Structure and Dynamics of Hydrated Ions

Hitoshi Ohtaki^{*}

Coordination Chemistry Laboratories, Institute for Molecular Science, Myodaljl-cho, OkazakI, **444** Japan

Tamás Radnai

Central Research Institute for Chemistry of the Hungarian Academy of Sciences, Budapest, P.O. Box 17, H-1525 Hungary

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Contents

/. Introduction

Ionic hydration is one of the most attractive subjects to chemists, especially to those who are interested in reactions in aqueous solutions, and numerous studies on ionic hydration, i.e., the determination of hydration numbers of ions, rates of exchange of coordinated water molecules around ions, and interaction energies between ions and water molecules, have been carried out since Arrhenius.

In the time from the 1930s to the 1960s, various classical methods have been applied to investigate these problems. However, confusion in understandings of the structure of hydrated ions has been noticed in this period because of scattered data obtained by different methods and the lack of information of static and dynamic properties of water molecules coordinated to ions. When new methods were applied to the field of solution chemistry beginning in the 1970s, which became possible owing to the development of modern electronics and high-speed electronic computers, a steep increase in publication of interesting papers was seen in studies on the structures of hydrated ions, dynamics of coordinated water molecules, and interaction energetics of ions with water molecules.

A fairly large number of monographs¹⁻⁵ have so far been published concerning ionic hydration phenomena. Frequent reviews on hydration of ions have been published in various scientific journals. An excellent review on radii of ions in the hydrated species in solution by Marcus⁶ has appeared in this journal in 1988. Although some overlaps with Marcus' paper are not avoidable in the description of this review because the sizes of hydrated ions is the most fundamental element in studies of ionic hydration, we intend to pay more attention to other structural, as well as dynamic properties of hydrated ions. Most works quoted in this review were published in the last 10 years, but some important articles and monographs which were published in the last 2 decades are also referred to.

Since we focus our interest on the structure and dynamics of hydrated ions, results of thermodynamic measurements and quantum mechanical calculations for hydration of ions are generally not included in the discussion.

Hitoshi Ohtakl was born in Tokyo, Japan, in 1932. He graduated from Nagoya University in 1955 and obtained M.Sc. and Dr.Sc. degrees in 1957 and 1961, respectively, from Nagoya University. He became a Research Associate of the Tokyo Institute of Technology in 1959. He studied complex equilibria in solution under the supervision of Prof. L. G. Sillen, Royal Institute of Technology, Stockholm, Sweden, as a postdoctoral research fellow from 1961 to 1964. After returning to Japan, he was appointed Lecturer of Nagoya University in 1965 and promoted to Associate Professor in 1967. In 1970 he moved back to Tokyo Institute of Technology as an Associate Professor and then became a Full Professor in 1973. On the appointment in 1988 as Professor of the Institute for Molecular Science of the Okazaki National Research Institutes, he has served as the Director of the Coordination Chemistry Laboratories. He is also the Dean of the School of Mathematical and Physical Science of the Graduate University for Advanced Studies, of which the Institute for Molecular Science is one of the parent institutes. His research interests cover various areas of solution chemistry, especially structural chemistry of solution, including studies on structures of solvents, solvated ions and complexes existing in solution by means of X-ray diffraction; EXAFS and neutron diffraction methods have been occasionally employed as well. Thermodynamic studies on solution equilibria have long been another major area of his study. Molecular dynamics simulation studies on dissolution and nucleation processes of crystals have thrown a new sight for chemistry of ionic solvation and crystal growth from the dynamical viewpoint. He has published more than 170 research papers and a book on isotope chemistry (with H. Kakihana), on solution equilibria and kinetics (with M. Tanaka and S. Funahashi, which has been translated into Chinese) and one S. Funahashi, which has been translated into Chinese) and one book on ionic hydration and two books on solution chemistry, all were written in Japanese. He was the editor of two English books and three Japanese books. He has translated two books (by V. Gutmann and by K. Burger) into Japanese. He was an author of chapters of more than 16 books. He has served as a visiting **has been extended on IUPAC bodies in various capacities since 1974. He is one of the three founders of Eurasia Conference on Chemical Sciences which was founded in 1988. He has received three awards: (Matsunaga Prize in 1976, Tejima Memorial Award in 1989 and Take) Prize in 1990).**

//. Methods for Determination of the Structure of Hydrated Ions

The concept of the structure of hydrated ions largely depends on the methods of observation. We may classify the structure of hydrated ions into three categories depending on the methods of investigations as follows:

(1) Static structure, in which the structure is discussed on the time and space averaged ion-water interactions. Results obtained by X-ray and neutron diffraction methods are included in this category.

(2) Structure discussed on the basis of dynamic properties of coordinated water molecules. Results obtained by NMR measurements may be the typical case. When an ion is surrounded by coordinated water molecules which move much faster than the time of

Tamas Radnai was born in Budapest, Hungary, in 1949. He finished his university course in the Eötvös Lorand University, Budapest in **1973 with receiving an equivalent degree of M.Sc. He received Doctor University Degree (corresponding to Ph.D.) from the Eötvös** Lorand University in 1977, and the academic title "Candidate for **Chemical Science", a post-Ph.D. degree in 1991. He has held a continuous appointment at the Central Research Institute for Chemistry of the Hungarian Academy of Sciences in Budapest as a researcher. In the meantime, he was employed by the Ministry of Education, Science and Culture of Japan as a Research Associate, a Government Officer of Japan, from 1987 to 1989. He was invited to the Institute for Molecular Science as a Guest Foreign Associate Professor for one year from 1991. He has worked as a member of several international research groups in Italy and Germany. He is a senior fellow of the department of solution chemistry in the Central Research Institute for Chemistry of Hungarian Academy of Sciences. His main interest includes structural studies of electrolyte solutions by using diffraction methods, as well as computer simulations. He has published more than 40 research** computer simulations. He has published more than 40 research papers.

frequencies of the magnetic field employed in the NMR measurement, we observe that the ion has practically no hydrated water molecule in the coordination shell.

(3) Energetic considerations lead to the discrimination of strongly combined water molecules with ions from loosely interacting water molecules to ions. Spectroscopic investigations, including frequency measurements by Raman and IR spectroscopies and thermodynamic studies on ionic hydration give us information on water molecules strongly coordinated to ions. Information on water molecules weakly interacting to ions is always ambiguous. The borderline between coordinated and noncoordinated water molecules may be drawn on the basis of the kinetic energy *kT* or *2kT* be drawn on the basis of the Kinetic energy κT or $2\kappa T$
according to Bierrum.⁷ However, it is obvious that the structure of hydrated ions discussed in terms of static and dynamic properties of water molecules coordinated also depends on interaction energies between ions and water molecules. Therefore, the classifications of the concept of ionic hydration is conventional.

A. Scattering Methods

The scattering methods may be separated into two. The one is the method to measure intensities of elastically scattered electrons, photons (X-ray), and neutrons by atoms and the other is the inelastic neutron scattering method. Since electrons and photons are so light, the quasi-elastic or inelastic scattering of them by atoms are not usually considered in the diffraction measurements as sources of structural information.

The former method provides the coordination number or frequency factor (n_{p-q}) of p-q atom pairs, the

distance between p and q atoms (r_{p-q}) , and the temperature factor (b_{p-q}) of the atom pairs, which relates to the root mean square displacement (rmsd, *I)* of the atom pairs as $b = \langle l^2 \rangle / 2$, *l* being also obtainable from frequency measurements. The distance between atoms is a unique quantity obtained by the diffraction methods and it is hard to directly measure by other methods.

The electron scattering method is included in the elastic scattering method in the present classification. However, since electrons are easily absorbed by air, electron scattering measurements should be performed in vacuum, and liquid samples should be cooled to reduce the vapor pressure or kept in a cell with extremely thin windows. In most cases the former method is employed.

The hydrogen atom is not well detected by the electron and X-ray diffraction methods due to its small scattering power. Since the neutron diffraction method can provide the interatomic distances between an atom and hydrogen atoms (hydrogen atoms are replaced with deuterium atoms in usual cases), the determination of the average orientation of coordinated water molecules around the central ion is possible. The isotopic substitution method, applicable to cases in which suitable isotopes are available for the central ions in the measurements, provides detailed information on the structure of hydrated ions. The isomorphous substitution method in the X-ray diffraction measurements, in which ions with essentially the same size and charge but different atomic numbers are used, may be compared with the isotopic substitution method in the neutron diffraction measurements.

X-ray beams are scattered by electrons around atoms, while neutrons are diffracted by nuclei of atoms, and electron beams are scattered in the Coulombic field formed by both shell electrons and nuclei, and therefore, these methods are complementary.

It should be noted here that the distribution of electrons in an atom is not always spherical due to interatomic interactions. In the X-ray diffraction method a spherical electron cloud is assumed around a nucleus and the location of all electrons of an atom is considered to be at the center of the atom. The same assumptions are made in the electron diffraction measurement. However, the center of the nucleus of the atom due to interatomic interactions, and thus, interatomic distances of a given atom pair determined by these methods under similar experimental conditions are sometimes different.

The quasi-elastic and inelastic neutron scattering methods provide dynamic information predominantly of the translational motion of protons.

Brief descriptions for each method of measurements will be presented in the following sections.

1. X-ray Diffraction (XD) Method

Scattered intensities *I* are measured as a function of the momentum transfer, which is often denoted by *k, s, q*, or Q. In the X-ray diffraction method s is frequently used, where $s = (4\pi/\lambda)\sin \theta$ (2 θ is the scattering angle and λ wavelength), and is also called a scattering variable or reciprocal space variable. However, in order to unify symbols with the same physical meaning, the letter *k* will be used for describing the quantity of the momentum transfer in all the scattering methods.

The intensities *I{k)* are measured as a function of *8* at a constant λ (the method is called as angular dispersive method) or as a function of λ or energy at a given θ (energy dispersive method). The former method is often used when the apparatus is equipped with a usual X-ray tube, from which a relatively strong characteristic X-ray is emitted. On the other hand, the latter is favorable when a strong white X-ray source such as synchrotron orbital radiation (SOR) is available. When sample solutions contain heavy atoms which emit fluorescence X-rays in the course of radiation with white X-rays, elimination of the fluorescence X-rays is sometimes difficult. X-rays emitted during the disturbing incoherent scattering process are more easily corrected in the former than in the latter. Therefore, for many electrolyte solutions which often contain atoms with various atomic numbers, the angular dispersive method is extensively employed. The energy dispersive method is used for the structural analysis of organic liquids in which atoms with relatively low atomic numbers are involved. Application of a multichannel detector tremendously shortens the measuring time in the energy dispersive method compared with the angular dispersive method.

Another way of classification is to separate the methods into the transmission method and the reflection method. The former allows the incident X-rays to pass through a sample cell, and the detector is placed behind the cell. On the other hand, reflected X-rays scattered at the surface of a liquid sample are measured in the latter technique. A free surface of liquids can be used as the reflection area in the latter method, so that corrections for a container of the sample are not necessary. On the other hand, the corrections are unavoidable in the former. The advantage of the transmission technique is that experimental conditions such as temperature and pressure can be changed over a much wider range than the reflection method can. Absorption, multiple scattering, and polarization of scattered X-rays should be corrected in both methods.

The scattering factor $f_M(k)$ of atom M is calculated from quantum mechanical theories.⁸ $f_M(0)$ is identical to the number of electrons in a given atom M. Sometimes an atomic group containing hydrogen atoms is regarded as one scattering unit with a group scattering factor $f_G(k)$. This treatment is convenient because it reduces the number of "atoms" in the system to be treated, but deviation from the assumed spherical distribution of electrons of the group should lead to errors in this treatment. For the X-ray analysis of pure water, however, the introduction of the group scattering factor for a water molecule does not result in significant errors in the structural information of intermolecular interactions of water.

The structure function *i(k)* multiplied by *k* as a weight can be calculated from the observed intensities after subtraction of coherent and incoherent self-terms of scattering of atoms as follows:

$$
ki(k) = k[Iabs(k) - \sum x_d f_a^{2}(k) - \sum x_d I_a^{inc}(k)] \quad (1)
$$

where $I_{\text{abs}}(k)$ is the corrected intensity converted to the absolute unit, x_{α} the atomic fraction of atom or atom group α . $I_{\alpha}^{\text{inc}}(k)$ represents the intensity caused by the incoherent scattering by α . The summation extends over each α type independent scattering units in one stoichiometric volume. In practice the *ki(k)* function

is often multiplied by an arbitrary chosen modification function or sharpening function *M(k)* in order to minimize the truncation effect of the measured intensities at the Fourier transform of the intensity data.

The radial distribution function *Dir)* is obtained by the Fourier transform of the structure function:

$$
D(r) = 4\pi r^2 \rho_0 + 2r\pi^{-1} \int_0^{k_{\text{max}}} ki(k)M(k) \sin(kr) \, dk \qquad (2)
$$

The function $D(r) - 4\pi r^2 \rho_0$ is conventionally used to emphasize peaks in the $D(r)$ function in the large r region.

Instead of the $D(r)$ function, the $G(r)$ and $H(r)$ functions are also often used:

$$
G(r) = D(r)/(4\pi r^2 \rho_0)
$$
 (3)

$$
H(r) = G(r) - 1 \tag{4}
$$

Analysis of the structure function $ki(k)$ is usually carried out on the assumption of various models of molecular structures and molecular arrangements in a given system. In the space area constructed with shortrange interatomic interactions, a discrete electron distribution can be reasonably assumed, while in the area of long-range interatomic interactions, ordering of atoms and molecules disappears and the region can be assumed to be constructed with a practically homogeneous electron distribution. Thus, the structure function may be separated into two parts, discrete $ki_d(k)$ and continuous $ki_c(k)$ parts:

$$
ki(k) = kid(k) + kic(k)
$$
 (5)

The discrete part can be calculated by the formula derived by Debye:⁹

$$
ki_{\rm d}(k) = \sum x_{\alpha} \sum n_{\alpha\beta} f_{\alpha}(k) f_{\beta}(k) \Delta_0(kr_{\alpha\beta},l_{\alpha\beta}) \qquad (6)
$$

and the continuous part can be formulated as

$$
ki_{c}(k) = -4\pi\rho_{0} \sum R_{\alpha\beta}^{2} x_{\alpha} x_{\beta} f_{\alpha}(k) f_{\beta}(k) \Delta_{1}(kR_{\alpha\beta},L_{\alpha\beta})
$$
 (7)

where x_{α} and x_{β} represent the atomic fractions of α and β -type particles, respectively, and Δ_m is associated with the $j_m(x)$ spherical Bessel function of the *mth* order

$$
\Delta_m(x, y) = j_m(x) \exp\{- (y^2/2)k^2\}
$$
 (8)

The structure function includes three kinds of shortrange structural parameters, the frequency factor or coordination number $n_{\alpha\beta}$, the average discrete interatomic distance $r_{\alpha\beta}$, and its root mean square deviation (rmsd) $l_{\alpha\beta}$ for the $\alpha-\beta$ atom pair, and moreover, two kinds of long-range parameters, $R_{\alpha\beta}$ for the boundary of the continuum distribution of electrons and $L_{\alpha\beta}$ for the sharpness at the boundary between discrete and continuum regions. Some authors prefer to use the *2* temperature factors $b_{\alpha\beta} = l_{\alpha\beta}^2/2$ instead of $l_{\alpha\beta}$.

In the course of discussion of the hydration structure of ions, the values of $n_{\alpha\beta}$ and $r_{\alpha\beta}$ are most interesting. The $l_{\alpha\beta}$ values are important to consider the reliability of the estimated $n_{\alpha\beta}$ values, because the $n_{\alpha\beta}$ and $l_{\alpha\beta}$ values are strongly correlated. When we have information from frequency measurements by using Raman and IR spectroscopies, the $l_{\alpha\beta}$ values may be compared with the vibrational amplitude of a given $\alpha-\beta$ atom pair in the system. It should be noted here that the introduction of the continuum electron distribution for long-range interactions usually also improves the accuracy of parameters in the short-range interactions, especially the $n_{\alpha\beta}$ values.

Interpretations of the radial distribution function and the structure function are usually started from a model of the first neighbors and the hydration number $n_{\alpha\beta}$ is estimated by eq 6 on the basis on the model proposed under the assumption that the shape of the peak in the radial distribution curve from which the $n_{\alpha\beta}$ value is evaluated is assumed to be Gaussian. This assumption is also applied to the data analysis in the neutron diffraction method.

The structural information for the second hydration shell is usually not clear when we directly analyze the radial distribution function and the structure function, because many interatomic interactions may overlap each other and it may be very difficult to extract the only peak due to the second hydration shell of ions.

In order to obtain more reliable structural information on the second hydration shell, the isomorphous substitution method is applied to some suitable systems, in which central metal ions with practically the same ionic radii but reasonably different atomic numbers are available. Details of the isomorphous substitution method will be given in the next section, because the method is compared with the isotopic substitution method in the neutron diffraction method.

2. Neutron Diffraction (ND) Method

The basic principles of the elastic neutron diffraction method are similar to those of the X-ray diffraction method except for some respects. One of the remarkable differences between the two methods is the difference in the scattering coefficients of atoms for X-rays and neutrons. In the former the scattering power of atoms is a function of *k* and depends on the number of electrons in the atom, i.e., the atomic number, and in the latter the scattering ability of atoms depends on the composition of the atomic nuclei and thus varied with isotopes of a given element.

Essentially two different techniques are employed for neutron scattering experiments. In one technique neutron beams emitted from a reactor are monochromatized by using the Bragg reflection. In the other pulsed neutrons are used and the time of flight after the production of pulsed neutrons is measured to determine energies of incident neutrons. The latter is called the "time-of-flight (TOF)" method.

An important advantage in the neutron diffraction method over the X-ray diffraction technique is the application of the isotopic substitution method, in which samples with the same atomic composition but different isotopes for a particular element are prepared and the difference in scattering intensities or $G(r)$ functions is used for the structural analysis of the sample solutions, in which no other interatomic pair correlations are included except for the functions related to the isotopic elements. For instance, in an MX aqueous solution the following 10 atomic pairs should be included in the *Gir)* function:

$$
G(r) = aG_{M-O}(r) + bG_{M-H}(r) + cG_{X-O}(r) + dG_{X-H}(r) +
$$

\ne
$$
G_{M-X}(r) + fG_{M-M}(r) + gG_{X-X}(r) + hG_{O-H}(r) +
$$

\n
$$
iG_{O-O}(r) + jG_{H-H}(r)
$$
 (9)

Coefficients a-j are either a function of scattering factors for X-rays or scattering lengths for neutrons. When the isotopic substitution method is applied, the difference $G(r)$ function, $\Delta G(r)$, of two identical solutions except for the isotopic composition of, say, M includes the only terms containing atom pairs with M:

$$
\Delta G(r) = (\mathbf{a} - \mathbf{a}')G_{\mathbf{M} - \mathbf{0}}(r) + (\mathbf{b} - \mathbf{b}')G_{\mathbf{M} - \mathbf{H}}(r) +
$$

$$
(\mathbf{e} - \mathbf{e}')G_{\mathbf{M} - \mathbf{X}}(r) + (\mathbf{f} - \mathbf{f}')G_{\mathbf{M} - \mathbf{M}}(r) \tag{10}
$$

In solutions where no polynuclear complex is formed, the last term drops in the structural analysis of ionic hydration in solution, and thus the $\Delta G(r)$ function contains only three terms. When one discuss the structure in the first coordination sphere in a solution in which M-X ion pairs do not practically exist, the $\Delta G(r)$ function becomes very simple.

Readers should notice that suitable isotopes available for measurements of difference radial distribution functions are rather limited, because the isotopes should have reasonably large differences in the scattering lengths among them. Combinations of ${}^{1}H^{-2}H(D)$, 6 , $\rm Li$ – $^7\rm Li$, $^{35}\rm Cl$ – $^{37}\rm Cl$, $^{40}\rm Ca$ – $^{44}\rm Ca$, $^{54}\rm Fe$ – $^{56}\rm Fe$ – $^{57}\rm Fe$, $^{58}\rm Ni$ – $^{60}\rm Ni$ – ⁶²Ni-⁶⁴Ni isotopes are often employed, but for other elements the isotopic substitution method may be hardly applicable. For the ¹H and D atoms corrections for inelastic collisions between the atoms and neutrons must be done.

A similar method can be applied to the X-ray diffraction method by using ions with isomorphous structures for suitable cases. In this case the atoms have different atomic numbers, but the structure around the relevant atoms is expected or known to be the same, and then, the difference in the radial distribution functions between the two sample solutions is calculated to eliminate the contribution of other atom pairs except pairs related to the isomorphic elements. However, the application of the isomorphous substitution method in the X-ray diffraction method is much more limited than the isotopic substitution method as anticipated, because the atoms or ions should have practically the same radius or the difference in the radii should be less than 2 pm . The $Y^{3+}-Er^{3+}$ pair has been used in the hydration structural analysis of the ions. 10.11 Similarly, the $CrO₄² - SeO₄²$ pair was recently applied in a combined study of X-ray with neutron methods.²⁶⁷ Some other pairs such as $Pt(II)-Pd(II)$, $Pt(IV)-Pd(IV)$, MoO_4^{2-} WO₄²-, etc. may be used, but less studies have been done so far.

Another advantage in neutron scattering measurements compared with X-ray scattering measurements is the ability to determine the M-H (or M-D) and X-H (or X-D) bond lengths. Since the M-O and X-O bond lengths are determinable by both X-ray and neutron diffraction measurements, the tilt angle of hydrated water molecules in the first coordination sphere can be measured by the neutron diffraction method.

Uncertainty of experimental intensity data in neutron diffraction measurements is usually more pronounced than that in the X-ray diffraction method.

Some reviews have previously been published for structural data obtained by the neutron diffraction method.^{5,6,12-14}

3. Electron Diffraction (ED) Method

Electrons are scattered both by electrons and nuclei of atoms due to Coulombic interactions. The accessible range in the momentum transfer in the course of electron scattering is usually much wider than that achieved in the course of X-ray scattering measurements. For instance, 68 keV electron beams corresponding to the wavelength of 5 pm cover the maximum k range of 0.25-0.30 pm^{-1} , in contrast to the upper limit of $k_{\text{max}} = 0.10$ (when Cu K α line is used) to 0.17 pm⁻¹ (Mo K_{α}) in the X-ray diffraction method and may be compared with that in the neutron diffraction method. A disadvantage in the electron diffraction method is that measurements must be carried out in vacuum, and this requirement is rather fatal for studies of liquids and solutions, although this serious problem may be partly avoided by using solutions with extremely low vapor pressure and by application of special techniques with short measuring times.

Since the scattering lengths of electrons, X-rays, and neutrons are different for a given atom, one can separate the total pair correlation function into contributions of each atom pair by combining electron, X-ray, and neutron diffraction methods for one sample.¹⁵

4. Small-Angle X-ray (SAXS) and Neutron Scattering (SANS) Methods

Small-angle scattering X-ray and neutron diffraction methods are essentially similar, in principle, to the scattering methods in the large-angle range described in previous sections, but they can provide information of the long-range structure of large molecules and clusters. Since the covered k range is very short (k_{max}) \approx several thousands per picometer), special measurement techniques are required. The SAXS-SANS techniques usually cover the region of distances well above 5-10 nm in the r space, and therefore, they should be distinguished from the "low-angle scattering", which has a somewhat confusing name and relates to distances comparable to those observed by the large-angle diffraction methods. In electrolyte solutions containing relatively simple ions and complexes, scattering data in a low-angle region (a small *k* range) sometimes show a small peak in the structure function. Some authors suggested that the peak position corresponds to the interionic distance arising from a quasi-lattice arrangement of ions in solution (super-arrangement).^{16,17} The distance between ions *Rc* is given approximately as

$$
R_{\rm c} = 7.73/k_{\rm m} \tag{11}
$$

where k_m denotes the k value of the position of the peak maximum. The distance thus calculated usually agrees with the size of the stoichiometric volume containing one cation in solution in which $[MX_n]^{(m-n)+}$ type complexes are formed. However, in aqueous solutions containing only hydrated ions, observation of such a peak in the *ki(k)* function is not clear, and moreover, alternative explanations can also be given.

5. Quasi-Elastic Neutron Scattering (QENS) Method

Intensity measurements of incoherent scattering of neutrons with protons, which provide a good probe due to its largest incoherent scattering contribution to the total scattering pattern of an aqueous solution, can elucidate dynamic properties of protons in aqueous solutions.

Intensities of scattered neutrons over a wide range of energy or frequency of a neutron beam ω at various k 's are measured. Under a special condition where the time scale of interactions between neutrons and ions is long enough, i.e., both the frequency of neutrons ω and *k* are small, the theory assumes that the translational and rotational diffusion can be decoupled. In this case the incoherent space-time correlation function for the hydrogen atoms, $G^sH(r,t)$, can be expressed in terms of the translational diffusion coefficient D for the hydrogen atoms as

$$
GsH(r,t) = (4Dt)-3/2 exp(-r2/4Dt)
$$
 (12)

and the corresponding scattering function $H^s_H(k,\omega)$ is described as

$$
H_{\mathrm{H}}^{\mathrm{s}}(k,\omega) = \frac{1}{\pi} \frac{Dk^2}{(Dk^2)^2 + \omega^2}
$$
 (13)

The *D* value obtained from the above equations for hydrogen atoms in water corresponds to the average diffusion coefficient of hydrogen atoms provided that the exchange rate of protons around an ion is fast; in other words, the bonding time τ_{RES} (called also as mean lifetime or residence time) of the protons in the hydration shell of an ion is much smaller than the characteristic observation time $\tau_{\rm meas}$ of neutron scattering experiment (approximately 5×10^{-9} s)

$$
\tau_{\rm RES} \ll \tau_{\rm meas} \tag{14}
$$

Equation 12 holds for many aqueous ionic solutions, and the diffusion coefficient of ions can thus be determined, the value being compared with those determined by usual diffusion methods.

In the slow exchange limit, however, the residence time of protons around an ion is much longer than the observation time $\tau_{\rm RES}$

$$
\tau_{\rm RES} \gg \tau_{\rm meas} \tag{15}
$$

and the scattering function can be decomposed into two Lorentzian type functions

$$
H^{\rm s}_{\rm H}(k,\omega) = \frac{1}{\pi} \bigg[c_{\rm b} \frac{D_{\rm b} k^2}{(D_{\rm b} k^2)^2 + \omega^2} + c_{\rm t} \frac{D_{\rm t} k^2}{(D_{\rm t} k^2)^2 + \omega^2} \bigg] \qquad (16)
$$

where $D_{\rm b}$ and $D_{\rm t}$ are the diffusion coefficients of protons in the hydrated water molecules (bound) and in the bulk, respectively, and c_b and c_t are constants characterizing the ratio of the population of the two kinds of protons. The average diffusion coefficient D_H is thus obtained as

$$
D_{\rm H} + c_{\rm b}D_{\rm b} + c_{\rm t}D_{\rm t} \tag{17}
$$

It is worth noting that the τ_{RES} value of oxygen atoms in water is usually the same order of magnitude as for the protons. Exceptions are observed for the $Cr³⁺$ and other inert cations.

The quasi-elastic neutron scattering method has recently been developed rather rapidly. However, only partial reviews of the results have been published until now.¹⁸⁴⁹

B. Spectroscopic Methods

Various spectroscopic methods have been employed in studies of ionic hydration. Among them, X-ray absorption spectroscopies have been developed very quickly in recent decades. The extended X-ray absorption fine structure (EXAFS) method is used for the determination of short-range intermolecular interactions and the X-ray absorption near edge structure (XANES) method is often employed to discuss the

electronic structure of the relevant element. NMR measurements have been used for determining the hydration structure of various ions, as well as structures of complexes in solutions. The method is preferably used for studies on dynamic properties of ions and hydrated water molecules rather than the static structure of hydrated ions after the appearance of the X-ray diffraction and X-ray absorption methods. Infrared and Raman spectroscopies are the traditional methods which have still been widely used in aqueous and nonaqueous solutions. Application of the Mossbauer spectroscopy to structural measurements of solutions is restricted to some special cases which contain Mössbauer active elements, but it is useful for studies on ionic hydration in solids, because the method can show different spectra for one element with different oxidation states. However, the method is not applicable to usual solutions without settling them by quick freezing.

/. Extended X-ray Absorption Fine Structure (EXAFS) and X-ray Absorption Near Edge Structure (XANES) **Spectroscopies**

Both methods are connected with X-ray absorption by atoms which are affected by their environmental structures. Separation of the EXAFS part from the XANES one is rather arbitrary.

