parameters, while for $\cos F_6$ ⁻ the data required $a_3^2k =$ 0.4. However, the latter result proved very sensitive to the values assumed for the free-ion parameters and was felt to be too low, while the estimation of a_5^2 is not possible in any case. Moreover, Gerloch and Miller²⁶ have shown that the presumed equivalence of *k* and *as2* is invalid and that the extent of delocalization of the $t_{2\alpha}$ set is normally appreciably underestimated by the orbital reduction factor.

Consequently the effective spin-orbit coupling constants for OsF_6 ⁻ and Ir Fr_6 ⁻ can probably be determined

more accurately from the electronic spectra, especially for Ir F_6 ⁻ where the spin-orbit splitting of the ground state is directly observable. Naturally the extent of the estimated nephelauxetic and relativistic effects depends closely on the free-ion parameters assumed, which are subject to an appreciable uncertainty. We are therefore attempting to assess the consistency of our parameterization by studying other 5d hexahalo complexes,²⁴ particularly those of $Os(IV)$ and $Ir(IV)$ in which the spin-orbit splitting of the ground levels allows good estimates of the effective spin-orbit coupling constants of the complexes to be made.

CONTRIBUTION FROM THE CHRISTOPHER INGOLD LABORATORIES, UNIVERSITY COLLEGE, LONDON, EXGLAND

Vapor-Phase Raman Spectra, Force Constants, and Values for Thermodynamic Functions of the Tetrachlorides, Tetrabromides, and Tetraiodides of Titanium, Zirconium, and Hafnium

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The Raman spectra of titanium tetrachloride (at 65°), titanium tetrabromide (at 125°), and the tetrachlorides, tetrabromides, and tetraiodides of zirconium and hafnium (at 380-420') have been recorded in the vapor phase, and values for all the fundamentals of these molecules are presented. The $v_2(e)$ and $v_4(t_2)$ modes have pronounced rotational structure in each case. Although the zirconium and hafnium tetrahalides are polymeric in the solid state, their Raman spectra are interpreted to indicate that all are tetrahedral monomers in the vapor phase and thus isostructural with the titanium tetrahalides (which are tetrahedral in all states of matter). Force constants have been calculated for all the halides, including titanium tetraiodide (by the use of solution data) on the bases of both the modified-valence force field as well as the Urey-Bradley force field. Values for various thermodynamic functions have also been computed for each halide.

The tetrahalides of zirconium and hafnium have closely similar physical properties on account of the near identity of the radii of the two group IV metal atoms. In view of the technological interest in the separation of zirconium from hafnium and the role of the tetraiodides in the van Arkel-de Boer process for the preparation of the pure metals, the thermodynamic properties of the tetrahalides in the vapor phase acquire considerable practical importance. The properties may be calculated from a knowledge of the frequencies of the various fundamentals for molecules of known structure. Although the tetrahalides of titanium are tetrahedral monomers in all states of matter, $¹$ this is</sup> not the case for those of zirconium and hafnium. In the solid state, the halides MX_4 ($M = Zr$ or Hf, $X =$ C1 or Br) consist² of zigzag chains, each metal atom being octahedrally coordinated to halogen atoms, two of which are terminal and two bridging, *viz.* $[MX_{1/2}X_2]_{\infty}$. The structures of the solid tetraiodides are uncertain. By contrast with the data for the solid state, electron diffraction data on zirconium tetrachloride^{$3-5$} and on hafnium tetrachloride⁵ indicate that in the vapor phase these molecules are monomeric and tetrahedral. Infrared data 6 on the same two halides in the vapor phase have been interpreted on this basis, one fundamental, $v_3(t_2)$, having been found directly in each case. Vapor density measurements on zirconium tetrabromide and zirconium tetraiodide have shown that these halides also are monomeric in the vapor phase.7 It remains to be shown that they, as well as the tetrabromides and tetraiodides of hafnium, are tetrahedral in the vapor phase. The present Raman data establish this conclusively.

Values for all the fundamentals of the group IVa tetrahalides (excepting the fluorides) have been determined, and appropriate force constants have been calculated on the modified-valence force field and Urey-Bradley force-field models. Values for various thermodynamic functions of each of the halides have also been calculated and are tabulated. A preliminary publication on this work has appeared.8

Experimental Section

⁽²⁶⁾ M. Gerloch and J. R. Miller, *Pvogv. Inorg. Chem.,* **10,** 1 (196s).

⁽¹⁾ R. J. H. Clark, "The Chemistry of Titanium and Vanadium," Elsevier, Amsterdam, 1968.

⁽²⁾ B. Krebs, *Z. Amrg. Allg. Chem.,* **378,** 263 (1970).

⁽³⁾ M. W. Lister and L. E. Sutton, *Trans. Faraday Soc.*, **37,** 393 (1941). **(4)** M. Kimura, K. Kimura, M. Aoki, and *S.* Shibata, *Bull. Chem. Soc.* Jap., 29, 95 (1956).

