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**Chlorine Trifluoride Oxide. III. Vibrational Spectrum, Force Constants, and Thermodynamic Properties**

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The infrared spectra of gaseous, solid, and matrix isolated  $\text{ClF}_3\text{O}$  and the Raman spectra of gaseous and liquid  $\text{ClF}_3\text{O}$  are reported. Nine fundamental vibrations were observed, consistent with symmetry  $C_s$ . The vibrational spectrum of  $\text{ClF}_3\text{O}$  agrees well with a trigonal-bipyramidal model with two fluorine atoms at the apexes and one fluorine atom, one oxygen atom, and one localized free electron pair at the remaining corners. A modified valence force field was computed for  $\text{ClF}_3\text{O}$ . These data indicated that the axial ClF bonds ( $f_r = 2.34$  mdyne/Å) are considerably weaker than the equatorial one ( $f_r = 3.16$  mdyne/Å) and that the bond order of the ClO bond is close to two ( $f_D = 9.37$  mdyne/Å). The Raman spectrum of the liquid and the infrared spectrum of the solid indicate association through the axial fluorine atoms. Thermodynamic properties were computed for  $\text{ClF}_3\text{O}$  in the range 0–2000°K.

**Introduction**

As part of our extensive study<sup>1–5</sup> of the novel chlorine oxyfluoride,  $\text{ClF}_3\text{O}$ , we have investigated the vibrational spectra of the molecule. In their recent paper, Bougon, *et al.*, have proposed<sup>6</sup> a structure of symmetry  $C_{3v}$  for  $\text{ClF}_3\text{O}$  on the basis of infrared and Raman spectra. In this paper, we report the complete vibrational spectrum, force constants, and thermodynamic properties of  $\text{ClF}_3\text{O}$ . The vibrational spectra, force constants, and structure of the two ions,  $\text{ClF}_2\text{O}^+$  and  $\text{ClF}_4\text{O}^-$ , derived from  $\text{ClF}_3\text{O}$  are discussed elsewhere.<sup>7,8</sup>

**Experimental Section**

The preparation of  $\text{ClF}_3\text{O}$  and its purification and handling have previously been described.<sup>3,8</sup> The  $\text{ClF}_3\text{O}$  samples used in this investigation had a purity of better than 99.8% and our spectra did not reveal any impurities. The infrared spectra were recorded either on a Beckman Model IR-7 spectrophotometer with NaCl and CsI interchange in the range 4000–200  $\text{cm}^{-1}$  or a Perkin-Elmer Model 457 spectrophotometer in the range 4000–250  $\text{cm}^{-1}$ . The instruments were calibrated by comparison with standard gas calibration points.<sup>9</sup> The gas cells were made of 304 stainless steel and had a path length of 5 cm. Silver chloride and CsI windows were used in the ranges 4000–450 and 450–200  $\text{cm}^{-1}$ , respectively. The inside of the CsI windows was protected against chemical attack by Teflon FEP sheet (1-mil thickness). To ensure close contact between the CsI window and the Teflon sheet, a thin film of Halocarbon grease (high-temperature grade from Halocarbon Products) was put on the CsI window prior to attaching the Teflon sheet. The Teflon sheet did not come loose or show blisters even upon evacuation of the cell provided all gas pockets between the CsI window and the thin Teflon sheet were carefully eliminated before assembling the cell. Compensation for bands due to the window material was achieved by placing an empty cell into the reference beam.

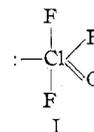
The apparatus, materials, and technique used for the matrix isolation study have previously been described.<sup>10,11</sup> Raman

spectra were recorded using a Coherent Radiation Laboratories Model 52 Ar laser as a source of  $\sim 1.3$  W of exciting light at 5145 Å. The scattered light was analyzed with a Spex Model 1400 double monochromator, a photomultiplier, and a dc ammeter. Polarization measurements were carried out using a Model 310 polarization rotator from Spectra-Physics. A stainless steel cell with Teflon O-rings and sapphire windows was used for obtaining the spectra of gases and liquids. The design of this cell was similar to that of a cell described previously.<sup>12</sup>