Both white X-rays and a monochromatized X-ray beam are used simultaneously. A detector scans the spectrum over a wide-angle range to cover a reasonable *k* range to the Fourier transform of the EXAFS spectra. The use of synchrotron orbital radiations (SOR) is the most effective for the measurements, but several laboratory-scale instruments have also been developed. The EXAFS method has a large advantage over a usual X-ray diffraction method because of its high selectivity of the central elements and its high sensitivity to be able to use sample solutions with fairly low concentrations which are not possible to measure by the X-ray diffraction method in which scattered intensities from solvents hide information from the solutes. Since longrange interatomic interactions are practically eliminated from EXAFS spectra, only the local structure around X-ray absorbing atoms can be determined. Disadvantages in the EXAFS method are the necessity for the phase-shift correction and less reliably for frequency factors (coordination numbers) because the latter ones strongly correlate to the Debye-Waller factor, which corresponds to the temperature factor in the X-ray diffraction method.

If the multiple scattering of back-scattered electrons from surrounding atoms of the X-ray absorbing central atom is neglected, the ratio $\chi(E)$ of the incident beam and the beam after passing through a sample at a given energy *E* is described as follows:

$$
\chi(E) = [\mu(E) - \mu_0(E)]/\mu_0(E) \tag{18}
$$

where μ and μ_0 denote the measured intensity and the intensity of the incident beam, respectively. The energy dependence of χ can be converted to the dependence on the photon wave vector *k.*

Under the plane-wave approximation, the EXAFS interference function for the photoexcitation of an atomic species α is given by a superposition of contributions from all back-scattering atoms *0* surrounding α , and each contribution is expressed in the sinusoidal form with a frequency which is a function of the interatomic distance $r_{\alpha\beta}$, the number of interactions (atom pairs) $n_{\alpha\beta}$, and the total phase shift $\phi_{\alpha\beta}(k)$. The $\chi(E)$ function is related to the pair-correlation function $g_{\alpha\beta}(r)$ through eq 19 as follows:

$$
k\chi(k) = 4\pi \rho_0 S_{\alpha}(k) \sum x_{\alpha} F_{\alpha}(k) g_{\alpha\beta}(r) \times
$$

exp $(-2r/\lambda) \sin[2kr + \phi_{\alpha\beta}(k)]$ dr (19)

where the summation runs over each species present in the system. χ_{α} is the atomic fraction of the α -type atoms, ρ_0 the bulk electron density, $F_\alpha(k)$ the absolute value of the complex scattering amplitude of atom α , $\phi_{\alpha\beta}(k)$ the total phase shift, λ the mean free path of the photoelectrons at *k,* and *Sa(k)* a phenomenological term which accounts for inelastic scattering. Introducing structural parameters of interatomic distance $r_{\alpha\beta}$, the coordination number $n_{\alpha\beta}$, and the root mean square displacement $l_{\alpha\beta}$ into eq 19, we obtain

$$
k\chi(k) = S_{\alpha} \sum F_{\alpha}(k) n_{\alpha\beta} r_{\alpha\beta}^{-2} \exp(-2r_{\alpha\beta} \lambda^{-1}) \times
$$

$$
\exp(-2k^2 l_{\alpha\beta}^2) \sin[2kr_{\alpha\beta} + \phi_{\alpha\beta}(k)] \quad (20)
$$

The parameter values are calculated by a least-squares fitting procedure. The values of S_{α} , F_{α} , and $\phi_{\alpha\beta}$ should be evaluated in advance by separate experiments or by an assumption. The value of $\mu_0(E)$ is also obtained in the course of the fitting process as an adjustable parameter. Therefore, we usually use a crystal with a known structure as a reference substance to determine some of the parameter values. When a central atom is coordinated with more than one kind of ligand atoms, the situation becomes more complicated.

The Fourier transferred EXAFS spectra usually contain many ripples arising from the truncation of the EXAFS k -space spectra at a finite value, and thus, a window function is usually introduced to cut off a main peak in the r space from Fourier transform (corresponding to the radial distribution curve), and the leastsquares procedure is carried out for the extracted $EXAFS\ \&$ -space spectrum to obtain structural parameters related to the interatomic interactions.

The principles of the method are given in a review by Lee et al.²⁰ and some recent results are summarized by Magini et al.⁵

2. Nuclear Magnetic Resonance (NMR) Spectroscopy

NMR spectroscopy is a very useful method for both studying the static and dynamic structures of hydrated ions. However, explanations of NMR spectra in terms of structures of species in solution are not straightforward compared with those of radial distribution functions obtained by X-ray, neutron, and electron diffraction methods and EXAFS spectra except for simple cases. The time scale accessible in NMR measurements should be longer than 10^{-9} s, because we can use frequencies of the magnetic field of some 100 MHz. Several effects are included in chemical shifts, and the separation of these factors in order to extract structural or dynamic properties from the spectra are usually complicated, especially when we discuss the behavior of protons in aqueous solutions. The static and dynamic behavior of water molecules in the coordination sphere can be separated from that in the bulk when we measure the NMR spectra of oxygen atoms (¹⁷O), but still deconvolution of a single peak into the components is

necessary except for limited cases containing inert ions such as $\mathrm{Cr^{3+}}$.

Principles of NMR measurements for aqueous electrolyte solutions and results obtained have been reviewed by various authors.^{3,21-32}

Interactions between the total spin *I* of an atomic nucleus in a solution system and the outer magnetic field H_0 applied result in a magnetic resonance spectra in which $2I + 1$ different energy levels are separated by

$$
h\omega L = h\gamma H_0 \tag{21}
$$

where γ represents the gyromagnetic ratio of the nucleus and ωL stands for the Larmor frequency. In a solution there are many spins which interact with the local electric and magnetic fields generated by the nuclear spins and electric charges of the other atoms. The most important interactions are those between dipoles, between an outer electric field and an electric quadrupole of a nucleus, and between the spin of an atomic nucleus and the magnetic field generated by the rotation of an adjacent molecule. The Hamiltonian is given as the sum of the Zeeman Hamiltonians of the eigenstates and other types of "lattice" Hamiltonians. Finally, the Hamiltonian *Hc* is described as

$$
H_c = -\sum \int d\mathbf{r} H(\mathbf{r}) \mu_{\alpha}^M \delta(\mathbf{r} - \mathbf{r}_{\alpha})
$$
 (22)

where $H({\bf r})$ is the external magnetic field and $\mu^{\rm M}{}_{\!\alpha}$ is the magnetic moment of an α type nucleus, and thus, eq 22 expresses the coupling between the applied magnetic field and the system. By using this Hamiltonian, the spectral densities and time-correlation functions can be calculated.

The experimental information is obtained from the position and the shape of the resonance maxima on NMR spectra. Moreover, the motion of molecules and ions can be monitored through the selection of the resonances characteristic to the nucleus in question. It is worth noting that the primary information is always the effect of the environment on the motion of the spins, and all other dynamic properties are deduced from it. The following properties are usually derived from NMR experiments:

(1) From the change in the *chemical shift* of peak positions, the change in the environment of the NMR nucleus, i.e., the coordination structure of the probe atom, can be estimated. In most cases the change in the coordination number causes a fairly large difference in the chemical shifts of a given central atom. A detailed analysis of chemical shifts of proton NMR spectra can elucidate ion-water interactions in solution in spite of proton exchange reactions much faster than water exchange reactions in aqueous solution on the NMR time scale. ¹⁷O NMR spectra are often measured in order to discuss the structure and dynamic properties of hydrated water molecules.

(2) *Coordination numbers,* which are usually called *dynamic hydration numbers,* are determined under the area of the resonance peaks. In case of very inert ions such as Cr^{3+} , Be²⁺, and A¹³⁺, two ¹⁷O peaks, arising from the hydrated water and bulk water, can be separated and the areas under the two peaks are characteristic of the ratio of the two kinds of water molecules. Despite the fact that the accuracy of the area integration is sometimes very limited (a high concentration of electrolytes and good separation of the shifts are required), reasonable results can be obtained for the coordination

structure of hydrated ions. Various technical modifications have been applied to derive structural and dynamic properties from measurements of relevant NMR ions, protons, and ¹⁷O atoms.³³⁻³⁵ Separation of peaks is usually tried by cooling down the sample solution to low temperatures by lowering the solvent substitution rate constants of the sample solutions. However, since we use aqueous solutions, only a limited range of temperature decrease is possible. Sometimes organic solvents are added to aqueous sample solutions assuming that the hydration structure is not varied by the addition of the organic solvents, but the freezing point of the solutions becomes low enough to separate the rate of solvent substitution reactions in the coordination sphere and in the bulk. Of course, preferential solvation of the relevant ions, if it occurs, can become a source of errors.

When dia- and paramagnetic ions are mixed, the residence time of surrounding molecules of the latter one is significantly shorter than that around the diamagnetic ion.

(3) *Spin-spin relaxation times and diffusion coefficients* can be derived because the width of the peaks in NMR spectra is a function of rates of molecular motions, and therefore, dynamic properties of the atom showing the NMR peak can be elucidated from the analysis of the peak shape.

In the case where the decay of the $C_{\mu\mu}(t)$ autocorrelation function with time is expressed in terms of an exponential function, the translational diffusion coefficient D_t and spin-spin relaxation times T_2 can be described by a Lorentzian law

$$
C_{\mu\mu}(t) = C_{\mu\mu}(0) \exp\left[\frac{t}{T_2} - \frac{\gamma^2 g^2 D_t^3}{12}\right]
$$
 (23)

where g is a field gradient and γ the gyromagnetic ratio. Two measurements have to be carried out (with and without the presence of field gradient *g),* and from the comparison the diffusion coefficient can be determined. In most cases the so called spin-echo method is used.

In comparison between the spin-spin relaxation time measurement and quasi-elastic neutron diffraction measurements for investigating dynamic properties of ions and water molecules, the former has an advantage that the accessible diffusion coefficients are smaller than 10^{-9} cm² s⁻¹. In the latter the limit of the diffusion coefficient measurable is in the order of 10^{-7} cm² s⁻¹. The former method can cover the interaction time scale I he former method can cover the interaction time scale
of $\leq 10^{-4}$ s, much longer than $\leq 10^{-9}$ s in the latter. A drawback of the spin-spin relaxation method is that *T²* itself must not be too short, and therefore, only a limited $\frac{1}{2}$ itself must not be too snort, and therefore, only a limited in the spin of $\frac{1}{2}$ can be used (1H, $\frac{19F}{2}$ set of nuclei with the spin of $\frac{1}{2}$ can be used $(4H, 4H)$.
31D, etc.). Moreover, the diffusion coefficient is de t_{t} , $\sigma(t)$, increase, the diffusion of the molecules or ions terminable when the diffusion of the molecules or ions can be described in terms of a single translational coefficient.

(4) *Spin-lattice relaxation.* Another relaxation time *Ti* can be assigned to the response for weakly interacting spins of a kind of nuclei in a solution with the "lattice" represented by the solution in a strong magnetic field. In the simplest case, the spin I is $\frac{1}{2}$ and the lattice is constructed from a random magnetic field, the corresponding autocorrelation function being exponential and the decay being characterized by the spin-lattice correlation time τ_c . The longitudinal magnetization decays, under some conditions, as an exponential

function of time with a characteristic T_1 spin-lattice relaxation time, which relates to the correlation time as

$$
T_1^{-1} = 2\gamma^2 H_x^{2} \tau_c \tag{24}
$$

where $H_{\rm x}$ denotes the strength of the random magnetic field. When the spin is larger than $\frac{1}{2}$, the exponential function becomes a sum of exponential functions. In real solutions H_x is not expressed in terms of the random magnetic field and the quantity becomes more complicated, and therefore, the extraction of the parameter τ_c is difficult. Calculations with preliminary models and a series of measurements are usually required for data analyses. The relaxation mechanism is often separated into a few typical relaxation mechanisms of different types of molecular motions such as translational and rotational motions. However, the separation is not simple, especially in the case of proton spinlattice relaxation measurements. The relaxation via quadrupole interactions gives information similar to that by the depolarized Raman scattering of light, and thus, it is sensitive to the rotational motions. An example of the application of this method is the determination of the anisotropic rotational diffusion with two different diffusion constants of parallel D_\parallel and perpendicular D_{\perp} to the symmetry axis.

The spin-lattice relaxation via spin-rotation mechanisms is also important for the determination of rotational motions of molecules. Since the rotational mechanism is different from that of the others discussed above, the relaxation time thus measured is different from the values determined by other methods. Moreover, the efficiency depends on the spin of nuclei. The ¹³C nucleus has a large efficiency in this measurement. It may be worth noting that there are some other relaxation mechanisms, but they are not included here because they are less important for aqueous solutions.

(5) *Application of NMR measurement to the elucidation of solvent substitution reaction mechanisms.* Since water molecules bound to ions to form the hydration shell have different relaxation times from that in the bulk solvent, attempts have been made to extract information about kinetics of substitution reactions of water molecules between the hydrate shell and the bulk from NMR data.

The rate constants for the solvent substitution reactions between the hydration shell and the bulk are measured by using a high-pressure NMR technique at various temperatures and pressures. From the variation of the rates with temperature the entropy of activation ΔS^* as well as the enthalpy of activation ΔH^* of the reaction are determined. The activation volume ΔV^* is determined from the variation of the rates with pressure. A simple assumption has been introduced for the exchange mechanisms, i.e., the reaction occurs *associatively* when ΔS^* or ΔV^* is negative, while the mechanism of reactions with positive ΔS^* and ΔV^* is *dissociative.* The temperature variation for the reaction rates can be observed by usual methods at different temperatures, and thus, the entropy of activation ΔS^* can be measured by many other experimental methods. However, the specialists believe²⁴ that the determination of the sign of ΔV^* is more reliable than that of ΔS^* because a long extrapolation is needed in the latter.

The distinction of the reaction mechanisms upon the sign of the activation volume thus measured has been done under a very simple assumption that in the

associative mechanism a water molecule in the bulk enters in the coordination sphere so that the entropy in the activation state of the complex may decrease and the volume of the complex may be smaller than the sum of the volume of the original hydrated ion plus the volume of one water molecule in the bulk, provided that no significant bond length variation within the complex occurs at the activation process. It is obvious that the assumption is too simple to explain the reaction mechanism from the signs of ΔS^* and ΔV^* , because in the associative reaction mechanism the coordination number of the central ion increases, and thus, it is expected that the bond length between the central atom and the ligand atom increases which may cause the expansion of the volume of the activated complex, while the reverse may be true for the dissociation mechanism. Therefore, further discussions with more careful considerations about changes in the bond length in activated complexes, and thus, volume changes at the activation step, may be needed to interpret the thermodynamic quantities in the solvent substitution reactions, as well as any other ligand substitution reactions.

3. Mössbauer Spectroscopy

Mössbauer spectroscopy is based on a recoilless nuclear γ -ray resonance phenomenon of atomic nuclei. The formal description of the Mössbauer scattering is essentially the same as that for the X-ray scattering, provided that the scattering factor $f_a(k)$ is characteristic to the γ -scattering substitutes. The Mössbauer effect is mostly applied to systems containing ⁵⁷Fe and ¹¹⁹Sn, and ⁹⁹Ru, ¹²¹Sb, ¹²⁵Te, ¹²⁷I, ¹²⁹I, ¹²⁹Xe, ¹³¹Xe, ¹⁵¹Eu, ¹⁵³ Eu, 191 Ir, 193 Ir, 197 Au, and 237 Np are also used for Mössbauer spectroscopic measurements.

Energy changes due to energy transitions are usually called the isomer shift, chemical isomer shift, or center shift, and they are influenced by the electron distribution around a nucleus, and thus, the local environment of the excited atom or ion. Moreover, since the size of the excited nucleus may be different in the excited state than in the ground state, the shift is related to the oxidation state of the ion. These two facts provide a sensitive tool for investigating the structural changes around an ion. However, the Mössbauer spectroscopy can be used only for solids where the whole crystal lattice absorbs the energy and the momentum of the recoil. Therefore, the use of this method is impossible, in principle, for liquids without quenching liquids and solutions. Some reports have been published that the Mössbauer effect is observable for solutions which are absorbed on porous glasses. However, a question remains if no structural change occurs when the bulk solution is absorbed on the glass, or if the quenched solution preserves its structure same as that at room temperature.

4. Infrared (IR), Raman, and Raleigh-Brlllouin **Spectroscopies**

Infrared and Raman spectroscopies played a pioneering role in structural and dynamic investigations at an early stage of investigation of solutions and they became dominant in solution chemistry in the 1960s and 1970s. Even now their role in solution chemistry is very important, of course. Methods such as the Brillouin scattering spectroscopy can also be used for investigating molecular coagulation in mixed solutions. Detailed studies and reviews on the results are referred to the literature. 3,32,36

The common feature of the spectroscopic methods of this group is that electromagnetic waves which are absorbed or scattered have a long wavelength compared to molecular dimensions. They interact with the electric field of ions and molecules and the interaction Hamiltonian can be written as

$$
H_c = -\int d\mathbf{r} \left[E(\mathbf{r},t)\mu(\mathbf{r},t) + \left(\frac{1}{2}\right)E(\mathbf{r},t)a(\mathbf{r},t)E(\mathbf{r},t) + o(\mathbf{r},t) \right] \tag{25}
$$

where $E(\mathbf{r},t)$ is the electric field vector of the incident electromagnetic field, $\mu(\mathbf{r})$ the dipole moment at position $\mathbf r$ in the liquid, $a(\mathbf r,t)$ the polarizability tensor, $a(\mathbf{r},t)E(\mathbf{r},t)$ the induced polarizability, and $o(\mathbf{r},t)$ the higher order terms, which can be ignored.

If we assume that only the permanent dipoles exist in the liquid, the dynamical variable may be given by

$$
\mu(\mathbf{r}) = \sum \mu_{\alpha} \delta(\mathbf{r} - \mathbf{r}_{\alpha})
$$
 (26)

where μ_{α} is the dipole moment of the α -type particle. The absorption of the electromagnetic radiation is proportional to the spectral density $C_{uu}(\omega)$, and the related autocorrelation function is

$$
C_{\mu\mu}(t) = (1/N)\langle \sum \mu_{\alpha}(0)\mu_{\beta}(t) \rangle \tag{27}
$$

If the electric field can be written in the form of E_0 cos *at,* then

$$
\mu = \mu_0 + \sum \mu_n \cos \omega_n t \tag{28}
$$

where ω_n are the normal frequencies. Dielectric absorption arises from the contribution of μ_0 , while IR absorption is originated from changes in μ brought by vibration or rotation. It can be shown that the second term of the right hand side of eq 28 is connected to the Raleigh-Brillouin and Raman scattering. In the simple case when the polarizability is scalar, the polarization *P* can be written as

$$
P = E\alpha_0 \cos \omega t + \left(\frac{1}{2}\right) E_0 \sum \alpha_n [\cos(\omega - \omega_n)t + \cos(\omega + \omega_n)t] \tag{29}
$$

where α_0 is the static polarizability. The first term of the right hand side of eq 29 corresponds to the origin of the Raleigh scattering, and the next two terms produce the Raman lines of the spectrum. The condition is that the rotation or vibration of a molecule changes the polarizability of the molecule. In case of Brillouin scattering, the situation is more complicated, because the normal frequencies are associated with acoustic phonons in the liquid.

C. Computer Simulations

Computer simulations are theoretical approaches to elucidate structure and dynamics of molecules and molecular ensembles in the condensed phases, as well as in the gas phase. The simulations may be classified into three categories: molecular dynamics (MD), Monte Carlo (MC) simulations, and molecular mechanics (MM) calculations. The first one can describe both structural and dynamic properties of a system, because the simulation includes time in the treatment, while the latter two have no time-dependent parameter in the course of the calculations, and thus, they give only

structural or energetic information. Molecular dynamics and Monte Carlo simulations are usually applied to systems containing a number of molecules, which may be assumed to be similar to bulk systems. Special techniques are usually used in the calculation, which include the periodical boundary condition, in which a small system containing a limited number of ions and molecules is repeatedly set up to build an infinite system. Such a periodical boundary condition is accepted in crystals, but an assumption that the condition is acceptable even in liquids and solutions, provided that the basic system (basic cell) is large enough in which long-range interatomic interactions are reasonably assumed to have no correlation, is introduced. On the other hand, the molecular mechanics simulation is a calculation usually applied to an isolated molecule. Combinations between MD or MC and MM in structural and dynamic calculations of liquid and solution systems have been examined in some solutions containing polyatomic solutes, although the computer program is still rather simple. Molecular dynamics simulations in combination with molecular orbital calculations have also been tried to treat some systems in order to correct effects of polarization of electrons in molecules on intermolecular interactions.

1. Molecular Dynamics Simulations (MD)

The basic assumption of MD simulations is that any kind of interatomic interactions can be described in terms of the sum of the interactions of the atomic pairs, and thus, the multibody problem is not taken into consideration in the course of the simulation procedures, except for a very limited cases.³⁷ An interatomic interaction is described in terms of a pair potential function and the total energy of the system calculated is reduced by changing positions of atoms in a given system until it reaches the minimum value, allowing fluctuations in energies due to thermal motions at a given temperature. Average dynamic properties of atoms and molecules of a system in the equilibrium state are deduced from the atomic and molecular motions around the minimum value of the total system. For a bulk system the periodic boundary condition is usually introduced, and long-range interatomic interactions such as Coulombic interactions which may affect the behavior of other atoms beyond the size of the cell are corrected by the application of the Ewald summation,³⁸ reaction field method,³⁹ or shifted force potential method.⁴⁰ Selection of a set of good pair-potential functions for each atom pair is an essential problem for the achievement of simulation calculations. If the potential functions are reasonable, acceptable reproduction of measurable properties should be done, and thus, other properties which have never been measured should also be predicted with reasonable certainties.

Statistical averages are taken over the ensemble of phase space points of the trajectories over a sufficient length of time in order to reproduce the structural and dynamical properties of the system. The most frequently used statistical ensemble is a modified microcanonical ensemble⁴¹ where the total momentum of the system is also fixed in addition to the *(N, V,E)* parameters.

Various types of models and potentials have been applied to describe behavior of ions and water molecules. Ions are usually characterized in terms of the point charges and radii. Interionic interaction potentials are written in various ways. A Lennard-Jones-type function with suitable parameter values derived from some theories or experimental values in combination with the Coulombic interaction is often used. Other polynomials and exponential type functions are also used which are fitted to the potential functions obtained from molecular orbital calculations. The Fumi-Tosi potential is often used for describing ion-ion interactions in the alkali halide systems of rock salt type.⁴²

The potential function describing water-water intermolecular interactions has been a big subject in the development of simulation calculations. A first attempt was done by Bjerrum⁴³ who represented a water molecule as a regular tetrahedral structure with the oxygen atom centered between two hydrogen atoms and two lone with pair electrons at the apexes of the tetrahedron. Atom-atom and electron-atom distances were fixed in the model. The charges were distributed on the atoms and electrons to reproduce the permanent dipole moment of an isolated water molecule. The water model was too simple to describe the behavior of molecules in the bulk water, so then, various modifications were applied to describe structural and dynamic properties of water. Water models proposed by Ben-Naim-Stillinger (BNS model)⁴⁴ and Stillinger-Rahman $(ST2 \text{ model})^{45}$ are typical ones among the four-point $\frac{1}{2}$ models. The model proposed by Rowlinson⁴⁶ is a modified four-point charge model. Matsuoka-Clem- μ modified from μ that charge model. The issues and μ is a three-point charge model of the water molecule.

Besides these rigid models, flexible models have also been introduced, in which the bond length and bond angle between the hydrogen and oxygen atoms are changeable due to vibrations. Typical ones are the central force (CF) model⁴⁸ and its modifications proposed by Bopp-Jancs6-Heinzinger (BJH).⁴⁹ Other, less frequently used models were discussed by Bopp.⁵⁰

The recent progress of MD simulations applied to aqueous solutions have been reviewed by Heinzinger and his co-workers.50-52

MD simulations can derive the following properties of solutions:

(1) Since distances of all kinds of atom pairs in the basic cell of a given system are calculable, the radial distribution function is derived by MD simulations, which can be compared with those obtained by X-ray and neutron diffraction measurements. Any kind of partial structures and pair-correlation functions is derived from the results which are often hardly obtained by the diffraction methods, because many partial functions are overlapped in experimentally obtained functions.

(2) Angular distributions of water molecules around ions in the system are obtained. The values may be obtained by neutron diffraction measurements. In MD simulations an angular correlation function $P(\cos \theta)$ is obtainable, where θ denotes either the cation-oxygenhydrogen or the anion-hydrogen-oxygen bond angle.

(3) The numbers of water molecules in the first and second hydration shells of each ion are obtained provided that a suitable definition of the hydration shells are set. By the diffraction methods only average hydration numbers are obtainable, which are often fractional numbers. However, the hydration numbers in the first and second hydration shells of individual

cations and anions in an electrolyte solution are determinable by MD simulations, and thus the average hydration numbers obtained by the diffraction measurements can be compared with the sum of the numbers of individual ions.

(4) The time evolution of interatomic distances and bond angles in the hydration shells is calculable, from which we can discuss the rigidity of the hydration shell and the symmetry of the hydration structure. If an extremely high-speed computer were available, the substitution reaction mechanism of the hydrating water molecules with bulk water molecules would be elucidated by computer simulations. However, unfortunately, the rates of water substitution reaction, as well as those of any kind of solvent substitution reactions, are too slow to compute by using presently existing computers.

(5) The time evolution of distances and angles in a specific atomic group can be seen, which cannot be determined by any kind of experimental techniques.

(6) The self-diffusion coefficient *D* can be derived from either mean square displacements

$$
D = \lim_{t \to \infty} \frac{\langle [\mathbf{r}(t) - \mathbf{r}(0)]^2 \rangle}{6t}
$$
 (30)

or through the velocity autocorrelation function

$$
D = \lim_{t \to \infty} \frac{1}{3} \int \langle \nu(0) \nu(t') \rangle \, \mathrm{d}t' \tag{31}
$$

The self-diffusion coefficient thus evaluated may be compared with those obtained by quasi-elastic neutron scattering and NMR measurements.

(7) Spectral densities of hindered translational motions of ions and water molecules are evaluated, which are obtained by the Fourier transform of the normalized velocity autocorrelation functions. From these spectra, important conclusions can be drawn directly for the collective motions of the ions and their hydrates.

(8) Flexible models allow the calculations of intramolecular vibrations, and the results are compared with infrared and Raman spectroscopic data. Information for the rotational spectra can be obtained even by rigid models.

It should be noted here that the hydration numbers $n_{\alpha\beta}$ determined by the diffraction methods and molecular dynamics simulations are not always the same even though the structure found in both methods are very similar. In the diffraction methods the hydration number (in more general, coordination number) is evaluated from the area under the relevant Gaussiantype peak. However, the concept of the coordination shell is not clear in MD, as well as MC, because no assumption is made in the simulations that all atoms in the nearest neighbor have the same interatomic distance from the central ion, and thus, the peak caused by the ion-water interactions in the nearest neighbor in the radial distribution curve drawn by the simulation calculations is not necessary to be a symmetrical Gaussian-type one. Of course, a similar shape of peaks may appear in the diffraction measurements and the asymmetrical peak shapes are analyzed under the assumption that two or more symmetrical peaks are involved in the peak. Since such an assumption is not made in the simulations, the first hydration shell is conventionally defined as the region from $r = 0$ up to

the first minimum $r_{1,\text{min}}$ or r_1 at $g(r_1) = 1$ after passing the first maximum. In many cases, especially in cases where ion-water interactions are weak, observation of the first minimum $(r_{1,\text{min}})$ is not clear, and thus, the second definition is preferably used. Therefore, the hydration numbers $n_{\alpha\beta}$ thus determined by the simulation calculations are sometimes larger or smaller than the value determined from the diffraction methods depending on the degree of asymmetry of the first peak. It is obvious that the hydration number thus estimated by the simulation methods are not necessarily integers. In systems in which ion-water interactions are strong, the peak shape calculated by the simulations become more symmetrical and the hydration number evaluated is close to an integer value. Such an asymmetrical distribution of coordinated ligand atoms reflect to the value of rmsd, $l_{\alpha\beta}$, in the diffraction methods.

2. Monte Carlo Simulations (MC)

Monte Carlo (MC) simulations differ basically from MD. In the former atomic configurations are created randomly, and thus, MC looses time-dependent information in the simulated system. Since the configurations are selected randomly, physically irrelevant configurations can also occur. The procedure distinguishes between acceptable and nonacceptable configurations which should be thrown away in the course of the calculations. The classical formulation of the of the calculations. The classical formation of the method was given by Metropolis et al.⁵³ The criterion for the acceptance of a configuration is fixed, and the total energy of the system decreases along the line of the criterion. If the configuration yields in a lower energy, it is accepted. If not, it is accepted upon a prefixed probability criterion. The accepted configurations are then collected and stored for further analysis.