⁽⁵⁾ V. P. Spiridonov, P. **A.** Akishin, and V. I. Tsirel'nikov, *J. Struct. Chem. (USSR),* **8,** 311, (1962).

Compounds.-The sources of the tetrahalides were as follows: British Drug Houses (TiCl₄, ZrCl₄); Alfa Inorganics (TiBr₄, TiI₄, ZrBr₄, ZrI₄, HfCl₄, HfBr₄, and HfI₄).

^{(6) (}a) J. K. Wilmshurst, *J. Mol. Speclrosc.,* **6,** 343 (1960); (b) **A.** Buchler, J. B. Berkowitz-Mattuck, and D. H. Dugre, *J. Chem. Ph>s.,* **84,** 2202 (1961).

⁽⁷⁾ W. Fischer and 0. Rahlfs, *Z. Elektrochem. Angew. Phys. Chem.,* **38, 592 (1932).**

⁽⁸⁾ R. J, H. Clark, B, K. Hunter, and D. M. Rippon, *Chenz. Ind. (Loudon),* **787** (1971).

The iodides were found to be insufficiently pure as supplied, there being oxide and/or oxyiodide as well as iodine impurity present. Whereas the presence of the oxo compounds was not serious, owing to their low volatility, the presence of even very small amounts of iodine rendered the recording of the Raman spectrum of a tetraiodide impossible owing to (a) absorption of the scattered beam and (b) the excitation of the very strong resonance fluorescence9 spectrum of iodine.

Titanium tetraiodide was purified by four consecutive sublimations. Zirconium and hafnium tetraiodides were purified by chemical transport *in vacuo* in a sealed tube *ca.* **18** in. long, the ends of which were held at *ca.* **440'** (end **A)** and **20'** (end B) and the center portion of which was held at *ca.* **170'.** The impure tetraiodide was placed at end **A** prior to the appropriate temperature gradient being set up. After *ca*. 24 hr, all the iodine was transported to the cold end of the tube, and the pure tetraiodide was transported to the center region of the tube, while the oxo impurities remained at the hot end of the tube. The pure tetraiodide was then sealed into a vacuum manifold.

Instruments.-The Raman spectra were recorded using a Spex **1401** spectrometer in conjunction with Coherent Radiation Laboratories Ar⁺ and Kr⁺ lasers. The scattered radiation was collected at 90° and focused by an $f/0.95$ lens onto the entrance slit of the monochromator after having been passed through a polarization scrambler. The **0.75-m** Czerny-Turner monochromator employed two **1200** lines/mm Bausch and Lomb gratings blazed at **5000 A.** The method of detection was photon counting (cooled FW **130** phototube) and the presentation of the spectra was linear in response. The power available at the exciting wavelengths used, 4880 and 5145 Å of the Ar⁺ laser, was ca . **1.3** W and *ca.* **1.6** W, respectively. The spectra were calibrated by reference to the emission lines of argon and neon. The spectral slit width was *ca.* 1 cm-l.

The samples were sealed into small Pyrex tubes *ca.* **2** cm long. The tubes were held inside a small furnace as shown in Figure 1,

Figure 1.—Hot cell used in conjunction with the recording of the Raman spectra of the tetrahalides in the gas phase.

the temperature inside of which could be controlled by a transformer. The spacings between the wire wound around the inner tube were arranged so that a temperature gradient was produced with the hottest part of the furnace being directly in the path of the laser beam. This was necessary in order that any sublimation of the sample would be away from the path of the laser beam and thus interference with the incident and scattered beam by the sublimate would be avoided.

Results and Discussion

The procedure was adopted of recording the spectrum of any tetrahalide at the lowest temperature for which the signa1:noise ratio of the Raman spectrum was acceptable. Although the volatility of the tetrahalides is in the order chloride $>$ bromide $>$ iodide, the order of metal-halogen bond polarizability derivatives is the reverse¹⁰ so that for any given metal atom the spectra of the chloride, bromide, and iodide were in fact recorded at rather similar temperatures. The (9) **I. R. Beattie,** *G.* **A. Ozin, and R.** *0.* **Perry,** *J. Chem. Soc. A,* **2071**

(1970). (10) R. J. **H. Clark and C.** J. **Willis,** *Inovg. Chem,* **10, 1118 (1971).** values found for the various fundamentals are listed in Table I, together with the known or estimated metal-

^a*Y.* Morino and H. Uehara, *J. Chem. Phys.,* **45, 4543 (1966). M.** W. Lister and L. E. Sutton, *Trans. Faraday* Soc., **37, 393** (1941). ^{*c*} Data refer to cyclohexane solutions: R. J. H. Clark and C. J. Willis, *J. Chem. Soc. A*, 838 (1971). *^d* Estimated value. **eV.** P. Spiridonov, P. **A.** Akishin, and V. I. Tsirel'nikov, *J. Struct. Chem. (USSR),* **3, 311 (1962).** f I. **A.** Cherkasov, quoted by *S.* S. Berdonosov, D. G. Berdonosova, **A.** V. Lapitskii, and L. G. Vlasov, *Russ. J. Inorg. Chem.*, 8, 277 (1963). *⁰* ν_1 and ν_2 both have sharp Q branches (with the exception of ν_2 in HfI₄) and are considered accurate to ± 1 cm⁻¹. ν_3 is broader and is considered accurate to ± 2 cm⁻¹. ν_4 is weaker than ν_2 and its Q branch is difficult to locate; its frequency is considered accurate to ± 3 cm⁻¹.

halogen bond lengths and the *B* values calculated therefrom.