**Results and Discussion**

**Vibrational Spectra.**—Figure 1 shows the infrared spectrum of gaseous  $\text{ClF}_3\text{O}$  recorded at various sample pressures. Figures 2 and 3 show the Raman spectra of gaseous and liquid  $\text{ClF}_3\text{O}$ , respectively. Figure 4 shows the infrared spectra of pure and  $\text{N}_2$ -matrix isolated  $\text{ClF}_3\text{O}$  at 4°K. Figure 5 shows the most intense infrared bands of matrix isolated  $\text{ClF}_3\text{O}$  under high resolution conditions allowing the determination of the <sup>35</sup>Cl–<sup>37</sup>Cl isotopic shifts. Table I lists the observed frequencies which can be attributed to fundamental vibrations. Table II lists the infrared frequencies attributed to overtones and combination bands.

In  $\text{ClF}_3\text{O}$ , the chlorine central atom possesses a non-bonding valence-electron pair. In related molecules (*e.g.*,  $\text{SF}_4$  and  $\text{ClF}_5$ ), such an electron pair appears to be sterically active, and similar behavior is anticipated here. The resulting structure can best be described as a trigonal bipyramid. In this type of structure, the apical positions are always occupied by the most electronegative ligands.<sup>13</sup> Hence, structure I of symmetry



$C_s$  (possessing one symmetry plane as the only symmetry element) might be expected for  $\text{ClF}_3\text{O}$ . The arrangement of the axial  $\text{FCIF}$  group may not be exactly linear owing to the different degrees of repulsion expected from the three equatorial ligands. However, possible deviations from 180° should be relatively small and symmetry  $C_s$  would be retained. A second structure of symmetry  $C_s$  (II) can be written for  $\text{ClF}_3\text{O}$  in which one fluorine and one oxygen atom occupy the axial positions. However, comparison with related

(1) D. Pilipovich, R. D. Wilson, and H. F. Bauer, U. S. Patent, to be issued, 1972.

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(3) D. Pilipovich, H. H. Rogers, and R. D. Wilson, *ibid.*, **11**, 2192 (1972).

(4) C. J. Schack, C. B. Lindahl, D. Pilipovich, and K. O. Christe, *ibid.*, **11**, 2201 (1972).

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(11) K. O. Christe, *Spectrochim. Acta, Part A*, **27**, 631 (1971).

(12) E. L. Gasner and H. H. Claassen, *Inorg. Chem.*, **6**, 1937 (1967).

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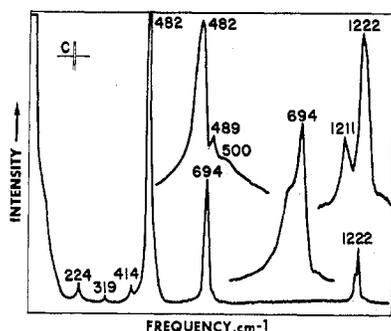


Figure 2.—Raman spectrum of gaseous  $\text{ClF}_3\text{O}$  at a sample pressure of 490 mm. The inserts show the three most intense bands at higher resolution and scale expansion. C indicates equivalent slit width.

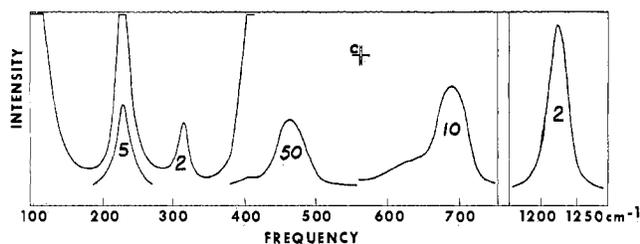


Figure 3.—Raman spectrum of liquid  $\text{ClF}_3\text{O}$ . C indicates equivalent slit width.