Since the acceptance criterion is fixed to the energy in the classical MC studies, the adoption of good interatomic potentials is very important to obtain reliable results. It is worth noting that although a very good reproduction of the bulk structural data is possible through the computation of the pair-correlation functions with a surprisingly small number of generated configurations, reasonable microscopic structures are sometimes hardly described in the MC simulations even after the generation of a huge number of configurations. A detailed review of the MC method has been given by Wood and Erpenbeck.⁴¹ Results on aqueous solutions are collected by Beveridge et al.⁵⁴

Recent MC techniques include significant modifications of the original MC procedures. One of such modifications is so-called reverse Monte Carlo (RMC) method, which has been invented by McGreevy and Pusztai.⁵⁵ In the RMC simulations pair potentials and energy criteria are completely omitted, and a new criterion is set for the fit between the *G(r)* function calculated from the randomly generated configurations and that from a diffraction experiment. Although the method requires extensive calculations, the advantage is clear: the often dubious pair potentials are no longer needed and the set of accepted configurations reproduces the experimental pair correlation functions to an incredibly good extent. The RMC simulations have widely been used for many monoatomic liquid systems, but only one attempt has been made for the determination of the hydration structure of ions in solution.⁵⁶

3. Molecular Mechanics Calculations (MM)

The molecular mechanics (MM) calculation is usually applied to the calculation of the total energies and energies of interatomic interaction in a molecule with different configurations or some isolated systems with a very limited number of molecules which interact each other. However, some attempts have been made to calculate optimum configurations of a polyatomic molecule in water. In the treatment, interactions between the solute and surrounding water molecules are calculated in a way similar to MD simulations and the optimum configuration of the solute is stepwise improved with the change in the orientation of water molecules. Such examinations have only been tried on the laboratory level, and therefore not yet published in scientific journals. However, in the near future elaborated programs may appear.⁴¹⁶

///. **Structural Aspects of Ionic Hydration**

In this section the static structural features of the ionic hydration shells will be discussed. Since a partial structural survey of ionic hydration has been done by Marcus in this journal,⁶ we aim to summarize structural data which are important to the discussion of dynamic properties of ionic hydration which will be done in the next section, along with time- and space-averaged structural data. The static structural information obtained by the X-ray diffraction method is believed one of the most reliable data. EXAFS data may be less accurate than the X-ray diffraction data especially for the coordination number. However, the structural data obtained by the EXAFS method are generally better than those derived by NMR measurements and any other classical methods. We summarize the structural data for the first hydration shell of various cations and anions determined by the X-ray diffraction and EXAFS methods in Tables 1-7.

The neutron diffraction method provides similar structural information as obtained by the X-ray diffraction and EXAFS methods (which are also included in Tables 1-7), but information about the orientation of hydrated water molecules in the first coordination sphere is obtainable by the former method. These structural data are separately summarized in Table 8.

In concentrated electrolyte solutions, ion pairs and complexes are often formed, which still include water molecules in the first coordination shell of the central ions. The diffraction and EXAFS methods can determine the hydration number of ion pairs and metal complexes in solution under favorable conditions, and thus, we summarize the number of water molecules in the first coordination shell of metal complexes determined by the methods in Tables 9-11. The structural information is directly connected to the total coordination number of the central ions in the ion pairs and complexes. The hydration structure of anions in ionpairs and metal complexes is difficult to measure due to their weak hydration. The structures of ion pairs and metal complexes themselves are not discussed in this review.

In Tables 12-14 the structural parameters, determined from the radial distribution functions derived by computer simulations (MD and MC), are tabulated, which can be compared with those found by the diffraction and EXAFS data. MD simulations can give

hydration numbers of individual ions so as to be able to calculate the distribution of the hydration numbers of ions, which are never obtainable by usual experimental observations. The average hydration number calculated from the distribution weighted with hydration numbers of individual ions should be compared with the hydration numbers obtained by experimental methods. Such distribution of the hydration numbers may be relatively wide in weakly solvated cations and anions and may be sharp in high valent small cations. The parameters that represent the distribution are tabulated in Table 15.

Under favorable experimental conditions or by using special techniques, we can estimate hydration numbers in the second hydration shell by the X-ray and neutron diffraction methods. The EXAFS method is less suitable than the diffraction methods for determining the structure of the second hydration shell, because the method is rather insensitive to long-range interatomic interactions. The structural data for the second hydration shell are summarized in Table 16.

The structure of the second hydration shell can certainly be determined by computer simulations with a less arbitrary definition for *the second shell.* The data are included in Table 17.

Although the NMR method is now believed a less reliable technique than the diffraction methods to determine the hydration numbers of substitution labile ions, the method can still give us useful information especially for the discussion of the dynamic behavior of hydrated water molecules. The static hydration numbers of ions determined by the NMR method are summarized in Table 18, which may be discussed by comparing those with the diffraction data given in preceding tables.

Dynamic parameters of hydrated water molecules determined by various methods are tabulated in Tables 19-27.

The methods applied are abbreviated as follows: X, X-ray diffraction; N, neutron diffraction; EX, EXAFS; MD, molecular dynamics; MC, Monte Carlo; E, electron diffraction; S, method explained in the footnote.

Sometimes various (usually two) methods are combined to discuss the hydration structure of ions. In such cases the first letter in the column of "method" in the tables refers to the source of the data.

Various models are usually adopted to interpret radial distribution curves obtained by the diffraction methods. The models used are referred to by the following abbreviations:

FNl (first neighbor model 1). A hydrated species is assumed to have a symmetrical structure such as an octahedron or a tetrahedron and the rest of the solution is assumed to have the bulk water structure, whose radial distribution function is separately measured by using pure water.

FN2 (first neighbor model 2). A hydration species is assumed to have a symmetry similar to FNl, but the water structure is also assumed to be describable in terms of some symmetrical molecular arrangement.

FN3 (first neighbor model 3). No symmetrical arrangement is a priori assumed for hydrated species and water structure.

SNM (second neighbor model). The structure of the second coordination shell is taken into account in the determination of structure of the cationic hydration shells. Water structure is subtracted by using the bulk data or assumed to have some symmetrical molecular arrangement as those assumed in FNl or FN2.

rdf. The structure is analyzed from the analysis of the radial distribution functions without least squares fitting.

iso. The isomorphous substitution method is applied by using ions with the same ionic charge and essentially the same ionic radii. The method is used for X-ray diffraction measurements.

fod (the first-order difference). The difference curve of two radial distribution functions obtained from isotopically substituted samples is analyzed. The method is only applicable to the neutron diffraction method.

Other abbreviations applicable to the tables are as follows: T, tetrahedral symmetry; O, octahedral symmetry; H, hexahedral (cubic) symmetry; L, linear configuration; Sp, special symmetry (explained in the footnote); n, no specific symmetry is assumed; n in the column "second shell" indicates that no second shell is found around the ion; i, the existence of the second hydration shell is indicated, but not analyzed; a, the structure of the second hydration is analyzed, and the structural parameters are shown in Table 16.

Most measurements were carried out at room temperature or at 25 ⁰C, unless otherwise stated in the footnote.

In the tables of computer simulation data, the structural parameters are obtained by the direct evaluation of the $g(r)$ functions. Since significant differences can arise from the potential models applied to describe water-water intermolecular interactions, the types of the potential functions are indicated in the tables: ST2, Stillinger-Rahman model;⁴⁵ MCY, Matsuoka-Clementi-Yoshimine model;⁴⁷ CF, central force model;⁴⁸ BJH, Bopp-Jancsó-Heinzinger model;⁴⁹ CI, $\frac{1}{2}$ configuration interaction model;⁵⁷ HF, potential functions derived from Hartree-Fock calculations are $\frac{1}{2}$ employed;⁵⁸ TIP4P, a version⁵⁹ of transferable interaction potential (TIP);⁶⁰ SPC-FP, an elastic potential for a water molecule with a harmonic oscillation of 0-H bonds.⁶¹

Various potential functions have so far been applied to describe ion-water and ion-ion interactions. Since variations of potential functions employed are numerous and practically dependent on authors in each work, no notation for the potential functions is indicated in the tables. Readers who are interested in the description of potential functions should consult the relevant reference. The characteristics of distributions of coordination numbers (Table 15) are calculated from the original histograms.

All values of distances (r) and root mean square deviations (rmsd) (l) are given in the picometer (pm) unit. The angles are represented by the degree unit.

The rmsd value reflects the statistical character of the distribution of interatomic ion-water distances in the hydration shell. They are usually determined from the theoretical $ki(k)$ (eq 1) or $G(r)$ (eq 3) values on the basis of a geometrical model fitted to the experimentally obtained structural data and treated as an independent variable. Under an assumption of the Gaussian shape of a peak, the rmsd value is related to the full width at half height (FWHH):

$$
FWHH = 2\sqrt{\ln 2} \times \text{rmsd} \tag{32}
$$

In the EXAFS analysis, the term Debye-Waller factor, is preferably used instead of rmsd (see eq 20). When the peak shape is not approximated to be a Gaussian type, the rmsd value does not coincide with the value of the Debye-Waller factor. It is a general trend that an increase in the size for ions with the same charge results in the increase in the rmsd value of an ion-water distance because of weakening of the ionwater bond.

Coordination structures around ions are expected to have symmetrical polyhedra according to the knowledge of coordination chemistry. This expectation is acceptable when the central ion is a transition metal ion in which the direction of the metal-water bonds is subjected by the direction of the d orbitals. However, in the hydration sphere in which ion-water interactions are essentially electrostatic, and therefore no specific direction is expected for the ion-water bonds, a polyhedral hydration structure may barely be constructed, and the number of water molecules accommodated in the first coordination sphere, and probably as well as in the second shell, of an ion may be restricted by the surface area of the central ion. The ion-water bond length may not necessarily be the same for all the bonds. Therefore an asymmetric structure of the hydration shell may result, and the average hydration number of the ion may not be an integer.

In the case where ion-water interactions are not very strong and a spherical electrostatic field is expected, the number of water molecules around ions of a given kind is not always the same, due to the fluctuation of the water structure in the bulk caused by thermal disturbance. The cases seem to be realized for alkali metal ions. MD and MC simulations show a distribution of the hydration numbers around ions, which also becomes a reason why the hydration number of a given ions is not an integer. In Table 15 the ranges of the coordination numbers found in the simulation calculations are given, together with the most frequently appearing coordination number for the ion and maximum percentage of the frequency. If the hydration structure is definitely concrete, only one value of the coordination number should be obtained with 100% frequency.

The orientation of water molecules around an ion is experimentally determinable by the neutron diffraction method. The distribution of the tilt angles is calculable by the simulations.

A. First Hydration Shell of Ions

1. Monovalent Cations

The hydration structure of ions is described, during the first stage of approach, in terms of the first neighbor ion-water distances and the number of nearest neighbor water molecules around the ion. Much attention has been devoted to the most accurate determination of these quantities from various measurement. Since 1930s widely spread numbers have been reported for a given ion under similar conditions by various experimental techniques which have different definitions of the hydration sphere, and thus a big confusion in the concept of the hydration structure of ions has been seen. Development of the diffraction methods removed

Table 1. Structural Parameters of the Hydration Shell of Monovalent Cations Derived from Diffraction and EXAFS Measurements (See Section III for the Full Definitions of the Abbreviations)

	H ₂ O/salt								
salt	molar ratio	r_{MO} , pm	l_{MO} , pm	n_{MO}	symmetry	second shell	method	model	ref
				H_3O^+					
					T		$\mathbf X$		
HCl	38.8	280		4		n		rdf	62
HCl	$4 - 96$	252		4	T	$\mathbf n$	X, N	FN1	63
DCI	55.6	288		4	T	n	N	FN3	64
DBr	55.6	288		4	T	n	N	FN ₃	64
$HCl + CoCl2$	4.3	244,290		3,1	Sp^a	$\mathbf n$	$\mathbf X$	FN3	65
$HCl + RhCl3$	75.9	270	8	4	т		$\mathbf X$	SNM	66
						n			
$HCl + Rh(ClO4)3$	20.5	275	10	4	T	n	$\mathbf X$	SNM	66
HNO ₃	25	280		4	n	n	X	FN3	67
$HClO4 + Rh(ClO4)3$	52.3	276	10	4	T	n	$\mathbf X$	SNM	68
				$NH4$ ⁺					
NH_4F	24,11	295		5.3	n	n	$\mathbf X$	FN1	69
NH_4F	3.6	288	38	4	\mathbf{T}^b	n	$\mathbf X$	FN1	70
NH ₄ Cl	8.5	280	14	4	\mathbf{T}^b	n	$\mathbf X$	FN1	70
NH ₄ Br	7.6	282	19	4	\mathbf{T}^b	n	$\mathbf X$	FN1	70
					\mathbf{T}^b				
NH ₄ I	8.2	291	8	4		$\mathbf n$	$\mathbf X$	FN1	70
NH ₄ Cl	25	306	$\bf{22}$	8	\mathbf{T}^c	n	X, MD	FN2	71
NH ₄ Cl	11.1	300		$10 - 11$	n	n	N	fod	72
NH_4ClO_4	37	304	12	8	н	n	$\mathbf X$	FN ₃	73
$(NH_4)_2SO_4$	15.8	301	15	8	н	$\mathbf n$	$\mathbf X$	FN ₃	73
				$Li+$					
LiCl	$3 - 8.2$	$195 - 225$	$25 - 31$	4	T	n	X, N	FN1	74
LiCl	$4.4 - 24.5$	210		4			X	FN1	75
					n	n			
LiCl	54.3	190		4	T	n	N	FN ₃	76
LiCl	13.9,27.8	208,217	51,25	4	T	$\mathbf n$	$\mathbf X$	FN2	77
LiCl	5.6,15.6	195,195		3.3,5.5	n	i	N	fod	78
LiCl	4.0	218		5	n	n	X, MD	FN3	79
$LiCl + CoCl2$	7.5	203	14	4	T	$\mathbf n$	x	FN1	80
$LiCl + CoCl2$	$8 - 17$	195-207	13	4	T	n	$\mathbf X$	FN1	81
LiI	25	210	12	6	$\mathbf 0$	a	X, MD	SNM	82
LiI	20,9.2	220,228	14	4	T	$\bf n$	X	FN1	83
LiBr	$8.4 - 25$	$215 - 225$	25	4	T	n	$\mathbf X$	FN1	84
LiBr	31.7	194		4.5	$\bf n$	$\bf n$	N	fod	85
				$Na+$					
NaCl	10.2	241		6	$\mathbf 0$	n	$\mathbf X$	FN1	86
NaCl	54.3	250		8	n	$\mathbf n$	$\mathbf N$	FN3	87
NaCl	13.9, 27.8	242,242	27,21	4	T		$\mathbf X$	FN2	77
						$\mathbf n$			
NaI	7	240	9	4	n	$\bf n$	$\mathbf X$	FN3	88
NaNO ₃ ^d	$6.1 - 9.3$	$244 - 248$	$19 - 22$	6	n	$\mathbf n$	$\mathbf x$	FN ₃	89
				K^+					
KOH	$27.5 - 15.9$	287,292		4	$\bf n$	n	X	rdf	90
KOH	3.2	292		4 or 6	n	$\mathbf n$	$\mathbf X$	rdf	91
KOH	38.8	270		6	$\mathbf n$	n		rdf	62
KF	$13.3 - 27$	295		5.3	$\mathbf n$	$\mathbf n$	$_{\rm X}^{\rm X}$	rdf	69
KCI	38.8	270		6	n	$\mathbf n$	X	rdf	62
KCI	10.1	292		$\overline{\mathbf{4}}$	n	n	$\mathbf x$	$\operatorname{\textbf{rdf}}$	90
KCl	53.7	270		8	n	n	N	FN3	87
KCI	13.9,27.8	280,280	19,17	6	$\mathbf 0$	n	$\mathbf x$	FN2	${\bf 77}$
KCI	12.5	260		weak	$\mathbf n$	n	N	fod	92
ΚI	108.4,19.5	290		3.2,2.2	$\mathbf n$	n	X	rdf	93
				$Cs+$					
CsF	$2.3 - 8.0$	$307 - 321$		6	\mathbf{O}^e	n	$\mathbf X$	FN ₃	94
CsCl	53.1	295		8	n	n	N	FN ₃	76
CsCl	13.9,27.8	315,309	27,14	6,8	0, H		$\mathbf X$	FN2	77
						n			
				Ag+					
AgNO ₃	14.2	245	9	2.45	L	a	x	FN ₃	95
AgNO ₃	16.6,5.5	236,234		3.9, 2.6			EX	FN3	96
					n	n			
AgNO ₃	16.6,4.3	242,243	4,4	4	T	a	X	SNM	97
$AgNO3 + AgI$	$2.7\,$	246	4	$\overline{\mathbf{4}}$	T	a	$\mathbf X$	SNM	97
AgClO ₄	16.8	231	8	2.9	$\mathbf n$	n	EX	FN ₃	98
AgClO ₄	10.4	241	9	$\bf{2}$	L	a	x	FN3	95
AgClO ₄	16.3,3.3	238,243	5,5	$\overline{\mathbf{4}}$	T	a	X	SNM	97
$AgClO4 + AgI$	$3.2\,$	248	5°	$\overline{\mathbf{4}}$	T	a	X	SNM	97
	10.6	241	17	3.9	T		N	FN3	99
$AgClO4 + HClO4$						i			

^a Trigonal pyramid. ^b A water molecule in the tetrahedral water network ^a Trigonal pyramid. ^{*b*} A water molecule in the tetrahedral water network is substituted with an NH₄+ ion. ^c Four water molecules are located at larger distances. ^{*d*} Outer-sphere ion pairs are formed. *^e* A sh -range order analogous to the crystal structure is assumed.

a major part of the confusion, because the coordination sphere is more clearly defined in terms of the radial distribution functions. Some of ambiguities that still remained in the explanation of peaks in the radial distribution curves are explained with supports of MD and MC simulations.

Ion-water distances can be determined with the rather high accuracy of a few tens picometers, independent of the experimental methods. However, the variation of the ion-water distance is sometimes significant and strongly depends on the methods. An extreme may be seen for the Li⁺ ion, the *rLi-o* distance reported varies between 194 and 228 pm as shown in Table 1. Another example may be the $Ni^{2+}-OH_2$ distance which converges in an 11-pm range from 204 to 215 pm (Table 2). On the other hand, the length for the $Mg^{2+}-OH_2$ bond distance varies within a 7.6-pm range, and four different measurements for the hydrated Al^{3+} ion reported the r_{Al-O} distances varying only by 3 pm (Table 3).

The hydration structures of monovalent cations are summarized as follows:

 H_3O^+ . The hydration number of H_3O^+ is 4.⁶²⁻⁶⁸ The hydrated species is reported to be the tetrahedral $[H_3O^+$ - $(\dot{H}_2O)_4$] complex for most cases, $62-64,66-68$ but a trigonalpyramidal structure with a very weak $HOH - [OH_3]$ ⁺- $(H₂O)₃$] interaction is also proposed, the structure being better fitted to the structure expected from molecular orbital calculations.⁶⁵ In the latter structure three short 0–0 bonds of 244 pm and a long and practically not interacting $O-O$ bond of 290 pm are concluded. Thus, a proton has four water molecules, one of the water molecules forms H_3O^+ and the remaining three molecules form a planar triangle structure around the H_3O^+ ion.⁶⁵ However, all other works concluded that one H3O + species is tetrahedrally surrounded by four water molecules, and thus, one proton *has five* water molecules including the core and the first coordination shell. The H_3O^+ $-H_2O$ bond length in the species is reported to be close to the H_2O-H_2O bond length in the bulk water, ϵ except for one case.⁶³ The hydration structure of H_3O^+ found by neutron diffraction experiments is also $tetrahedral^{63,64}$ (see Table 8).

NH⁴ + . An ammonium ion is believed to be easily embedded in the bulk water structure because its molecular structure is similar to the hydrogen-bonded water structure. The reported hydration number of an NH⁴ + ion is widely spread from 4 to 11. The hydration number of 4 may be easily concluded from the expected structure of hydrogen-bonding interactions between each proton in the $NH₄⁺$ ion and a water molecule. However, $NH₄$ ⁺ is a structure-breaking ion and thus it may not form a strong hydration shell like Li⁺ and Na⁺. A large hydration number and a long r_{MO} distance as reported in refs $71-73$ may indicate weak NH_4 ⁺- H_2O interactions. Similar results are obtained from computer simulations.²⁷⁸²⁷⁹ However, it should be noted that it is extremely difficult to distinguish between oxygen and nitrogen atoms by the X-ray diffraction method and the method is practically blind for the detection of protons around an atom. Neutron diffraction with isotopic substitution is more promising.⁷²

Li⁺ . Since a water molecule is regarded as a hard ligand, the ion-water bond length is well approximated as the sum of a half of the size of water molecule plus the characteristic ionic radius of an ion. The definition of the size of molecules and ions is more or less arbitrary because the electron cloud diffuses without a sharp boundary. Nevertheless, the radius of a water molecule, which is conventionally assumed to be spherical, is about 140 pm which is estimated from the water-water bond length in the bulk water. The distance of a water

molecule from the central ion depends on the coordination number and the spin state of the ion according to Shannon.³⁴¹ The ion-water bond lengths in the first coordination sphere of most ions have been determined with deviations of \pm (2-3) pm from the simply averaged measured values with various methods for ions with rather definite coordination structures.

However, the ion-water bond length in the first coordination shell is rather largely dispersed in alkali metal ions. As already mentioned above, the Li^+ - H_2O bond length has been reported to be 194-228 pm, the variation is certainly beyond the uncertainties of measurements.74-85 The bond length measured by the neutron diffraction method seems to be shorter than that by X-ray diffraction (see Table 1). An interpretation of the shorter bond is that the neutron diffraction method measures the distance between the nuclei of lithium and oxygen, while the X-ray diffraction method measures the distance between the centers of electron clouds of the atoms, which may be deformed due to polarization. An assumption of a point electron distribution at the center of an atom at the analysis of X-ray diffraction data may cause another error for such a largely polarized case of water molecules. Another and more probable reason why the $Li⁺-OH₂$ bond length disperses by observers is the concentration dependence of the coordination numbers of the lithium ion, which causes the change in the bond length. It is generally accepted that the ion-ligand bond length in a complex with a higher coordination number is longer than that with a lower coordination number. As we see from MD and MC simulations (see Table 15), the hydration number of the lithium ion distributes over the range 3-7, and the most frequent coordination number estimated by the simulations depends on the case, probably on the concentration of the ions. As expected, the most frequently appearing coordination number $(MFACN)$ of Li⁺ is smaller (e.g., in systems with $H₂O/$ salt ratio of 4, the MFACN becomes 4-5) in a system with a high concentration than that with a low concentration (MFACN $= 6$). The concentration dependence of the hydration number of Li⁺ is not clear from the data in Table 1, however. The average hydration number of Li⁺ determined from the area of the first peak of the radial distribution function derived by computer simulations is also not conclusive, but the Li⁺ -H2O bond length seems to be converged in a range of 200-213 pm (see Table 12). The hydration number of 200-210 pm (see Table 12). The hydration number
of Lit estimated from the chemical shift of proton NMR of L₁ estimated from the chemical shirt of proton NWIN
spectra was 3³¹³ (Table 18), and the value seems to be too small.

Na⁺ . Sodium ions have various numbers of water molecules in the first hydration shell according to computer simulations^{60,281,300,301} (see Table 15), as seen in the case of Li⁺. The hydration number of Na⁺ measured by the diffraction methods is distributed over the range 4-8.77,86 - 88 The average hydration number of Na⁺ estimated from the area under the peak of the radial distribution function calculated by computer simulations is 6^{60,281–284,297,299} or close to 6.^{58,285,289,398,300–302} The Na^{+} – H_2O bond length determined by the diffraction methods falls in a relatively narrow range of 240- $250 \text{ pm.}^{77,86-89}$ A shorter length of the Na⁺-H₂O bond **of 230-240** $\text{pm}^{57,58,60,280-285,289,295-302}$ **is given in the** estimation by computer simulations (Table 12).

Table 2. Structural Parameters of the Hydration Shell of Divalent Cations Derived from Diffraction and EXAFS Measurements (See Section III for Full Definitions to the Abbreviations)

	H ₂ O/salt								
salt	molar ratio	r_{MO} , pm	l_{MO} , pm	n_{MO}	symmetry	second shell	method	model	ref
				$Be2+$					
BeCl ₂	10	167	3.5	4	$\mathbf T$	\mathbf{i}	X, MD	FN2	100
				Mg^{2+}					
Mg^{2+}	55	206.6		6	n	i	X	iso	101
	25	200			\mathbf{o}		$\mathbf X$		
MgCl ₂				6		n		rdf	102
MgCl ₂	$9.8 - 13.0$	210		$8.1 - 7.9$	n	n	$\mathbf X$	rdf	103
MgCl ₂	$27.1 - 55.5$	210-212	$4 - 12$	6	o		$\mathbf x$	SNM	104
						a			
MgCl ₂	11.5	211	$\mathbf{12}$	6	$\mathbf O$	a	$\mathbf x$	SNM	105
MgCl ₂	50	212	4	6	\mathbf{o}	a	X, MD	SNM	106
$MgCl2 + CaCl2$	25	204.4	6	6	o	n	X	FN1	107
$MgCl2 + KClo$	12.3,73.6	209.2	8.4	6	$\mathbf o$	n	$\mathbf X$	FN2	108
$MgCl2 + CsClb$	13.2,41.5	208.1	9.3	6	$\mathbf 0$	n	$\mathbf X$	FN2	108
$Mg(NO3)$ ²	$10.8 - 24.8$	211	u	6	\mathbf{o}	a	$\mathbf X$	SNM	109
Mg(ClO ₄) ₂	18	212	12.9	6	$\mathbf 0$	a	$\mathbf X$	SNM	110
MgSO ₄	20.5	209.4	12	6	0	a	$\mathbf X$	SNM	m
$Mg(BF_4)_2$	$14 - 37$	215		$6.4 - 6.2$	0	i	$\mathbf X$	rdf	112
				Ca^{2+}					
$Ca2+$	55	233		6		\mathbf{i}	$\mathbf X$	iso	101
					n				
CaCl ₂	38.8	240		6	$\mathbf n$	n	$\mathbf X$	rdf	62
CaCl ₂	$12.3 - 55.8$	$241 - 242$	$13 - 15$	6	\mathbf{o}	n	$\mathbf X$	FN1	113
$CaCl2 + MgCl2$	25	242.8	14	6	o	n	$\mathbf x$	rdf	107
CaCl_{2}	8.8	236,270	9.	$3 + 6$	0	n	$\mathbf X$	rdf	105
CaCl ₂	12.4	240		5.5	n	n	${\bf N}$	fod	114
CaCl ₂	55,20,12	246,239,241		10,7.2,6.4	n	n	${\bf N}$	fod	115
CaCl ₂	50	239	13	6.9	$\mathbf n$	n	$\mathbf X$	FN3	116
CaCl ₂	8.6	244		6.0	n	a	$\mathbf x$	FN3	117
							$\mathbf X$		
CaBr ₂	26,44.1	240,249	12,20	6.6	\mathbf{o}	n		FN1	118
				Sr^{2+}					
SrCl ₂	21.5	260		7.9	n	n	X	rdf	103
SrCl ₂	26.5,34.6	264	12,14	8	Sp ^d	a	$\mathbf X$	SNM	119
Sr(CIO ₄) ₂	18.5	265		15	n	$\mathbf n$	${\bf N}$	fod	120
				$Ba2+$					
BaCl ₂	36	290		9.5	n	$\mathbf n$	$\mathbf X$	rdf	103
				$V^{2+}e$					
		221		6	n	$\mathbf n$	EX	FN3	344
				Cr^{2+}					
CrSO ₄	55	208'		$\boldsymbol{\psi}$	n	$\mathbf n$	EX	FN3	344
				Mn^{2+}					
Mn(ClO ₄) ₂	25.4	220		6	0	$\mathbf n$	X	rdf	121
Mn(CIO ₄) ₂	55.5	217.7	7.8	6	O	n	\mathbf{EX}	FN3	122
				6					
MnSO.	20,50,100	220			n	n	X	rdf	123
				$Fe2+$					
FeCl ₂	14-140	218		6	o	n	EX	rdf	124
$\text{FeCl}_2 + \text{HCl}$	52.2	213	28	6	n	i	N	fod	125
Fe(CIO ₄) ₂	25.5	212		6	o	n	$\mathbf X$	rdf	121
		209.5		6	0		EX	FN3	
Fe(CIO ₄) ₂	55.5		8.1			n			122
$\mathbf{Fe}(\mathbf{ClO}_4)_2$	55.5	228	$\mathbf{12}$	6.1	0	n	\mathbf{X}^g	SNM	126
				$Co2+$					
$Co2+$		209		6					
	55				n	i	$\mathbf X$	iso	101
CoBr ₂	$\bf 17$	211		5.9	$\mathbf n$	n	$\mathbf x$	rdf	127
$Co(NO_3)_2$	550	208					\mathbf{EX}	rdf	128
					n	n			
Co(CIO ₄) ₂	$25\,$	208		6	\mathbf{o}	n	$\mathbf X$	rdf	121
Co(CIO ₄) ₂	15	210	12	6	$\mathbf O$	a	$\mathbf X$	SNM	129
							$\mathbf X$		
$\cos \! \! \alpha_{\scriptscriptstyle 4}$	$25 - 100$	215		$6.1 - 6.3$	$\mathbf n$	n		rdf	130
				$Ni2+$					
$Ni2+$	55	206.5		6	n	\mathbf{i}	$\mathbf X$	iso	101
NiCl ₂	25	210		6	$\mathbf 0$	n	$\mathbf X$	rdf	102
NiCl ₂	12.6	205		5.8	$\mathbf n$	n	${\bf N}$	fod	131
$\rm NiCl_{2}$	12.6,27.4	205,206	10,13	6	\mathbf{o}	n	$\mathbf X$	FN1	132
NiCl ₂	12.6-645	$207 - 210$		$5.8 - 6.8$	$\mathbf n$	n	${\bf N}$	fod	133
NiCl ₂	14,18	206,207	4,4	7,6.2		n	\mathbf{EX}	FN3	134
					n				
NiCl ₂	28	206	12	6	0	a	X	SNM	105
NiCl ₂	14	206	13	6	$\mathbf 0$	a	$\bar{\mathbf{x}}$	SNM	105
NiCl ₂	13.8	205.9	7.4	6.4	$\mathbf n$	n	\mathbf{EX}	FN3	135
NiCl ₂	11.5	207	13	5.8	$\mathbf n$	n	${\bf N}$	fod	136
				$5.9 - 5.4$					
NiCl ₂	$28 - 550$	206			n	n	$\mathbf N$	fod	137
NiCl ₂	25	206		5.9	$\mathbf n$	n	${\bf N}$	rdf	138
		207		6.0			${\bf N}$	rdf	
NiCl ₂	25				$\mathbf n$	n			138
\mathbf{Nic}_{12}	$\bf 25$	207		5.5	n	n	N	rdf	138
$Ni(NO_3)_2$	550	205			n	n	EX	rdf	139
$Ni(NO_3)_2$	13.8	205.5	6.7	6.6	$\mathbf n$	$\mathbf n$	EX	FN3	135
$Ni(NO_3)$	$13 - 110$	205-206		$6.4 - 7.1$	n	n	${\bf EX}$	FN ₃	140
Ni(ClO ₄) ₂	26.6	204		6	0	n	$\mathbf X$	rdf	121
Ni(CIO ₄) ₂	14.6	207		5.8	n	n	${\bf N}$	fod	133
Ni(ClO ₄) ₂	15.0	205	10	6	$\mathbf 0$	$\mathbf a$	$\mathbf X$	SNM	129
NiSO ₄	$25 - 100$	215		$6.15 - 6.32$	$\mathbf n$	n	$\mathbf X$	rdf	130