Although the Raman spectra of titanium tetrachloride and titanium tetrabromide have long been known in the liquid state,¹¹ there has been no report of their vapor-phase Raman spectra. On the other hand, the vapor-phase infrared spectrum of titanium tetrachloride has been reported,¹² the values for the fundamentals being as follows: ν_1 , 388; ν_2 , 119; ν_3 , 498.5; ν_4 , 139 cm⁻¹ (ν_3 was observed directly, whereas ν_1 , ν_2 , and ν_4 were deduced from combination bands). The values for ν_1 and ν_3 are in excellent agreement with those reported herein, but the values for ν_2 and ν_4 are 3-5 cm⁻¹ higher than those reported herein. A later report¹³ of the vapor-phase infrared spectrum of titanium tetrachloride places v_4 at 136 cm⁻¹, in exact agreement with the value reported herein. The liquid state Raman values for the fundamentals of titanium tetrabromide (ν_1 229.5; ν_2 74; ν_3 384; ν_4 90.5 cm⁻¹)¹⁴ differ by up to 9 cm⁻¹ from the present vapor-phase data. The vapor-phase Raman spectra of these two halides are shown in Figures **2** and 3. All attempts to obtain the vapor-phase Raman spectrum of titanium tetraiodide were unsuccessful due to its decomposition in the laser beam (6471-A excitation, 700 mW). Accordingly, subsequent calculations were carried out by use of data pertaining to cyclohexane solutions.

The only previous reports of the Raman spectra of the zirconium or hafnium tetrahalides refer to spectra of zirconium tetrachloride as a melt in phosphorus oxytrichloride and in phosphorus pentachloride.ls Unfortunately, it is not certain what the scattering

(12) N J. Hawkins and D. **R Carpenter,** *J Chem. Phys* , **23, 1700 (1955).**

- **(13) H Burger and A.** Ruoff, *Spectrochzm Acta, Part A,* **24, 1863 (1968).**
- **(14) F A Miller and** *G.* L. **Carlson,** *zbzd* , **16, 6 (1960) (15) R.** J. **H. Clark and** C. J. **Willis,** *J. Chenz. SOL A,* **838 (1971).**
- **(16) M. L. Delwaulle and** M. **F. Fransois,** *Bull* Soc *Chzm. Fv,* **205 (1946),** *J. Phys. Radzum,* **7, 15 (1946).**

⁽¹¹⁾ P. Pascal, Ed, "Nouveau Trait6 de Chimie Minerale," Vol 9, Masson, Paris, 1963.

Figure 2.-Raman spectrum of titanium tetrachloride at 65° in the vapor phase (5145-A excitation). Instrument settings: scanning speed 10 cm^{-1}/min, gain 500 counts/sec, time constant 1 sec, slits 70/200/70 *p;* inset for *vg,* settings identical except for gain 200, time constant 2 sec, slits l00/200/100 *p.*

Figure 3.—Raman spectrum of titanium tetrabromide at 125° in the vapor phase (5145-A excitation). Instrument settings: scanning speed 10 cm⁻¹/min, gain 500 counts/sec, time constant 1 sec, slits $30/200/30 \mu$; inset for ν_3 , settings identical except for time constant 2 sec, slits $80/200/80 \mu$; inset for ν_2 and ν_4 , settings as for main scan except for time constant *2* sec, slits 35/200/35 *p.*

species are in these media.¹⁷ The Raman spectrum of zirconium tetrachloride in the solid state has also been reported,¹⁸ these data referring therefore to the sixcoordinate polymeric form.

The vapor-phase spectra of the zirconium and hafnium tetrahalides are new, and the assignments therefore require comment. One band in the spectrum of each compound is completely polarized $(\rho < 0.005)$, and must therefore be assigned to the totally symmetric stretching mode of a cubic *(i.e.,* in this case tetrahedral) molecule. Accordingly, all the fundamentals have been assigned on the basis of T_d symmetry for each molecule. The $v_1(a_1)$ band is very strong and comparatively sharp in each case; it has no rotational structure in agreement with the appropriate selection rules.¹⁹ By contrast, the fundamental of highest frequency for each molecule is very weak and is accordingly assigned to the $\nu_3(t_2)$ mode. Its frequency agrees well with the corresponding infrared value in the two cases for which this comparison is possible, *viz.*, $ZrCl_4$ 421 cm⁻¹ (infrared),⁶ 418 cm⁻¹ (Raman); $HfCl₄ 393 cm⁻¹ (infrared),⁶ 390 cm⁻¹ (Raman). Al$ though the expected rotational structure of this band

could not be resolved owing to the weakness of the band, the latter is considerably broader than the band associated with the $v_1(a_1)$ band, presumably on this account.