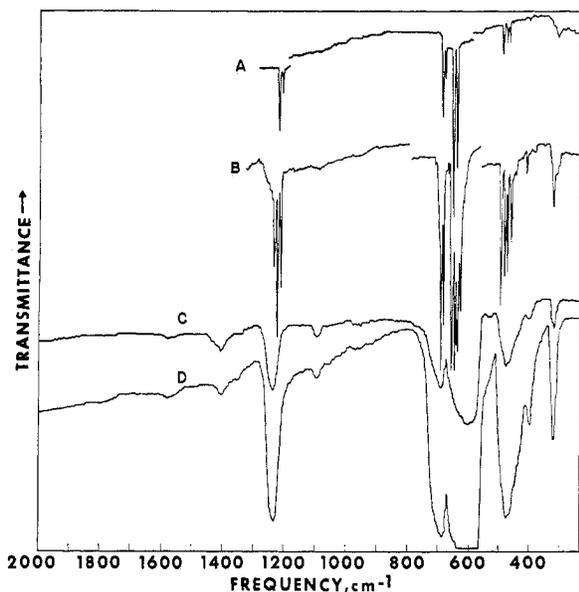
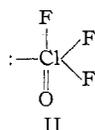


Figure 4.—Infrared spectra of solid  $\text{ClF}_3\text{O}$  at 4°K: (A) 1.0  $\mu\text{mol}$  of  $\text{ClF}_3\text{O}$  in 10,000  $\mu\text{mol}$  of  $\text{N}_2$ ; (B) 10  $\mu\text{mol}$  of  $\text{ClF}_3\text{O}$  in 10,000  $\mu\text{mol}$  of  $\text{N}_2$  after controlled diffusion at 27°K; (C and D) pure  $\text{ClF}_3\text{O}$  at two different concentrations.



molecules and the observed  $^{35}\text{Cl}$ – $^{37}\text{Cl}$  isotopic shifts (see below) eliminate this model. If the free electron pair of the chlorine atom in  $\text{ClF}_3\text{O}$  would not be sterically active, a tetrahedral molecule of symmetry  $C_{3v}$  would result, as proposed<sup>6</sup> by Bougon, *et al.*

For structures I and II, one would expect nine fun-

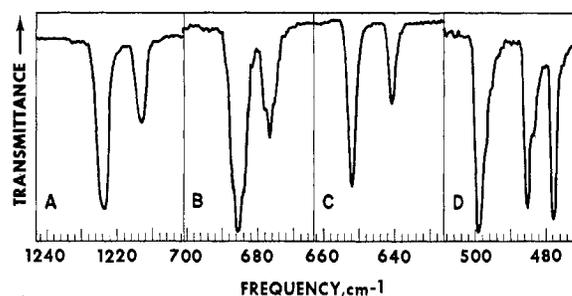
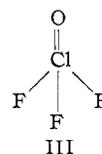


Figure 5.—Infrared spectra of  $\text{ClF}_3\text{O}$  in  $\text{N}_2$  matrix at 4°K recorded at tenfold scale expansion under high resolution conditions: (A) 2.9  $\mu\text{mol}$  of  $\text{ClF}_3\text{O}$  (matrix ratio (mr) 1:1000); (B) 2.0  $\mu\text{mol}$  (mr 1:10,000); (C) 0.24  $\mu\text{mol}$  (mr 1:10,000); (D) 10  $\mu\text{mol}$  (mr 1:1000).

TABLE II  
OVERTONES AND COMBINATION BANDS OBSERVED IN  
THE INFRARED SPECTRUM OF  $\text{ClF}_3\text{O}$  GAS

Frequencies, $\text{cm}^{-1}$			Assignment
Obsd	Calcd		
800 vvw, br	$475 + 318 = 793$	$(\nu_8 + \nu_5)$	
899 vw	$408 + 491 = 899$	$(\nu_4 + \nu_3)$ or $(\nu_4 + \nu_8)$	
980 w	$2 \times 491 = 982$	$2\nu_3$ or $2\nu_8$ or $(\nu_3 + \nu_3)$	
1161 mw	$491 + 680 = 1171$	$(\nu_3 + \nu_7)$ or $(\nu_3 + \nu_7)$	
1372 mw	$2 \times 680 = 1360$	$2\nu_7$	
1899 vw	$680 + 1221 = 1901$	$(\nu_7 + \nu_1)$	
2436 w	$2 \times 1221 = 2442$	$2\nu_1$	