 a KCl ion pairs are formed. b CsCl ion pairs are formed. c The short-range order analogous to the crystalline hydrate structure is assumed. ^d Square antiprism. Composition of the salt used is not reported. No detail information is given in the literature except for the bond length. $/\overline{A}$ xial Cr²⁺-OH₂ bond length is not observable. I Mössbauer spectra of frozen solutions are also investigated.
^h Outer-sphere complex formation occurs. In H₂O. In D₂O. * In "null mixtur lines indicate the parameters values of the equatorial and axial bonds, respectively, in the distorted octahedral structure. *^m* Ab initio calculations and IR spectroscopy are also applied. " Distorted (elongated) octahedron. *°* Jahn-Teller effect is not found, *P* Anomalous X-ray scattering investigation. « In the glassy state. ^r Parameters refer to the hydrated Zn2+ ion. The structural parameters of hydrated water molecules of the mono(glycinato)zinc(II) and bis(glycinato)zincate(II) complexes are given in Table 11. The structural parameters of the tris(glycinato)zincate(II) complex is also given in the same reference. ^s Parameters refer to the hydrated Zn2+ ion. The structural parameters of hydrated water molecules of the mono(α -alaninato)zinc(II) and bis(α -alaninato)zincate(II) complexes are given in Table 11. The structural parameters of the tris(α -alaninato)zincate complex are also given in the same reference. α Two different M-OH₂ bond lengths are reported.

K + . The potassium ion is one of the most difficult ions to determine the hydration structure by the X-ray and neutron diffraction methods, because the ionic size of K^+ is evaluated to be 152 pm for the case of $CN =$ 6 and 165 pm for $CN = 8,341$ and thus, the length of the

 K^+ -H₂O bond is very close to that of the H₂O-H₂O bond in the bulk. Therefore, the hydration number of K + is not possibly determined by the diffraction methods without an assumption for the water structure, and thus, the hydration number estimated by the

^a Four equatorial H₂O molecules at 198.5 pm and two axial ones at 230 pm are also proposed. ^b Three different isotopic mixtures were used. *^c* Highly hydrolyzed solution. *^d* The second sample is heated to 90 ⁰C during the preparation.

diffraction methods is not very reliable, although the K^+ - H_2O bond length may be determinable by the diffraction method within a reasonable uncertainties. The length of the K⁺ -H2O bond is reported to be 260- $295 \,\mathrm{pm}^{62.87,90-93}$ (Table 1). On the other hand, computer s imulations can evaluate the K^+ - H_2O bond length and the average hydration number of the K^+ ion without an assumption of the bulk water structure. The results μ assumption of the burk water structure. The results summarized in Table 12 show that the K⁺-H₂O bond length is 271-286 pm^{57,58,280-284,285,303} except for one case (265 pm^{295}) and the hydration number of K⁺ is 6.3-7.8.57,281-285,303 The distribution of hydration numbers around individual K⁺ ions in water shifts toward larger numbers $(4-8)^{281}$ (see Table 15) than that for Li^+ and $Na⁺$ as expected from the larger size of the $K⁺$ ion than $Li⁺$ and $Na⁺$ ions. However, since the $K⁺-H₂O$ interactions are weaker than the Na^+ - H_2O interactions, the hydration number of the K^+ ion estimated by NMR measurements is smaller than that of $Na⁺ ³¹³$ (Table 18).

Rb⁺ . The rubidium ion has not been investigated by means of X-ray diffraction, because it emits strong fluorescent X-rays upon irradiation of X-ray beams on samples by the most frequently used Mo characteristic wavelength. Since no diffraction data are available, no computer simulation has been carried out because of the lack of data to be compared with the results. An NMR measurement shows the hydration number of

Rb⁺ to be 3.5,³¹³ but the result has no reason to be believed.

Cs⁺ . There are some results for the hydration structure of Cs⁺ . The diffraction methods show a hydration number of either 6 or 8 for Cs⁺ and the Cs⁺- $\dot{\mathrm{H_2O}}$ bond length of 295–321 pm^{76,77,94} (Tables 1 and 8). The results by computer simulations show a spread of values for the average hydration number of 5.3- $8.2^{289,304,305}$ (Table 12). The Cs⁺-H₂O bond length estimated from the peak position of the radial distribution curve determined by computer simulations is 303-320 pm, which is close to the value determined by the diffraction methods.

Ag⁺ . The hydration structure of Ag⁺ is interesting from the viewpoint of coordination chemistry, because the ion can form both 4- and 2-coordination structures. The final conclusion is that the Ag⁺ ion has a regular tetrahedral structure with the Ag⁺– $\rm H_2O$ distance of 238– $248 \text{ pm}^{97,99}$ (Table 1).

Orientation of water molecules in the first hydration shell has been investigated by means of neutron diffraction and computer simulations. Molecular orbital calculations of hydrated ions in the gas phase show that the hydrated water molecules are oriented so as to direct their dipoles toward the central ion, i.e., the M + -O-H angle is close to 127° and the ion, oxygen and two hydrogen atoms in a hydrated water molecule lay in one plane with the tilt angle $\theta = 0^{\circ 342}$ (see figure in

Table 4. Structural Parameters of the Hydration Shell of Rare Earth Ions Derived from Diffraction and EXAFS Measurements (See Section III for the Full Definition of the Abbreviations)

	H_2O /salt								
salt	molar ratio	r_{MO} , pm	$\iota_{\textsf{MO}}$, pm	n_{MO}	symmetry	second shell	method	model	ref
LaCl ₃ ^a	$20.8 - 31.9$	248		8.0	H	n	X	FN3	184
LaCl ₃ ^a	14.6	258		9.1	n	n	$\mathbf X$	rdf	185
\mathbf{LaBr}_{3}^a	$20.9 - 18.8$	248		$8.0 - 7.9$	n	n	$\mathbf X$	rdf	186
La(CIO ₄) ₃	12.0	257.0	8.9	8.0	n	a	$\mathbf X$	SNM	10
PrCl ₃	14.6	254		9.2	n	n	$\mathbf X$	rdf	185
PrCl ₃	14.6	248		10	n	n	${\bf N}$	iso	187
NdCl ₃	32	241		$_{\rm 8.0}$	н	n	$\mathbf X$	rdf	188
NdCl ₃	16.5	251		8.9	n	n	$\mathbf X$	rdf	185
NdCl ₃	17.5	248		8.5	n	n	${\bf N}$	rdf	189
NdCl ₃	17.5	248	9	8.4	n	n	${\bf N}$	fod	190
NdCl ₃	14.6	248		10	n	$\mathbf n$	${\bf N}$	iso	187
$Nd(CIO4)3a$	28	251	9.1 ^b	9.5	n	n	EX	rdf	191
$Nd(ClO4)3c$	28	251	5.0^b	9.6	n	n	EX	rdf	191
Nd(ClO ₄) ₃	55,165	250,252	$30,28^d$	8.9,8.8	n	n	N	fod	192
SmCl ₃ ^a	17.2	247		8,8	n	$\mathbf n$	$\mathbf X$	rdf	193
Sm(ClO ₄)	15.2	245.5	8.9	8,0	n	n	$\mathbf X$	SNM	10°
Sm(ClO ₄) ₃	28	245	7.7 ^b	9.3	n	n	EX	rdf	191
$Sm(ClO4)3c$	28	245	6.0 ^b	8.6	n	n	EX	rdf	191
Sm(ClO ₄) ₃	55	246	31 ^d	8.5	n	n	${\bf N}$	fod	192
EuCl ₃ ^a	17.2	245		8.3	n	n	X	rdf	193
Eu(CIO ₄) ₃	28	243	7.4 ^b	8.6	n	n	EX	rdf	191
$Eu(CIO4)3c$	28	243	5.7 ^b	8.9	n	n	EX	rdf	191
GdCl ₃	26.8	240		9.9	n	n	$\mathbf x$		194
GdCl ₃	30	237		8.0	н	n	$\mathbf x$	rdf	195
Gd(CIO ₄) ₃	28	241	6.6 ^b	7.6	n	n	EX	rdf	191
$Gd(CIO4)3c$	28	241	5.5 ^b	7.5	n	n	EX	rdf	191
TbCl ₃	15.9	241		8.2	n	n	X	rdf	196
TbCl ₃	18.4	239		9	$\bf n$	n	${\bf N}$	iso	187
Tb(CIO ₄) ₃	39.4,13.7	240	7.7,10	8.0	n	n	X	SNM	10
Tb(CIO ₄) ₃	13.7	240		8.0	n	n	X	iso	180
Tb(CIO ₄) ₃	28	239	6.6 ^b	3.9	n	$\mathbf n$	EX	rdf	191
Tb(CIO ₄) ₃ c	28	238	7.5^{b}	8.0	n	$\mathbf n$	EX	rdf	191
DyCl ₃	6.9	240		7.9	n	n	$\mathbf x$	rdf	196
DyCl ₃	21.0	237	15	7.4	n	n	${\bf N}$	fod	197
\mathbf{DyCl}_3	18.4	239		9	n	n	${\bf N}$	iso	187
DyCl ₃	55,165 28	238 237	25 ^d 6.8 ^b	7.8	n	$\mathbf n$	N EX	iso rdf	198
Dy(CIO ₄) ₃ Dy(C1O ₄) ₃ c	28	237		8.1	n	n	EX	rdf	191 191
Dy(C1O ₄) ₃	55	239	4.2 ^b 25 ^d	6.8	n	n	N		198
$\mathrm{Er} \mathrm{Cl}_{2}^{a}$	$58.5 - 18.2$	230		7.9 $6.3 - 6.5$	n н	$\mathbf n$ i	$\mathbf x$	iso rdf	199
ErCl_3	15.7	237		8.2			$\mathbf X$	rdf	196
ErCl_3	53.9	235	9,5	8.0	n	$\mathbf n$		SNM	11
ErBr ₃	52.9	235	9.5	8.1	n n	a a	$\frac{\text{X}}{\text{X}}$	SNM	11
$\mathbf{Er} \mathbf{I}$ ^a	41.7	230		6.3	н	\mathbf{i}	X	rdf	199
Er(CIO ₄) ₃	12.1	236	10	8.0	n	$\mathbf n$	$\mathbf x$	SNM	10
Er(CIO ₄) ₃	12.1	236		8.0	n	$\mathbf n$	X	rdf	180
Er(CIO ₄) ₃	28	234	6.4^{b}	7.8	n	n	EX	rdf	191
$Er(CIO4)3c$	28	234	4.3 ^b	7.5	n	$\bf n$	\mathbf{EX}	rdf	191
$Er(CIO4)3e$	$49.2 - 12.2$	$235 - 236$	9	$8.0 - 7.9$	n	a	X	SNM	11
TmCl ₃	15.3	236		8.2	n	n	$\mathbf X$	rdf	196
TmCl ₃	15.6	233		8	n	$\mathbf n$	$\mathbf N$	iso	187
Tm(CIO ₄) ₃	28	233	6.6 ^b	8.0	n	$\bf n$	\mathbf{EX}	rdf	191
$Tm(CIO4)3c$	28	233	4.2^{b}	7.3	n	$\mathbf n$	EX	rdf	191
YbCl ₃	15.6	233		8	n	$\mathbf n$	${\bf N}$	iso	187
Yb(CIO ₄) ₃	55	232	24^d	7.8	n	$\bf n$	N	iso	198
LuCl ₃	15.3	234		8.0	n	$\mathbf n$	$\mathbf X$	rdf	196
Lu(ClO ₄) ₃	28	231	6.4 ^b	7.7	n	n	\mathbf{EX}	rdf	191
$Lu(ClO4)3c$	${\bf 28}$	231	4.5 ^b	6.8	n	$\bf n$	\mathbf{EX}	rdf	191

" The formation of outer-sphere complexes is probable. temperature. ation of outer-sphere complexes is probable. ⁵ Debye–Waller factor. ^c
^d The value of the full width of half height (FWHH). ^c Complex for. Valler factor. ^c Quenched to the glassy state at a liquid nitrogen e Complex formation is not excluded at higher concentrations.

^a In the glassy state at -148 °C. b KCl ion pairs are formed. ^c CsCl ion pairs are formed. ^d 25% [Zn(H₂O)₆]²⁺, 25% [Zn(H₂O)₂Br₂], 35% [Zn(H₂O)Br₃]-, and 15% [ZnBr₄]²⁻ complexes coexist. e Parameters at –40, 25, 40, 80, and 100 °C are given. / At 70 °C. s LiI ion pairs are formed. ^h [ZnI_n]⁽²⁻ⁿ⁾⁺ complexes are formed. ⁱ Supercooled at -5 °C. ^j In the glassy state at -20 °C.

Table 8). In the solution phase where the hydrated ions are surrounded by water molecules in the second coordination shell and in the bulk, on the other hand, the water molecules in the first hydration sphere are hydrogen bonded with water molecules outside the sphere, and thus, the water molecules in the first hydration sphere have the tilt angle θ of 30-70° (Tables 8 and 12).

2. Divalent Cations

Be2+ . The beryllium ion is unique among the cations in the 2 group due to its high charge and small ionic size. Different from other ions, the beryllium ion has only four water molecules in the first coordination sphere according to the results of the X-ray diffraction measurements.¹⁰⁰ The same structure has been confirmed by MO calculations.³⁴³ MD simulations with modified potential functions, in which charge transfer from water molecules to the beryllium ion is more or less taken into consideration, show the hydration number of the beryllium ion to be 4^{306} instead of 6^{100} The conclusion may be obvious for most coordination chemists, because the coordination number of beryllium ions in many complexes is 4, and no 6-coordinated beryllium complex has been found. The $Be^{2+}-OH_2$ bond length has been determined to be 169 pm by the solution X-ray diffraction method, 100 which is close to the sum of the reported ionic radius of the beryllium $\frac{1}{100}$ and the size of a water molecule. The bond length evaluated by molecular dynamic simulations^{100,306} is slightly longer than the experimental value, probably due to a weaker $Be^{2+}-H_2O$ interaction in the simulation calculations than the real one in spite of the modification of the potential function 306 from the original one.¹⁰⁰ NMR data also tell us the hydration number of 4 for the beryllium ion as expected $35,315-317$ (Table 18).

Mg2+ . The magnesium ion is one of the most widely investigated ions by the X-ray diffraction method and computer simulations, and their results agree fairly well each other.^{101-112,132,136} The Mg²⁺-OH₂ bond length is evaluated to be 200-215 pm by the X-ray diffraction method, and many values fall in the range of 200-212 pm.^{103-106,109,110} An MD calculation shows an octahedrally coordinated hydration structure of Mg^{2+} with

the Mg2+-OH2 bond length of 200 pm.³⁰⁸ In fact, the distribution of hydration numbers around Mg²⁺ ions is very sharp at $CN = 6$, although the distribution covers the range 1-6³¹² (Table 15). NMR data also show the hexahydration structure of Mg²⁺.^{313,318,319} The tilt angle of the hydrated water molecules θ is 14°.³⁰⁸

 Ca^{2+} . The hydration structure of the Ca^{2+} ion may be the 6-coordinated according to most X-ray diffraction data,^{62,101,107,113,117} although some other hydration numbers have been reported. However, molecular dynamics simulations show a much larger value for the hydration number of 9.2-9.6^{116,285} than that obtained by the diffraction methods. Since the hydration number of the first shell of Ca2+ is evaluated from the area under the peak ranged from $r = 0$ to $r_{1,\text{min}}$ in the radial distribution curve obtained by the MD calculation, the large hydration number may be obtained due to an asymmetric shape of the first peak of the radial distribution curve. The tilt angle of the hydrated water molecules is estimated to be 26-43°.²⁸⁵

 Sr^{2+} . The hydration number of Sr^{2+} may be larger than that of Ca^{2+} because of the larger ionic radius of the former than that of the latter. \check{X} -ray and neutron diffraction data^{103,115,119,120} give a value more than 6. A result by MD simulations also gives a value of 9.8.³⁰⁹ A value given by an NMR measurement 317 is even smaller than that of Ca^{2+} , which is probably due to an erroneous evaluation of the peak area of the NMR spectrum for such an ion with a fast exchange rate for the solvent substitution reaction.

The $Sr^{2+}-OH_2$ bond length has been determined to be 260-265 $pm^{103,119,120,309}$ by the diffraction and simulation methods.

Ba2+ . Since barium extensively absorbs X-rays, barium salt solutions for the structural studies by the X-ray diffraction methods are not favorable. Since Ba2+ has so many electrons, reliable functions for the cationanion, cation-water pair potentials are hard to obtain by molecular orbital calculations. Therefore, very few investigations have been done for the hydration structure of Ba2+ . An X-ray diffraction study shows us the $Ba^{2+}-H_2O$ bond length of 290 pm.¹⁰³ The hydration number evaluated by the diffraction work is 9.5 , 103 which accidentally agrees with an NMR data $(9.7).^{317}$

Table 7. Structural Parameters of the Hydration Shell of Oxyanions (XO_y²) **(See Section III for the Full Definitions of the Abbreviations) ~) Derived from Diffraction Measurements**

salt	H ₂ O/salt molar ratio	$r_{\text{XO(H}_2O)}$, pm	$\iota_{\text{XO(H}_2\text{O})},\, \text{pm}$	$n_{XO(H_2O)}$	symmetry	method	$_{\rm model}$	ref
			NO ₃					
HNO ₃	25	325		9	n	$\mathbf X$	FN ₃	67
NH ₄ NO ₃	4.4	350	31	3	n	$\mathbf X$	FN ₃	223
NaNO ₃ ^a	6.1,9.3	314,322	15,28	6	n	$\mathbf X$	FN ₃	89
NaNO ₃	7.1	340		1.3,2.4	n	${\bf N}$	fod	224
AgNO ₃	14.2	317	27	4.3	$\mathbf n$	$\mathbf X$	FN ₃	95
AgNO ₃	6.6.4.3	350,340	18,16	11,7,2	n	$\mathbf X$	FN ₃	97
$\text{Zn}(\text{NO}_3)_2$	9.0	344	19	17.7	n	$\mathbf X$	SNM	153
Cd(NO ₃) ₂	$9 - 54$	349	16	8.8	n	$\mathbf X$	SNM	225
Al(NO ₃) ₃	14.5	339	22	5.9	n	$\mathbf X$	SNM	174
Rh(NO ₃) ₃	61,145	350	16	8.2	n	$\mathbf X$	SNM	226
Ce(NO ₃) ₃	25.4	340	28	7.0	n	X	SNM	227
			ClO ₄					
NH ₄ ClO ₄	37	400	51	12.1	n	$\mathbf X$	FN ₃	69
NaClO ₄	17.1	360-380		$4 - 5$	n	${\bf N}$	fod	228
AgClO ₄	10.4	357	50	25.6	n	$\mathbf X$	FN ₃	95
AgClO ₄	16.3,3.3	375,370	32,28	7.5,7.9	n	$\mathbf X$	FN ₃	97
Mg(CIO ₄) ₂	18	374	31	10.1	$\mathbf n$	$\mathbf X$	SNM	110
Sn(ClO ₄) ₂	16.8	380	5	8	n	$\mathbf X$	SNM	164
Y(ClO ₄) ₃	40,12.4	380	30	8	n	$\mathbf X$	SNM	10
$Rh(CIO4)3 + HCl$	277	376	26	12	n	$\mathbf X$	SNM	66
La(CIO ₄) ₃	12.0	380	30	8	n	$\mathbf X$	SNM	10
$Ce(CIO4)3 + HClO4$	38	374	31	10.1	n	$\mathbf X$	SNM	110
Sm(ClO ₄) ₃	15.2	380	30	8	$\mathbf n$	$\mathbf X$	SNM	10 [°]
Tb(CIO ₄) ₃	39.4,13.7	380	30	8	n	$\mathbf X$	SNM	10
Er(CIO ₄) ₃	12.1	380	30	8	n	$\mathbf X$	SNM	10
			SO_4^{2-}					
$(NH_4)_2SO_4$	15.6	379	24	7.6	n	X	FN1	229
$(NH_4)_2SO_4$	15.8	393	30	11.2	n	$\mathbf X$	FN ₃	73
$(NH_4)_2SO_4$	14.3	370	20	8	n	$\mathbf X$	FN3	230
MgSO ₄	20.5	370	$\bf{22}$	7.7	n	$\mathbf X$	SNM	111
$MnSO4$ ^a	28	382	20	8	n	$\mathbf X$	SNM	141
N iSO ₄	27.5	381	29	8.1	n	$\mathbf X$	SNM	111
N iSO ₄ + H_2 SO ₄	48	380	25	7.6	n	$\mathbf X$	${\bf SNM}$	111
$NiSO_4^a$	28	392	10	9.6	n	$\mathbf X$	SNM	141
N iSO ₄	18,27	379	28	$\bf 8.2$	$\mathbf n$	$\mathbf X$	SNM	142
N iSO ₄ + Li ₂ SO ₄	37	380	20	8.2	n	$\mathbf X$	SNM	142
CuSO ₄	40.4	378	17	8.2	$\mathbf n$	$\mathbf X$	SNM	150
ZnSO ₄	18.6	383	27	8.2	$\mathbf n$	$\mathbf X$	SNM	231
ZnSO ₄	27.1	387	19.3	6.95	n	$\mathbf X$	SNM	154
ZnSO ₄	$100 - 16.9$	383-391	$17 - 28$	$14.3 - 7.6$	n	$\mathbf X$	SNM	155
CdSO ₄	17.6	374	28	12	n	$\mathbf X$	SNM	160
CdSO ₄	25	389,383	15,13	6.9,6.0	$\mathbf n$	$\mathbf X$	SNM	232
$Cr_2(SO_4)_3$	18.7	381	21	7.8	n	$\mathbf X$	SNM	233
$In_2(SO_4)_3$	14.2	389	23	6.4	n	X	SNM	234
			$SO32-$					
	14.3	380	36	9		$\mathbf X$	FN3	230
$(NH_4)_2SO_3$					n			
			$S_2O_5{}^{2-}$					
$(NH_4)_2S_2O_5$	6.1, 4.7	376,377	25	9.2,8.9	n	$\mathbf X$	FN3	230
			$\text{SeO}_4{}^{2-}$					
Y_2 (SeO ₄) ₃	69,58.2	395	30	8	n	X	SNM	10
$La_2(SeO_4)_3$	76.9,47.0	395	30	8	n	$\mathbf X$	SNM	10
$Tb_2(SeO_4)_3$	44.3	395	30	8	n	$\mathbf X$	SNM	10
$Er2(SeO4)3$	68.7,53.1	395	$30\,$	8	n	$\mathbf X$	SNM	$10\,$
			$CrO42-$					
Na ₂ CrO ₄	30.7	396	25	$12\,$	n	X, N	iso, fod	267
			$MoO42-$					
$Na_2(MoO_4)_3$	25	406	34	$12\,$	n	$\mathbf X$	iso	235
			$WO42-$					
$Na_2(WoO_4)_3$	$25\,$	406	34	12	n	$\mathbf X$	iso	235
			$PO43-$					
H_3PO_4	25.1,7.2	373,360	18,16	4 or 8	n	X	FN1	236
			$H_2PO_4^-$					
$Mg(H_2PO_4)_2$	34.0	387	33	8.8	n	X	SNM	237
$Ni(H_2PO_4)_2$	45.6	375	$\bf{22}$	4.2	n	$\mathbf X$	SNM	238
$Cd(H_2PO_4)_2$	41.7	391	28	4.4	n	X	SNM	238
			$CH_3CO_2^-$					
$MgCH_3CO_2)_2$	33.1,14.9	360,370	17,14	4.2, 6.1	n	X	SNM	239
$Mn(CH_3CO_2)_2$	33.5	363	19	5.3	n	$\mathbf X$	SNM	239

Divalent Transition Metal Ions in the First Row. For a long time, the divalent transition metal ions have attracted much attention of traditional coordination chemists, and the ions are still being extensively studied by coordination chemists. The hydration structure of the ions are believed to be regular octahedral except for the Cr^{2+} and Cu^{2+} ions which have d^4 and d^9 electron configurations, respectively. These two ions should have a distorted octahedral structure with either elongated or shortened axial bonds perpendicular to the square plane. Investigations for the hydration structure of the divalent transition metal ions have been done by many authors using the visible spectroscopic method, and the octahedral structure predicted by the ligand field theory has been well accepted. However, information about the length of the $M^{2+}-OH_2$ bonds have been scarcely available, because crystals containing hexahydrated metal ions are hardly obtainable and some of the water molecules in the coordination sphere are often replaced with other ligands. Since the water molecules at the apexes of the elongated axis in the distorted octahedral structure of Cr^{2+} and Cu^{2+} are easily replaced with other ligands, no reliable data for the metal-water(ax) bond length have been reported. The diffraction and EXAFS methods have an advantage of determining the bond length of the elongated axis, in principle, provided that the hydrated ions with a distorted octahedron exist in the solution. In the following we will discuss the hydration structure of divalent transition metal ions of the first row in the periodic table.

Ti2+ . The titanium(II) ion can exist in water, but due to its very strong reducing properties, it is extremely difficult to handle as solution samples containing pure Ti(II) salts. Therefore, no structural data has been reported for Ti²⁺.