For all tetrahedral molecules, the $\nu_2(e)$ mode lies at a lower frequency than the $\nu_4(t_2)$ mode, this having in general been established by the existence of an infrared-Raman coincidence for the latter but not for the former. Accordingly we have assigned these two fundamentals of the present molecules in the same relative frequency order. The band envelopes of the two bending modes of each molecule tend to overlap, in some cases seriously, owing to pronounced rotational structure of each band. On intensity grounds, 19 this structure is expected to arise from OQS as well as PQR branches. The Q branch of the $\nu_2(e)$ mode shows up clearly in all cases except for hafnium tetraiodide. However, that of the $\nu_4(t_2)$ mode is frequently difficult to locate accurately; this situation arises partly because of the close proximity of the two fundamentals to each other and partly because of Coriolis interactions between the two modes. The detailed form of these interactions and the procedure used to decide upon the value of the $v_4(t_2)$ mode in such cases will be discussed elsewhere in conjunction with a study of the vaporphase spectra of the group IVb tetrahalides.²⁰ However, the situation is typified by that for titanium tetrachloride, for which the separation of the branch maxima at 65° is *ca*. 13 cm⁻¹ for the $\nu_2(e)$ mode and *ca.* 12 cm⁻¹ for the $\nu_4(t_2)$ mode (Figure 2). The two Q branches are readily recognized by the fact that their separation (22 cm^{-1}) is close to that in cyclohexane solution. The situation is more difficult with the zirconium and hafnium tetrachlorides, because in these cases, the intensity of the $\nu_4(t_2)$ mode appears to be considerably less than that of the $\nu_2(e)$ mode. Values for the *Q* branches of each fundamental only are quoted in Table I.

The vapor-phase infrared spectrum of zirconium tetrachloride has been reported⁶ in the range $200-900$ cm^{-1} . In addition to the $v_3(t_2)$ fundamental, a large number of combination bands and overtones were observed. The values for the four fundamentals deduced therefrom $(\nu_1$ 388, ν_2 102, ν_3 421, ν_4 112 cm⁻¹) agree moderately well with those found directly herein (377, 98, 418, and 113 cm-l, respectively).

Force Constant Calculations.---Force constants for the molecules have been calculated 21 both on the basis of the modified valence force field (MVFF) as well as on the basis of the Urey-Bradley force field (UBFF) and the results are given in Tables I1 and 111. As commonly occurs in force constant calculations, two mathematically correct sets of solutions are obtained for the force-field equations. In the MVFF case, the solution has been taken which contains the smaller values of the interaction force constants. The discarded solution contains an interaction force constant which is nearly as large as the stretching force constant. In the UBFF case the solution has been taken which contains the larger stretching force constant. The discarded solution contains a stretching force constant of less than 0.1 mdyn A^{-1} and large negative interaction constants. The force constants are not, of course, chemi-

⁽¹⁷⁾ **A** Suvarov and E K. Krzhizhanovskaya, *Russ. J. Inovg. Chem* , **14, 434** (1969).

⁽¹⁸⁾ J Weidlein, U Muller, and K Dehnicke, *Specfuochim. Ada,* Part *A, 24,* 253 **(1968).**

⁽¹⁹⁾ G Herzberg, "Infraxed arid Ranian Spectra of Polyatomic *Mole*cules," Van Nostrand, Princeton, N. J., 1945.

⁽²⁰⁾ R J **H** Clark and D. 11. Rippon, *Chem. Commun.,* 1297 (1971), and **work** in preparation.

⁽²¹⁾ **A** Muller and B. Krebs, *J. Mol. Spectuosc.,* **24,** 180 **(1907).**

TABLE II MVFF FORCE CONSTANTS^a (MDYN A^{-1}) FOR \mathbf{m} IV Trep \mathbf{m} in \mathbf{m}

THE GROUP IV TETRAHALIDES					
Halide	fr.	fα	frr	Íαα	
TiCl ₄	2.663	0.101	0.165	0.005	
TiBr ₄	2.196	0.083	0.109	0.005	
TiI.	1.640	0.074	0.107	0.004	
ZrCl ₄	2.514	0.081	0.151	0.007	
ZrBr ₄	2.157	0.061	0.078	0.002	
ZrI ₄	1.664	0.052	0.067	0.003	
HfCl ₄	2.631	0.093	0.138	0.011	
HfBr4	2.268	0.069	0.114	0.003	
HfL	1.849	0.078	0.005	0.002	

^{*a*} 1 mdyn Å⁻¹ = 10² N m⁻¹: A. Müller and B. Krebs, *J. Mol.* Spectrosc., 24, 180 (1967). In this approximation, $f_{\alpha\alpha'} = 0$ and $f_{r\alpha} = f_{r\alpha}'$.