damental vibrations active in both the infrared and Raman spectrum. For structure III, one would expect only six fundamental vibrations, again, active in both the infrared and Raman spectrum. Since nine fundamental vibrations were unambiguously observed for  $\text{ClF}_3\text{O}$  (see Table I), structure III can be ruled out. Consequently, assignments for  $\text{ClF}_3\text{O}$  were made in point group  $C_s$  (see Table I) contrary to the conclusions<sup>6</sup> reached by Bougon, *et al.*, on the basis of incomplete experimental data. Further evidence against symmetry  $C_{3v}$  is based on the band shapes observed for the infrared spectrum of the gas showing that  $\text{ClF}_3\text{O}$  cannot be a symmetric top molecule. For example, in the case of  $C_{3v}$ , the  $\text{Cl}=\text{O}$  stretching mode at about 1220  $\text{cm}^{-1}$  should show well-separated PQR branches. Since the geometry of  $\text{ClF}_3\text{O}$  of symmetry  $C_s$  can be estimated (see below), the three principal moments of inertia were computed for structure I resulting in  $A = 0.216$ ,  $B = 0.129$ , and  $C = 0.0934 \text{ cm}^{-1}$ . Based on these values, the infrared band contours were estimated for  $\text{ClF}_3\text{O}$ , according to the method of Ueda and Shimanouchi.<sup>14</sup> The observed band contours were consistent with those predicted. However, they cannot be used to distinguish structures I and II due to the interference by the  $^{37}\text{Cl}$  isotope bands, the double and triple coincidences of bands in the infrared spectrum occurring at about 680 and 490  $\text{cm}^{-1}$ , respectively, and the fact that most bands are blends of the A, B, and C type.

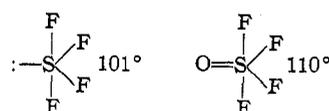
The assignments (see Table I) for  $\text{ClF}_3\text{O}$  in point group  $C_s$  (structure I) are based on the following argu-

(14) T. Ueda and T. Shimanouchi, *J. Mol. Spectrosc.*, **28**, 350 (1968).

ments. The band at about  $1220\text{ cm}^{-1}$  occurs at a frequency much too high for any Cl—F fundamental vibration and, hence, must be assigned to the stretching vibration involving the Cl=O double bond. Comparison with related Cl—F containing compounds<sup>15-18</sup> indicates that of the three Cl—F stretching vibrations, at least two should show a frequency higher than  $600\text{ cm}^{-1}$  (see Table I). Indeed two bands were observed at about  $690$  and  $650\text{ cm}^{-1}$ . These were assigned to the equatorial Cl—F and the antisymmetric axial FCIF stretching vibrations, respectively. These assignments were made on the basis of the observed  $^{35}\text{Cl}$ — $^{37}\text{Cl}$  isotopic shifts (see below), the relative infrared and Raman intensities, and the Raman polarization measurements. The fourth stretching vibration, *i.e.*, the symmetric FCIF mode, should be polarized and of the highest relative intensity in the Raman spectrum. Furthermore, it should be relatively weak in the infrared spectrum and show no detectable Cl isotopic shift if the F—Cl—F group is approximately linear. Hence, this mode must be assigned to the band at about  $480\text{ cm}^{-1}$ .

The remaining five bands must be assigned to deformation modes. Of these, three involve a motion of the doubly bonded oxygen atom and, consequently, should occur at higher frequencies. The two lowest frequency bands are assigned to the two F—Cl—F bending modes,  $\nu_5$  and  $\nu_6$ , by analogy with  $\text{ClF}_3$ .<sup>15</sup> Of the remaining modes, the torsional motion should result in a band of very low intensity in both the infrared and the Raman spectrum. Obviously, only the  $414\text{-cm}^{-1}$  band fulfills these requirements. The two highest frequency bands should belong to the F'—Cl=O scissoring and the rocking mode. On the basis of the observed and computed (see below) Cl isotopic shifts, we prefer to assign the scissoring mode,  $\nu_8(A')$ , to the  $480\text{-cm}^{-1}$  band and the rocking mode,  $\nu_8(A'')$ , to  $500\text{ cm}^{-1}$ . These assignments are supported by simple valence force field arguments and by comparison with the vibrational spectra of the related species,  $\text{ClF}_3$ ,<sup>15</sup>  $\text{ClF}_2^-$ ,<sup>16</sup>  $\text{ClF}_2\text{O}_2^-$ ,<sup>17</sup> and  $\text{ClO}_2\text{F}$ <sup>18</sup> (see Table I).