 V^{2+} . The V^{2+} ion is also a strong reducing agent and hard to use as sample solutions for diffraction studies. An attempt appears to have been made by the EXAFS method according to ref 344, in which the $V^{2+}-H_2O$ bond length is reported. However, no detailed information, including the composition of the salt used, is given in the literature. Therefore, the structural data given in ref 344 may not be very reliable.

 Cr^{2+} . The Cr^{2+} ion has 4 d electrons, and thus, the ion is expected to show the Jahn-Teller distortion when the ion has an octahedral structure with a high-spin electron configuration. The structure of the hydrated Cr2+ ion is studied by the EXAFS method, and the $Cr^{2+}-OH_2$ (equatorial) is reported to be 208 pm.³⁴⁴ The $Cr^{2+}-H_2O(axial)$ bonds are not observed in the experiment, although the authors expected the distorted octahedral structure for the Cr^{2+} ion.

Mn2+ to Zn2+ . The structure of hydrated ions from Mn^{2+} to Zn^{2+} has thoroughly been studied by various experimental and theoretical methods. Accumulated experimental data for the static structure of the ions in water have been reviewed by $Marcus^6$ and they are

also summarized in Table 2. AU the ions have six water molecules in the first coordination sphere, and the M^{2+} - $H₂O$ bond length varies from 220 pm in $Mn^{2+121,123}$ to $208 - 210$ pm in $\text{Zn}^{2+101,121,154,156-158}$ passing through the shortest $M^{2+}-OH_2$ bonds in the square plane of the $[Cu(H₂O)₆]$ ²⁺ complex having a distorted octahedral structure. The lengths of the $M^{2+}-OH_2$ bond, as well as the dissociation energies of the hydrated water molecules, have recently been calculated by using the local spin density function method for ions from Cr^{2+} to $\sum n^{2+}$, 345 According to the results the bond lengths and the number of the bonds are as follows: Cr^{2+} [2 \times 202.5 pm + 2 \times 202.8 pm + 2 \times 226.8 pm], Mn²⁺ [6 \times 212.1 pm], Fe²⁺ $[2 \times 207.0$ pm + 2 \times 207.2 pm + 2 \times 207.5 pm], Co^{2+} [$2 \times 202.8 \text{ pm} + 2 \times 203.8 \text{ pm} + 2 \times$ 205.6 pm], Ni²⁺ [$6 \times 200.2 \text{ pm}$], Cu²⁺ [$2 \times 197.3 \text{ pm}$ + $2 \times 199.0 \text{ pm} + 2 \times 219.3 \text{ pm}$, and Zn^{2+} [6 \times 205.7 pm]. Since the calculations have been done for the isolated $[M(H₂O)₆]^{2+}$ ions, the negative end of the dipoles of the hydrated water molecules orient toward the central ions, which do not fit to the experimental results found by the neutron diffraction measurements^{125,131,133,136,146,240} (Table 8). The trend in the bond length variation of the ions corresponds to the change in the hydration energies of the ions, which is known as the Irving-Williams series and has been explained in terms of the change of the ligand field stabilization energies.¹²¹ The variation of the dissociation energy of a water molecule from the hydration sphere of the ions thus calculated is similar to the experimental²⁰³ and calculated^{121,346} values reported in previous studies. Results for the $M^{2+}-H_2O$ bond lengths obtained by simulation calculations slightly disagree with experimental results (Table 8), due probably to incomplete potential functions to describe ion-water interactions. The NMR method is not quite suitable for studies of the hydration μ is not quite suitable for studies of the hydration
structure of the ions, except for Zn^{2+} due to the paramagnetic property of the ions.

Cd2+ and Hg2+ . Cadmium(II) and mercury(II) ions have the regular octahedral structure and the M^{2+} - H_2O bond lengths are 229 pm^{101,160} and 233 pm,¹⁶³ respectively. NMR data³¹⁷ (Table 11) give smaller coordination numbers than those found by the diffraction methods, due probably to fast water substitution reaction rates of the ions. No simulation calculation has been undertaken for such ions having a large number of electrons.

 Sn^{2+} . Sn²⁺ is a peculiar ion, having a less symmetrical hydration structure than other d¹⁰ ions. In the 6-coordinated hydration structure, two kinds of bonds have been observed, one has the bond length of 233-234 pm and the other of 238–290 pm.^{164,165} The asymmetry may arise from the electronic configuration of the ion having $5p^2$ electrons, but since no other ions in the 14 group has been investigated for their hydration structures, we cannot draw a reliable conclusion for the origin of the asymmetric hydration structure of the ion.

6.0

317.5

5.4

222.5

5.4

n

n

fod

207

Table 8. Structural Parameters of the Orientation of Water Molecules in the First Hydration Shell of Ions Derived from Neutron Diffraction Measurements (See Section III for the Full Definitions of the Abbreviations)

Table 8. (Continued)

^{*a*} H denotes hydrogen or deuterium. *b 8* is used for cations, ψ for anions (see below). C The X-ray diffraction method is also applied. *d* intervalses the method is also applied. invar = orientation of water molecules is invariant and independent of concentrations of the salt. ϵ In H₂O. ^{*I*} In D₂O. ϵ The Jahn-Teller effect is not observed. * Three different isotopic mixtures.' The second sample is heated to 90 ⁰C during the preparation.*'* Full width of half height (FWHH) values. * bifurc = water molecules combine with the ion in a bifurcated way. ^Tundet = bond angle is not determined. ^{*m*} In the glassy state at -148 °C.

Table 9. Structural Parameters of the Incompletely Hydrated First Shell of Metal Complexes $[M^I A_m(H_2O)_n]^s$ of **Monovalent Cations Derived from Diffraction and EXAFS Measurements (See Section III for the Full Definitions of the Abbreviations)**

salt	H ₂ O/salt molar ratio	r_{MO} , pm	l_{MO} , pm	n_{MO}	$r_{\rm MA}$, pm	$l_{\rm MA}$, pm	$n_{\rm MA}$	symmetry	second shell	method model		ref
					$\rm Li^+$							
LiCl	2.0	195		2.3	276		1.5	n	n	N	fod	241
LICI	3.35	195		2.3	275			n	$\mathbf n$	N	fod	242
LiCl	3.0	200		4	280		1	$\overline{0}$	$\mathbf n$	X, MD	FN3	243
					$Na+$							
NaCl	9.0	241	23	4.6	282	30	0.3	$\mathbf n$	$\mathbf n$	X	FN3	244
					K^+							
KC _I	12.2	281	13	5.8	315	15	0.6	n	n	x	FN3	244
$KCl + MgCl2a$	12.3,73.6	277	18	3.7	320	17	2.4	n	n	X	FN2	108
KI	8.4	290		1.7	370		$\mathbf 1$	n	n	$\mathbf X$	rdf	93
KF	9.0	280	17	3.3	269	26	$2.3\,$	n	n	X	FN ₃	244
					$Cs+$							
CsCl	$5.5 - 22.2$	315		$1.9 - 6.2$	330-360		$2.7 - 3.5$	n	n	X	rdf	245
$CsCl + MgCl2a$	13.2,41.5	315	29	4.7	339	23	2.0	n	n	X	FN2	108
CsI	22.2	315		$8.8 - 8.9$	390-410		$1.4 - 2.4$	n	$\mathbf n$	$\mathbf x$	rdf	245
CsBr	11.1,22.2	315		$7.2 - 8.9$	350-390		$2.0 - 2.5$	n	n	X	rdf	245
$\operatorname*{CsI}\nolimits_{\mathbf{C}\mathbf{s}\mathbf{I}^{b}}$	20	301	18	5.8	388	35	0.8	n	n	X	FN3	83
	20,10	306,302	21,15	4.7,3.0	385,385	36,48	0.8, 1.2	n	n	$\mathbf x$	FN3	83
CsF	1.7	321	25	$3.6\,$	312	33	3.3	n	n	X	FN3	244
					$Ag+$							
AgNO ₃	$5 - 16.8$	236-231	$6 - 10$	$1.9 - 2.9$	315-326		1	n	n	EX	FN3	98
AgCIO ₄	3.3	243	16	3	352	20	$\mathbf{1}$	$\mathbf T$	a	$\mathbf X$	SNM	97
AgNO ₃	4.3	243	14	3	313	21		T	a	$\mathbf X$	SNM	97
^{<i>a</i>} The Mg ²⁺ ions do not form ion pairs. ^{<i>b</i>} At 70 °C.												

 Pb^{2+} . The Pb^{2+} ion is not suitable for study by the diffraction method, because it extensively absorbs X-rays. The NMR method may be more adequate because of its diamagnetic property. According to the NMR measurement, ${}^{317}Pb^{2+}$ has six water molecules in the first coordination sphere as expected.

3. Trivalent Cations

Al3+ . The Al3+ ion is an inert ion and is one of the ions with the strongest ion-water interactions because of its high charge and small ionic radius. The structure of the hydrated Al3+ ion is a 6-coordinated one with the $Al^{3+}-H_2O$ bond length of 187-190 pm.^{169,173,174} Since

the ion is very inert, splitted NMR signals of ¹⁷O in H₂¹⁷O water are observed, which are the signal from the bulk water in the lower field and that from the hydrated water in the high field, the two peaks being separated by 2330 Hz (430 ppm).³¹⁵ From the area under the peaks, the hydration number of Al^{3+} is calculated to be 6. In most NMR measurements it is difficult to observe separate peaks of the bulk and the hydrated water, but the Al^{3+} case is a typical example for the NMR determination of the hydration structure of ions.

 Cr^{3+} . The Cr^{3+} ion is an extremely inert ion. The hydration structure of the ion is regular octahedral with

Table 10. Structural Parameters of the Incompletely Hydrated First Shell of Metal Complexes $[M^H A_m(H_2O)_n]^2$ of **Divalent Cations Derived from Diffraction and EXAFS Measurements (See Section III for the Full Definitions of the Abbreviations)**

salt	H_2O /salt molar ratio	r_{MO} , pm	l_{MO} , pm	n_{MO}	r_{MA}, pm	l_{MA} pm	$n_{\rm MA}$	symmetry	second shell	method model ref		
$Mg(H_2PO_4)_2$ $MgCH_3CO_2$ ₂	34 33.1,14.9	211 209,219	12.6 8.8,14.1	6 5.2,5.2	Mg^{2+}		1.0 0.8, 0.8	0 o	a a	X X	SNM SNM	237 239
CaCl ₂ ^a $CaCl2$ ^b CaCl ₂ ^c CaCl ₂ ^d CaCl ₂ e	6.2 6.0 6.2 5.6 4.0	244 245 245 244 246		5.9 5.6 5.3 4.8 3.9	Ca^{2+} 275 274 274 275 280		0.9 1.0 1.4 1.7 2.1	n n n n n	a a a a a	X $\mathbf X$ X $\mathbf X$ X	FN3 FN3 FN3 FN ₃ FN ₃	117 117 117 117 117
MnCl ₂ MnCl ₂ MnBr ₂ Mn(NO ₃) ₂ MnSO ₄ MnSO ₄ $Mn(CH_3CO_2)_2$	9.7.17.3 550-11 8.3,16.2 12.6 13.2,26.8 28 33.5	218.4 220 218.1 219 222,220 220 220	10 11 9. 13 10,12 10 11	4.65,4.45 6-5.3 5.0, 4.8 5.4 5,2,5,1 5.5 5.2	$\mathbf{M} \mathbf{n}^{2+}$ 250,250 249 264,260 350	10,10 31 10,10 12 10.13 11 19	1.3, 1.5 1.1, 1.2 0.6 0.75,0.92 $0.5\,$ 0.8	o o o o o o o	n a n a a a a	x EХ X X $\mathbf X$ X X	FN3 SNM FN3 SNM SNM SNM SNM	246 247 246 248 249 141 239
$_{\rm{FeCl_2}}$ FeBr ₂ FeSO ₄	55.5 17.8.10.3 55.5	224 212 228	8 9,7 8	5.6 5.5, 5.1 5.8	\mathbf{Fe}^{2+} 258 260,262 349	18 11,12 26	0.4 0.3, 0.7 0.2	0 n 0	n n n	\mathbf{X}^f X X/	SNM FN ₃ SNM	126 250 126
CoCl ₂ $CoCl2 + HCl$ $CoCl2 + LiCl$ $CoCl2 + LiCl$ $CoCl2 + LiCl$ $CoCl2 + NiCl2 +$	17.9 23.2 7.5 23.1 8–17 110	210 214 210 214 208-209 210,211	10 11 12 11 $10 - 12$ 10	5.5 5 5.5 5.0 5 5.9,5.7	$Co2+$ 247 235 249 235 $245 - 242$ 239,229	11 11 16 11 13–12 14	0.5 1 1 ı 1 0.8, 1.2	0 o o 0 o 0, n	a n n n a a	$\mathbf X$ $\mathbf X$ $\mathbf X$ $\mathbf x$ X X, EX	SNM FN ₃ FN ₃ FN ₃ SNM SNM	251 65 182 65 81 252
6LiCl CoBr ₂ $Co(CH_3CO_2)_2$	18.2,10.8 53.7	210 214	8 11	6,5.3 5.2	258	12 21	0.3, 0.6 0.8	n o	n a	$\mathbf X$ X	FN3 SNM	250 239
NiCl_2 $\rm NICl_2$ $NiCl2 + HClg$ NiCl_2 + 2LiCl ^h $NiCl2 + 3LiCl7$ NiCl_2 + 4LiCl $NICl2 + 6LiCl$ $\text{NiCl}_2 + \text{CoCl}_2 +$	18 25 24.6 24.8 23.2 25 25 110	207 206 205.6 206.4 207.2 205 205 210,206	9 7 8 9 10	5.5 6.5 5.1 4.8 4.6 6.0 5.3 5.9,5.2	$Ni2+$ 244 247 243 244 238 238 239,229	10 14 13 12 $14,-$	0.5 0 0.7 1.0 0.81,0.67	o $\mathbf n$ 0 o 0 n n o	a n a a a n n a	X EX X $\mathbf X$ x EX EX X, EX	SNM FN3 SNM SNM SNM FN3 FN3 SNM	251 253 254 254 254 253 253 252
6LiCl NiBr ₂ NIBr ₂ NIBr ₂ $Ni(en)_2(NO_3)_2$ Ni(gly) ₂ $Ni(H_2PO_4)$	$9.6 - 16.5$ 24.7 11.9,24.7 22.9 15.7 45.6	$204 - 205$ 206.5 207.9,206.6 210 208 203	$8 - 9$ 11 10,8 7.7 10 11	$5.7 - 5.4$ 5.7 5.6,5.7 2 4 5.25	258–252 262 262,261 210 209(O), 209(N) 11, 11	13–12 9 13,13 7.4 10	$0.2 - 0.7$ 0.3 0.44, 0.29 4.0 1(0),1(N) 0.75	o o o 0 $\mathbf 0$ $\mathbf O$	n a a n n a	x X X X $\mathbf X$ $\mathbf X$	SNM SNM SNM FN3 $\mathbf{FN}3$ SNM	214 215 216 268 269 238
CuCl ₂ ^k	10.9	195		2.4 $\overline{}$	Cu^{2+j} l 243		3.6	$\mathbf 0$	n	$\mathbf x$	rdf	255
CuCl ₂ ^k	15.7	195		2.7 $\overline{}$	- 243		3.3	$\mathbf 0$	n	$\mathbf X$	rdf	255
CuCl ₂ ^l	18	195 263	$\{^{8}_{16}$	4 $\,2\,$	225 \rightarrow	$\begin{pmatrix} 8 \\ - \end{pmatrix}$	1.2	\mathbf{o}	n	$\mathbf X$	FN1	251
CuCl ₂	11.6	205 250		$2.3\,$ \overline{a}	256 l –		4.2 or 2.8	\mathbf{o}	n	N	fod	256
CuCl ₂ $CuCl2 + nNH3m$	12.8 $18.9 - 20.5$ \sim	196 (233	16 $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$	3.4 $\begin{smallmatrix} 1 & 1 \ 2 & \end{smallmatrix}$				$\mathbf 0$ O	n n	N X	fod FN1	257 258
CuBr ₂ CuBr ₂	>1000 12,100	197 193		3.5 ι _	239		$\{\bar{}}_{1.5}$	n $\mathbf 0$	n n	EX \mathbf{EX}	rdf FN3	217 259
CuBr ₂	>100	$\qquad \qquad -$ 193		$\boldsymbol{7}$ \overline{a}	235		$\frac{-}{0.5}$	n	n	EX	FN3	260
CuBr ₂ CuBr ₂	sat. $10.8 - 53.1$ ($196 - 199$	193 l 237–251		3.5 $2.5 - 3.7$ $\boldsymbol{2}$	(240) $246 - 241$		$1.3 - 0.3$	H o	n n	EX X	FN3 FN3	260 261
Cu(NO ₃) ₂	23.5,34.7 (200		13	4 \ շ	303,307	14,11	1.2, 1.7	$\mathbf 0$	a	X	SNM	262
$Cu(en)_2(NO_3)_2$	28.1	l 212,222 292)	20	$\big\{ \frac{-}{1.8}$	194,194		7.5,7.5 3.8,3.8	$\mathbf 0$	$\mathbf n$	X	FN3	270

a At 15 °C. ^{*b*} At 33 °C. ^{*c*} At 80 °C. ^{*d*} At 72 °C. *^e* At 120 °C. *i* Mössbauer spectra of frozen solutions are also investigated. ^{*g*} 29% $\rm [Ni(H_2O)_6]^{2+}$, 50% $\rm [Ni(H_2O)_5Cl]^+$, 21% $\rm [Ni(H_2O)_4Cl_2]$ complexes are present. ^h 17% $\rm [Ni(H_2O)_6]^{2+}$, 50% $\rm [Ni(H_2O)_5Cl]^+$, 33% $\rm [Ni(H_2O)_4Cl_2]$ complexes are present. 1 8% $\rm [Ni(H_2O)_6]^{2+}$, 40% $\rm [Ni(H_2O)_5Cl]^+$, 52% $\rm [Ni(H_2O)_4Cl_2]$ complexes are present. j Parameters in the first and second lines refer to equatorial and axial bonds, respectively. * Complexes share ligands in the first hydration shell of the cations. ${}^{\prime}40\%$ [Cu(H₂O)_{4eq}(H₂O)_{2ax}]²⁺, 60% [Cu(H₂O)_{2eq}Cl_{2eq}(H₂O)_{2ax}] complexes are present. ${}^{\pi}$ π = 4.8,5,11.2; [Cu(NH₃)₄(OH₂)₂]²⁺, (Cu(NH₃)₄(OH₂)₂]²⁺, and [Cu(N P 25% $[Zn(H_2O)_6]^{2+}$, 25% $[Zn(H_2O)_2Br_2]$, 35% $[Zn(H_2O)Br_3]$, and 15% $[ZnBr_4]^{2-}$ complexes are present. ^q The MO parameters refer to those of the hexaaqua complex. The data given in this table are the averaged one for the mono-, di-, tri-, and tetrahalogeno complexes. 'Parameters refer to those measured at -40, 25, 40, 80, and 100 °C. 'Supercooled at -5 °C. 'In the glassy state at -20 °C.
"60% [Zn(H₂O)₆]²⁺, 40% [Zn(H₂O)₅SO₄] complexes are probably formed. 'At 9 °

the $Cr^{3+}-H_2O$ bond of about 198–200 pm^{166–172} (Table 3).

 Fe^{3+} . The Fe³⁺ ion has also rather small rate constant for the water substitution reaction due to its strong interaction with hydrated water molecules. The hydration structure of the ion is octahedral with the $Fe^{3+}-$ H₂O bond length of 201-205 pm,^{125,175-179} which is slightly longer than the $Cr^{3+}-H_2O$ bond. The lengths of the $Fe^{3+}-H_2O$ and $Cr^{3+}-H_2O$ bonds are slightly shorter than the sum of the radii of the ions $(Cr^{3+}, 76)$ pm; Fe^{3+} (high spin), 79 pm³⁴¹) and water (140 pm). The hydration energy of $Fe³⁺$ should be larger than that of Cr^{3+} , although the data so far reported do not fit to this consideration.203,347 MD simulations also show the hexahydrated structure of Fe³⁺ with the Fe³⁺-H₂O bond length of 203-205 pm,³¹⁰ which fairly well agrees with the experimental result. The tilt angle of the Fe³⁺⁻ O-H bond is 7.7° according to MD simulations.³¹⁰

Rh3+ . The rhodium(III) ion is an extremely inert ion in water, and the hydrated Rh3+ ions cannot be obtained by simple dissolution of rhodium(III) salts in water and must be prepared by a special synthetic way. The hydrated ion is 6-coordinated with the $Rh^{3+}-H_2O$ bond

length of 204–206 pm,^{66,68} which is much shorter than the sum of the ionic radius of Rh^{3+} (82 pm³⁴¹) and that of water. The result indicates that the $\bar{R}h^{3+}-H_2O$ bond has an extensively covalent character.

 In^{3+} and T^{13+} . The In^{3+} ion has an octahedral structure with the $In^{3+}-H_2O$ bond length of 215 pm, which is much shorter than the sum of the sizes reported for In³⁺ (94 pm³⁴¹) and water. NMR data also show the 6-coordinated structure of the hydrated In³⁺ ion. The $T13+$ ion has also an octahedral structure and the $T13+$ -H2O bond length is 223 pm.¹⁸²¹⁸³ Again the measured bond length is much shorter than the sum of the sizes of the TI^{3+} ion (104 pm^{341}) and water molecule by 20 pm, the result indicating a significant covalent-type bonding between the ions and water molecules.

Y 3+ *and Lanthanoid Ions, Ln3+ .* The yttrium(III) ion is often compared with the lanthanoid(III) ions in their reactions and structures in solution because of their similar ionic charge and size. In fact, the isomorphous substitution method can be applied to the analysis of the hydration structure of the ions by using the X-ray diffraction method.^{10,180} The Y^{3+} ion is octahydrated, according to the X-ray diffraction meth-

Table 11. **Structural Parameters of Incompletely Hydrated First Shell of Metal Complexes [M¹¹¹An(H2O)1,]' and [M,vAjn(H2O)n]* of Tri- and Tetravalent Cations, Respectively, Derived from Diffraction and EXAFS Measurements (See Section III for the Full Definitions of the Abbreviations)**

	H ₂ O/salt molar								second			
salt	ratio	r_{MO} , pm	l_{MO} , pm	n_{MO}	r_{MA} , pm	$l_{\rm MA}$, pm	$n_{\rm MA}$	symmetry	shell	method model ref		
					Cr^{3+}							
CrCl ₃ ^a	17.9,26.5	196,198	25,21	4.5,4.9	230,232	10,9	1.5.1.1	0	a	X	SNM	209
$CrCl3 + HClb$	24.5	197	23	5.35	233	9	0.65	$\mathbf 0$	a	X	SNM	209
$Cr_2(SO_4)_3$	18.7	194	10	5.17	318	12	0.83	$\mathbf 0$	a	X	SNM	233
					\mathbf{Fe}^{3+}							
FeCl ₃		207		4	220		2	$\mathbf 0$	n	x	rdf	274
\mathbf{FeCl}_{3}^{c}	$6 - 21.5$	$194 - 208$	$12 - 15$	$1.8 - 4.2$	230-237	$13 - 17$	$3.2 - 1.8$	$\mathbf 0$	a	X	SNM	210
$\text{FeCl}_3{}^d$	$12 - 6.7$	$201 - 203$	$9 - 12$	$1.7 - 2.2$	$221 - 235$	$5 - 10$	$3.2 - 3.4$	0, T	n	$\mathbf x$	FN3	211
FeCl ₃	55	210		6	213			0	n	EX	rdf	124
$FeCl3 + NaOHe$	55	199-201	$7 - 10$		$4.6 - 6.1$ 230-231 7-5		$1.8 - 0.3$	$\mathbf n$	n	EX	rdf	275
$Fe(NO3)3 + NaOHe$	55	$192 - 208$	$5.7 - 2.8$	$10 - 7$				n	n	EX	rdf	275
$Fe2(SO4)3$	$14.4 - 23.8$	$201 - 202$	$7 - 9$	6			$1.2 - 1.0$	$\mathbf 0$	i	X.	FN3	276
					\mathbf{V}^{3+}							
Y_2 (SeO ₄) ₃	69,58.2	233	10	8.0	375	8,12	0.35,0.60	n	a	x	SNM 10	
					Rh^{3+}							
$RhCl3 + HCl$	141	209	13	3	233	5	3	$\mathbf 0$	a	$\mathbf X$	SNM	66
Rh(NO ₃) ₃	61,145	202,203	10,6	4		24	$\overline{2}$	Ω	a	X	SNM	226
					In^{3+}							
InCl ₃	8.8	232	6	5	250	6	$\bf 2$	S^g	a	X	SNM	212
InBr ₃	9.2	260	6	6	260	6	$\overline{2}$	S^g	a	X	SNM	212
$In_2(SO_4)_3$	58.9	216	11	5.12	334	12	0.88	Ω	a	X	SNM	181
					Ln^{3+}							
$La2(SeO4)3$	76.9	256.0	11	8.0	393	12	0.60	$\mathbf n$	n	x	SNM	10
Ce(NO ₃) ₃	25.4	255	12	7.5	344	28	0.9	n	a	x	SNM	227
$Ce(CIO4)3 + HClO4$	38	257.4	14.6	7.5	359	18	2.09	$\mathbf 0$	a	X	SNM	110
$NdCl3 + 3HCl$	30	241		7.0	278		1	н	n	$\mathbf X$	rdf	188
$GdCl3 + 3HCl$	30	237		6.0	280		$\overline{2}$	$\mathbf H$	n	$\mathbf X$	rdf	195
$\text{Tb}_2(\text{SeO}_4)_3$	44.3	238	11	8.0	378	12	0.6	n	n	X	SNM	10
ErCl ₃	21.5	233	8.4	7.8	265	10	0.25	n	a	$\mathbf X$	SNM	11
$\text{ErCl}_3 + 7 \text{LiCl}$	55.3	235	7.7	7.8	267	11	0.8	n	a	$\mathbf x$	SNM	11
$ErBr3 + 7LiBr$	52.4	235	7	7.9	287	11	0.3	n	a	x	SNM	11
$Er(NO_3)_3^h$	$52.4 - 14.9$	$232 - 237$	6	$5.8 - 3.3$	286,278	14.8	$0.9 - 2.2$	n	a	X	SNM	277
$Er2(SeO4)3$	68.7,53.1	234	10	8.0	375	8,12	0.35,0.6	n	n	X	SNM	10
					$Th4+$							
$\rm ThCl_4$	30	2.48	10	7.0	284	22	1.6	n	n	X	FN3	201
$\text{Th}(\text{NO}_3)_4$ ^h	$49.3 - 17$	$245 - 251$	$4 - 10$	$3.0 - 7.4$	$300 - 310$	10	$3.5 - 1.0$	n	i	X	FN ₃	201

^a 20–25% [Cr(H₂O)₆]³⁺, 37–50% [Cr(H₂O)₆Cl]²⁺, 38% [Cr(H₂O)₄Cl₂]-, 30% [Cr(H₂O)₃Cl₃] complexes are present. ^b 35% [Cr(H₂O)₆]³⁺, 65% [Cr(H₂O)₅Cl₂]⁺ complexes are present. ^c Beside [FeCl4]- complexes also coexisted. *^d* Only the [Fe(H2O)3Cl3] and [FeCl4]- complexes were present.^e Extended polymerization occurs due to hydrolysis. / Hydrated cations are bridged through anions. «Bipyramidal structure with the anions in axial positions. * Bidentate nitrato complexes are formed.

od,^{10,180} like the heavier lanthanoid ions are. The Y^{3+} -H2O bond length is determined to be 236-237 pm. The hydration number of the Y^{3+} ion measured by NMR is 2.4,³²⁸ which is fairly small compared with the value determined by the diffraction method. The difference may be too large to attribute to experimental uncertainties in both X-ray diffraction and NMR measurements and suggests that two water molecules in the $\frac{1}{2}$ coordination sphere of Y^{3+} may be substituted much more slowly than the other six.

