TABLE VII

RMS RADIAL THERMAL DISPLACEMENTS AND COMPONENTS
ALONG THE PRINCIPAL AXES R_1, R_2, AND R_3
OF THE VIBRATIONAL ELLIPSOIDS

Atom	Rms radial displacements, Å	Principal axes, Å		
		R_1	R_2	R_3
U(1)	0.228	0.124	0.132	0.138
U(2)	0.245	0.156	0.137	0.159
O(1)	0.293	0.136	0.174	0.193
O(3)	0.311	0.160	0.185	0.192
N(1)	0.263	0.138	0.155	0.162
O(11)	0.284	0.140	0.153	0.194
O(12)	0.327	0.140	0.179	0.235
O(13)	0.299	0.142	0.163	0.207
N(2)	0.311	0.145	0.164	0.222
O(21)	0.309	0.138	0.156	0.229
O(22)	0.440	0.142	0.213	0.358
O(23)	0.343	0.152	0.160	0.262
O(2)	0.291	0.131	0.159	0.205
H(21)	0.400	0.180	0.237	0.268
H(22)	0.413	0.180	0.206	0.310
O(4)	0.314	0.131	0.166	0.233
H(41)	0.401	0.172	0.226	0.283
H(42)	0.493	0.182	0.212	0.406

participate in hydrogen bonding since it is much larger than the motion of the corresponding oxygen of the other molecule.

Conclusion

All attempts to resolve the differences between this study and the X-ray study of Vdovenko, *et al.*,[2] have failed. The neutron study does indicate a face-centered symmetry with respect to the uranium atoms, and the β angles are similar when the usual convention is applied to the X-ray results, but the large differences with respect to the unique "*b*" axis cannot be reconciled. Also, the bond lengths are considerably different. It must be concluded, in light of the neutron study, that the abnormal U–O distances found by Vdovenko, *et al.*,[2] are in error and that the correct distances for such a structure are near those predicted by Zachariasen and Plettinger.[4]

The Crystal and Molecular Structure of the Polymeric Complex Mercuric O,O'-Diisopropylphosphorodithioate, Hg[(*i*-C₃H₇O)₂PS₂]₂

BY STEPHEN L. LAWTON

Mercuric O,O'-diisopropylphosphorodithioate, Hg[(*i*-C₃H₇O)₂PS₂]₂, crystallizes in the monoclinic space group $C2/c$ with unit cell parameters $a = 31.70 \pm 0.02$, $b = 8.66 \pm 0.01$, $c = 22.34 \pm 0.01$ Å, and $\beta = 129.34 \pm 0.02°$. The observed and calculated densities are 1.74 ± 0.02 and 1.758 ± 0.004 g/cm³ (for $Z = 8$ molecules/unit cell), respectively. The structure was refined by least-squares methods to a final conventional R index of 0.057 using three-dimensional X-ray diffraction counter data. Unlike the zinc(II) and cadmium(II) O,O'-diisopropylphosphorodithioate complexes which exist as dimers, this complex polymerizes into helical chains around the unit cell twofold screw axis aligned parallel with the crystal needle axis. One of the (RO)₂PS₂ groups functions as a chelating group and the other as a bridging group linking two mercury atoms together. Each mercury atom is surrounded by five sulfur atoms: two at relatively short distances of 2.388 ± 0.005 and 2.391 ± 0.006 Å, two at intermediate distances of 2.748 ± 0.005 (the polymeric bridge) and 2.888 ± 0.005 Å, and one making a van der Waals contact at 3.408 ± 0.006 Å. The –S–Hg–S– angle involving the two short Hg–S bonds is $149.7 \pm 0.2°$. There are no Hg–S interactions with neighboring chains. The P–S bonds within each ligand are unequal in length, averaging 2.01 ± 0.01 and 1.94 ± 0.01 Å; sulfur atoms forming the shorter bonds are associated with the longer (and weaker) Hg–S interactions. The P–O bonds within the nonbridging ligand average 1.62 ± 0.03 Å and those within the bridging ligand average 1.54 ± 0.02 Å.

Introduction

Recently the structures of zinc O,O'-diisopropylphosphorodithioate, Zn₂[(*i*-C₃H₇O)₂PS₂]₄, and of cadmium O,O'-diisopropylphosphorodithioate, Cd₂[(*i*-C₃H₇O)₂-PS₂]₄, were reported.[1] These two complexes were found to exist in the crystalline state as isomorphous dimers in which each metal atom is coordinated with four sulfur

(1) S. L. Lawton and G. T. Kokotailo, *Inorg. Chem.*, **8**, 2410 (1969).

atoms in a somewhat distorted tetrahedral environment. The two phosphorodithioate groups per metal were found to serve two independent functions: one acting as a chelating group and the other as a bridging group linking two metal atoms together to form the dimer. The four phosphorus–sulfur bonds were found to be essentially equivalent.

It is well known that mercury, the third member of

the group IIb elements, exhibits a strong tendency to form complexes with considerably different coordination characteristics than are found for zinc and cadmium. For example, in the divalent state mercury often forms compounds with linear HgX_2 coordination as well as with tetrahedral HgX_4 coordination (X = monovalent ligands) whereas compounds with ZnX_2 and CdX_2 coordination are quite rare. Mercuric compounds with the HgX_2 coordination, however, often exhibit weak $Hg\cdots X$ interactions in addition to the two strong short Hg–X bonds, resulting in polymeric structures. It was of interest to extend our investigation of the metal O,O'-dialkylphosphorodithioates to mercury(II) to determine which type of coordination, MX_2 or MX_4, this metal atom would assume when chelated by the $(RO)_2PS_2$ ligand. Crystals of the mercury(II) diisopropyl derivative of empirical formula

$$H_7C_3\text{—O} \quad S \quad S \quad O\text{—}C_3H_7$$
$$P \quad Hg \quad P$$
$$H_7C_3\text{—O} \quad S \quad S \quad O\text{—}C_3H_7$$
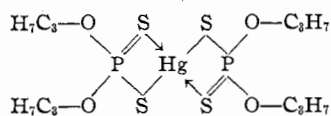
were therefore synthesized. The crystal morphology and X-ray photographs of this complex indicated that the structure of this material differed considerably from that of its zinc(II) and cadmium(II) analogs. Its crystal and molecular structure was subsequently determined and the results of this investigation are reported here.

Experimental Section

Preparation.—Mercuric O,O'-diisopropylphosphorodithioate, $Hg(dtp)_2$, where dtp = $(i\text{-}C_3H_7O)_2PS_2$, was prepared by a method described elsewhere.[2] Soft, light gray needles for use in the X-ray investigation were obtained by recrystallization of the compound from warm absolute ethanol.