^{*a*} 1 mdyn $A^{-1} = 10^2$ N m⁻¹: A. Müller and B. Krebs, *J. Mol.* Spectrosc., 24, 180 (1967).

cally significant to the number of figures quoted, on account of the possible uncertainty in the frequencies of the fundamentals; however, they are quoted as given because in this way they will regenerate the fundamentals to within 0.1 cm^{-1} . The only force constant sufficiently large to warrant discussion is the bond stretching force constant $(f_r$ in the MVFF, K in the UBFF). This constant behaves in the same way in the two force fields. It is very little dependent on the central metal atom, but it decreases in the order $MC1 > MBr >$ MI for both force fields, the average values being 2.60, 2.21, and 1.72, respectively, for the MVFF case, and 2.40, 2.12, and 1.70 mdyn A^{-1} , respectively for the UBFF case. This behavior appears to be typical of halide complexes in this part of the periodic table; in particular, the same dependence of metal-halogen bond stretching force constants on the halogen has been found for the $MX_6{}^{2-}$ ions.²²

Thermodynamic Functions.—Values for a number of thermodynamic functions of the nine tetrahalides have been calculated in each case for 1 mol of the ideal gas at 1 atm pressure and at a number of different temperatures. The functions are $H^{\circ}{}_{T} - H^{\circ}{}_{298}$, $-(G^{\circ}{}_{T} - H^{\circ}{}_{298})/T$, $S^{\circ}{}_{T}$, and C_{p}° . The calculations were based on the ¹²C atomic weight scale and the N. A. S. tables of physical constants, the rigid-rotor harmonic oscillator approximation, a singlet ground electronic state, and a symmetry number of 12. Contributions from nuclear spin, anharmonicity, centrifugal distortion, and isotope mixing to the functions were omitted. The results are given in Tables IV-XII and are based on the values for the fundamentals and bond lengths given in Table I.

Less complete tables of similar data have previously been published for three of the tetrahalides, TiCl₄,

(22) W. van Bronswyk, R. J. H. Clark, and L. Maresca, Inorg. Chem., 8, 1395 (1969).

TABLE IV CALCULATED THERMODYNAMIC PROPERTIES OF TITANIUM TETRACHLORIDE IN THE VAPOR STATE

			$-(G^o r -$	
	C_n °, cal	H° T $ H^{\circ}$ 298, a	H° ₂₀₈ $/T$, cal	S° <i>T</i> , cal
T , \mathcal{C} K	deg ⁻¹ mol ⁻¹	cal mol ⁻¹	deg ⁻¹ mol ⁻¹	deg ⁻¹ mol ⁻¹
100	15.98	$-3,991$	103.03	63.12
150	18.59	-3.123	90.94	70.12
200	20.54	$-2,142$	86.46	75.75
250	21.91	$-1,079$	84.80	80.49
298.15	22.83	0	84.43	84.43
300	22.86	42	84.43	84.57
350	23.52	1,203	84.71	88.15
400	24.00	2,392	85.35	91.33
450	24.34	3,601	86.17	94.17
500	24.60	4.825	87.10	96.75
600	24.96	7.304	89.10	101.27
700	25.18	9,812	91.12	105.14
800	25.32	12.338	93.09	108.51
900	25.43	14,876	94.97	111.50
1000	25.50	17,423	96.76	114.18
		$^a H^{\circ}_{\text{298}} - H^{\circ}_{\text{0}} = 5142 \text{ cal mol}^{-1}.$		

TABLE V CALCULATED THERMODYNAMIC PROPERTIES OF TITANIUM TETRABROMIDE IN THE VAPOR STATE

 $^a H^{\circ}{}_{298} - H^{\circ}{}_{0} = 5714 \text{ cal mol}^{-1}.$

TABLE VI CALCULATED THERMODYNAMIC PROPERTIES OF TITANIUM TETRAIODIDE IN THE VAPOR STATE

			$-(G^{\circ}T -$	
	C_p °, cal	H° T $ H^{\circ}$ ₂₉₈ , ^{<i>a</i>}	$H^\mathrm{o}_{-293})/T$, cal	S° <i>T</i> , cal
$T, \, {}^{\circ}\mathbf{K}$	deg ⁻¹ mol ⁻¹	\degree cal mol $^{-1}$	deg^{-1} mol ⁻¹	$\text{deg}^{-1} \text{ mol}^{-1}$
100	19.82	$-4,554$	123.90	78.35
150	$22.07\,$	$-3,503$	110.20	86.85
200	23.38	$-2,363$	105.21	93.40
250	24.14	$-1,174$	103.40	98.70
298.15	24.59	∘ 0	103.00	103.00
300.	24.60	46	103.00	103.15
350	24.91	1.284	103.30	106.97
400	25.11	2.534	103.97	110.31
450	25.26	3.794	104.84	113.27
500	25.36	5,059	105.82	115.94
600	25.50	7.603	107.90	120.58
700	25.59	10.158	110.00	124.51
800	25.64	12.720	112.03	127.93
900	25.68	15,286	113.97	130.96
1000	25.71	17,856	115.81	133.66
a H° .		$H^0 = 6050$ and maluk		