The hydration structure of the lanthanoid(III) ions has been extensively discussed by many authors after the works by Habenshuss and Spedding,^{185,193,196} and the discussions seem to converge to the conclusion that the lighter lanthanoid(III) ions from La^{3+} to Nd^{3+} are nonahydrated and the heavier ions after Gd^{3+} or Tb^{3+} have eight water molecules in the first coordination sphere, and intermediate ions such as Sm^{3+} , Eu^{3+} (and $probability Gd^{3+}$, too) have both nona- and octahydrated structures which are in equilibrium. The conclusion has been supported from the Raman spectroscopic measurements³⁴⁸ and EXAFS measurements of the hydrated structure of the ions in water and in the aqueous glassy state.¹⁹¹

The octa- and nonahydrated structures may be depicted in the way that six water molecules are arranged in a prism or antiprism shape around the ion and the rest two to three water molecules, depending on the sizes of the $Ln³⁺$ ions, are located at the middle of the prism structure. If we imagine such a structure, slowly exchangeable water molecules found in Y^{3+} by the NMR method may be explainable. In fact, such a small hydration number is observed by the NMR method for Er^{3+} $(n = 1)$, which is one of the heavier lanthanoid(III) ions with eight water molecules in the first coordination shell. On the other hand, the hydration number of lighter lanthanoid(III) ions determined by the NMR method is much larger than those found for Y^{3+} and Er^{3+} (Table 11), and the results suggest that the water molecules in the middle part of the prismatically (or antiprismatically) 6-coordinated structure of a lighter lanthanoid(III) ion may be more weakly bound to the central ion than those of the heavier lanthanoid(III) ions so that the six water molecules at

Table 12. Structural Parameters of the Hydration Shell of Monovalent Cations Derived from Radial Distribution Functions Obtained by Computer Simulations (See Section III for the Full Definitions of the Abbreviations)

	H_2O /salt								
salt	molar ratio					θ , deg ^a	water model	method	ref
		r_{MO} , pm	$n_{\rm MO}$	$r_{\rm MH},$ pm	$n_{\rm MH}$				
$NH4$ +	215	290	14	$NH4$ ⁺			MCY	MC	278
NH ₄ Cl	25	305	8.1	350		15.7	ST2	MD	279
NH ₄ Cl	25	306	8.1				ST2	MD, X	71
				$Li+$					
$Li+$	200	195	5	259			CI	MC	57
$Li+$	64	190					HF	MC	280
$Li+$	215	210	6	270		$20 - 40$	MCY	MС	281
$Li+$	64,125	198	5.3	257	10.6	40	MCY	MD	282
$Li+$	79	192	4	262	8	55.4	CF	MD	283
$Li+$	215	213	6 4.9				MCY TIP ₄ P	MC MC	284
$Li+$ $Li+$	125 64	195 204	6	260 260	10.6 >12	65.8 45	TIP4P	MD	60 285
LiF	50	200	4				HF	MC	286
LiF	100	200	$\bf 5$	260			HF	MC	287
LiF ^b	31	200	\leq 3	350		45	SPC-FP	мс	61
LiCl	10	190-200	4				HF	cluster	58
LICI	100	200	5	260			HF	MC	288
LiCl	25	206	7.1			49	ST2 CF	MD MD	289
LiCl LiCl ^b	4 4	211 200	5.3 5.2	280 270	10.7	30 41	BJH	MD	79 290
LiCl ^b	3	200	4	280	8	0	BJH	MD, X	243
LiI	25	210	7.3			46.4	ST2	MD	289
LiI	25	213	6.1	268	13.1	42.1	ST2	MD	291
LiI	25	213 $\ddot{}$	6.1	267	47		ST2	MD	292
LiI ^c	100	213	6.0				ST2	MD	293
LiI ^d	25	212	6.1	268	15.4		ST2	MD	294
				$Na+$					
$Na+$	10	$230 - 240$	$5 - 6$				HF	cluster	58
Na^+ ^e	200	233	5.4	300			$_{\rm CI}$	MC	57
$Na+$ $Na+$	64 200	230 240					HF	MC MC	280 295
$Na+$	215	235	6	289		$20 - 40$	MCY	MC	281
$Na+$	64,125	229	6.0	295	12	20	MCY	MD	282
Na^+ ^e	79	224	6.0	292	12	55.9	CF	MD	283
$Na+$	215	237	6.0				MCY	MC	284
$Na+$	125	233	6.0	290	13.2	62.4	TIP4P	MC	60
$Na+$	64	235	6.2	291	>12.4	43	TIP4P	MD	285
NaCl NaCl	25 100	231 232	6.6 7.3			46.4	ST2 ST2	MD MD	289 296
NaCl	25	236	6.0				ST2	MD	297
NaCl	25	230	5,9				CF	MD	298
NaCl	100	230	6.0				MCY	MC	299
NaCl	25	230	5.8	295	13.9	37.5	BJH	MD	300
NaCl	$25\,$	230	6.3	293	16.2	42.2	BJH	MD	300
NaCl NaClO ₄	$100 - 7.9$ 25	$231 - 230$ 236	$6.1 - 4.9$ 6.5				MCY ST2	MC MD	301 302
				\mathbf{K}^+					
$\rm K^+$ K^{+e}	10 200	285 283	$5 - 7$ 6.8				HF CI	cluster MC	58 $57\,$
K^+	64	280					HF	MC	280
K^+	200	265						MC	295
\mathbf{K}^+	215	271	6.3	319		$20 - 40$	MCY	MC	281
K^+	64,125	276	7.5	335	15	36	MCY	MD	282
\mathbf{K}^{+} e	79	276	7.0	344	14	54.4	CF	MD	283
\mathbf{K}^+	215	286	6.6				MCY	MC	284
$\rm K^+$ KC1	64 25	286 280	7.6 7.8	332 332	>15.2 23.4	55 45.5	TIP4P ST2	MD MD	285 303
CsF	$\bf 25$	310	7.3	$Cs+$		56.6	ST2	MD	289
CsF	25	320	7.9	372	23	50.2	ST2	MD	304
CsCl	25	310	8.2			56.6	ST2	MD	289
$\mathbf{C}\mathbf{s}\mathbf{I}^b$	20	303	5.8				CF	MD	305
$CsI^{b,g}$	$20\,$	303	5.3				CF	MD	305
$CsI^{b,h}$	10	305	6.0				CF	MD	305

^a The orientational angle of the M–O–H bond (see Figure in Table 8). ⁵ Ion-pair formation occurs. ^c At 235 °C and 3 kbar pressure.
^d The solution is placed between two Lennard-Jones walls of 1231-pm separation. ^e kbar pressure. ^{*s*} At 68 °C. ^{*h*} At 76 °C.

Table 13. Structural Parameters of Hydration Shells of Di- and Trivalent Cations Derived from Radial Distribution Functions Obtained by Computer Simulations (See Section III for the Full Definitions of the Abbreviations)

	H ₂ O/salt								
salt	molar ratio	r_{MO} , pm	n_{MO}	r_{MH} , pm	$n_{\rm MH}$	θ , deg ^a	water model	method	ref
				$Be2+$					
BeCl ₂	50	175	6	249	12.7		BJH	MD	100
BeCl ₂	50	175	4	252			BJH	MD	306
				\mathbf{Mg}^{2+}					
$Mg^{2+} + 37NH_3$	164	205	4.0	270	8.1	0	MCY	MC	307
MgCl ₂	50	200	6.0	275	12.0	14	CF	MD	308
				Ca^{2+}					
Ca^{2+}	64	254	9.3	313	>18.6	35	TIP ₄ P	MD	285
$Ca^{2+ b}$	64	254	9.6	307	>19.2	43	TIP ₄ P	MD	285
$Ca^{2+ c }$	64	254	9.3	317	>18.6	26	TIP4P	MD	285
CaCl ₂	50	239	9.2	313	18.7	0	CF	MD, X	116
				Sr^{2+}					
SrCl ₂	50	263	9.8	335	20.1	$\bf{0}$	BJH	MD	309
				\mathbf{Fe}^{2+}					
$\mathbf{Fe}^{2+ d}$	100	215	6.0	289	6.0		BJH	MD	310
\mathbf{Fe}^{2+} e	100	220	6.0	288	6.0		BJH	MD	310
				$Ni2+$					
$Ni2+$	64	217	8.0	276	>16.0	38	TIP4P	MD	285
				\mathbf{Zn}^{2+}					
\mathbf{Zn}^{2+}	200	205	7.0	275	$14 - 15$	dipolar	MCY	МC	311
				$Fe3+$					
$\mathbf{Fe}^{3+ d}$	100	203	6.0	281	6.0	7.7	BJH	MD	310
\mathbf{Fe}^{3+e}	100	205	6.0	281	6.0	7.7	BJH	MD	310
							^{<i>a</i>} The orientational angle of the M-O-H bond (see Figure in Table 8). ^b At 10 kbar pressure. ^c At 81 °C, ^{<i>d</i>} In H ₂ O, ^{<i>e</i>} In D ₂ O.		

breaking point such as so called "gadolinium break" has been observed.

4. Tetravalent Cations

 Sn^{4+} . Since the Sn^{4+} ion is easily hydrolyzed in water, no diffraction experiment has been attempted. NMR measurements carried out in highly acidic solutions show that the ion has six water molecules in the first hydration shell.

Th4+ . Experimental difficulties due to the emission of fluorescence X-rays from Th⁴⁺ ions in the course of the X-ray diffraction measurements must be overcome with suitable corrections. Experiments must be done in highly acidic solutions, because the Th⁴⁺ ions are easily hydrolyzed in water to form various polynuclear $complexes.$ The hydrated Th⁴⁺ ion has about eight water molecules in the first coordination sphere of the ion.²⁰¹ The small hydration number found for a Th- $(NO₃)₄$ solution shows the coordination of an $NO₃$ ion with Th4+ 20° (Table 5). An NMR measurement carried out with H_2 ¹⁷O also shows a large hydration number of about 10.³³⁵

 U^{4+} . The U⁴⁺ ion is also easily hydrolyzed in water to form polynuclear complexes, and the ion emits fluorescence X-rays upon irradiation under X-ray beams. In spite of these difficulties, an X-ray diffraction study has been done for the structural analysis of the hydrated U4+ ion. The ion has more than six water molecules (7.9-9.2) and a square antiprism structure has been suggested.²⁰² The U⁴⁺-H₂O bond length has been determined to be 250-251 pm.²⁰²

No monoatomic ion with a charge higher than +4 can exist as a hydrated ion in water.

The hydration energies of various cations, as well as anions, have been reviewed by Rosseinsky²⁰³ in this journal.

5. Halide Ions

Water is a moderate donor and a rather strong acceptor solvent, and thus, water is expected to solvate anions more strongly than cations. However, hydration of anions is usually much weaker than that of cations, except for limited cases such as OH⁻ and F⁻, because most anions have much larger ionic radii than cations. In fact, the hydration enthalpy of F- is more negative than that of Na⁺ which has a similar or even smaller ionic radius than F⁻²⁰³ Therefore, information on the hydration structure of anions is not very conclusive compared with that of cations in spite of fairly large number of measurements have been carried out for studies on the hydration of anions.

F~. The F- ion is a structure-making anion, and it may be coordinated with six water molecules in the first coordination sphere, $60,70,87,282,283,289,304$ although the results so far reported for the hydration structure of Fare rather scattered.^{57,69,94,281,284,287} The F-O bond length is determined to be 262-269 pm by the X-ray diffraction method^{69,70,94} (Table 6) and 260-267 pm^{87,281-284,304} by computer simulations (Table 14). A sharp distribution of hydration numbers of F-60,281 suggests the formation of a well-ordered arrangement of water molecules around F^- as is expected from its structure-making properties.

Ct⁻. The Cl-ion is the most extensively investigated one among anions for the hydration structure and has been studied by various methods. Many X-ray and neutron diffraction measurements show that Cl- is hydrated with six water molecules (Tables 6 and 8), but results of computer simulations suggest that the hydration structure around Cl- is not so definite and water molecules of such numbers as $1-8^{300,312}$ or $4-13^{79}$ seem to surround Cl- ions (Table 15). The peak corresponding to the Cl- $-H_2O$ interactions found in the radial distribution function is very asymmetric and the

Table 14. Structural Parameters of the Hydration Shell of Anions (X-) Derived from Radial Distribution Functions Obtained by Computer Simulations (See Section III for the Full Definitions of the Abbreviations)

 a The angle is defined as ϕ instead of ψ of the Figure shown in Table 8. b Ion pair is formed. c At 10 kbar pressure. d The solution is placed between two Lennard-Jones walls of 1231-pm separation. ^e No preferential orientation is found. / At 68 °C. *s* At 76 °C. ^h The 0 -O distance between an oxygen atom in $ClO₄$ and that in $H₂O$.

running coordination number has no plateau. Since C_l-is regarded as a structure-breaking ion and rotational and translational correlation times of water molecules around Cl-are even smaller than those of the bulk water, the structural information of the hydration of Cl- is hardly obtained by NMR measurements, and the hydration number of the Cl- ion is often arbitrarily assumed to be 0 in many NMR studies.

The Cl-O bond length is determined to be 310-320 pm. In some cases a much longer bond length has been reported (see Table 6).

The orientation of water molecules around Cl- has been investigated by the neutron diffraction measurements. Most neutron diffraction studies reveal that the tilt angle of the Cl⁻⁻O-H bond is 0-12°.^{114,131,136,206} The small angle value has also been obtained in computer simulation studies.^{281-283,285,308} In some cases bifurcatedly coordinated water molecules to Cl⁻ are reported.64,87

Br. The hydration number of Br determined by the X-ray and neutron diffraction methods is 6 for most cases. $64,70,84,118,214,215$ However, it is obvious from the

Table 15. Characteristic Values for the Distribution of Coordination Numbers (CN) of Ions Estimated from Monte Carlo (MC) and Molecular Dynamics (MD) Simulation Data (See Section III for the Full Definitions of the Abbreviations)

salt	$H2O/salt$ molar ratio	range of CN value	most frequent CN value	frequency, %	water model	method	ref
			$NH4$ +				
NH_4 ^{+ a}	215	$0 - 4$	3	45	MCY	MC	278
			\mathbf{Li}^+				
$Li+ b$	215	$5 - 6$	6	95	MCY	MC	281
$Li+ b$	125	$4 - 6$		62	TIP ₄ P	MC	60
LiCl ^c	4	$3 - 6$	$\begin{array}{c} 5 \\ 5 \end{array}$	42	CF	MD	79
LiCl ^b	4	$3 - 6$	4,5	38	BJH	MD	290
LiCl ^c	$\boldsymbol{4}$	$3 - 7$	6	45	BJH	MD	290
			$Na+$				
$Na+b$	215	$5 - 6$	6	95	MCY	MC	281
$Na+b$	125	$5 - 7$	6	95	TIP ₄ P	MC	60
NaCl ^e	25	$1 - 6$	$\frac{6}{5}$	35	BJH	MD	300
NaCle./	25	$1 - 6$		25	BJH	MD	300
NaCl ^b	100	$4 - 7$	6	60	MCY	MC	301
NaCl ^b	16.5	$3 - 8$	5	35	MCY	MC	301
NaCl^b	7.9	$2 - 7$	$4 - 5$	30	$_{\mathrm{MCY}}$	MC	301
			\mathbf{K}^+				
K^+ ^b	215	$4 - 8$	$6 - 7$	40	MCY	MC	281
			\mathbf{Mg}^{2+}				
MgCl ₂ ^e	50	$1 - 6$	6	60	CF	MD	312
			\mathbf{Zn}^{2+}				
$\mathbf{Zn}^{2+~b}$	200	$4 - 7$	6	53	MCY	MC	311
\mathbf{Zn}^{2+d}	200	$4 - 7$	7	98	MCY	MC	311
			\mathbf{F}^\perp				
$F-b$	215	$4 - 6$	$\bf 5$	90	MCY	MC	281
$F-b$	125	$6 - 7$	6	90	TIP ₄ P	MC	60
			Cl^-				
$Cl - b$	215	$7 - 10$	8	45	MCY	MC	281
Cl^{-b}	125	$6 - 10$	$7 - 8$	40	TIP4P	MC	60
LiCl ^c	4	$4 - 13$	8	21	CF	MD	79
LiCl ^b	$\boldsymbol{4}$	$3 - 9$	6	36	BJH	MD	290
LiCl ^c	4	$6 - 12$		31	BJH	MD	290
$\rm NaCl^e$	25	$1 - 7$		30	BJH	MD	300
NaCl ^{e,f}	25	$1 - 8$	$\begin{smallmatrix}8\3\3\2\end{smallmatrix}$	32	BJH	MD	300
MgCl ₂ ^e	50	$1 - 6$		35	CF	MD	312

 α Definitions of coordination numbers (CN): $\operatorname{first\,minimmum}, r_{1,\min}, \operatorname{in} g(r).$ minimum, $r_{\text{ab,min}}, \operatorname{in} g(r).$ e F .^c Integrated Integrated over the all pairs with $r_{\text{H(N)},\dots}$ over the range between the two $r_{\rm i}$..o < 210 *i* values where $g(r_i) = 1$. Frequency of simultaneous occupation of symmetry sites.' pm. 6 Integrated from r = 0 up *d* Integrated from r = ' At 10 kbar pressure. to the = 0 to the absolute

discussion given in the case of Cl- that the hydration structure of Br- is not uniform and the number of water molecules around Br- ions may be distributed over a wide range.

The Br- -O bond length determined by the diffraction and EXAFS methods is 329-340 pm.^{70,84,118,132,214,216,217} Some values smaller and larger than the range of the hydration number given above are also reported.

 I . The hydration structure of the $I⁻$ ion has been studied rather extensively among those of anions, because it has been used as a good counteranion of various cations, since it scarcely forms ion pairs with hard cations even in concentrated electrolyte solutions. Since the ion is most weakly hydrated among the halide ions, a clear conclusion for the hydration number of the ion may not be obtainable by the diffraction method and a variety of values have been reported (see Table 6). The ambiguity in evaluation of the hydration number arises from difficult extraction of the peak for the I- -H2O interactions from the total radial distribution function obtained by the diffraction method, because the peak may be overlapped with various longrange interatomic interactions in the systems under study. In spite of a large variety of reported hydration numbers of I- , the I- -O bond length determined by the diffraction measurements converges to values in the

range of 355-370 pm, and values between 360 and 370 pm7o,82,88,93,i9o,222 *mos\,* frequently appear in the literature, because the position of the I--O peak can easily be seen in the radial distribution curve owing to a large scattering factor of iodine. Similar bond lengths have also been reported from computer simulation studies.^{291,294,305} The average hydration number of I⁻ estimated by computer simulations is around 9, but it does not mean that the ion has a definite hydration structure, and the arrangement of water molecules around I-may be disordered very much due to its weak interactions with water molecules.

6. Oxyanions

The hydration structure of oxyanions XO_y^2 - has been investigated by the diffraction methods and the X-O (within H₂O) distances (noted as $r_{XO(H_2O)}$ in Table 7), rmsds. $(l_{XO(H_2O)})$ and the numbers of water molecules around the $XO_y^{\nu^2}$ ions $(n_{XO(H_2O)})$ have been determined. Some of the oxyanions interact rather strongly with water molecules in solution.

 $NO₃$. Hydration of $NO₃$ is not very strong and the hydration number $n_{XO(H_2O)}$ reported in the literature spreads over a wide range from 1.3^{224} to 17.7^{153} (Table 7). If we simply assume that each oxygen atom in the $NO₃$ ion contacts with $1-2$ water molecules, the number

of water molecules around an $NO₃⁻$ ion becomes 3 to 6, and therefore, $n_{XO(H,0)} = 1.3$ may be too small. The positively charged N atom in $NO₃$ may interact with 1 or 2 water molecules above and/ or below the planar triangle structure of the ion. If so, the hydration number of an $NO₃$ ion may become a maximum of 8. Therefore, $n_{XO(H_2O)} = 17.7$ may be too large. The values $5.9-967,8997,174,225-227$ seem to be acceptable as the average number of water molecules around an $NO₃$ ion.

The $N-O(H₂O)$ distance has been reported to be 314 pm as the minimum⁸⁹ and 350 pm as the maximum.^{97,223,226} Since the N-O bond length in NO_3 ⁻ is known to be about 120 pm and the $O(NO₃-)-O(H₂O)$ distance may be close to 280 pm, the $N-O(H₂O)$ bond length may become to be about 320-350 pm, depending on the N-O-O($H₂O$) angle.

 $ClO₄$ ⁻. The regular tetrahedral $ClO₄$ ⁻ ion may be surrounded by 8-12 water molecules if we simply assume that each oxygen atom in the $ClO₄$ ion interacts with 2-3 water molecules. Of course, the hydration structure around $ClO₄$ may not be so definite as described above, but many of the analyses of radial distribution functions of solutions containing ClO₄-ions have been done on such a simple and rather unacceptable assumption.^{10,66,164} Some other values between 4^{228} and 25.6^{95} have also been reported for the hydration number of $ClO₄$. A computer simulation study³⁰² gives an average hydration number of 6.5 for $ClO₄$ ⁻ (Table 14).

 SO_4^2 . The SO_4^2 -ion can strongly combine with water molecules and still keeps some water molecules even in crystals of metal sulfates. According to the results by X-ray diffraction studies of metal sulfate aqueous solutions, an SO_4^{2-} ion is surrounded by about $7^{111,154,229,233,234}$ to $12^{73,160}$ water molecules. The S-O(H₂O) length is estimated to be 370-393 pm. Since the S-O bond length in $\text{SO}_4{}^{2-}$ is about 150 pm, the values around 380 pm may be expected to be the $S-O(H₂O)$ length.

SO3² - and S2Os² -. Intramolecular S-O bond length of the ions may be close to that of $\text{SO}_4{}^{2-}$, and each oxygen atom in the ions may interact with 1-2 water molecules in solution. An X-ray diffraction study²³⁰ shows the number of water molecules around SO_3^2 and $S_2O_5^2$ ions to be 9 and 8.9-9.2, respectively.

 $SeO₄²$. The $SeO₄²$ ion has a similar structure to SO_4^2 -, although the length of the Se–O bond (165 pm) is longer than that of the S-O bond. The hydration number of the ion is assumed to be 8 in the structural analysis of the selenates of yttrium(III) and some lanthanoid(III) ions.¹⁰ The value has been consistent with the diffraction data, but some other values, somewhat smaller or larger than that, may also fit to the data. The Se- $O(H_2O)$ length is assumed to be 395 pm from the structural model of $[SeO_4(H_2O)_8]^{2-10}$

 $CrO₄²$. The hydration structure of the chromate-(VI) ion has not been reported. However, in a recent study with a combined X-ray and neutron diffraction method, 267 the $CrO₄²$ ion was used as the reference ion for the isomorphous substitution of SeO_4^2 , and the lengths of 163, 266, and 396 pm were determined for the $M-O(MO_4^2)$, $O-O(MO_4^2)$, and $M-O(H_2O)$ bonds, respectively, where M denotes either Se or Cr. Although the last distance coincides with the previously determined one $(395 \text{ pm})^{10}$ within the limit of error, the mined one (see pm) which the finite of effect, $i_{\text{O(MO₄²)-O(H₂O)}} = 12^{267}$ instead of 8^{10}). Therefore, further investigations may be needed for the hydration structure of both anions.

 $MoO₄²$ and $WO₄²$. Since $MoO₄²$ and $WO₄²$ are isomorphous in their structures, the isomorphous substitution method has been employed in the structural analysis of their sodium salts solutions by means of X-ray diffraction.²³⁵ Therefore, no difference should be found in the hydration structures of the ions. The hydration number of 12 and the metal- $O(H_2O)$ interatomic distance of 406 pm are introduced in the course of the structural analysis of the solutions.

 PO_4^3 and $H_2PO_4^-$. The hydration structures of solutions containing $H_3PO_4^{236}$ and dihydrogen phosphates of divalent metal ions^{237,238} have been studied by means of X-ray diffraction. The $P-O(H_2O)$ distance has been estimated to be 373 and 360 pm for PO₄³⁻ under assumptions of the tetra- and octahydrated structures, respectively, of the ion.²³⁶ The $P-O(H₂O)$ distance in the hydrated H_2PO_4 - ion has been estimated to be 375-391 pm with the hydration number of 4.2- 8,8.237,238 The hydration structures of the ions are not conclusive, however.

CH3COO-. Divalent metal acetate solutions have been investigated by means of X-ray diffraction.²³⁹ The hydration structure of the acetate ion has been estimated by using the second neighbor model around the acetate and metal ions. The number of water molecules around an acetate ion has been evaluated to be 4.2 to 6.1 with the $C(CO_2^-)$ -O(H₂O) distance of 352-372 pm.²³⁹

B. Second Hydration Shell of Ions

Information on the second hydration shell of ions is much poorer than that of the first hydration shell. The quantitative analysis of radical distribution curves in the long r range is difficult and much less reliable. Reliable information in the structure of the second shell may be obtained by the isomorphous substitution method in X-ray diffraction measurements and the isotopic substitution method in neutron diffraction measurements. Of course, such methods cannot be applied to every electrolyte solution. In most works in which some attempts have been done to determine the structure of the second hydration shell of ions, suitable structural models have been introduced. In some other cases fitting procedures between experimental and theoretical *G(r)* functions have been continued until a more or less satisfactory agreement is obtained between them. The distance between a central ion and water molecules in the second hydration shell, $r_{\text{MO}(2)}$, the rmsd, $l_{\text{MO}(2)}$, and the number of water molecules in the second hydration shell of the ion, $n_{MO(2)}$, are summarized in Tables 16 and 17. Data in Table 17 are obtained by computer simulations. The $r_{O(1)-O(2)}$ and $l_{O(1)-O(2)}$ values between water molecules in the first and second hydration shells are also listed in Table 16.

