Ti_2C_{80} is more likely a titanium carbide endohedral metallofullerene $(Ti_2C_2)@C_{78}^\dagger$

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We show by means of density functional calculations that the previously synthesized metallofullerene Ti_2C_{80} does not take the form of $Ti_2@C_{80}$, but is a titanium carbide endohedral metallofullerene, $Ti_2C_2@C_{78}$, that has a $C_{78}^{-6-}(D_{3h})$ cage which follows faithfully the stable closed-shell electronic rule.

Endohedral metallofullerenes that encapsulate a metal atom or cluster inside hollow carbon cages have received extensive attention in the past two decades owing to their fascinating structural and electronic properties.¹ Electron transfer occurs from the encapsulated metal or clusters to the fullerene cages, resulting in negatively charged carbon cages adopting preferentially stable closed-shell electronic configurations.^{2,3} As such, the number of electrons transferred is a key factor that controls the stability of the otherwise unstable fullerene cages. This is exemplified by $Sc_2@C_{66}$, $4Sc_3N@C_{68}$ and $La_2@C_{72}$, 6 where the carbon cages even violate the well-known isolated pentagon rule (IPR)³ and are certainly unstable in their empty neutral forms. However, exceptions⁷ to the aforementioned stable closed-shell electronic rule seem available, e.g., the recently synthesized Ti₂C₈₀.⁸ In Ti_2C_{80} , the two Ti atoms were believed to be encapsulated in a C_{80} (D_{5h}) or C₈₀ (I_h) cage with a total of four electrons transferred to the carbon cages,⁸ although clear cut evidence is available that both C80 cages prefer to accept six electrons to attain a closed-shell electronic configuration.9 Here we show by means of density functional calculations that the Ti₂C₈₀ is not in the form of Ti₂@C₈₀ (Fig. 1), but that it is actually a titanium carbide endohedral metallofullerene, $Ti_2C_2@C_{78}$ (Fig. 2), that has a $C_{78}^{\ \ 6-}(D_{3h})$ cage^{9,10} that follows faithfully the stable closed-shell electronic rule.

All calculations were carried out with the hybrid density functional theory at the B3LYP level¹¹ using the Gaussian 98 program.¹² Two basis set–RECP (relativistic effective core potential) combinations were used. The first, denoted DZ, is the combination of the split-valence 3-21G basis set¹³ for C with the small core RECP, plus the valence double- ζ basis set (denoted LanL2DZ)¹⁴ for Ti. The second, denoted DZP, combines the split-valence d-polarized 6-31G* basis set¹⁵ for C with the LanL2DZ set for Ti. The geometries of all Ti₂C₈₀ isomers concerned were first optimized at the B3LYP/DZ level; the geometries of key structures

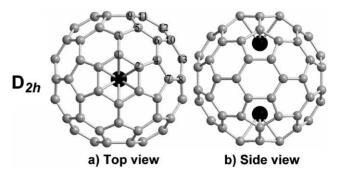


Fig. 1 (a) Top view and (b) side view of the B3LYP/DZP-optimized D_{2h} -symmetric geometry of singlet state $Ti_2@C_{80}$, in which the carbon cage is the $I_h(7:7)$ IPR isomer of C_{80} . The shortest C–Ti distances are 2.20 Å.

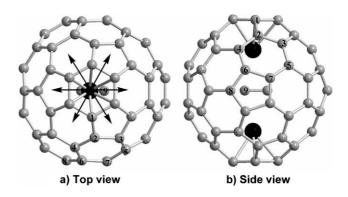


Fig. 2 (a) Top view and (b) side view of the B3LYP/DZP-optimized static C_{2v}-symmetric geometry of singlet state Ti₂C₂@C₇₈, in which the carbon cage is the D_{3h}(5 : 5) IPR isomer of C₇₈. The average Ti–C1 and Ti–C9 distances are 2.20 and 2.38 Å, respectively. The optimal C–C bond length for the central acetylide C₂^{2–} group is 1.27 Å.