 $- H^{\circ}{}_{0} = 6050 \text{ cal mol}^{-1}.$ H^{\bullet} 298

TiBr₄, and ZrCl₄. In the case of titanium tetrachloride,¹¹ the data were then based on a slightly inaccurate titanium-chlorine bond length (2.18 Å) and on values for v_1 , v_2 , and v_4 as determined from infrared-active combination bands. The original data on titanium tetrabromide were based on liquid-state values for the fundamentals, which, as discussed earlier, differ significantly from those reported herein. The originally

 $^a H^{\circ}_{198} - H^{\circ}_{0} = 5392 \text{ cal mol}^{-1}.$

TABLE VIII CALCULATED THERMODYNAMIC PROPERTIES OF ZIRCONIUM TETRABROMIDE IN THE VAPOR STATE

			$-(G^{\circ}T -$	
	C_p° , cal	H° T $ H^{\circ}$ ₂₉₈ , ^{<i>a</i>}	H° ₂₉₈ $/T$, cal	S° r, cal
$T,~^{\circ}K$	deg ⁻¹ mol ⁻¹	cal mol ⁻¹	deg ⁻¹ mol ⁻¹	deg^{-1} mol ⁻¹
100	19.37	$-4,521$	119.81	74.59
150	21.84	$-3,486$	106.20	82.96
200	23.25	$-2,355$	101.23	89.45
250	24.06	$-1,171$	99.42	94.74
298.15	24.53	0	99.02	99.02
300	24.55	45	99.02	99.17
350	24.87	1,281	99.32	102.98
400	25.08	2,530	99.99	106.31
450	25.23	3,788	100.86	109.28
500	25.34	5,053	101.84	111.94
600	25.49	7,595	103.92	116.58
700	25.58	10.149	106.01	120.51
800	25.64	12,710	108.05	123.93
900	25.68	15.275	109.98	126.95
1000	25.71	17.845	111.82	129.66

 $^a H^{\circ}_{198} - H^{\circ}_{0} = 5962 \text{ cal mol}^{-1}.$

published tables of thermodynamic functions of zirconium tetrachloride⁶ were also based on somewhat different values for its four fundamentals from those reported herein, values determined (except for $v_3(t_2)$) from combination bands as discussed previously; the zirconium-chlorine bond length was not specified. In none of these three cases, however, do the present

TABLE X CALCULATED THERMODYNAMIC PROPERTIES OF HAFNIUM TETRACHLORIDE IN THE VAPOR STATE

			$-(G^{\circ}T -$	
	C_p °, cal	H° T $ H^{\circ}$ ₂₉₈ a	$H^\circ{}_{^{298}})/T$, cal	S° <i>T</i> , cal
T , \mathbf{K}	deg ⁻¹ mol ⁻¹	cal mol ⁻¹	deg^{-1} mol ⁻¹	$\rm deg^{-1} \ mol^{-1}$
100	17.05	$-4,208$	108.57	66.49
150	19.84	$-3,282$	95.84	73.96
200	21.71	$-2,239$	91.14	79.95
250	22.90	$-1,122$	89.42	84.93
298.15	23.65	0	89.03	89.03
300	23.67	44	89.03	89.18
350	24.18	1,241	89.32	92.87
400	24.54	2,459	89.97	96.12
450	24.79	3.693	90.82	99.03
500	24.98	4,937	91.77	101.65
600	25.23	7,448	93.81	106.23
700	25.38	9.980	95.87	110.13
800	25.49	12,524	97.87	113.53
900	25.56	15,076	99.78	116.53
1000	25.61	17,635	101.59	119.23
TTTO	YTO	F 4 0 F 4 - 1 - 1 - 1 - 1		

 $^{a}H^{o}{}_{298}-H^{o}{}_{0}=5435$ cal mol⁻¹.