 Li^{+} . The $r_{MO(2)}$ value of 441 pm has been evaluated from a model of the octahedral structure of the hydrated Li⁺ ion with additional 12 water molecules around the hydrated ion,⁸² in which each water molecule in the first hydration shell is hydrogen-bonded with two water molecules with the C_{3v} arrangement. On the other hand, the water molecules in the first hydration shell of Li⁺ is found to have an orientation which is favorable to form three hydrogen bonds with water molecules in the second hydration, the orientation helping to construct a more stable hydration structure than that forms only two hydrogen bonds.^{74,75,78} Therefore, the $n_{MO(2)}$

Table 16. Structural Parameters of the Second Hydration Shell of Cations Derived from Diffraction Measurements (See Section HI for the Full Definitions of the Abbreviations)

	H ₂ O/salt								
salt	molar ratio	$r_{MO_{(2)}}$, pm	$l_{MO_{(2)}}$, pm	$n_{\text{MO}_{(2)}}$	$r_{O_{(1)}-O_{(2)}},$ pm	$l_{\mathbf{O}_{(1)}\text{-}\mathbf{O}_{(2)}$, pm	method	model	ref
				$Li+$					
LiI	25	441	13	12	272	10	X, MD	SNM	82
				$Ag+$					
AgClO ₄	10.4	444	45	9.5			x	FN3	95
AgNO ₃	14.2	429	45	17.3			$\mathbf X$	FN ₃	95
AgClO ₄	16.3,3.3	440,440	35,28	9,8	285,299 ^a	15,17	$\mathbf X$	FN ₃	97
AgNO ₃	16.6,4.3	436,430	31,45	4.9,5	286,289 ^a	101,5	X	FN ₃	97
				Mg^{2+}					
MgCl ₂	27.1.55.5	410,423	68,61	12,12	276,275	1,5	X, MD	SNM	106
MgCl ₂	50	410	5	12	279	$\boldsymbol{2}$	MD, X	SNM	106
Mg(NO ₃) ₂	10.8.24.8	420,420	60,67	12,12	275,278	7,11	x	SNM	109
Mg(CIO ₄) ₂	18	428	23	12	278	11	$\mathbf X$	SNM	110
MgSO ₄	20.5	421	24	12	281	9	$\mathbf x$	SNM	111
$Mg(H_2PO_4)_2$	34	428	47	9.5	276	11.8	$\mathbf X$	FN3	237
				12.5,14			X	SNM	111
$MgCH_3CO_2$ ₂	33.1,14.9	424,415	10,10						
				Ca^{2+}					
CaCl ₂	6.0	460		5.9	292		$\mathbf X$	SNM	117
$CaCl2$ ^b	6.2	450		5.0	291		$\mathbf X$	SNM	117
							$\mathbf X$		
CaCl ₂ ^c	6.0	460		5.0	294			SNM	117
CaCl ₂ ^d	6.2	460		3.8	294		$\mathbf X$	SNM	117
CaCl ₂ ^e	5.6	448		2.8	295		$\mathbf X$	SNM	117
CaCl ₂	4.0	454		1.0	293		$\mathbf x$	SNM	117
				$\mathbf{Sr^{2+}}$					
SrCl ₂	26.5,34.6	494,491	28,33	9.3,14.9	287,281	10,9	X	SNM	119
				Mn^{2+}					
MnCl ₂	550-11	443	33	$14 - 18$			EX	SNM	247
Mn(NO ₃) ₂	12.6	425	29	8.9	271	6	$\mathbf X$	SNM	248
MnSO ₄	13.2,26.8	417,422	31,29	10.4,9.8	278,274	4,3	$\mathbf X$	SNM	249
MnSO ₄	${\bf 28}$	434	26	10.7	273	5	$\mathbf X$	SNM	141
$Mn(CH_3CO_2)_2$	33.5	436	22	10	274	5	$\mathbf x$	SNM	239
				\mathbf{Fe}^{2+}					
FeCl ₂	55.5	451	25	12	284	14	$\mathbf X$	SNM	126
Fe(CIO ₄) ₂	55.5	430	26	12	287	18	$\mathbf X$	SNM	126
FeSO ₄	55.5	443	26	12	285	18	$\mathbf X$	SNM	126
				$Co2+$					
CoCl ₂	17.9			5.7	275	5	$\mathbf X$	SNM	251
$CoCl2 + LiCl$	$8 - 17$	428-393	$40 - 27$	$14.8 - 10$	$275 - 273$	$1 - 10$	$\mathbf X$	SNM	81
$CoCl2 + NiCl2 + 6LiCl$	110			7.3	277	1	X, EX	SNM	252
		420		6.3	279				
Co(ClO ₄) ₂	15.0		29			4	X	SNM	129
$CoCH_3CO_2$ ₂	53.7	428	18	10	279	5	$\mathbf x$	SNM	239
				$Ni2+$					
NiCl ₂	28	399	56	12	272	6	X	SNM	105
NiCl ₂	14	405	32	6.2	268	18	X	SNM	105
NiCl ₂	18			4.1	274		\mathbf{X}	SNM	251
NiBr ₂	24.7	401	21	10.9	279	$\begin{array}{c} 5 \\ 9 \end{array}$	$\mathbf X$	SNM	215
NiBr ₂	11.9,24.7			4.0,10.2	275,280	11,5	$\mathbf X$	SNM	218
NiSO ₄	28	426	$\bf 27$	9.2	273	8	$\mathbf X$	SNM	141
N i $SO4$	27.5	420	33	12	274	8	X	SNM	111
		420	31	11.8	274	$\bf 5$	$\mathbf x$	SNM	
N iSO ₄ ^{$\#$}	27								142
$NiSO_4 + H_2SO_4$	48	433	30	12	277	10	$\mathbf X$	SNM	111
N i SO_4 + Li_2SO_4	37	428	25	12.1	270	$\bf 5$	$\bar{\mathbf{x}}$	SNM	142
$Ni(H_2PO_4)_2$	45.6	414	15	$7.3\,$	270	6		SNM	238
$\text{NiCl}_2 + \text{HCl}$	24.6			4.3	276	4		SNM	254
$NiCl2 + 2LiCl$	24.8			2.7	273	4		SNM	254
$NiCl2 + 3LiCl$	23.2			$2.8\,$	272	$\overline{\mathbf{4}}$		SNM	254
			$\bf{21}$		277	10	X X X X X X		
Ni(ClO ₄) ₂	15.0	410		5.5				SNM	129
				$Cu2+$					
Cu(NO ₃) ₂	23.5,34.7	408	41,37	9.1,8.3	272,274	3	X	SNM	262
Cu(CIO ₄) ₂	15.4,25.1	396,395	38,33	11.1,11.6	280,273	13,13	$\mathbf X$	SNM	147
CuSO ₄	40.4	420	20	7.6	279	3	$\mathbf X$	SNM	150
				${\bf Zn^{2+}}$					
$\mathbf{ZnBr}_{2}h$	10.1	396-395	14	$11.8 - 7.5$			$\mathbf X$	SNM	219
\mathbf{ZnBr}_2	$7.57 - 1.0$	395-410	50	$12 - 14$				SNM	220
\mathbf{ZnI}_{2}	$18 - 55$	410-417	34–44	$3.7 - 12.5$			$_{\rm X}^{\rm X}$	SNM	221
$\text{Zn}(\text{NO}_3)_2$	9.0	402	38	10.8	279	31		SNM	153
							$\frac{\mathbf{X}}{\mathbf{X}}$		
ZnSO ₄	18.6	426	37	12	270	11		SNM	113
ZnSO ₄	27.1	424	28	9.9	279.6	2.5	$\mathbf x$	SNM	154
ZnSO ₄	$100 - 16.9$	421-426	$24 - 40$	$13.2 - 7.6$	$276 - 280$	$2 - 4$	$\mathbf X$	SNM	155
$Zn(CH_3CO_2)_2$	34	423	14	$10\,$	269	$\overline{\mathbf{4}}$	X	SNM	239

Table 16. (Continued)

^a The O₍₁₎-O₍₂₎ distance was not distinguishable from the O-O distance in the bulk water. ^b At 15 °C, $n_{O(1)-O(2)} = 2.2$. ^c At 33 °C, $n_{Q(1)-Q(2)} = 1.7.$ 4 At 80 °C, $n_{Q(1)-Q(2)} = 2.0.$ 4 At 72 °C, $n_{Q(1)-Q(2)} = 1.9.$ At 120 °C, $n_{Q(1)-Q(2)} = 0.9.$ 4 Outer-sphere complexes are formed.
A Parameters refer to those measured at -40, 25, 40, 80, and 100 °C. Para other complexes are also present in the solutions.^{*'*} At 9[']°C and 62 °C. ^{*k*} Parameter values of the second shell of cations and those of the second shell of cations and those of the hydrotography of π . π the first shell of anions are correlated. Parameters refer to those of the hydrated complexes. ${}^m n_{0(1)-0(2)} = 5-10$, ${}^n n_{0(1)-0(2)} = 4-6$. $p_{\text{O}(1)}$ - $\text{O}(2)$ = 12. P Oxygen atoms of water molecules in the second shell of cation and those of anion are not distinguishable.

value of 12 may be better explained, in spite of $n_{\text{MO}(1)}$ value of 6 in ref 82, in terms of 3 water molecules bonding to each water molecule in the tetrahedrally hydrated Li⁺ ion, which has been reported by most of the authors given in references in Table 1. The $r_{O(1)-O(2)}$ distance of 272 pm, which is shorter than the water-water intermolecular distance in the bulk, fit the intensity data.

Table 17. Structural Parameters of the Second Hydration Shell of Cations Derived from Radial Distribution Functions Obtained by Computer Simulations (See Section III for the Full Definitions of the Abbreviations)

	$\rm H_2O/salt$ molar ratio	$r_{MO_{(2)}}$		water model		
salt		pm	$n_{MO_{(2)}}$		method	ref
$Li+$	200	400	14.2	СI	мс	57
LiCl	25	450		$_{\mathrm{ST2}}$	MD	289
LiI	25	430		$\mathrm{ST}2$	MD	289
LiI	25	419		$_{\mathrm{ST2}}$	MD	291
LiI	25	441	12	$_{\mathrm{ST2}}$	MD. X	82
$Na+$	200	450	12.4	СI	MС	57
NaCl	25	450		$_{\mathrm{ST2}}$	MD	289
NaCl	100	444		$\mathrm{ST}2$	MD	296
NaCl	25	447		BJH	MD	300
NaCl ^a	25	480		BJH	MD	300
K+	200	530	17.4	СI	МC	57
KCl	25	460		ST2	MD	303
CsF	25	530		ST2	MD	289
CsF	25	530		ST2	MD	304
CsCl	25	490		ST2	MD	289
NH4Cl	25	506		ST2	MD	279
BeCl_2	50	373		BJH	MD	100
${MgCl}_{2}{}^{b}$	50	447	12	CF	MD	106
$Mg^{2+} + 37NH_3$	164	430	9	MCY	MD	307
CaCl ₂	50	453		$_{\rm CF}$	MD. X	116
SrCl ₂	50	500		BJH	MD	309
$\mathrm{Zn^{2+}}$	200	420	$16 - 18$	MCY	MС	311
^{<i>a</i>} At 235 °C and 3 kbar pressure. ^b Analyzed by the SNM method.						

Results of MD and MC simulations show the number of water molecules in the second hydration shell of 12- 14.2^{57,82} with the $r_{\text{MO}(2)}$ value of 400–450 pm.^{57,82,289,29}

Na⁺ . The number of water molecules in the second hydration shell of Na⁺ estimated by MD and MC simulations is 12.4,⁵⁷ and the $r_{\text{MO}(2)}$ value is ranged from 441 to 480 pm.^{57,289,296,300} The evaluation of the structural parameters of the second hydration shell of Na⁺ from the radial distribution curve may be difficult.

K + . The structure of the second hydration shell of K + may be diffused, because the ion-water interactions of K⁺ is not so strong. An MC simulation calculation reveals that the number of water molecules in the second hydration shell of K^+ is 17.7 and the $r_{\text{MO}(2)}$ distance is 530 pm.⁵⁷ On the other hand, another MD calculation shows that the $r_{\text{MO}(2)}$ value is 460 pm,³⁰³ which is much shorter than the value obtained by the MC calculation. It may be difficult to conclude now which value is more reliable.

 $Cs⁺$. No value has been proposed for $n_{MO(2)}$ of Cs⁺. The $r_{\text{MO}(2)}$ value estimated by MD simulations is 490- $530 \text{ pm.}^{289,304}$

 NH_4 ⁺. It may be impossible to experimentally determine the second hydration structure of NH₄⁺ by the X-ray diffraction method, because the method can hardly distinguish between oxygen and nitrogen atoms in ammonium salt aqueous solutions. An attempt made by MD simulations for the estimation of the structure of the second hydration structure of NH₄+ shows that the N-O(2) distance is 506 pm.²⁷⁹

Ag⁺ . The structure of the second hydration shell of Ag⁺ has been explored from the peak shape of the radial distribution curves of silver(I) perchlorate and nitrate solutions.⁹⁵⁹⁷ Since the linear structure of the hydrated Ag⁺ ion is not accepted, the values in ref 95 may not be reliable. The Ag⁺ -0(2) distance of 430-440 pm and the number of water molecules in the second hydration

shell of Ag⁺ of about 5-9 have been estimated.⁹⁷ The $r_{O(1)-O(2)}$ value is about 285 pm.

Be²⁺. The structure of the second hydration shell of Be2+ has been investigated by an MD simulation with a modified central force model for water.¹⁰⁰ However, since the result obtained for the first hydration structure of the ion in the study is wrong, which is corrected later³⁰⁶ to fit the experimental result,¹⁰⁰ the information on the second hydration shell may be less reliable.

Mg2+ . The number of water molecules in the second hydration shell of Mg²⁺ is estimated to be about 12 in most studies^{106,109-111} and the $r_{\text{MO}(2)}$ value is 410-428 $pm^{106,109-111,237}$ according to the results by X-ray diffraction measurements, but the $r_{\text{MO}(2)}$ value estimated by MD simulations is $430-447$ pm, $106,307$ which is longer than the experimental value. The $O(1)-O(2)$ distance estimated by the X-ray diffraction studies is 275-281 pm, which is shorter than the water-water distance in the bulk water.

Ca2+ . The number of water molecules in the second coordination sphere of Ca²⁺ should be larger than that of Be^{2+} and Mg^{2+} , because Ca^{2+} has a larger ionic surface than the other two. In fact, the $r_{\text{MO}(2)}$ value of Ca²⁺ $(448-460 \,\mathrm{pm}^{116,117})$ is larger than those of Be²⁺ and Mg²⁺. Nevertheless, the estimated value of $n_{MO(2)}$ for Ca²⁺ $(1.0-5.9^{117})$ is smaller than those found for Be^{2+} and Mg2+ , and the result may be difficult to be reasonably explained.

Sr2+ . An X-ray diffraction study shows the result that $r_{\text{MO}(2)}$ is 492-494 pm and $n_{\text{MO}(2)}$ is 9.3-14.9;¹¹⁹ the latter value is much larger than the corresponding value of Ca²⁺. The $r_{MO(2)}$ distance estimated by an MD simulation is 500 pm.³⁰⁹

Mn2+ to Zn2+ . The structural parameters of the second hydration shell of the divalent transition metal ions in the first row of the periodic table from Mn²⁺ to Zn²⁺, summarized in Table 16, does not show clear conclusions for the structural change of the ions with the atomic number as discussed in section III.A.l. An optimistic observation of the data in Table 16 may draw a conclusion that the $r_{\text{MO}(2)}$ changes in a similar manner to the change in $r_{\text{MO}(1)}$, but the conclusion is not definitive, because the values of $r_{\text{MO}(2)}$ reported are so spread even for a given metal ion and are strongly dependent on the models adopted in the structural analyses.

 Cd^{2+} , Hg^{2+} , and Sn^{2+} . It seems to be strange that the $r_{\text{MO}(2)}$ value of Hg²⁺ is smallest¹⁶¹ among the values of the three ions^{160,161,164,225,232,238,239} and even smaller than the values of ions from Mn^{2+} to Zn^{2+} . More investigation may be necessary for the $r_{MO(2)}$ value of Hg²⁺.

Al3+ , Cr3+ , and Fe3+ . The strongly hexahydrated trivalent ions may have a stronger second hydration shell than mono- and divalent ions. Structural models with 12 water molecules in the second hydration shell of the trivalent ions have been adopted by many authors. However, the tilt angle of water molecules in the first hydration shell of $Cr³⁺$ and $Fe³⁺$ ions is in favor of the construction of three hydrogen bonds with water molecules in the second hydration shell. Thus, for the evaluation of orientation of water molecules and the $n_{\text{MO}(2)}$ value, more detailed investigations may be $n_{\text{MO}(2)}^{1000}$ values may be shortest in Al^{3+} and longest in Fe^{3+} , although the difference is not clearly seen from the data in Table 16.

 Y^{3+} , Rh^{3+} , In^{3+} , Ln^{3+} , and U^{4+} . In the course of the structural analysis of the first hydration shell of the ions, the structure of the second hydration shell has also been taken into consideration to reduce uncertainties in the structural parameters of the first hydration shell. Therefore, the structural parameters obtained for the second hydration shell are treated as adjustable parameters which may include various kinds of uncertainties in the structure parameters of interatomic interactions except for those of the parameters in the first hydration shell. Thus, we do not intend to discuss the parameter values of the second hydration shell in detail and just summarize the values in Table 16.

C. Influences of Ion-Pair and Complex Formation on Hydration of Ions

When a part of water molecules in the hydration shell of M^{z+} or X^{z-} is replaced with another ligand, say A, the r_{MO} (or r_{XO}), n_{MO} (n_{XO}) and l_{MO} (l_{XO}) values may be influenced by electrostatic, donor-acceptor, and other interactions between M^{z+} or X^{z-} and A. The values of r_{MO} , n_{MO} , and l_{MO} are summarized together with the values of r_{MA} , n_{MA} , and l_{MA} , in Tables 9-11.

Since the structural data obtained by the solution diffraction measurements are much less accurate than those of crystallographic and the frequency data, no clear bond length variations upon the introduction of A (or A's) to the coordination sphere of M^{z+} is seen.

It has been observed that the $M^{z+}-H_2O$ bond length significantly changes when the structure of an octahedrally solvated divalent transition metal ion is changed to a tetrahedral structure at the complex formation reactions with halide ions. In aqueous solutions the complex formation reactions of divalent transition metal ions do not proceed to form highly ligated halogeno complexes, different from nonaqueous systems, and thus, we do not see such variations of the $M^{2+}-H_2O$ bond in the data of Table 10.

When a large ligand is introduced to the coordination sphere of a small metal ion, a structural change also occurs due to a steric hindrance between ligands, although good examples are not given in the tables.

In nonaqueous solutions in which we can choose solvents with various donor-acceptor properties and sizes, we can see such bond length variations upon the complex formation. However, we do not discuss problems in nonaqueous solution systems in this review.

D. Influences of Temperature and Pressure on the Hydration Structure of Ions

This interesting subject has not been well studied because ranges of temperature examined for and pressure applied to aqueous solutions are narrow in structural studies of hydrated ions and complexes. Results so far reported in the literature show that the effects of temperature and pressure on the metal-ligand bond lengths is very small and not detectable within accuracies of the present experimental techniques. Under extremely high temperatures and pressures, aqueous electrolyte solutions may become electrolyte solutions in water under supercritical conditions. It may be very interesting to study structures of hydrated ions in the supercritical conditions.

IV. Dynamic Aspects of Ionic Hydration

Dynamic properties of hydrated ions are more sensitive than structural properties to environmental conditions such as temperature, pressure, and concentration. The isotopic effect which can be neglected in the structural consideration of hydrated ions, becomes significant in some cases. Information on dynamic behavior of hydrated water molecules and the hydrated ions themselves is still rather limited, and in this review we discuss translational and rotational motions including reorientational motions of coordinated water molecules, diffusional motions of hydrated ions and rates of substitution of water molecules in the first hydration sphere with the bulk water. Various techniques have been adopted in studies on dynamic properties of water molecules in the hydration shells and hydrated ions. The NMR method is one of the most powerful and convenient techniques. The quasielastic neutron scattering (QENS) method has sometimes been used to investigate dynamics of hydrated water molecules by using protons as the probe. Other relaxation techniques of various kinds have been examined to determine translational and rotational correlation times of water molecules in the hydration shell. The tracer diffusion method is important to determine diffusion coefficients of hydrated ions and water molecules. The rate of solvent substitution reactions, techniques of investigation of fast substitution reactions having been first explored by Eigen (see ref 349 as a review of his work), is an interesting kinetic parameter to understand behavior of water molecules in the hydration shell and the mechanism of water substitution reactions in aqueous solutions. Recently developed MD simulations are very useful to elucidate dynamic properties of water in pure and salt solutions.

A. Self-Diffusion Coefficients of Water Molecules and Ions

Translational, rotational, and reorientational correlation times of pure water are important quantities to the investigation of the behavior of hydrated water molecules, and the corresponding correlation times in the hydration shell are compared with those in pure water. The terms of "structure-making" and "structurebreaking" are originated from the comparison of the correlation times of water molecules in the two different situations. The translational (τ_T) and rotational (τ_R) correlation times are given in Table 19. The dielectric relaxation times (τ_{DE}) of pure water at various temperatures are also given in the table. The reorientational correlation time of water molecules in pure liquid (τ_{RO}) is obtained as $\tau_{\text{RO}} = 3\tau_{\text{R}}$, which is correlated with τ_{DE} and should be between $\tau_{\text{DE}}/2$ and $2\tau_{\text{DE}}/3$. Measured τ_{RO} values are given in Table 19 for comparison.

The self-diffusion coefficient of water in the bulk increases with the increase in temperature. The translational and rotational correlation times of water molecules in the bulk shorten with temperature, as expected from the breaking of hydrogen bonds in water. The shortening of the reorientational (τ_{RO}) and dielectric relaxation (τ_{DE}) times are accompanied by the change in $\tau_{\rm R}$.

Self-diffusion coefficients experimentally obtained by the quasi-elastic neutron scattering (QENS) method are summarized in Table 20, in which the diffusion

Table 18. (Continued)

" Estimated value from the concentration and temperature dependences over a broad temperature range. *^b* Upon hydrolysis, the predominant species formed are $[A](H_2O)_6]^{3+}$, $[A](OH)(H_2O)_5]^{2+}$, $[A]_2(OH)_2(H_2O)_8]^{4+}$, $[A]_1_3O_4(OH)_{24}(H_2O)_{12}]^{7+}$, and probably $[{\rm Al}_8({\rm OH})_{20}({\rm H}_2{\rm O})_x]^{4+}.$

Table 19. Self-Diffusion Coefficient *D,* **Translational Correlation Time** τ _T, **Rotational Correlation Time** τ _R, **Reorientational Correlation Time TRO, and Dielectric Relaxation Time** TDE **of Pure Water at Various Temperatures**

$t, \degree C$	D , 10^{-5} cm ² s ⁻¹ (ref 353)	τ _T , ps (ref 30)	τ_R , ps (ref 30)	τ_{RO} , ps (ref 354)	τ_{DE} , ps (ref 354)
0	1.12	11.2	$5.2\,$	15.6	17.8
5	$_{1.31}$	9.6	4.3	12.9	15.0
10	1.54	8.2	$_{3.8}$	11.4	12.7
15	1.78	7.1	3.3	9.9	10.8
20	2.02	6.2	2.9	8.7	9.55
25	$2.3\,$	5.5	$2.5\,$	7.5	8.25
30	2.55	4.9	$^{2.2}$	6.6	7.37
35	2.92	4.3	2.0	6.0	6.60
40	3.23	3.9	$1.8\,$	5.4	5.94
45	3.58	3.5	1.7	5.1	5.30
50	3.90	$3.2\,$	1.5	4.5	4.84
55	4.20	3.0	1.4	4.2	4.40

coefficients of hydrated *(Dh)* and nonhydrated water molecules (D_t) and the average diffusion coefficient (D_H) are given, these quantities being related through eqs 16 and 17. However, accuracies of the values determined by MD, QENS, and NMR are still less than those of the values determined by the classical tracer technique. Diffusion coefficients are determinable by using tracers and by the NMR methods. A very thorough and critical review of self-diffusion coefficients have recently been made by Mills et al.³⁵⁰ A part of the data which can be compared with those by QENS are also given in Table 20.

Self-diffusion coefficients of water molecules in pure solvent (D^w) , in the hydration shell of cations (D^{w+}) and anion (D^{w-}) , together with self-diffusion coefficients of hydrated cations (D^+) and anions (D^-) themselves, determined by MD simulations are summarized in Table 21. The data experimentally obtained are shown in the same table for comparison. Since the origins of the data are sometimes different from those given in Table 20, the values in Tables 20 and 21 are occasionally different. In general, an agreement between MD and experimental results is satisfactory. Times elapsed for the simulation calculations are also given in the table to see how long the simulations have been carried out. The data may be important to judge the reliability of the results. Experimental results are compared with the simulation results when the data are available.

B. Rotational Correlation Time of Hydrated Water Molecules

Rotational correlation times of water molecules in the hydration shell of ions can be determined by the NMR method. Intensive work has been devoted to the relaxation studies on aqueous solutions by the NMR method. A review published by Hertz³⁵¹ is worth citing as an important one. In Table 22 the rotational correlation times τ_R^* , given as a time of rotation of a vector from a proton in a hydrated water molecule to the center of an ion, and the rotational correlation times τ_R^+ and τ_R^- of the dipolar axis of water molecules around cations and anions, respectively, are given, together with hydration numbers of the ions assumed in the works. For paramagnetic ions such as transition metal ions, proton relaxation times burden that τ_R ** can only be determined instead of τ_R^* . Here τ_R^{**} is defined as

$$
(1/\tau_{\rm R}^{**}) = (1/\tau_{\rm R}^{*}) + (1/\tau_{\rm RES}^{i})
$$
 (33)

where τ_{RES} ⁱ is the residence time of water molecules in the hydration shell of ion i. Therefore, the residence time τ_{RES} must be known to obtain τ_R^* . Most cations, however, do not strictly follow the dipole-dipole mechanism expressed by eq 33, except for the cases such as Cu^{2+} and Gd^{3+} . Since the τ_R^* indicates the rotational correlation time of ion-proton vector, the rotational relaxation time of water molecules in the hydration shell can be obtained, in the strict sense, only for ions such as Al^{3+} and Cr^{3+} , which very strongly combine with water molecules, and thus, the rotation of the ion-proton vector can be regarded as the rotation of a water molecule itself.

For systems containing diamagnetic ions the assumption of the rigid body of solvated ions should be rejected. The average rotational correlation time τ_R ^w of water molecules in a given system is measured and is separated into values of individual τ_R^+ and τ_R^- after extrapolation of the τ_R ^w values to 0 concentration. The τ_{R} ⁺ and τ_{R} ⁻ values, the rotational correlation times of water molecules around a cation and an anion, respectively, without specification of bond vector, are usually given in the literature as the ratio τ_R^+/τ_R^0 or τ_R^-/τ_R^0 , where τ_R^0 denotes the rotational correlation time of pure water in the system, which is given to be 2.5 \times 10^{-12} s at 25 °C. At the separation of the τ_R ^w value into τ_{R} ⁺ and τ_{R} ⁻, the hydration numbers of the cation and anion must be known. In Table 22 the τ_{R} ⁺ and τ_{R} ⁻ values so far reported are summarized.

The values of τ_R ⁺ for Li⁺ and Na⁺ are larger than τ_R ⁰ and the τ_R^+ values of other alkali metal ions are even smaller than $\tau_{\textrm{R}}$ ⁰. The former are called the structuremaking ions and the latter the structure-breaking ones. Rotational motions of water molecules are enhanced around structure-breaking cations due to the breaking

^a Reference 358. ^b Reference 359. ^c Reference 360. ^d Reference 361. ^e Reference 362. *[D*ⁱ = D_b_is assumed. ^g Hydration number n^i = ⁴ Reference 358. *P* Reference 359. *CREFERGE 360. P* Reference 361. *P* Reference 362. *I* $D^* = D_b$ is assumed. *P* Hydration number $n^i = 4$ is assumed. n^i Reference 423. *P* $n^i = 6$ is assumed. However, due to n^{i} may be reduced to about 5. $i n^{i} = 8$ is assumed.

of hydrogen bonds between hydrated and bulk water molecules. Similar trends are seen in halide ions. $\tau_{R}^$ of F⁻ is larger than $\tau_{\rm R}^0$, and thus, $\tau_{\rm R}^-/\tau_{\rm R}^0 > 1$, while the values for Cl-, Br-, and I- are less than unity. The τ_{R} ⁺ values of divalent and trivalent cations are larger than τ_R^0 , as expected from the strong hydration structure of the ions. The $\tau_{\rm R}$ value of $\rm SO_4^{2-}$ is larger, while that of ClO₄- is smaller than $\tau_{\rm R}^0$. The results show that the hydration of SO_4^2 is strong, while ClO_4 breaks the water structure by the ion-water interaction. The $\tau_{\rm R}^$ value of BF_4 - is obtained to be almost equal to τ_R^0 when the hydration number of the ions is assumed to be 0.

In Table 23 rotational correlation times $\tau_{\rm R, D}$ and $\tau_{\rm R, OH}$ are listed, which are determined by using proton-, are usted, which are determined by using proton-,
deuteron-, and ^{17}O -waters. τ_{DD} is defined as the correlation time of rotation of a position vector at D in D_2O in the hydration shell of ions, and $\tau_{R,OH}$ is that of rotation of an O-H axis connecting in a water molecule. These values are determined by the NMR method. The I nese values are determined by the INNIR method. I he
isotope effect between ¹H and ²H (D) has been seen. isolope effect between Tr and Tr (D) has been seen.
The τ_{ext} ($\tau_{\text{o}}{}^{0}$ in the preceding section) is determined The $\tau_{R,H}$ (τ_R° in the preceding section) is determined
to be 2.5 X 10⁻¹² s, while $\tau_{R,R}^{0}$ is 1.95 X 10⁻¹² s. The $\frac{1}{100}$ be 2.0 \land 10 s, while $\frac{1}{100}$ is 1.00 \land 10 s. The rotational correlation time of ions, $\tau_{\text{R}}r^{\text{th}}$. The $\tau_{\text{R}}r^{\text{th}}$ values rotational correlation time of ions, $\tau_{R,D}$ ⁱ. The $\tau_{R,D}$ ⁱ values are very close to the $\tau_{R,OH}$ ⁱ, the rotational correlation time of O-H bond in the hydrated water of the ion i (i $= +$ for cations and $i = -$ for anions). Both $\tau_{R,D}$ and

 $\tau_{R,OH}$ ⁱ decrease with an increase in the ionic radius in the series of alkali metal ions, and $\tau_{\rm R,D}$ ⁺ and $\tau_{\rm R,OH}$ ⁺ of alkali metal ions are close to the $\tau_{\textrm{R}}{}^+$ values in Table 22. The same situation is seen for halide ions.

C. Reorientational Time of Hydrated Water Molecules

MD simulations reveal molecular movements of water in the hydration shell at the molecular level. The reorientational $(\tau_{\text{RO,X}})$ times of hydrated water molecules are summarized in Table 24 for rather limited systems. The reorientational time of the dipole (μ) of a water molecule, which is denoted as $\tau_{\text{RO},\mu}$ for water molecules around an ion i ($i = +$ for cation and $i =$ - for anion; the reorientational time of pure water or water in the bulk is denoted with the superscript 0 in Table 24) may reflect the strength of the ion-water bond in the hydration shell of ions. The reorientation times of the H-H axis in a water molecule $(\tau_{\text{RO,HH}})$ around ions i are close to the values of $\tau_{\text{RO}, \mu}$. The reorientational anisotropy of hydrated water molecules can be seen from the different values between $\tau_{\text{RO},\mu}$ ¹ and $\tau_{\text{RO,HH}}$.