Crystal Data.—Mercuric O,O'-diisopropylphosphorodithioate, $Hg[(i\text{-}C_3H_7O)_2PS_2]_2$ (formula weight 627.15), is monoclinic, with space group $C2/c$ (C_{2h}^6). The lattice parameters at 24° are $a = 31.70 \pm 0.02$, $b = 8.65 \pm 0.01$, $c = 22.34 \pm 0.01$ Å, $\beta = 129.34 \pm 0.02°$, $V = 4739 \pm 12$ Å³; $d_{obsd} = 1.74 \pm 0.02$ g/cm³ (by flotation), $d_{calcd} = 1.758 \pm 0.004$ g/cm³ for $Z = 8$ formula units/unit cell. The crystal habit is acicular and elongated along the b direction.

The crystal symmetry was determined from Weissenberg photographs which yielded systematic extinctions hkl, $h + k = 2n + 1$, and $h0l$, $l = 2n + 1$, consistent with the space groups Cc and $C2/c$. The centrosymmetric choice, $C2/c$, was initially indicated on the basis of intensity statistics.[3] The Patterson synthesis also revealed a set of vectors consistent with $C2/c$. The centrosymmetric space group was subsequently confirmed by the successful refinement of the derived structure. The lattice parameters were determined by a least-squares fit[4] to 55 independent unresolved Cu Kα reflections (λ 1.5418 Å) measured on zero-level Weissenberg zones in the 2θ range 77–120° using crystals mounted in three different orientations; each film was calibrated with superimposed aluminum powder lines ($a_0 = 4.0330$ Å). The Nelson–Riley extrapolation function was employed in

the refinement. Indicated errors in the cell parameters are 2σ. The observed lattice constants were checked with the computer program TRACER II[5] to verify that no symmetry higher than C-centered monoclinic was present.

Collection and Reduction of X-Ray Intensity Data.—Complete three-dimensional X-ray diffraction intensity data were taken at room temperature with zirconium-filtered molybdenum Kα radiation from a needlelike crystal 0.230 mm long and of cross section 0.092 × 0.130 mm, parallel with b^*, c^*, and approximately $[\bar{2}01]^*$, respectively. Owing to its fragile character, the crystal was mounted in a 0.3-mm Lindemann glass capillary with the needle axis coincident with the ϕ axis of the diffractometer. The half-widths of a number of narrow-source, open-counter 2θ scans[6] ranged from 0.07 to 0.09° in θ, indicating that the crystal mosaicity was acceptably small. A General Electric quarter-circle Eulerian cradle mounted on a Siemens diffractometer equipped with a Siemens scintillation detector and pulse height discriminator was used to collect the intensity data with the moving crystal, moving-counter measurement technique (θ, 2θ coupling) and a 3.5° takeoff angle. The receiving aperture size selected to minimize extraneous background was 4.2 mm wide by 2.5 mm high. The counter angle, 2θ, was scanned over 2° at a speed of 1°/min. Background counts of 24 sec were taken at each end of the 2θ scan. A total of 3093 independent reflections in the range $2\theta \leq 45°$ were collected. No attenuators other than the β filter were necessary. Three standard reflections were measured periodically as a check on crystal decomposition and a total 7% decrease in intensities was observed.

The raw intensity of each reflection was corrected for background, crystal decomposition, Lorentz, polarization, and absorption effects. The transmission factors were computed by the program ACACA.[7] Owing to the moderately large linear absorption coefficient (71.1 cm⁻¹), refined crystal dimensions normal to the needle axis, rather than those obtained by optical measurement, were used in the absorption correction to reduce possible errors in the calculated transmission factors. Pursuant to this refinement, the shape, or profile, of the crystal normal to the needle axis was first carefully examined by measuring the intensity of a strong low-order $0k0$ reflection, in this case 020, at $\chi = 90°$ in 10° intervals of ϕ from 0 to 360°. These intensities were used to prepare a plot of I_n/I_{av} *vs.* ϕ, where I_n (corrected for background) was the intensity at $\phi = n$ and I_{av} was the average intensity over the full range of ϕ. By systematically adjusting the shape and optically measured crystal thickness, a "best fit" of the transmission factor curve A_n/A_{av} with that of I_n/I_{av} plotted against ϕ was obtained by defining the crystal as an ellipsoidal cylinder with *minor* diameter 0.1034 mm and *major* diameter 0.1150 mm. The maximum deviation in F_o from the mean, defined as $(|F_{max} - F_{min}|/F_{av}) \times 100$, for the 36 measurements was subsequently reduced from 6.3% with no applied correction to 0.3% with this correction. (The corresponding factor based on an absorption correction using the optically measured shape was 7.1%. This value was sufficiently large to indicate that refraction problems arising from the capillary walls and crystal adherent interfered strongly with the optical measurement of the crystal, thereby rendering the observed dimensions unreliable.) The calculated ratio of the minimum and maximum transmission factors for the full set of three-dimensional data was 0.867. Standard deviations in the intensities were computed as described elsewhere,[8] using a value of 0.05 for p. The estimated standard deviation in each F_o was determined by employing the

(2) J. J. Dickert and C. N. Rowe, *J. Org. Chem.*, **32**, 647 (1967).

(3) (a) H. Lipson and W. Cochran, "The Determination of Crystal Structures," G. Bell and Sons, London, England, 1957, pp 32–41; (b) L. Guggenberger, "WSTAT, a Fortran Crystallographic Intensity Statistical Analysis Program," Experimental Station, E. I. du Pont de Nemours and Co., Wilmington, Del., 1967.

(4) D. E. Williams, "LCR-2, a Fortran Lattice Constant Refinement Program," Report IS-1052, Ames Laboratory, Iowa State University, Ames, Iowa, 1964.

(5) S. L. Lawton, "TRACER II, a Fortran Lattice Transformation–Cell Reduction Program," Mobil Research and Development Corp., Research Department., Paulsboro, N. J., 1967.

(6) In the region examined, $2\theta < 11°$, the 2θ-scan and ω-scan techniques yield comparable results: T. C. Furnas, "Single Crystal Orienter Instruction Manual," General Electric Co., Milwaukee, Wis., 1966.

(7) (a) B. J. Wuensch and C. T. Prewitt, *Z. Kristallogr.*, **122**, 24 (1965); (b) C. T. Prewitt, "ACACA, a Fortran Polyhedral Absorption Correction Program," Experimental Station, E. I. du Pont de Nemours and Co., Wilmington, Del., 1966.

(8) P. W. R. Corfield, R. J. Doedens, and J. A. Ibers, *Inorg. Chem.*, **6**, 197 (1967).

method of finite differences.[9] A total of 1649 reflections were observed above the background level, of which 1533 had $F_o{}^2 \geq \sigma(F_o{}^2)$. Those with $F_o{}^2 < \sigma(F_o{}^2)$ were considered as unobserved and thus omitted from the refinements.