**Force Constants and Bonding.**—The potential and kinetic energy metrics for chlorine trifluoride oxide were computed by a machine method.<sup>19</sup> The geometry assumed for this computation was  $D(\text{ClO}) = 1.42\text{ \AA}$ ,  $R(\text{ClF}') = 1.62\text{ \AA}$ ,  $r(\text{ClF}) = 1.72\text{ \AA}$ ,  $\alpha(\text{OCIF}') = 120^\circ$ ,  $\beta(\text{OCIF}) = \gamma(\text{FCIF}') = 90^\circ$ , based on the observed geometry for  $\text{ClF}_3$ <sup>20</sup> and a correlation between ClO bond length and stretching frequency.<sup>21</sup> In the absence of structural data, we assumed the oxygen atom, the fluorine atom, and the lone pair to require about the same space and, hence, to be about  $120^\circ$  apart. However, comparison with the known geometry of the trigonal-bipyramid type molecules  $\text{SF}_4$ <sup>22</sup> and  $\text{SF}_4\text{O}$ <sup>23</sup> indicates that the repulsion from a double bonded



oxygen atom should be intermediate between that of a free electron pair and that of a fluorine ligand. Consequently, the true structure of  $\text{ClF}_3\text{O}$  might be expected to deviate somewhat from that used for our computations.

The symmetry coordinates used for  $\text{ClF}_3\text{O}$  are given in Table III. The bending coordinates were weighted

TABLE III  
SYMMETRY COORDINATES FOR  $\text{ClF}_3\text{O}$

A'	$S_1$	$D(\text{ClO})$
	$S_2$	$R(\text{ClF}')$
	$S_3$	$\alpha(\text{OCIF}')$
	$S_4$	$(1/\sqrt{2})[r_1(\text{ClF}) + r_2(\text{ClF})]$
A''	$S_5$	$(1/\sqrt{2})[\beta_1(\text{OCIF}) + \beta_2(\text{OCIF})]$
	$S_6$	$(1/\sqrt{2})[\gamma_1(\text{FCIF}') + \gamma_2(\text{FCIF}')]$
	$S_7$	$(1/\sqrt{2})[r_1(\text{ClF}) - r_2(\text{ClF})]$
	$S_8$	$(1/\sqrt{2})[\beta_1(\text{OCIF}) - \beta_2(\text{OCIF})]$
	$S_9$	$(1/\sqrt{2})[\gamma_1(\text{FCIF}') - \gamma_2(\text{FCIF}')]$

by unit ( $1\text{ \AA}$ ) distance so the stretching force constants have units of  $\text{mdyn}/\text{\AA}$ , the deformation force constants units of  $\text{mdyn \AA}/\text{radian}^2$ , and the stretch-bend interaction constants  $\text{mdyn}/\text{radian}$ . The  $G$  matrix and  $Z$  transformation were found by the computer and, hence, are not given here.

The force constants were adjusted by trial and error with the aid of a computer to give an exact fit between the observed and computed frequencies. The computed force constants are given in Table IV. An un-

TABLE IV  
INTERNAL FORCE CONSTANTS OF  $\text{ClF}_3\text{O}^a$

$f_D$	9.37	$f_{rr}$	0.26
$f_R$	3.16	$f_{\beta\beta}$	0.11
$f_r$	2.34	$f_{\gamma\gamma}$	0.13
$f_\alpha$	1.84	$f_{r\beta} = -f_{r\beta'}$	0.25
$f_\beta$	1.69	$f_{\beta\gamma} = f_{\beta\gamma'}$	0.22
$f_\gamma$	1.87		

<sup>a</sup> Stretching constants in  $\text{mdyn}/\text{\AA}$ , deformation constants in  $\text{mdyn \AA}/\text{radian}^2$ , and stretch-bend interaction constants in  $\text{mdyn}/\text{radian}$ .