TABLE XI CALCULATED THERMODYNAMIC PROPERTIES OF HAFNIUM TETRABROMIDE IN THE VAPOR STATE

			$-(G^{\circ}r -$	
	C_v° , cal	H° T $ H^{\circ}$ 298, a	$H^{\circ_{298}}/T$, cal	S° <i>T</i> , cal
T , \circ K	deg ⁻¹ mol ⁻¹	cal mol ⁻¹	deg ⁻¹ mol ⁻¹	deg ⁻¹ mol ⁻¹
100	19.84	$-4,597$	121.20	75.23
150	22.33	-3.536	107.37	83.80
200	23.63	$-2,384$	102.34	90.42
250	24.34	-1.183	100.51	95.78
298.15	24.75	0	100.10	100.10
300	24.76	46	100.10	100.26
350	25.03	1,291	100.41	104.09
400	25.21	2.548	101.08	107.45
450	25.34	3.812	101.96	110.53
500	25.43	5,081	102.94	113.10
600	25.55	7.630	105.03	117.75
700	25.62	10.190	107.14	121.69
800	25.67	12,755	109.18	125.12
900	25.71	15,324	111.12	128.15
$1000\,$	25.73	17.896	112.96	130.86

 $^{a} H^{\circ}_{198} - H^{\circ}_{0} = 6049 \text{ cal mol}^{-1}.$

 $$ CALCULATED THERMODYNAMIC PROPERTIES OF HAFNIUM TETRAIODIDE IN THE VAPOR STATE

			$-(G^{\circ}r -$	
	C_p ^o , cal	$H^{\circ}{}_{T} \, - \, H^{\circ}{}_{^{298}}{}^{a}$	H° ₂₉₈)/T, cal	S° <i>T</i> , cal
$T, \ ^{\circ}K$	deg ⁻¹ mol ⁻¹	cal mol ⁻¹	$\text{deg}^{-1} \text{ mol}^{-1}$	deg ⁻¹ mol ⁻¹
100	21.28	-4754	127.62	80.09
150	23.38	-3630	113.37	89.17
200	24.34	-2434	108.21	96.04
250	24.84	-1203	106.35	101.53
298.15	$25.12\,$	0	105.94	105.94
300	25.13	46	105.94	106.09
350	25.31	1,308	106.24	109.98
400	25.43	2.577	106.93	113.37
450	25.51	3.850	107.81	116.37
500	25.57	5.127	108.80	119.06
600	25.65	7.689	110.91	123.73
700	25.70	10.257	113.03	127.69
800	25.73	12.828	115.09	131.12
900	25.75	15,402	117.04	134.15
1000	25.77	17,978	118.89	$136.87\,$
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 $^{a}H^{\circ}_{198} - H^{\circ}_{0} = 6307 \text{ cal mol}^{-1}.$

tables differ greatly from those mentioned above. In addition to the above, however, similar data are available in the "JANAF Thermochemical Tables"²³ for the tetrahalides of titanium and zirconium. The data therein for titanium tetrachloride and titanium tetra-

(23) "JANAF Thermochemical Tables," Air Force Contract AF 04(611)-7554, 1965.

bromide are based on the fundamental frequencies reported by Hawkins and Carpenter¹² and by Miller and Carlson¹⁴ which have already been discussed; the maximum difference between the literature value and the present value for any fundamental is 5 cm^{-1} for the tetrachloride and 10 cm^{-1} for the tetrabromide. The data in the JANAF tables for zirconium tetrachloride are in error owing to their being based on a grossly inaccurate value²⁴ for $v_1(a_1)$ (315 cm⁻¹ *vs.* the correct value of 377 cm^{-1}) and on estimated values for $\nu_2(e)$ and $\nu_4(t_2)$ (these being in error by up to 6 cm⁻¹). The data in the JANAF tables for the halides TiI4, Zr-Br4, and ZrI4 are entirely based on estimated values for

(24) Ya. \$3. **Bobavich, Ogl.** *Speklvosk.,* **11, 85 (1961).**

the four fundamentals in each case; these estimates are in error to the maximum extent of 63 cm^{-1} for titanium tetraiodide, 9 cm-' for zirconium tetrabromide, and **17** cm-I for zirconium tetraiodide. Tables VII-XI1 therefore place, for the first time, the very close similarities between the tetrahalides of zirconium and hafnium on a quantitative basis.25

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(25) NOTE ADDED IN PROOF.-Revised values for the fundamentals and the thermodynamic properties of lead tetrachloride have recently been published; see R. J. **H. Clark and B. K. Hunter,** *J. Mol. Stvuct.,* **9, 364 (1971).**

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Kinetics and Mechanism of the Reduction of Hydrazoic Acid by Chromium(II) and Vanadium(II)

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The rates of reduction of HN_3 by $V(II)$ and $Cr(II)$ are reported. Both reactions produce 1 mol of N_2 and 1 mol of NH_3 and consume 2 mol of reductant. Both reductions proceed by mixed second-order kinetics with the rate constant independent of [H⁺] over the range 0.1–1.0 *M*. For the rate law defined by $-d[HN_3]/dt = k[HN_3]$ [reductant], the rate param
are: for Cr²⁺ $k = 18$ *M*⁻¹ sec⁻¹, $\Delta H^{\pm} = 7.1$ kcal mol⁻¹, and $\Delta S^{\pm} = -29$ cal mol⁻¹ deg⁻ a Cr-N bond in the first step. The fate of the unstable species produced thereafter is discussed. For the rate law defined by $-d[HN_3]/dt = k[HN_3]$ [reductant], the rate parameters The reaction of Cr^{2+} with HN_3 produces