Solution and Refinement of the Structure

An unsharpened Patterson function was first calculated using all observed reflections, from which the mercury and phosphorus atoms were located.[10] The positional and isotropic thermal parameters of these seven atoms plus the scale factor were refined by least-squares to an R index $[R_1 = \Sigma||F_o| - |F_c||/\Sigma|F_o|]$ of 0.17 in the centrosymmetric space group $C2/c$. A series of three-dimensional electron density functions (F_o and $F_o - F_c$) and isotropic least-squares refinements led to the location of all remaining nonhydrogen atoms in the molecule.

The structure was refined by using a full-matrix least-squares procedure. The function minimized was $\Sigma w(|F_o| - |F_c|)^2$, where $|F_o|$ and $|F_c|$ are the observed and calculated structure amplitudes, respectively, and w is the weight defined as $1/\sigma^2(F_o)$. Atomic scattering factors for neutral S, P, O, and C tabulated by Hanson, *et al.*,[11] and those for neutral Hg tabulated by Doyle and Turner[12] were used. Effects of anomalous scattering of Hg, P, and S were included in the structure factor calculations,[13] using values for $\Delta f'$ and $\Delta f''$ given in ref 14. The best fully weighted isotropic refinement of all 23 atoms plus the scale factor resulted in $R_1 = 0.117$ and $R_2 = (\Sigma w||F_o| - |F_c||^2/\Sigma w|F_o|^2)^{1/2} = 0.083$ for the reflections above $\sigma(F_o{}^2)$. Further refinement with anisotropic thermal factors introduced for Hg, S, and P reduced R_1 and R_2 to 0.064 and 0.040, respectively. A subsequent three-dimensional difference map revealed vibrational patterns of the isopropyl groups which seemed physically reasonable. Extending anisotropic thermal refinement to the oxygen and carbon atoms reduced the R values still further. All thermal ellipsoids were well-behaved during all stages of the refinement. A second difference map was then computed to locate hydrogen atoms, but because of the large thermal motions of the carbon atoms, none could be readily located.

Convergence was reached with $R_1 = 0.058$ and $R_2 = 0.034$ for the 1533 observed reflections. Corresponding values for all 1649 reflections were $R_1 = 0.066$ and $R_2 = 0.037$. The final standard deviation for an observation of unit weight (*i.e.*, the "error of fit") based on the 208 variables was 1.40. Variation in the value of $w\Delta^2$ ($\Delta = |F_o| - |F_c|$) for different ranges of $|F_o|$ and $(\sin \theta)/\lambda$ was small and was considered to be satisfac-

(9) D. E. Williams and R. E. Rundle, *J. Amer. Chem. Soc.*, **86**, 1660 (1964).

(10) In addition to various local programs for the CDC 1604B computer, local modifications of the following programs were used in the solution and refinement of this structure: Guggenberger's FOUR Fourier program; Busing, Martin, and Levy's ORFLS crystallographic least-squares program; Busing and Levy's ORFFE function and error program; and Johnson's ORTEP thermal ellipsoid plotting program.

(11) H. P. Hanson, F. Herman, J. D. Lea, and S. Skillman, *Acta Crystallogr.*, **17**, 1040 (1964).

(12) P. A. Doyle and P. S. Turner, *ibid.*, **22**, 153 (1967).

(13) J. A. Ibers and W. C. Hamilton, *ibid.*, **17**, 781 (1964).

(14) "International Tables for X-Ray Crystallography," Vol. III, Kynoch Press, Birmingham, England, 1962, pp 215, 216.

tory.[15] On the final cycle the shift in each positional and thermal parameter averaged 0.004 times its own σ. A final difference map calculated on an absolute scale showed no peaks greater than 0.3 e$^-$/Å^3, consistent with good refinement. Comparison of the final values of $|F_o|$ and $|F_c|$ indicated that a correction for secondary extinction was not necessary.

The question of the correct space group arose in the refinement with the appearance of generally large thermal parameters for all atoms, particularly those of the light atoms in the outer coordination sphere. Derived magnitudes of these parameters indicated the possible existence of *two* crystallographically independent molecules of slightly dissimilar molecular geometry and group orientations such as would be allowed in the noncentrosymmetric space group Cc. This seemed unlikely, however, because of the low R values obtained in the $C2/c$ refinement, the lack of any convergence difficulties, the featureless difference Fourier synthesis at the conclusion of this refinement, and a statistical analysis of the observed intensities. It was therefore concluded that the $C2/c$ space group adequately describes the structure.

The final refined positional and anisotropic thermal parameters are given in Table I, along with the associated standard deviations in these parameters as estimated from the inverse matrix. Table II gives the root-mean-square amplitudes of vibrations of the atoms. Table III lists the observed and calculated structure factors, excluding contributions in F_c due to hydrogen. Four reflections—200, 002, $\bar{1}11$, and $\bar{2}02$—were too close to the direct beam to be measured accurately and so only their calculated structure factors are indicated. Data with $F_o{}^2 < \sigma(F_o{}^2)$ are not included in Table III; of these, none had $F_o{}^2 > 2\sigma(F_o{}^2)$.

Description of the Structure

Crystals of mercuric O,O'-diisopropylphosphorodithioate contain eight Hg(dtp)$_2$ molecules per unit cell. These molecules polymerize into infinite zigzag chains around the crystallographic screw diad axes parallel with the needle axis of the crystal, as illustrated in Figures 1 and 2. Tables IV and V list the final interatomic
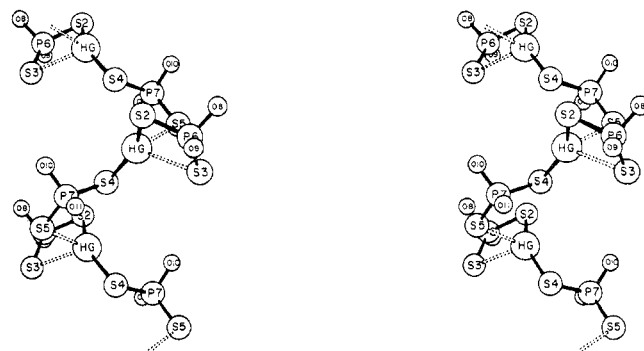
Figure 1.—A stereographic view of three Hg[(i-C$_3$H$_7$O)$_2$PS$_2$]$_2$ molecules within the infinite chain. Dotted bonds represent the "long" Hg–S interactions. The carbon atoms are not shown.