certainty estimate is difficult to make. However, numerical experiments with this and similar molecules indicate the uncertainties in the valence force constants to be 0.2 to 0.3 for Cl=O and  $0.2\text{ mdyn}/\text{\AA}$  for Cl—F. The force constants of greatest interest are the stretching force constants since they are a direct measure for the strength and covalent character of the various bonds. The value of  $9.37\text{ mdyn}/\text{\AA}$  obtained for  $f_{\text{Cl=O}}$  is similar to those computed for  $\text{ClO}_2\text{F}$  ( $9.07\text{ mdyn}/\text{\AA}$ )<sup>18</sup> and  $\text{ClO}_2^+$  ( $8.96\text{ mdyn}/\text{\AA}$ )<sup>24</sup> indicating double bond character. The value of  $2.34\text{ mdyn}/\text{\AA}$  computed for the axial Cl—F stretching force constant  $f_r$  is almost identical with that of  $2.34\text{ mdyn}/\text{\AA}$ , previously calculated<sup>16</sup> for  $\text{ClF}_2^-$ . The corresponding interaction constant,  $f_{rr}$ , is also very similar for both species. The relatively low value of  $f_r$  in  $\text{ClF}_2^-$  has previously been interpreted<sup>16</sup> in terms of semiionic three-center four-electron bonds. The same rea-

(15) H. Selig, H. H. Claassen, and J. H. Holloway, *J. Chem. Phys.*, **52**, 3517 (1970).

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(19) E. C. Curtis, *Spectrochim. Acta, Part A*, **27**, 1989 (1971).

(20) D. F. Smith, *J. Chem. Phys.*, **21**, 609 (1953).

(21) E. A. Robinson, *Can. J. Chem.*, **41**, 3021 (1963).

(22) K. Kimura and S. H. Bauer, *J. Chem. Phys.*, **39**, 3172 (1963).

(23) J. L. Hencher, D. W. Cruickshank, and S. H. Bauer, *ibid.*, **48**, 518 (1968).

(24) K. O. Christe, C. J. Schack, D. Pilipovich, and W. Sawodny, *Inorg. Chem.*, **8**, 2489 (1969).

soning<sup>16</sup> holds for the axial ClF bonds of ClF<sub>3</sub>O and, hence, will not be repeated here. It should be pointed out, however, that in ClF<sub>3</sub>O, enhancement of the ionic character of the axial ClF bonds is due to oxygen substitution, whereas in ClF<sub>2</sub><sup>-</sup> it is due to the formal negative charge. Since the influence of oxygen substitution on the remaining ClF bonds has previously been discussed<sup>17</sup> at length for ClO<sub>2</sub>F<sub>2</sub><sup>-</sup>, it will not be reiterated. The value of 3.16 mdyn/Å computed for the equatorial ClF bond of ClF<sub>3</sub>O is considerably larger than that of the axial bonds indicating<sup>16</sup> predominantly covalent bonding. In summary, the results from the force constant analysis are in excellent agreement with a generalized bonding scheme previously discussed<sup>25</sup> for a large number of halogen fluorides and suggest that the overall bonding in ClF<sub>3</sub>O might be described by the following approximation. The bonding of the three equatorial ligands (including the free electron pair on Cl as a ligand and ignoring the second bond of the Cl=O double bond) is mainly due to a sp<sup>2</sup> hybrid, whereas the bonding of the two axial ClF bonds involves mainly one delocalized p-electron pair of the chlorine atom for the formation of a semiionic three-center four-electron pσ bond.

The deformation force constants computed for ClF<sub>3</sub>O are as expected. The values for interaction constants cannot be unique; those not given in Table IV were assumed zero. The values for  $f_{rr}$ ,  $f_{\beta\beta}$ ,  $f_{\gamma\gamma}$ , and  $f_{r\beta}$  are in accord with those for similar molecules.<sup>26</sup> Fitting  $\nu_5$  and  $\nu_6$  required nonzero values for  $f_{\beta\gamma}$  and  $f_{\beta\gamma'}$ . Similarly,  $\nu_7$  and  $\nu_8$  required nonzero values for  $f_{r\beta}$  and  $f_{r\beta'}$ . While these four interaction constants are strictly indeterminate, numerical experiments showed that the observed frequencies could not be fitted with significantly different values.