were reoptimized at the B3LYP/DZP level. Similar combinations of theoretical method and basis sets have been shown to be reasonable at predicting the geometries of metallofullerenes such as $La_2@C_{72}$, $^6La_2@C_{80}$ ⁹ and $Sc_3N@C_{80}$. 16 NMR chemical shielding tensors were evaluated by employing the gauge-independent atomic orbital (GIAO) method. 17 Based on the computed chemical shielding tensors, theoretical 13 C NMR chemical shifts were calculated relative to C_{60} and converted to the TMS (tetramethyl-silane) scale using the experimental value for C_{60} (142.5 ppm). 18 The GIAO-B3LYP/DZP theory proved to be sufficiently accurate at reproducing 13 C NMR chemical shifts of fullerenes such as C_{60} , C_{70} , C_{76} and C_{78} , 19 and metallofullerenes such as $Sc_3N@C_{80}$. 16 and $Sc_2@C_{66}$. 20

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E-mail: xinlu@xmu.edu.cn; Fax: 86 592 2183047; Tel: 86 592 2181600 † Electronic Supplementary Information (ESI) available: Cartesian coordinates and energies of the optimized $Ti_2@C_{80}$ and $Ti_2C_2@C_{78}$, and electron isodensity surfaces of selected molecular orbitals of $Ti_2@C_{80}$ and $Ti_2C_2@C_{78}$. See http://dx.doi.org/10.1039/b507855e

In the experimental ¹³C NMR characterization of pure Ti_2C_{80} ,⁸ Cao et al. found a total of eight ¹³C NMR chemical shifts ranging from 130 to 145 ppm, and proposed that their sample was a mixture of two $Ti_2@C_{80}$ isomers of I_h and D_{5h} symmetric C₈₀ cages in a ratio of 1 : 3. Hence, we chose the I_h (7 : 7) IPR isomer of C₈₀ to investigate the structures of Ti₂@C₈₀. The encapsulation of two equivalent metal atoms in the C80 (Ih) cage leads to several highly symmetric structures, including those of D_{5d}, D_{2h} and D_{3d} symmetries.9 Geometry optimizations of these high symmetry isomers in a singlet state were performed at the B3LYP/DZ level. The singlet D_{2h} structure (Fig. 1) appears to be the most stable by 20.1 kcal mol⁻¹, compared to the singlet D_{5d} structure. The shortest C-Ti bond length in the singlet D_{2h} structure (2.24 Å) is much shorter than that (2.35 Å) in titanocenes,²¹ implying stronger Ti-cage covalent bonding in Ti2@C80. Even shorter C-Ti bond lengths (~ 2.20 Å) were predicted by the more sophisticated B3LYP/DZP optimizations. A detailed analysis of its molecular orbitals (Fig. S1, ESI^{\dagger}) indicates that the Ti₂@C₈₀ can be described as $Ti_2^{6+} @C_{80}^{6-}$, but that the presence of strong covalent dative bonding between Ti^{3+} and C_{80}^{6-} results in back-donation of charge from the negatively charged cage to the cations. The relatively strong Ti-cage covalent bonding should prohibit free rotation of the Ti₂ unit in the carbon cage at room temperature, and consequently Ti2@C80 with a C80 (Ih) cage would show thirteen ¹³C NMR signals, not the expected two lines.⁸ In addition, single-point B3LYP/DZP calculations revealed that the triplet state of D_{2h} symmetric $Ti_2 @C_{80}$ is 14.0 kcal mol⁻¹ lower in energy than its singlet state, with one spin-unpaired electron localized on each Ti atom. Hence Ti₂@C₈₀, preferentially adopting an open-shell electronic configuration, should be paramagnetic in nature and consequently should not have detectable signals in its NMR spectrum at all. Instead, the other structural model, which is diamagnetic, should be proposed to account for the experimental ¹³C NMR spectrum of Ti₂C₈₀.

We then inferred that the synthesized Ti₂C₈₀ should adopt a different structure model, *i.e.* Ti₂C₂@C₇₈, with a C₇₈ cage of D_{3h} symmetry (Fig. 2). This structural model is based on the following considerations. Firstly, a D_{3h}-symmetric C₇₈ cage would show eight ¹³C NMR signals.¹⁰ Secondly, the C₇₈ cage can accept up to six electrons,¹⁰ whereas the encapsulated Ti₂C₂ cluster can be hexavalent consisting of two Ti⁴⁺ cations and a C₂²⁻ group; the whole structure would thus have a diamagnetic closed-shell electronic configuration.

