of  $Cu@C_{60}$  in B88-PW91 DFT geometry are shown.

show that ESR for Cu@C<sub>60</sub> is shifted upfield by ~3.5% and ~7% when compared to diethanoamine–Cu<sup>2+</sup>and exohedral CuC<sub>60</sub>. respectively. The upfield shift can be explained by considering the ring current effect of the fullerene cage in a magnetic field. Compared to diethanoamine–Cu<sup>2+</sup> (without  $\pi$  conjugated electrons), Cu<sup>2+</sup> inside C<sub>60</sub> requires a higher magnetic field (+3.5%) for ESR due to the diamagnetic ring-current effect from  $C_{60}$ , while exohedral doping requires a lower magnetic field  $(-3.5%)$  due to the paramagnetic ring-current effect. These results imply that the local magnetic field induced by the ring-current effect is about 11.6 mT in a magnetic field of  $\sim$  330 mT, a value that cannot be ignored when dealing with the magnetic properties of fullerenes and endofullerenes.

It is worthwhile noting that the extent of diamagnetic shielding of the fullerene cage observed in ESR is much larger (5470 times) than that observed in NMR for  ${}^{3}He@C_{60}$ , in which the resonance is shifted upfield only by 6.4 parts per million.14 This difference is also in excellent agreement with the ratio (5509 times) between the Bohr magneton  $(\mu_B)$  and the nuclear magneton of <sup>3</sup>He.

Pure  $Cu@C_{60}$  was obtained through HPLC, as is shown in Fig. 2. Appearing around 5.6 min is a mixture of pristine  $C_{60}$  and N@C<sub>60</sub>. Cu@C<sub>60</sub> elutes around 18 min., much later than C<sub>60</sub>, indicating an enhanced interaction between  $Cu@C_{60}$  and the column fixed-phase material. The ratio of  $Cu@C_{60}$  to  $C_{60}$  is about 0.1%, estimated from ESR and HPLC results;  $\sim$  20 µg of purified  $Cu@C_{60}$ was obtained from 100 mg initial C<sub>60</sub>.

The TOF-mass spectrum of purified  $Cu@C_{60}$  is shown in Fig. 3. Besides the prominent peak  $(m/z)$  783) from Cu@C<sub>60</sub>, peaks attributable to  $CuC_{58}$  ( $m/z$  759) and  $CuC_{60}O$  ( $m/z$  799) can also be seen. The presence of the CuC<sub>58</sub> peak and the absence of a  $C_{60}$  peak clearly confirm the endohedral nature of this metallofullerene, otherwise Cu instead of  $C_2$  will be dissociated from its mother compound under laser irradiation.

In summary,  $Cu@C_{60}$ , a new endofullerene, has been generated in rf-plasma. The well-resolved ESR spectrum of  $Cu@C<sub>60</sub>$ 



**Fig. 2** HPLC trace of the soluble portion of the rf-plasma-treated  $C_{60}$ deposit. The sample was dissolved in toluene after replacing  $CS<sub>2</sub>$ . HPLC conditions: semipreparative Buckyclutcher column,  $3mL$  min<sup>-1</sup> 7:3 toluene/hexane mobile phase. The peak around 5.6 min is due to a mixture of  $C_{60}$  and N@C<sub>60</sub>, and the almost invisible peak around 18 min comes from  $Cu@C_{60}$ . Other small peaks following  $C_{60}$  are ESR silent and not yet clarified. ESR spectra of N@C<sub>60</sub> and Cu@C<sub>60</sub> are shown in the insets.



facilitated comparison with theoretical data on  $Cu^{2+}$  ions. Ringcurrent-induced diamagnetic interior and paramagnetic exterior of  $C_{60}$  have been quantitatively described using the Cu<sup>2+</sup> probe.

Purified  $Cu@C_{60}$  has been obtained, and further investigation of

## **Notes and references**

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