Introduction

The role a reducible ligand plays in redox reactions between metal ions has become a topic of considerable interest since Nordmeyer and Taube produced convincing evidence of a radical ion intermediate in the Cr^{2+} reduction of isonicotinamidepentaamminecobalt-(111) complex.' One feature of this interest involves the study of the rate and mechanism of reduction of free ligands. A detailed study of the V^{2+} reduction of cobalt(III)-azide complexes² made information about the V2+ reduction of hydrazoic acid necessary. Several preliminary experiments indicated that the Cr²⁺ reduction of HN_3 , which has been examined previously, $3,4$ was not satisfactorily understood. In this report a description of the kinetics and mechanism of the reduction of $HN₃$ by these two reagents is given. Of interest in these noncomplimentary reactions is the fate of unstable, intermediate, oxidation states ; these are discussed.

Experimental Section

Reagents and Analyses.-The sodium azide was an Eastman product and was recrystallized from water as described by Browne.6 A solution of this solid was analyzed by adding the $\text{Na}\,\text{N}_3$ solution to excess $\text{Ce}(IV)$ and titrating that excess with standard Fe(II) solution to a ferroin end point.⁶

Solutions of Cr(I1) were derived from Cr(II1) solutions by reduction with amalgamated Zn. The Cr(II1) stock solution was prepared by H_2O_2 reduction of CrO_3 in HCO_4 . Total $Cr(III)$ content was determined by oxidation with peroxide of an aliquot made basic using ϵ 4815 at 3730 Å.⁷ The $[ClO_4^-]$ of the stock Cr(II1) solution was determined by ion exchange of an aliquot on Dowex $50W-X8$ in the H⁺ form and titration of the acid released with standard base. Reducing power of the $Cr(II)$ solution was measured by adding an aliquot to deaerated $Ce(IV)$ solution in H_2SO_4 medium⁸ followed by titration of the excess $Ce(IV)$ with standard $Fe(II)$. The concentration of total cations (exclusive of H^+) was determined by heating an aliquot of the Cr(I1) solution with excess EDTA buffered at pH 5 for 5-15 min at 80° ;⁹ after raising the pH to 10 with NH_a-NH_4Cl buffer, the excess EDTA was titrated with standard **Zn2+** to an Erio-Black-T end point.

Solutions of VO²⁺ were prepared either by electrolytic reduction
of V₂O₅ in HClO₄ acid or by treatment of a solution of VOSO₄ with an anion-exchange resin in the $ClO₄$ form. No differences in rate behavior between the two solutions were observed. Reduction was accomplished with amalgamated Zn. When not in use this solution of $\rm V(II)$ was stored at -78° in order to prevent reduction of ClO_4^- . Total V was determined spectrophotometrically as the peroxide complex;¹⁰ [ClO₄⁻] was determined as outlined above. Analysis for total reducing power utilized the ceriometric technique described for $[Cr(II)]$, except that the excess $Ce(IV)$ was titrated with $Fe(II)$ in a solution >5 *M* in H₂SO₄ in order to reduce VO_2 ⁺ before $Fe(phen)_3^{3+11}$ The $[V^{2+}] + [V^{3+}] + [Zn^{2+}]$ was determined by an EDTA titration after oxidation of V^{2+} with ClO₂⁻ or NO₃⁻.

(7) G. Haupt, *J. Res. Nat. Bur. Stand.*, **48**, 414 (1952). $ClO₂$ – or $NO₃$ –.

⁽¹⁾ F. Nordmeyer and H. Taube, *J. Amev. Chem.* Soc., **90, 1163 (1968).**

⁽²⁾ K. Hicks, D. **L. Toppen, abd** R. **G. Linck, unpublished observations. (3) M. Ardon and B. E. Mayer,** *J. Chem.* **Soc., 2819 (1962).**

⁽⁴⁾ *C.* **F. Wells and M. A. Salem,** *ibid., A,* **1568 (1968).**

⁽⁵⁾ A. W. **Browne,** *Inorg. Syn.,* **1, 79 (1939).**

⁽⁶⁾ J. W. **Arnold,** *Ind.* **Eng.** *Chem., Anal. Ed., 11,* **215 (1945).**

⁽⁸⁾ J. Y. Tong and E. L. King, *J. Amev. Chem.* Soc., *8%* **3805 (1960).**

⁽⁹⁾ G. Doppler and R. **Patzak,** *Z. Anal. Chem.,* **162, 45 (1956).**

⁽¹⁰⁾ P. R. Guenther and R. *G.* **Linck,** *J. Amer. Chem. Soc.,* **91,3769 (1969). (11) G. P. Walden,** Jr., **L.** P. **Hammett, and** S. **M. Edmonds,** *ibid.,* **66, 57 (1934).**