Neon is regarded as a Lennard-Jones solute. The reorientational time of water molecules around neon is close to that of $\mathrm{Li^{+}}$ and larger than those of $\mathrm{NH_{4}^{+}}$ and

Table 21. Self-Diffusion Coefficients Determined by MD Simulations for Cations *(D⁺),* **Anions** *(D-),* **Solvent Water** $(D^{\mathbf{w}})$, and Water Molecules in the Hydration Shell of the Cations $(D^{\mathbf{w}+})$ and Anions $(D^{\mathbf{w}-})$ and in the Bulk Water $(D^{\mathbf{w}0})$ $\mathbf{a} \times \mathbf{b} = \mathbf{c} \times \mathbf{c} = (10^{-5} \text{ cm}^2 \text{ s}^{-1})^{\frac{1}{2}}$

salt	H_2O /salt	D^+	D-	D»	D^{w+}	D ^{w-}	D^{w0}	simulation time, ps	water model	ref(s)
NH ₄ Cl	$25\,$	1.0 (1.58)	1.3 (2.66)	2.80	2.90 (2.78)	2.98	2.72	3.5	ST2	363 353,364
$Li+ b$	64,125	0.3 (0.60)						30	MCY	282 365
$Li^{+ c}$	64,125	1.8 (3.10)						30	MCY	282 365
LIC1	3	0.20 (0.21)	0.17 (0.21)	0.35 (0.41)				7.5	BJH	366 367
L _i Cl	4	0.70	0.90	0.96				2.4	CF	79
LiI	25	0.7 (1.0)	1.40 (1.47)	2.48 (2.35)	1.33	2.67	2.85	10	ST2	368 362
Na^{+d}	64,125	1.0 (0.91)						30	MCY	282 365
NaCl	25	0.82 (1.12)	1.08 (1.60)	1.60 (1.92)	0.90	1.51	1.86	5	BJH	300 365,369
NaCl ^e	25	0.55 (1.30)	1.25 (1.85)	1.51 (1.60)	0.89	1.44	1.76	$\bf 5$	BJH	300 365,369
K^{+} /	64,125	2.8 (0.86)						30	MCY	282 $365 -$
KC1	25	1.1 (1.50)	1.1 (1.85)	2.60 (1.86)	2.40	2.50	2.70	8.7	ST2	370 371,372
SrCl ₂	50	0.6	1.2		0.9	1.4	1.4	$\boldsymbol{4}$	BJH	309
$F-b$	64,125		0.9 (0.86)					30	MCY	282 365
$Cl \beta$	64,125		2.3 (1.60)					30	MCY	282 365
HS ^h	107	2.91		3,46				70	ST ₂	373

 a Values in parentheses are those experimentally obtained. b At 5 °C. c At 95 °C. d At 9 °C. e At 30 °C and 10 kbar pressure. f At 1 °C. ^{*k*} At 14 °C. ^{*h*} Lennard-Jones (Hard Sphere) solute at 30 °C.

I", the result suggesting the hydrophobic structuremaking effect around neon atoms. A similar trend can also been seen for the rotational correlation times of τ_R ^{LJ} for neon.

The rotational correlation $(\tau_{R,X}^i)$ times of the dipole $(X = \mu)$ and proton-proton $(X = HH)$ vectors of water molecules in the first hydration shell of cations $(i = +)$, anions $(i = -)$, and bulk water $(i = 0)$ have been **investigated by MD simulations, and the results are summarized in Table 24. The rotational anisotropy of water molecules around the ions examined is not conclusive.**

O. Residence Time of Water Molecules In the First Hydration Shell of Ions

A question how long water molecules stay in the hydration shell is interesting, but is very difficult to answer. Molecular dynamics simulations may reply to the question to some extent. The residence time τ_{RES} **is, of course, different from the rate of solvent substitution process of ions, the data being summarized in** the last table. The residence times (τ_{RES}^i) of water **molecules in the first hydration shell of cations and** anions and at a given position in the bulk $(i = +, -, and)$ **0, respectively) are listed in Table 25.**

The dynamic hydration number n_{hyd} ^{*i*} is newly defined **as follows:**

$$
n_{\text{hyd}}^{i} = n^{i} \exp(-\tau_{\text{RES}}^{0}/\tau_{\text{RES}}^{i})
$$
 (34)

where *ri* **denotes the static hydration number of a cation** $(i = +)$ and an anion $(i = -)$. τ_{RES}^0 and τ_{RES}^i stand for **the residence times of water molecules in the bulk and in the hydration shell of ions, respectively.²⁸² From**

Table 22. Rotational Correlation Times τ_R^* , τ_R^+ , and $\tau_R^$ **for Some Paramagnetic and Diamagnetic Ions Obtained** by the NMR Method at 25 °C

ion	nucleus	$\tau_{\rm R}$ *, ps	$\tau_{\rm R}{}^+$, $\tau_{\rm R}{}^-$, a ps	ref
Li ⁺	⁷ Li	12	4.5	374
$Li+$	ıΗ		7.3	375
$Li+$	'H		5.8	375
$Na+$	ıН		4.5	375
Na ⁺	ŀΗ		4.0	375
K+	'H		2.3	375
Rb^+	'H		1.5	362
$Cs+$	1H		1.3	362
$Cs+$	'H		1.8	362
Mg^{2+}	¹ H		13.0	375
$Ca2+$	١H		8.8	362
Ca^{2+}	'H		7.3	362
$\mathbf{Sr^{2+}}$	ıΗ		6.5	362
Ba^{2+}	'H		5.5	362
V^{2+}	'H	21		376
\mathbf{Mn}^{2+}	ıΗ	31		376
Mn^{2+}	'H	31		377
$\mathbf{M} \mathbf{n}^{2+}$	'H	32		378
$Cu2+$	ıн	25		378
$Cu2+$	17 _O	29.5		351
Cr^{3+}	'H	45		378
Cr^{3+}	'H	83		376
Al^{3+}	27 Al	53	44	374
Gd^{3+}	¹ H	20		376
F-	'H		5.8	362
F-	¹⁹ F	$11 - 17$		360
$Cl-$	ıн		2.3	362
$Cl-$	'H		2.3	362
Br^-	ıΗ		1.5	362
Br^-	ıн		1.8	362
I-	'H		1.0	375
SO_4^2	ıΗ		4.3	375
ClO ₄	ŀН		1.3	375
BF_{4}^-	^{19}F		2.4	360
			^{<i>a</i>} Recalculated from the $\tau_{\mathbf{R}}^i/\tau_{\mathbf{R}}^0$ ratios with $\tau_{\mathbf{R}}^0 = 2.5 \times 10^{-12}$ s	

for pure water at $25 °C$.

Table 23. Rotational Correlation Times $\tau_{R,D}$ and $\tau_{R,OH}$ for Ions Obtained by the NMR Method at 25 °C

jon	salt	$H_2O/salt$	nucleus	$\tau_{\text{R,D}}$ ⁱ , ps	$\tau_{\rm R,OH}$, ps	ref
$Li+$ ^a	LiCl	14.0	${}^{1}H, {}^{2}H, {}^{17}O$	4.1	4.6	379
$Li+$ ^a	LiCl	30.3	${}^{1}H, {}^{2}H, {}^{17}O$	3.9	3.7	379
Li^{+b}	LiCl	14.0	${}^{1}H, {}^{2}H, {}^{17}O$	5.0	5.0	379
$Li+ b$	LiCl	30.3	$^{1}H, ^{2}H, ^{17}O$	3.0	3.0	379
Li^{+a}	LiBr	$270 - 55$	$^{1}H, ^{17}O$	4.53	4.54	380
Li^{+c}	LiBr	$270 - 55$	$^{1}H, ^{17}O$	4.83	4.81	380
Li^{+d}	LiBr	$270 - 55$	$^{2}H, ^{17}O$	4.72		381
$Na+ a$	NaCl	13.9	${}^{1}H, {}^{2}H, {}^{17}O$	3.1	3.5	379
$Na+b$	NaCl	13.9	${}^{1}H, {}^{2}H, {}^{17}O$	3.7	3.7	379
$Na+ a$	NaBr	$270 - 55$	$^{1}H, ^{17}O$	3.36	3.34	380
$Na+b$	NaBr	$270 - 55$	${}^{1}H, {}^{17}O$	3.58	3.54	380
Na^{+d}	NaBr	$270 - 55$	$^{2}H, ^{17}O$	3.49		381
K^+ ^a	KCl	13.9	${}^{1}H, {}^{2}H, {}^{17}O$	1.8	1.7	379
K^+ ^b	KC1	13.9	${}^{1}H, {}^{2}H, {}^{17}O$	1.6	1.6	379
\mathbf{K}^+ a	KCl	$270 - 55$	$^{1}H, ^{17}O$	1.64	1.63	380
K^+ ^b	KCl	$270 - 55$	¹ H, ¹⁷ O	1.75	1.73	380
K^{+d}	KC1	$270 - 55$	$^{2}H, ^{17}O$	1.96		381
Cs ^{+ a}	CsBr	$270 - 55$	$^{1}H, ^{17}O$	1.52	1.50	380
Cs^{+b}	CsBr	$270 - 55$	1H, 17O	1.62	1.59	380
Cs^{+d}	CsBr	$270 - 55$	$^{2}H, ^{17}O$	1.67		381
$\mathbf{Mg}^{2+~a}$	MgCl ₂	$13.9 - 55.5$	${}^{1}H, {}^{2}H, {}^{17}O$	2.02	1.96	382
$\mathbf{Mg}^{2+~b}$	MgCl ₂	$13.9 - 55.5$	$^{1}H, ^{2}H, ^{17}O$	1.94	1.94	382
\mathbf{F} ^{-a}	ΚF	13.9	1H, 2H, 17O	3.9	4.7	379
\mathbf{F} - b	KF	13.9	${}^{1}H, {}^{2}H, {}^{17}O$	5.3	5.3	379
\mathbf{F}^-	ΚF	$270 - 55$	$^{2}H, ^{17}O$	6.47		383
$Cl-$	KCl	$270 - 55$	$^{2}H, ^{17}O$	1.70		383
Br^-	KBr	$270 - 55$	$^{2}H, ^{17}O$	1.33		383
\mathbf{I}^\perp	ΚI	$270 - 55$	$^{2}H, ^{17}O$	0.90		383

 a The intramolecular relaxation rates are calculated by assuming that the O–H bond length is same as that in pure water. b The rotational isotropy is taken into account. The bond lengths in a water molecule are slightly changed. $^{\circ}$ The intramolecular relaxation rates are calculated with OH distance longer than that in pure water. ^d In D₂O.

Table 24. Reorientational Characteristics of Water Molecules Determined by the MD Simulation Method: Reorientation Times $\tau_{RO,X}$ ['] and Rotational Correlation Times $\tau_{R,X}$ ['] of Dipole Vector (X = μ) and of Proton-Proton **Vector (X = HH) of the Water Molecules in the First Hydration Shells of Cations (i = +), Anions (i = -) , and in the** Bulk ($i = 0$) at 25 °C (10^{-12} s)

salt	$\rm H_2O/salt$				$\tau_{\text{RO}_{\mu}}{}^+$ $\tau_{\text{RO}_{\mu}}{}^ \tau_{\text{RO}_{\mu}}{}^0$ $\tau_{\text{RO}, \text{HH}}{}^+$ $\tau_{\text{RO}, \text{HH}}{}^ \tau_{\text{RO}, \text{HH}}{}^0$ $\tau_{\text{R}_{\mu}}{}^+$ $\tau_{\text{R}_{\mu}}{}^ \tau_{\text{R}_{\mu}}{}^0$ $\tau_{\text{R}, \text{HH}}{}^+$ $\tau_{\text{R}, \text{HH}}{}^-$								$\tau_{\rm R,HH}$ ⁰	simulation time, ps	water model	ref
NH _a Cl LiI Ne^a	25 25 107	4.6 6.8 6.59	7.9 4.6	5.8 4.4 3.82	4.5 4.03	3.3	3.1 3.42	1.7 2.5 2.55	2.4 1.5	$1.6\,$ 1.3 1.61	2.1 2.4 2.65	2.0 1.4	1.4 $1.5\,$ 2.12	3.5 10 5.4	ST2 ST2 ST2	363 368 384
	^a A Lennard-Jones solute (neon) at 30 °C.															

Table 25. Residence Times of Water Molecules in the First Hydration Shells of the Cations τ_{RES}^+ and Anions $\tau_{RES}^$ and in the Bulk τ_{RES} ⁰, and Dynamic Hydration Numbers n_{hyd} ⁺, n_{hyd} ⁻ for Cations and Anions, Respectively, Determined by the MD Simulation Method at 25 °C

the definition, the dynamic hydration number is 0 when the residence time of hydrated water molecules is very small compared with that in the bulk and increases with the increase in the residence time of water molecules to approach the static hydration number *ri^l* The dynamic hydration number may be an interesting quantity to discuss dynamic behavior of hydrated water molecules around ions, but the usefulness of the quantity is not yet well established in solution chemistry.

The residence times of oxygen and hydrogen atoms

in the hydration shell are usually different, because a proton can move faster than an oxygen atom to the next water molecule in the second hydration shell through a hydrogen bond. Therefore, the residence time of protons should be shorter than that of oxygen atoms. The residence time of an oxygen atom is regarded as the residence time of hydrated water molecules. Various methods have been used to determine the residence times of proton and oxygen atoms in hydrated water molecules. The results are summarized in Table 26.

Table 26. Residence Time of Oxygens TR ES° or Protons TRES^P for Water Molecules in the First Hydration Shells of Ions Determined by Various Methods at 25 ⁰C (See Sections III and IV for the Full Definitions of the Abbreviations)

ion	salt	$\tau_{\rm RES}$ ^O , s	$\tau_{\rm RES}$ ^P , s	method	ref
Li ⁺	LiI	2.5×10^{-12}		MD	368
Li ⁺			\sim 3 \times 10 ⁻¹¹	NMR	351
Li+			\leq 10–10	QENS	355
$\rm K^+$		${<}10^{-10}$		QENS	415
Cs^+			$\leq 10^{-10}$	QENS	355 315
\mathbf{Be}^{2+}	Be(CIO ₄) ₂	\geq 3 \times 10 ⁻⁴	$\geq 5 \times 10^{-9}$	NMR QENS	355
Mg^{2+} $\overline{\text{Mg}}^{2+}$		2×10^{-6}		NMR	415
$\overline{\text{Mg}}^{2+}$		\sim 10 ⁻⁵		NMR	415
$Ca2+$			$\leq 10^{-10}$	QENS	355
$Ca2+$		$< 10^{-10}$		NMR	415
Mn^{2+}	MnSO ₄	3.2×10^{-8}		NMR	385
$Fe2+$	$Fe(NH_4)_2(SO_4)_2$	3.13×10^{-7}		NMR	385
$Co2+$	$\cos \! \! \alpha_4$	8.8×10^{-7}		NMR	385
$Ni2+$			$\geq 5 \times 10^{-9}$	QENS	355
$Ni2+$	Ni(NO ₃) ₂	3.7×10^{-5}		NMR	385
$Ni2+$		3×10^{-6}	$^{\sim 10^{-6}}$	NMR NMR	415 386
Ni^{2+} $Ni2+$	Ni(ClO ₄) ₂		3.3×10^{-5}	NMR	387
$Cu2+$	Cu(NO ₃) ₂	1×10^{-4}		NMR	385
		5×10^{-9}			
$Cu2+$			$≤10^{-10}$	QENS	355
Zn^{2+}			$>10^{-10}$	QENS	388
Ti^{3+}		1.0×10^{-5}		NMR	389
$_{\mathrm{V}^{3+}}$		6×10^{-4}		NMR	390
$\mathrm{Cr^{3+}}$		1.8×10^5 to		NMR	391
		2.2×10^{5}			
$\mathrm{Cr^{3+}}$ Cr^{3+}			2×10^{-6} $\geq 5 \times 10^{-9}$	NMR QENS	351 356
Cr^{3+}	Cr(CIO ₄) ₃		1.1×10^{-5}	NMR	387
\mathbf{Fe}^{3+}	Fe(ClO ₄) ₃		3.6×10^{-7}	NMR	387
$Fe3+$			$\geq 5 \times 10^{-9}$	QENS	357
Fe^{3+}		5×10^{-5}		NMR	389
\mathbf{Fe}^{3+}		6×10^{-3}		NMR	392
\mathbf{Fe}^{3+}		6×10^{-3}		NMR	393
$Co3+$		10 ⁵		NMR	394
Al3+	AICl ₃	1		NMR	394
Al^{3+} Al^{3+}			2×10^{-3}	NMR	395
Al^{3+}		6.0		NMR NMR	389 396
Al^{3+}		0.78	$\geq 5 \times 10^{-9}$	QENS	357
Ga^{3+}			$\geq 5 \times 10^{-9}$	QENS	357
\mathbf{Ga}^{3+}		5.5×10^{-4}		NMR	389
Nd3+			$> 10^{-10}$	QENS	355
Th^{3+}		1.3×10^{-8} to		NMR	397
		3.9×10^{-8}			
$\mathrm{Th^{3+}}$		1.4×10^{-8} to 4.8×10^{-8}		NMR	397
$\mathbf{D}\mathbf{y}^{3+}$		3.1×10^{-8} to		NMR	397
		5.9×10^{-8}			
$\mathbf{D} \mathbf{y}^{3+}$		3.2×10^{-8} to		NMR	397
		7.1×10^{-8}			
$\mathbf{D} \mathbf{y}^{3+}$			$>10^{-10}$	QENS	357
Ho^{3+}		1.6×10^{-8} to 9.1×10^{-8}		NMR	397
Ho^{3+}		1.8×10^{-8} to		NMR	397
		1.1×10^{-7}			
Er^{3+}		7.4×10^{-9} to 1.5×10^{-7}		NMR	397
Er^{3+}		7.7×10^{-9} to 1.9×10^{-7}		NMR	397
\mathbf{Tm}^{3+}		1.5×10^{-7} to 2.5×10^{-7}		NMR	397
Tm^{3+}		1.6×10^{-7} to 3.0×10^{-7}		NMR	397
F-			\leq 10- 10	QENS	355
Cŀ-	NaCl	3.8×10^{-12}		MD	298
Cŀ-			$\leq 10^{-10}$	QENS	355
I-	LiI	5×10^{-11}		MD	368
I^-			$\leq 10^{-10}$	QENS	355
ClO ₄			\leq 10 ⁻¹⁰	QENS	355

E. Rates of Water Substitution Reactions of Ions

In the last table, Table 27, kinetic parameters for water substitution reactions of various hydrated ions are summarized. The geometries and the static hydration number *(n)* of the hydrated ions taken into consideration of the water substitution reactions are also shown in the table, which are not absolutely necessary to discuss the thermodynamic quantities, but the geometries and the hydration numbers are important to discuss the reaction mechanisms from the thermodynamic quantities of activation state of the reaction. The rates (k) , and enthalpies (ΔH^*) , entropies (ΔS^*) , and volumes (ΔV^*) of activation are listed in the table. It is supposed that a positive value of the entropy of activation and a negative value of the volume of activation indicate the associative mechanism, and the reverse shows the dissociative mechanism at the substitution reaction of hydrated water molecules. From the criteria, reaction mechanisms of the solvent substitution reactions are proposed for various cases. However, it should be noted that no sound theoretical base has been established to say that the negative (positive) entropy of activation or the negative (positive) activation volume indicates the associative (dissociative) mechanism. The criteria might be reasonable as long as the bond length in the hydrated ion, and thus the radius of the spherical shape of the hydrated ion, were preserved during the solvent substitution process. These requirements may hardly be fulfilled in most cases. Reliable theoretical considerations should be done to discuss reaction mechanisms from the kinetic parameters.

Rate constants of solvent substitution reactions decrease with increasing ionic radii of alkali metal and alkaline earth metal ions due to the decrease in the ion-water interaction energy. The irregularity of the rate constants of transition metal ions has been explained in terms of the stabilization of activated complexes by using the ligand field theory.³⁵²

Ruthenium(H1III), palladium(II), and platinum(II) ions are very inert, which may be compared with chromium(III).

V. Concluding Remarks

Structural studies of hydrated ions seem to be established by the development of the diffraction techniques and simulation procedures. Of course, more reliable and precise measurements will be done in the future for the static structures of hydrated ions. The methods are extended to nonaqueous solutions. Structures of various complexes formed in solution have also been determined by the diffraction methods, and the bond length variations with solvents and composition of complexes have been discussed in the field of solution chemistry. However, studies on dynamic properties of ions and water molecules in the bulk and the hydration sphere have not been preformed as much as those on static properties and appears to be an important subject for more intensive investigation in the beginning with the 1990s. Reaction mechanisms of solvent substitution reactions and complex formation reactions will be discussed in terms of behavior of individual ions and molecules with time. Thermodynamics may still provide a sound base in solution chemistry, but thermodynamic properties must be connected with molecular behavior of each species in solution. According to recent MD investigations, the water structure is not so rigid as seen in pictures drawn in various texts and water molecules are very quickly moving at every moment.

Table 27. The Assumed Geometry and Hydration Number *n* for the $[M(H_2O)_n]^{\tau+}$ Ion, the Rate Constant *k* at 25 °C, the **Enthalpy** ΔH^* , **Entropy** ΔS^* , and Volume of Activation ΔV^*

ion	salt	geometry	n	k, s^{-1}	ΔH^* , kJ mol ⁻¹	$\Delta S^*,$ J K ⁻¹ mol ⁻¹	ΔV^* , cm ³ mol ⁻¹	proposed mechanism ^a	ref(s)
Be^{2+}		T	$\overline{\mathbf{4}}$	730	59.2	8.4	-13.6	\mathbf{A}	24
V^{2+}		0	6	87	61.8	-0.4	-4.1	I_a	398
Mn^{2+}	MnSO ₄	$\mathbf 0$	6	3.1×10^{7}	33.9	12.2			385
Mn^{2+}		$\mathbf 0$	6	2.2×10^{7}	32.2	3.8			399
Mn^{2+}	Mn(C1O ₄) ₂	$\mathbf 0$	6	2.1×10^{7}	32.9	5.7	-5.4	\mathbf{I}_a	400
$\mathbf{Fe^{2+}}$		0	6	3.2×10^6	32.2	-12.5			385
\mathbf{Fe}^{2+}	Fe(CIO ₄) ₂	$\mathbf 0$	6	4.39×10^6	41.4	21.2	3.8	$I_a - I_d$	400
$Co2+$	CoSO ₄	$\mathbf 0$	6	1.35×10^6	33.4	-17.4			385
$Co2+$		$\mathbf 0$	6	2.38×10^{6}	43.5	22.1			401
$Co2+$		$\mathbf 0$	6	2.45×10^{6}	49.7	44.3			402
$Co2+$	Co(C1O ₄) ₂	$\mathbf 0$	6	3.18×10^{6}	46.9	37.2	6.1	$I_a - I_d$	400
$Ni2+$	Ni(NO ₃) ₂	$\mathbf 0$	$\bf 6$	2.7×10^{4}	48.5	2.5			385
$Ni2+$	Ni(ClO ₄) ₂	0	6		58.1	41.8			326
$\mathbf{Ni^{2+}}$	Ni(ClO ₄) ₂	$\mathbf 0$	$\bf 6$	3.14×10^{4}	56.8	32.2			403
$Ni2+$	Ni(ClO ₄) ₂	0	6	3.37×10^{4}	52.3	17.2	7.1	I_d	404
$Cu2+$	Cu(NO ₃) ₂	0	(4)	1×10^4	192	-16.8			385
			ខ្មែ	12×10^8	l 87.8				
Ru^{2+}		$\mathbf 0$	6	1.8×10^{-2}	87.8	16.1	-0.4	T	24
$\mathbf{Pd^{2+}}$	Pd(CIO ₄) ₂	Sp ^b	4	560	49.5	-26	-2.2	Ia or A	405
Pt^{2+}	Pt(ClO ₄) ₂	\mathbf{Sp}^b	4	3.9×10^{-4}	89.7	-9	-4.6	I_a or A	24,405
Pt^{2+}	Pt(ClO ₄) ₂		4	5.8×10^{-4}	100	29		I_d	406
\mathbf{Al}^{3+}	$\rm AlCl_3$		6	0.17	113	117			389
Al^{3+}	$\text{Al}(\text{ClO}_4)_2$	$\mathbf 0$	6	1.29	84.7	41.6	5.7	I_d	396
$Ga3+$	Ga(CIO ₄) ₃	$\mathbf 0$	6	400	67.1	30.1	5.0	I_d	407
$Ti3+$		0	6	1.8×10^{5}	43.4	$1.2\,$	-12.1	I_a or A	407
V^{3+}	$V(CF3SO3)3$	Ω	6	500	49.4	-27.8	-8.9	I_a	408
V^{3+}	\rm{VCl}_3		6	1.67×10^3	25.9	-96			408
Cr^{3+}		$\mathbf 0$	6	2.4×10^{-6}	108.6	11.6	-9.6	I_a	409
$Fe3+$	Fe(CIO ₄) ₃	$\mathbf O$	6	167	77.3	57.7			392
\mathbf{Fe}^{3+}	Fe(CIO ₄) ₃	$\mathbf 0$	6	160	64.0	12.1		\mathbf{I}_a	393
$\rm Fe^{3+}$		0	6	160	64.0	12.1	-5.4	I_a	24
\mathbf{Fe}^{3+}	Fe(CIO ₄) ₃	0	6	6.8×10^{4}	40.3	80	-5.4	I_{a}	410
Ru^{3+}		0	6	3.5×10^{-6}	89.8	-48.3	-8.3	I_a	24
Rh^{3+}		0	6	2.2×10^{-9}	131.2	29.3	-4.2	I_{a}	24
Gd^{3+}	Gd(CIO ₄) ₃		9	10.6×10^{8}	12.0	-31.9		I_{a}	411
Gd^{3+}	Gd(CIO ₄) ₃		8	11.9×10^{8}	12.0	-30.9		I_a	192
$\mathbf{T} \mathbf{b}^{3+}$	Tb(CIO ₄) ₃		9	4.96×10^{8}	12.08	-37.9	-5.7	I_a	411
\mathbf{Th}^{3+}	Tb(CIO ₄) ₃		8	5.58×10^8	12.1	-36.9	-5.7	I_a	412
Dy^{3+}	Dv(CIO ₄) ₃		9	3.86×10^8	16.57	-25.0	-6.0	I_a	411
$\mathbf{D} \mathbf{y}^{3+}$	$Dy(CIO_4)_3$		8	4.34×10^{8}	16.6	-24.0	-6.0	I_a	412
Ho^{3+}	Ho(CIO ₄) ₃		9	1.91×10^8	16.36	-31.5	-6.6	\mathbf{I}_a	411
Ho^{3+}	Ho(CIO ₄) ₃		8	2.14×10^{8}	16.4	-30.5	-6.6	I_a	412
Er^{3+}	Er(CIO ₄) ₃		9	1.18×10^8	18.37	-28.8	-6.9	\mathbf{I}_a	411
Er^{3+}	Er(CIO ₄) ₃		8	1.33×10^{8}	18.4	-27.8	-6.9	I_a	412
$\mathrm{Tm^{3+}}$	Tm(CIO ₄) ₃		9	0.81×10^{8}	22.68	-17.4	-6.0	I_a	411
\mathbf{Tm}^{3+}	Tm(CIO ₄) ₃		8	0.47×10^{8}	22.7	-16.4	-6.0	I_a	412
Yb^{3+}	Yb(CIO ₄) ₃		9	0.41×10^8	23.29	-21.0		I_a	411
Yb^{3+}	Yb(CIO ₄) ₃		8	0.47×10^8	23.3	-21.0		I,	192
						^a A, associative; D, dissociative; I, interchange; I _a , associative interchange; I _d , dissociative interchange. ^b Square planar.			

This behavior of water molecules is well understood by chemists, but such pictures drawn from dynamic properties of ions and water molecules do not coincide well with structural pictures which are too much static in the drawings. The difference in the two kinds of pictures are so different that our understandings of *solution* are still confused. Besides studies under equilibrium conditions, studies on nonequilibrium processes in complex formation reactions have scarcely been investigated so far. Dynamic studies of ions and water molecules in solution may be a first step of investigation of nonequilibrium processes in solution. Therefore, the studies on dynamic processes of ions and water molecules should be investigated more extensively. No information is available for the structure of activated complexes. Discussions on the associative (A-) and dissociative (D-) mechanisms for solvent substitution reactions shown in Table 26 should disappear when we could *see* the structure of the activated appear when we could see the structure of the activated
complexes of $[M(H_2O)_{n-1}]^{z+}$ of $[M(H_2O)_{n+1}]^{z+}$ in the solution. Some attempts have been started to determine structures of reaction intermediates in solution by the EXAFS method combined with a stopped-flow technique,⁴¹⁷ but the goal to the direct structural analysis of activated complexes in solution is still far.

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