The computed potential energy distribution is given in Table V. The results were normalized, but the sums

TABLE V  
POTENTIAL ENERGY DISTRIBUTION FOR ClF<sub>3</sub>O

Assignment	Frequency, cm <sup>-1</sup>	Potential energy distribution
A' $\nu_1$	1223	0.96 $f_D$
$\nu_2$	694	0.76 $f_R$ + 0.16 $f_\alpha$
$\nu_3$	490	0.76 $f_\alpha$ + 0.17 $f_r$
$\nu_4$	482	0.90 $f_r$ + 0.10 $f_{rr}$
$\nu_5$	319	0.61 $f_\gamma$ + 0.53 $f_\beta$ - 0.14( $f_{\beta\gamma}$ + $f_{\beta\gamma'}$ )
$\nu_6$	227	0.41 $f_\beta$ + 0.30 $f_\gamma$ + 0.09( $f_{\beta\gamma}$ + $f_{\beta\gamma'}$ )
A'' $\nu_7$	652	0.89 $f_r$ + 0.14 $f_\beta$ - 0.09( $f_{r\beta}$ + $f_{r\beta'}$ )
$\nu_8$	500	0.97 $f_\beta$
$\nu_9$	414	0.80 $f_\gamma$

do not add up to 1.0 since the less important terms are not listed. As can be seen from Table V, most vibrations are highly characteristic except for  $\nu_5$  and  $\nu_6$ , which are almost an equal mixture of the symmetry coordinates, S<sub>5</sub> and S<sub>6</sub>. The  $\nu_5$  and  $\nu_6$  modes can be visualized as a symmetric bending motion of the two axial fluorine atoms in a plane perpendicular to the direction of the lone pair, and  $\nu_6$  as a symmetric bending motion in the plane containing the lone pair.

The <sup>35</sup>Cl-<sup>37</sup>Cl isotopic shifts computed for ClF<sub>3</sub>O from the above force constants are listed in Table VI.

TABLE VI  
OBSERVED AND COMPUTED <sup>35</sup>Cl-<sup>37</sup>Cl ISOTOPIC SHIFTS FOR ClF<sub>3</sub>O

Normal mode	Isotopic shift, cm <sup>-1</sup>	
	Computed	Obsd
A' $\nu_1$	11.8	10.9
$\nu_2$	9.0	8.8
$\nu_3$	2.7	2.0
$\nu_4$	0	0
$\nu_5$	1.4	
$\nu_6$	0.8	
A'' $\nu_7$	11.2	11.2
$\nu_8$	0.5	1.2 or less
$\nu_9$	0	

As can be seen, the agreement between the observed and computed shifts is good and lends further support to the assignments suggested above. Thus, the possibility of assigning the bands at 686, 652, and 478 cm<sup>-1</sup> to  $\nu_{\text{sym}}(\text{ClF}_2(\text{eq}))$ ,  $\nu_{\text{asym}}(\text{ClF}_2(\text{eq}))$ , and  $\nu(\text{ClF}(\text{ax}))$ , respectively, of structure II can be ruled out based upon the lack of observing any <sup>35</sup>Cl-<sup>37</sup>Cl isotopic splitting for the 478-cm<sup>-1</sup> band.

**Association in the Liquid and Pure Solid.**—The relatively high boiling point and Trouton constant<sup>2</sup> of ClF<sub>3</sub>O imply its association in the liquid phase. More specific evidence about the nature of this association can be obtained from the vibrational spectra recorded for the liquid and the solid and from a controlled diffusion experiment carried out for matrix isolated ClF<sub>3</sub>O. Principally, association could involve either oxygen or fluorine bridges. In the case of fluorine bridges, distinction might be made between equatorial and axial fluorine bridges. The stretching frequencies of the bonds involved in the bridging should decrease when going from the gas or matrix isolated solid to the liquid or pure solid. On the other hand, the bonds not involved in the bridging should show no decrease or possibly a very small increase in frequency. Comparison between the infrared spectra of matrix isolated (trace A, Figure 4) and pure solid ClF<sub>3</sub>O (trace C, Figure 4) reveals that the Cl=O stretching mode frequency increases by about 10 cm<sup>-1</sup>, that the equatorial ClF stretching mode shows no detectable change, but that the antisymmetric axial FCIF stretching frequency decreases by about 50 cm<sup>-1</sup>. Similarly, the controlled diffusion experiment carried out for matrix isolated ClF<sub>3</sub>O (trace B, Figure 4) shows the appearance of new bands on the high-frequency side of  $\nu(\text{Cl}=\text{O})$  and  $\nu(\text{Cl}-\text{F}')$  and on the low frequency side of  $\nu_{\text{as}}(\text{FCIF})$ . Furthermore, a weak band observed at 468 cm<sup>-1</sup> in the infrared spectrum of matrix isolated ClF<sub>3</sub>O must be due to associated ClF<sub>3</sub>O owing to its strong increase in relative intensity upon controlled diffusion. In addition, the Raman spectrum of the liquid (Figure 3) shows a pronounced frequency decrease only for  $\nu_{\text{sym}}(\text{FCIF})$ . Consequently, association appears to involve exclusively the axial fluorine atoms. This finding agrees with the association proposed<sup>27</sup> by Frey, Redington, and Aljibury for the structurally related, trigonal-bipyramidal molecules SF<sub>4</sub> and ClF<sub>3</sub>.