(15) D. W. J. Cruickshank and D. E. Pilling, "Computing Methods and the Phase Problem in X-Ray Crystal Analysis," R. Pepinsky, J. M. Roberts, and J. C. Speakman, Ed., Pergamon Press, New York, N. Y., 1961.
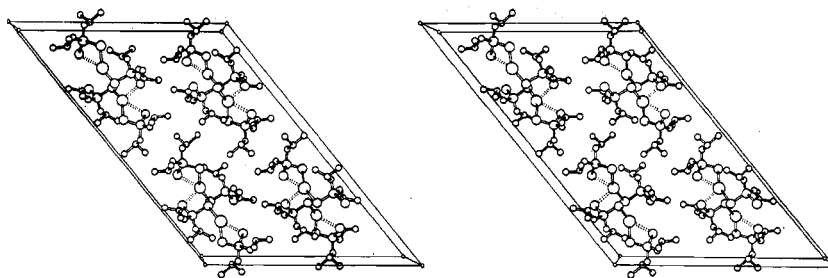
Figure 2.—Stereographic illustration of the molecular packing of the Hg[(i-C₃H₇O)₂PS₂]₂ helical chains in the unit cell. View is along the b axis with the origin in the lower left corner below the plane of the paper. Dotted bonds represent the "long" Hg–S interactions.

TABLE I

FINAL POSITIONAL AND ANISOTROPIC THERMAL PARAMETERS FOR Hg[(i-C₃H₇O)₂PS₂]₂[a]

Atom	x	y	z	β_{11}[b]	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}	Iso equiv, B, Å²
Hg(1)	0.31437 (3)	0.53349 (8)	0.26431 (5)	37.1 (0.2)	200.7 (1.3)	80.0 (0.4)	6.8 (0.6)	39.7 (0.3)	28.8 (0.9)	7.55
S(2)	0.3973 (2)	0.6786 (6)	0.3350 (3)	39 (2)	304 (12)	67 (3)	−22 (4)	35 (2)	−25 (5)	8.51
S(3)	0.3679 (2)	0.4199 (6)	0.2060 (3)	40 (2)	258 (13)	73 (3)	24 (4)	39 (2)	14 (5)	8.10
S(4)	0.2545 (2)	0.3637 (5)	0.2633 (3)	28 (1)	226 (14)	46 (3)	−9 (3)	22 (2)	11 (4)	6.49
S(5)	0.2570 (2)	0.1495 (5)	0.3834 (3)	35 (1)	182 (9)	45 (2)	−17 (3)	28 (2)	−5 (4)	6.14
P(6)	0.4182 (3)	0.5812 (8)	0.2747 (4)	31 (2)	346 (17)	70 (3)	15 (4)	30 (2)	26 (6)	8.71
P(7)	0.2972 (2)	0.3053 (6)	0.3749 (3)	26 (1)	132 (9)	39 (2)	−2 (3)	21 (2)	−1 (4)	4.92
O(8)	0.4274 (7)	0.7186 (20)	0.2365 (11)	59 (5)	451 (41)	105 (10)	−19 (12)	66 (7)	18 (17)	11.5
O(9)	0.4825 (7)	0.5309 (25)	0.3347 (11)	51 (5)	411 (43)	111 (11)	24 (14)	39 (6)	24 (20)	13.8
O(10)	0.3104 (4)	0.4525 (13)	0.4238 (5)	22 (2)	165 (19)	47 (5)	3 (7)	16 (3)	16 (10)	5.8
O(11)	0.3563 (5)	0.2676 (13)	0.4131 (7)	30 (3)	152 (21)	70 (7)	8 (7)	30 (4)	5 (10)	6.5
C(12)	0.3873 (13)	0.8191 (40)	0.1800 (17)	67 (10)	532 (95)	52 (15)	32 (26)	33 (11)	90 (31)	13.8
C(13)	0.4128 (15)	0.9717 (35)	0.2144 (22)	120 (15)	318 (67)	247 (33)	−48 (28)	133 (20)	−56 (41)	20.2
C(14)	0.3823 (14)	0.7958 (39)	0.1121 (16)	97 (13)	804 (111)	82 (17)	8 (29)	67 (13)	30 (37)	18.5
C(15)	0.5001 (15)	0.4065 (39)	0.3726 (20)	74 (13)	432 (83)	111 (22)	73 (26)	64 (15)	146 (35)	14.0
C(16)	0.5406 (11)	0.4306 (52)	0.4590 (17)	49 (8)	1340 (173)	97 (16)	88 (34)	38 (10)	143 (54)	21.8
C(17)	0.5382 (15)	0.3356 (47)	0.3625 (18)	76 (12)	1220 (163)	104 (19)	142 (36)	53 (14)	65 (43)	22.9
C(18)	0.2704 (9)	0.5627 (28)	0.4107 (12)	44 (6)	268 (53)	74 (12)	33 (16)	37 (8)	−14 (22)	9.1
C(19)	0.3002 (11)	0.7179 (21)	0.4369 (14)	86 (10)	70 (33)	133 (18)	4 (15)	73 (12)	6 (20)	12.3
C(20)	0.2625 (11)	0.5243 (31)	0.4711 (17)	97 (10)	353 (59)	207 (23)	−17 (23)	120 (14)	−91 (34)	15.9
C(21)	0.3694 (9)	0.1230 (29)	0.3999 (13)	36 (6)	324 (56)	78 (13)	16 (15)	40 (8)	42 (23)	8.4
C(22)	0.3982 (13)	0.1383 (26)	0.3700 (17)	125 (13)	325 (55)	204 (24)	94 (22)	148 (17)	95 (29)	15.8
C(23)	0.4063 (13)	0.0379 (34)	0.4747 (17)	98 (12)	499 (79)	139 (20)	143 (28)	80 (13)	86 (36)	17.7

[a] All atoms are in the general symmetry position (8f). Numbers in parentheses in all tables and in the text are estimated standard deviations occurring in the least significant digit of the parameter. [b] The form of the anisotropic temperature factor expression is exp $[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. The β's and their estimated standard deviations have been multiplied by 10^4. [c] Calculated from the anisotropic thermal parameters and unit cell parameters by the equation $B \cong (4/3)(\beta_{11}a^2 + \beta_{22}b^2 + \beta_{33}c^2 + 2\beta_{12}ab \cos\gamma + 2\beta_{13}ac \cos\beta + 2\beta_{23}bc \cos\alpha)$: W. C. Hamilton, *Acta Crystallogr.*, **12**, 609 (1959).