Fig. 2 depicts the static C_{2v} -symmetric geometry of $Ti_2C_2@C_{78}$, optimized at the B3LYP/DZP level. The averaged Ti–C1 and Ti–C9 distances are 2.20 and 2.38 Å respectively, indicating that the Ti–cage bonding is much stronger than the Ti–acetylide bonding. The $Ti_2C_2@C_{78}$ can be viewed as $Ti_2^{8+}C_2^{-2}@C_{78}^{6-}$ with covalent dative bonding between the Ti^{4+} cations and the C_{78}^{6-} cage, as well as ionic Ti^{4+} –acetylide interactions (Fig. S2, ESI†). The weaker ionic Ti–acetylide interaction allows rapid (free) rotation of the acetylide group around the C₃ axis of the C₇₈ cage with a small rotation barrier of about 0.1 kcal mol⁻¹, estimated at the B3LYP/DZ level. Thus at room temperature, the C₇₈ cage in $Ti_2C_2@C_{78}$ should maintain a D_{3h} symmetry over a long timescale due to the constant rotation of the encapsulated acetylide group.

Table 1 lists the ¹³C NMR chemical shifts of $Ti_2C_2@C_{78}$, predicted by GIAO-B3LYP/DZP calculations, along with the experimental data⁸ for Ti_2C_{80} . The good agreement between the

Table 1 Theoretical and experimental $^{13}\mathrm{C}$ NMR chemical shifts (in ppm) of Ti_2C_2@C_{78}

Theory/ppm ^{<i>a,b</i>}	Experiment/ppm ^c	Site ^d
128.6	129.8	C1
130.5	130.5	C2
133.5	131.6	C5
137.4	136.4	C4
139.8	138.9	C3
141.0	140.4	C6
142.5	141.6	C7
143.6	144.3	C8
^a The chemical shif	ts of equivalent isomeric si	tes were averaged to

^a The chemical shifts of equivalent isomeric sites were averaged to obtain the value. ^b The predicted ¹³C NMR chemical shift for the acetylide C atoms is 288.5 ppm. ^c Data extracted from ref. 8. ^d See Fig. 2 for the definition of the carbon sites.

theoretical ¹³C NMR chemical shifts and the experimental data convinces us that the metallofullerene Ti_2C_{80} , synthesized by Cao *et al.*, is not $Ti_2@C_{80}$, but is most likely in the form $Ti_2C_2@C_{78}$. For the acetylide carbon atoms in $Ti_2C_2@C_{78}$, the calculated ¹³C NMR chemical shift is 288.5 ppm, much higher than that (92 ppm) observed for $Sc_2C_2@C_{84}$.²² Unfortunately, this signal was not observed experimentally,⁸ due probably to the spin–rotation interaction^{9,23} at room temperature.

In summary, we have shown theoretically that the previously synthesized metallofullerene Ti₂C₈₀ does not take the form of Ti₂@C₈₀, but is most likely a titanium carbide endohedral metallofullerene, Ti₂C₂@C₇₈. Hence, Ti₂C₈₀ is the smallest metal carbide endohedral fullerene disclosed so far; other metal carbide endohedral fullerenes including Sc₂C₂@C₈₄²² and Y₂C₂@C₈₂.²³ The present work shows that care should be taken during the structural determination of metallofullerenes as Sc₂C₈₂,²⁴ Sc₂C₈₄,^{24,25} Ti₂C₈₄²⁶ and Hf₂C₈₀,²⁷ in that their simple endohedral forms, M₂@C_n, may not fulfil the stable closed-shell electronic rule. To attain stable closed-shell electronic configurations, it is likely that some of these dimetallofullerenes will form M₂C₂@C_{n-2} cages, as implied by high-resolution ion mobility measurements on Sc metallofullerenes.²⁴

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