TABLE II

FINAL ROOT-MEAN-SQUARE THERMAL
AMPLITUDES OF VIBRATION (Å)
IN Hg[(i-C₃H₇O)₂PS₂]₂

Atom	Min	Med	Max
Hg(1)	0.249 (1)	0.306 (1)	0.362 (1)
S(2)	0.296 (8)	0.309 (7)	0.375 (7)
S(3)	0.269 (7)	0.311 (7)	0.372 (7)
S(4)	0.247 (7)	0.278 (7)	0.330 (7)
S(5)	0.226 (7)	0.256 (7)	0.342 (6)
P(6)	0.298 (9)	0.316 (8)	0.377 (9)
P(7)	0.224 (8)	0.238 (7)	0.284 (7)
O(8)	0.25 (2)	0.42 (2)	0.44 (2)
O(9)	0.36 (3)	0.42 (2)	0.47 (2)
O(10)	0.24 (2)	0.26 (2)	0.31 (1)
O(11)	0.23 (2)	0.30 (2)	0.32 (2)
C(12)	0.21 (4)	0.48 (4)	0.50 (4)
C(13)	0.33 (4)	0.49 (3)	0.64 (4)
C(14)	0.29 (4)	0.55 (4)	0.56 (4)
C(15)	0.23 (4)	0.43 (4)	0.54 (4)
C(16)	0.33 (4)	0.42 (3)	0.74 (5)
C(17)	0.38 (4)	0.41 (4)	0.74 (5)
C(18)	0.24 (3)	0.34 (3)	0.42 (3)
C(19)	0.16 (4)	0.42 (3)	0.51 (3)
C(20)	0.26 (4)	0.42 (3)	0.60 (3)
C(21)	0.26 (3)	0.31 (3)	0.39 (3)
C(22)	0.23 (3)	0.32 (3)	0.66 (3)
C(23)	0.27 (4)	0.43 (3)	0.64 (4)

distances, angles, and corresponding standard deviations. Important distances and angles within the inner coordination sphere are summarized in Figure 3. Owing to the complexity of thermal motion within the molecule, no corrections in the bond lengths were made.

The Mercury(II) Coordination.—The three ions Zn^{2+}, Cd^{2+}, and Hg^{2+} each have a filled d^{10} shell. All three form stable compounds with covalent bonds when the metal (M) is attached to four ligands (X) to form an MX₄ group of tetrahedral symmetry. However, unlike Zn^{2+} and Cd^{2+}, Hg^{2+} exhibits a *strong* tendency also to form MX₂ compounds in which the two bonds are shorter than the four found in the MX₄ compounds. These two classes, MX₂ and MX₄, are summarized in Table VI for known mercury(II) complexes with sulfur ligands.

In Table VI the following observed structural features may be noted. (i) In the MX₂ group the two short mercury–sulfur bonds vary in length from 2.36 to 2.49 Å. In the MX₄ group the mercury–sulfur bond varies in length from 2.49 to 2.66 Å. (ii) In the MX₂ group, the two short bonds are linear in some cases but significantly nonlinear in others; the X–M–X angle ranges from 129 to 180° with no apparent effect on the

TABLE III

OBSERVED AND CALCULATED STRUCTURE FACTORS (IN ELECTRONS × 10)
FOR MERCURIC O,O'-DIISOPROPYLPHOSPHORODITHIOATE, $Hg[(i\text{-}C_3H_7O)_2PS_2]_2$

[Table of observed and calculated structure factors; columns labeled H, L, FO, FC repeated across the page.]

length of the two bonds. The MX_4 group, on the other hand, exhibits tetrahedral or near-tetrahedral symmetry. Metacinnabarite, for example, crystallizes in the zinc blende (wurtzite) structure in which each atom is surrounded tetrahedrally by four other atoms.[16] (iii) In the MX_2 group mercury often exhibits a tendency to form *weak* interactions with additional atoms, thus increasing its sphere of nearest neighbors in some cases to as high as six. These additional interactions often result in the formation of polymers. Cinnabar, for example, is hexagonal with mercury surrounded by six sulfur atoms: two at 2.36 Å and making an angle

with mercury of 172 (2)°, two at 3.10 Å, and two at 3.30 Å.[17] The structure consists of zigzag –Hg–S–Hg– helical chains which are bound to neighboring chains through weak Hg···S interactions to produce a pseudopolymeric compound.

Molecules of $Hg(dtp)_2$ in the present investigation exhibit features fully consistent with those found in the MX_2 group. Each mercury atom is surrounded by five sulfur atoms: two at relatively short distances of 2.388 (5) and 2.391 (6) Å, two at intermediate distances of 2.748 (5) and 2.888 (5) Å, and one at 3.408 (6) Å. The two short distances are well within the range 2.36–2.49 Å observed for the structures of this

(16) A. F. Wells, "Structural Inorganic Chemistry," Oxford at the Clarendon Press, 3rd ed, 1962, pp 528, 895.

(17) K. L. Aurivillius, *Acta Chem. Scand.*, **4**, 1413 (1950).

TABLE IV
INTERATOMIC DISTANCES IN $Hg[(i\text{-}C_3H_7O)_2PS_2]_2$

Intrachelating ligand		Interchelating ligand	
Atoms[a]	Dist, Å	Atoms[a]	Dist, Å
Hg(1)–S(2)	2.391 (6)	Hg(1)–S(4)	2.388 (5)
Hg(1)–S(3)	2.888 (5)	Hg(1)–S(5')	2.748 (5)
S(2)···S(3)	3.284 (8)	S(4)···S(5)	3.221 (6)
P(6)–S(2)	2.023 (8)	P(7)–S(4)	2.006 (6)
P(6)–S(3)	1.935 (8)	P(7)–S(5)	1.946 (6)
P(6)–O(8)	1.60 (2)	P(7)–O(10)	1.55 (1)
P(6)–O(9)	1.64 (2)	P(7)–O(11)	1.52 (1)
O(8)–C(12)	1.39 (3)	O(10)–C(18)	1.46 (2)
O(9)–C(15)	1.26 (3)	O(11)–C(21)	1.41 (2)
C(12)–C(13)	1.48 (4)	C(18)–C(19)	1.53 (3)
C(12)–C(14)	1.44 (3)	C(18)–C(20)	1.56 (3)
C(15)–C(16)	1.51 (4)	C(21)–C(22)	1.44 (3)
C(15)–C(17)	1.49 (4)	C(21)–C(23)	1.49 (3)
		Hg(1)···S(4')	3.408 (6)

[a] Bonds signified by dashes and nonbonds by dots.

TABLE V
INTERATOMIC ANGLES IN $Hg[(i\text{-}C_3H_7O)_2PS_2]_2$

Atoms	Angle, deg	Atoms	Angle, deg
Mercury Coordination			
S(2)–Hg(1)–S(3)	76.3 (2)	S(3)–Hg(1)···S(4')	134.1 (1)
S(2)–Hg(1)–S(4)	149.7 (2)	S(3)–Hg(1)–S(5')	81.6 (1)
S(2)–Hg(1)···S(4')	89.4 (2)	S(4)–Hg(1)···S(4')	95.8 (1)
S(2)–Hg(1)–S(5')	109.0 (2)	S(4)–Hg(1)–S(5')	99.7 (2)
S(3)–Hg(1)–S(4)	118.5 (2)	S(4')···Hg(1)–S(5')	62.0 (1)
Sulfur Coordination			
Hg(1)–S(2)–P(6)	91.6 (3)	Hg(1)–S(5')–P(7')	103.2 (2)
Hg(1)–S(3)–P(6)	79.7 (3)	Hg(1)···S(4')–Hg(1')	158.7 (2)
Hg(1)–S(4)–P(7)	103.7 (3)		
Phosphorus Coordination			
S(2)–P(6)–S(3)	112.1 (4)	S(4)–P(7)–O(10)	109.6 (5)
S(4)–P(7)–S(5)	109.2 (3)	S(4)–P(7)–O(11)	111.9 (5)
S(2)–P(6)–O(8)	107.1 (7)	S(5)–P(7)–O(10)	113.6 (5)
S(2)–P(6)–O(9)	109.2 (7)	S(5)–P(7)–O(11)	116.5 (5)
S(3)–P(6)–O(8)	116.8 (7)	O(8)–P(6)–O(9)	94.6 (11)
S(3)–P(6)–O(9)	115.4 (8)	O(10)–P(7)–O(11)	95.5 (6)
Oxygen Coordination			
P(6)–O(8)–C(12)	126 (2)	P(7)–O(10)–C(18)	126 (1)
P(6)–O(9)–C(15)	125 (2)	P(7)–O(11)–C(21)	120 (1)
Carbon Coordination			
O(8)–C(12)–C(13)	102 (3)	O(11)–C(21)–C(22)	112 (2)
O(8)–C(12)–C(14)	107 (3)	O(11)–C(21)–C(23)	109 (2)
O(9)–C(15)–C(16)	113 (4)	C(13)–C(12)–C(14)	109 (3)
O(9)–C(15)–C(17)	108 (3)	C(16)–C(15)–C(17)	98 (3)
O(10)–C(18)–C(19)	104 (2)	C(19)–C(18)–C(20)	105 (2)
O(10)–C(18)–C(20)	106 (2)	C(22)–C(21)–C(23)	107 (2)

Figure 3.—Selected bond distances and angles in $Hg[(i\text{-}C_3H_7O)_2\text{-}PS_2]_2$. The carbon atoms are not shown.

type. They are equal in length within one standard deviation. If the van der Waals radius of mercury(II) is assumed to be at least half of the interatomic distance in crystalline mercury,[18] 1.5025 Å, and if that of sulfur[19] is taken as 1.85 Å, then their sum (3.35 Å or greater) suggests that an appreciable interaction also exists in the ~2.8 Å set. In this set the bridging Hg–S interaction (2.748 (5) Å) appears to be slightly stronger than the nonbridging interaction (2.888 (5) Å). These sulfur atoms form essentially double bonds with phosphorus. The fifth sulfur atom (that which also forms the Hg–S bridge) occurs at a distance slightly greater than 3.35 Å, revealing no appreciable interaction. The angle formed by the two short Hg–S covalent bonds is 149.7 (2)° which, being considerably less than 180°,

lends further evidence that in mercury(II) compounds of this type the bicovalent angle is quite flexible. As with most other members of this group a pseudopolymeric structure exists. As shown in Figures 1 and 2 the molecules arrange themselves into helical chains with no Hg–S linkages existing between the chains. Cleavages of the crystals parallel with the needle axis are readily explained by this molecular packing.

In a nonpolar solvent such as benzene (38°) Hg-(dtp)$_2$ dissociates into monomers. In a study of divalent zinc, cadmium, mercury, and lead O,O'-diisopropylphosphorodithioates, vapor pressure lowering data have shown[20] that the mercury(II) derivative has the smallest association constant (order: Cd > Zn > Pb > Hg), indicating that the metal–sulfur interaction linking the molecules together is the weakest of the four compounds studied.

The Phosphorodithioate Ligands.—As was found in the molecules of $Zn_2(dtp)_4$ and $Cd_2(dtp)_4$ previously determined,[1] the two phosphorodithioate groups per metal atom serve the *same two* independent functions with one acting as a chelating group and the other as a bridging group linking two metal atoms together. However, unlike those in the zinc and cadmium derivatives, no two bridging groups in Hg(dtp)$_2$ are associated with the same two metal atoms; consequently, polymeric chains, rather than dimers, result. This same type of chain was recently shown to form in crystals of $Zn[(C_2H_5O)_2PS_2]_2$.[21]

The lengths of the two P–S bonds within each ligand vary inversely with the lengths of the adjoining Hg–S bonds, as required by bond order conservation around sulfur. The two sets of P–S distances average 1.94

(18) L. E. Sutton, Ed., *Chem. Soc., Spec. Publ.,* **No. 11**, S6 (1958).

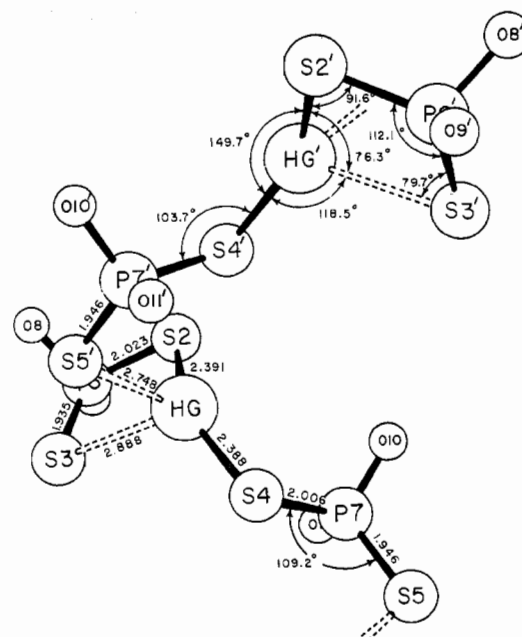(19) L. Pauling, "The Nature of the Chemical Bond," Cornell University Press, Ithaca, N. Y., 1960, p 260.

(20) I. J. Heilweil, *Amer. Chem. Soc., Div. Petroleum Chem., Preprints,* **10**, No. 4-D, 19 (1965).

(21) T. Ito, T. Igarashi, and H. Hagihara, *Acta Crystallogr., Sect. B,* **25**, 2303 (1969).

TABLE VI

SUMMARY OF KNOWN MERCURY(II) COMPOUNDS WITH MERCURY–SULFUR INTERACTIONS[a]

Compound	Mercury environment	Hg–S contacts,[b] Å	S–Hg–S angles,[b,c] deg	Ref
A. Two Short Bonds (MX$_2$)				
HgS (cinnabar)	6S	2 at 2.36 2 at 3.10 2 at 3.30	1 at 172.4 (17)	d
Hg(SCH$_3$)$_2$	5S	2 at 2.36 (5) 3 at 3.25 (5)	1 at 180.0 (25)	e
Hg(SC$_2$H$_5$)$_2$	6S	2 at 2.45 2 at 3.53 2 at 3.56	1 at 180	f
Hg$_3$S$_2$Cl$_2$	2S, 2Cl, 2Cl	2 at 2.45 (2)	1 at 166 (1)	g
2HgCl$_2$·(C$_2$H$_5$)$_2$S	S, 3Cl, 2Cl$^-$	1 at 2.41 (2)	...	h
HgCl$_2$·C$_4$H$_8$S	2S, Cl, 3Cl$^-$	1 at 2.40 (2)	...	i
Hg[SC(NH$_2$)$_2$]$_2$Cl$_2$	Not available	2 at 2.37	1 at 128.75	j
Hg[SC(NH$_2$)$_2$]$_2$(SCN)$_2$	6S	1 at 2.360 (6) 1 at 2.488 (6) 4 at 3.100–3.392 (6)	1 at 180 (8)	k
Hg(C$_{13}$H$_{11}$N$_4$S)$_2$·2C$_2$H$_5$N	4S, 2N	2 at 2.4 2 at 2.8	1 at 155	l
Hg[(C$_6$H$_5$)NC(OCH$_3$)S]$_2$	2S	1 at 2.317 (7) 1 at 2.343 (7)	1 at 174.7 (2)	m
CoHg$_2$(SCN)$_6$·C$_6$H$_6$	4S, 2C	1 at 2.424 1 at 2.455 1 at 2.719 1 at 2.855	1 at 140	n
Hg[(i-C$_3$H$_7$O)$_2$PS$_2$]$_2$	5S	1 at 2.388 (5) 1 at 2.391 (6) 1 at 2.748 (5) 1 at 2.888 (5) 1 at 3.408 (6)	1 at 149.7 (2)	This study
B. Four Short Bonds (MX$_4$)				
HgS (metacinnabarite)	4S	4 at 2.54	6 at 109.47	o
Cu[Hg(SCN)$_4$]	4S	2 at 2.520 (14) 2 at 2.620 (14)	6 at 93.33–143	p
Hg[Co(NCS)$_4$]	4S	4 at 2.558 (4)	6 at 105.1–118.7 (1)	q
[Hg(SCN)$_4$][Cu(C$_2$H$_8$N$_2$)$_2$]	4S	4 at 2.49–2.61	6 at 100.017–120.2	r
Hg[SC(NH$_2$)$_2$]$_4$Co(SCN)$_4$	4S	4 at 2.590 (7)	6 at 109.47 (150)	s
Hg(tert-C$_4$H$_9$S)$_2$	4S	2 at 2.590 (25) 2 at 2.660 (25)	6 at 87.0–121.0 (15)	t
Hg(SC$_4$H$_8$O)$_2$Cl$_2$	2S, 2Cl	1 at 2.540 (22) 1 at 2.590 (22)	1 at 115.0 (8)	u
C$_8$H$_{12}$S$_2$Hg$_2$Cl$_4$	2S, 2Cl	2 at 2.53	1 at 109	v
Hg(S$_2$C$_3$H$_6$)Cl$_2$	2S, 2Cl	2 at 2.61 (1)	1 at 113.0 (3)	w

[a] Structural features of a number of these compounds have been reviewed by C.-I. Bränden, *Ark. Kemi*, **22**, 501 (1964). [b] Numbers given in parentheses represent standard deviations occurring in the least significant digits of the parameters as quoted by the respective authors. [c] The indicated angles refer only to those formed by the two short Hg–S bonds in the MX$_2$ group and to all Hg–S bonds in the MX$_4$ group. [d] See ref 17. [e] D. C. Bradley and N. R. Kunchur, *J. Chem. Phys.*, **40**, 2258 (1964). [f] D. C. Bradley and N. R. Kunchur, *Can. J. Chem.*, **43**, 2786 (1965); a reference to earlier work is cited therein. [g] A. J. Frueh and N. Gray, *Acta Crystallogr., Sect. B*, **24**, 156 (1968). [h] C.-I. Bränden, *Ark. Kemi*, **22**, 83 (1963). [i] C.-I. Branden, *ibid.*, **22**, 495 (1964). [j] K. K. Cheung, R. S. McEwen, and G. A. Sim, *Nature (London)*, **205**, 383 (1965). [k] A. Korczynski, *Rocz. Chem.*, **40**, 547 (1966). [l] M. M. Harding, *J. Chem. Soc.*, 4136 (1958); a reference to earlier work is cited therein. [m] R. S. McEwen and G. A. Sim, *J. Chem. Soc. A*, 1552 (1967). [n] R. Grønbaek and J. D. Dunitz, *Helv. Chim. Acta*, **47**, 1889 (1964). [o] See ref 16. [p] A. Korczynski, *Rocz. Chem.*, **36**, 1539 (1962). [q] J. W. Jeffery and K. M. Rose, *Acta Crystallogr., Sect. B*, **24**, 653 (1968). [r] H. Scouloudi, *ibid.*, **6**, 651 (1953). [s] A. Korczynski and M. A. Porai-Koshits, *Rocz. Chem.*, **36**, 1567 (1965). [t] N. R. Kunchur, *Nature (London)*, **204**, 468 (1964). [u] R. S. McEwen and G. A. Sim, *J. Chem. Soc. A*, 271 (1967). [v] K. K. Cheung and G. A. Sim, *ibid.*, 5988 (1965). [w] W. R. Costello, A. T. McPhail, and G. A. Sim, *ibid., A*, 1190 (1966).

(1) and 2.01 (1) Å. The P–O bonds are also nonequivalent. Those within the nonbridging ligand average 1.62 (3) Å in length whereas those within the bridging ligand average 1.54 (2) Å in length, the averages differing by at most three standard deviations. Both deviate from the average length 1.57 (2) Å observed within this series.[22] The relative importance of π bonding within the

PS$_2$O$_2$ groups is difficult to assess directly because of the paucity of data on M←SPO bond lengths. We therefore offer no formal explanation for the P–O bond length differences at this time. The S–P–S and O–P–O bond angles in the distorted PS$_2$O$_2$ tetrahedron are consistent

(22) The weighted average, plus associated rms deviation, was based on the lengths of P–O bonds present in –PS$_2$(OR)$_2$ ligands reported in ref 1 and 21 and the following: (a) V. Kastalsky and J. F. McConnell, *Acta Crystallogr., Sect. B*, **25**, 909 (1969); (b) J. F. McConnell and V. Kastalsky, *ibid.*,

22, 853 (1967); (c) S. Ooi and Q. Fernando, *Inorg. Chem.*, **6**, 1558 (1967); (d) S. Husebye, *Acta Chem. Scand.*, **20**, 24 (1966); (e) S. L. Lawton, *Inorg. Chem.*, **9**, 2269 (1970). (f) Bond lengths used include corrections for thermal motion of the atoms when available. The 1.62 (3) Å average bond length in Hg(dtp)$_2$ is within one standard deviation of the P–O bond length 1.64 (5) Å reported in ionic KS$_2$P(OCH$_3$)$_2$: P. Coppens, C. H. MacGillavry, S. G. Hovenkamp, and H. Douwes, *Acta Crystallogr.*, **15**, 765 (1962).
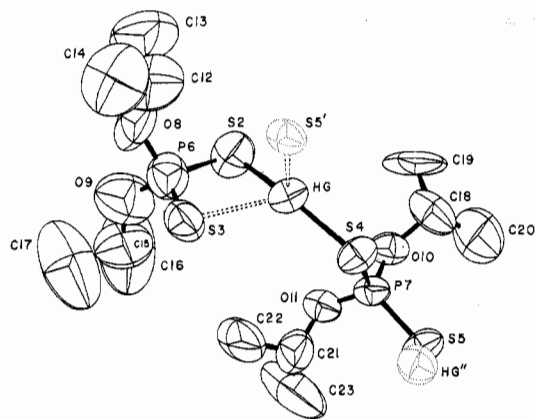
Figure 4.—A parallel projection of the $Hg[(i\text{-}C_3H_7O)_2PS_2]_2$ molecule illustrating the principal root-mean-square thermal displacements. Ellipsoidal boundaries are at the 60% probability level. Atoms S(5') at $1/2 - x, 1/2 + y, 1/2 - z$ and Hg'' at $1/2 - x, -1/2 + y, 1/2 - z$ are related by symmetry to the parent molecule at x, y, z and are indicated by dotted lines.

with those in other phosphorodithioates.[1,21,22a−22f,23]

The two multidentate PS_2 ligands each form a nearly planar array of linkages with the mercury atom, the maximum deviation from planarity not exceeding 0.08 Å. Expressed in unit cell coordinates, the equation of the weighted plane[24] for the chelating ligand with mercury is $0.844x - 6.028y + 12.011z - 0.225 = 0$; distances of the defining atoms from this plane are as follows (Å): Hg(1), −0.0006 (8); S(2), 0.044 (6); S(3), 0.028 (5); P(6), −0.076 (7). The equation of the weighted plane for the bridging ligand with mercury is $19.836x - 6.634y - 11.351z - 0.305 = 0$; distances of the defining atoms from this plane are as follows (Å): Hg(1), 0.0012(8); S(4), −0.048(5); S(5), 0.059(5);

(23) J. R. Knox and C. K. Prout, *Acta Crystallogr., Sect. B*, **25**, 2281 (1969).
(24) W. C. Hamilton, *ibid.*, **14**, 185 (1961).

P(7), −0.080 (5). The mercury atom Hg(1) is substantially out of the plane defined by atoms S(2), S(3), and S(4), however, by an amount 0.480 (1) Å. This deviation may be due to the presence of the Hg(1)−S(5') bridge as well as to crystal packing.

Thermal Motion.—As can be seen in Figure 4, atoms within the inner coordination sphere appear to undergo normal vibrational motions. The relative sizes and orientations of the derived thermal ellipsoids for the isopropyl groups, however, are so excessively large that it is doubtful the ellipsoidal model adequately describes actual motions of the atoms in these groups. It is more reasonable to expect that these groups undergo extensive libration. Significantly, the principal rms thermal displacements of the Hg, S, and P atoms range from 0.22 to 0.38 Å (Table II), indicating the presence of large lattice vibrations as well; such vibrations would be anticipated in a low-melting material (99.5°).

As was found in crystals of polymeric $Zn[(C_2H_5O)_2\text{-}PS_2]_2$, the bridging group exhibits generally smaller thermal parameters than does the chelating group, reflecting the tight binding of atoms within the chains.

van der Waals Contacts.—The laterally stacked chains are held together by methyl···sulfur, methyl···oxygen, and methyl···methyl interactions, the latter being the most prevalent. All the distances are normal when compared with the expected values of 3.85, 3.4, and 4.0 Å, respectively.

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The Crystal and Molecular Structure of Chlorotris(N,N'-dimethylthiourea)copper(I)

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The crystal structure of chlorotris(N,N'-dimethylthiourea)copper(I) has been determined from counter diffractometer data at room temperature from 240 crystallographically independent reflections. The compound crystallizes in the rhombohedral space group $R3m$ with $a = b = c = 8.867 \pm 0.002$ Å, $\alpha = \beta = \gamma = 112.43 \pm 0.08°$ (hexagonal cell: $a = b = 14.741$ (3) Å, $c = 7.471$ (2) Å, $d_m = 1.44$ (3) g/cm³, and $d_c = 1.46$ g/cm³ with one molecule in the rhombohedral cell. The structure was solved by conventional heavy-atom techniques and refined by full-matrix least squares to a final conventional R of 0.030. The structure is composed of discrete "tetrahedral" molecules separated by ordinary van der Waals distances. The space group requires the molecule to have $3m$ (C_{3v}) symmetry. The Cu–S and Cu–Cl distances are 2.360 ± 0.001 and 2.406 ± 0.005 Å, respectively, while the S–Cu–S and S–Cu–Cl angles are 106.4 ± 0.1 and $112.4 \pm 0.1°$, respectively. The Cu–S–C angle is $111.1 \pm 0.3°$ and the metal–sulfur bond involves the sulfur sp² orbital. The analogous Ag complex has a cell six times as large.

Introduction

It has been shown that the Cu(I) and Ag(I) complexes with thiourea, $SC(NH_2)_2$, hereafter tu, give rise

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to a multiplicity of stoichiometries and some unusual and unexpected bonding as well. $Cu(tu)_3Cl^2$ is a chain structure made up of almost regular tetrahedrally co-

(2) (a) Y. Okaya, C. B. Knobler, and R. Pepinsky, *Z. Kristallogr., Kristallgeometrie, Kristallophys., Kristallchem.*, **111**, 385 (1959); (b) Y. Okaya and C. Knobler, *Acta Crystallogr.*, **17**, 928 (1964).