A decrease of the volatility of ClF<sub>3</sub>O due to extensive self-ionization (*i.e.*, 2ClF<sub>3</sub>O  $\rightleftharpoons$  ClF<sub>2</sub>O<sup>+</sup> + ClF<sub>4</sub>O<sup>-</sup>) in the liquid or solid phase can be ruled out since it should result in more dramatic changes in the spectra upon phase change. Furthermore, no evidence for the presence of

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TABLE VII  
 THERMODYNAMIC PROPERTIES FOR ClF<sub>3</sub>O

T, °K	C <sub>p</sub> <sup>o</sup> , cal/mol	H <sup>o</sup> - H <sub>0</sub> <sup>o</sup> , kcal/mol	-(F <sup>o</sup> - H <sub>0</sub> <sup>o</sup> )/T, cal/(mol deg)	S <sup>o</sup> , cal/ (mol deg)
0	0	0	0	0
100	9.721	0.837	49.255	57.624
200	14.932	2.072	55.613	65.971
298.15	18.593	3.732	60.159	72.675
300	18.646	3.766	60.237	72.790
400	20.875	5.751	64.108	78.486
500	22.260	7.913	67.478	83.305
600	23.160	10.187	70.470	87.448
700	23.771	12.536	73.159	91.067
800	24.200	14.936	75.602	94.271
900	24.512	17.372	77.838	97.141
1000	24.744	19.835	79.900	99.736
1100	24.921	22.319	81.813	102.103
1200	25.059	24.818	83.595	104.277
1300	25.168	27.330	85.265	106.288
1400	25.256	29.851	86.834	108.156
1500	25.328	32.380	88.314	109.901
1600	25.387	34.916	89.715	111.538
1700	25.437	37.458	91.044	113.078
1800	25.479	40.003	92.309	114.533
1900	25.514	42.553	93.516	115.912
2000	25.545	45.106	94.668	117.221

the bands characteristic for ClF<sub>2</sub>O<sup>+</sup> and ClF<sub>4</sub>O<sup>-</sup> could be detected in the spectra of liquid or solid ClF<sub>3</sub>O.

**Thermodynamic Properties.**—The thermodynamic properties were computed with the molecular geometry and vibrational frequencies given above assuming an ideal gas at 1 atm pressure and using the harmonic-oscillator rigid-rotor approximation.<sup>28</sup> These properties are given for the range 0–2000°K in Table VII.

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## Chlorine Trifluoride Oxide. IV. Reaction Chemistry

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Reactions of chlorine trifluoride oxide, ClF<sub>3</sub>O, with Cl<sub>2</sub>, Cl<sub>2</sub>O, ClOSO<sub>2</sub>F, N<sub>2</sub>F<sub>4</sub>, HNF<sub>2</sub>, and NF<sub>2</sub>CFO are reported. In these reactions, ClF<sub>3</sub>O either fluorinates, oxygenates, or both oxygenates and fluorinates the substrates. The interaction of ClF<sub>3</sub>O with PtF<sub>6</sub> is also described. In this reaction, ClF<sub>3</sub>O liberates F<sub>2</sub> and forms the new complex, ClF<sub>2</sub>O<sup>+</sup>PtF<sub>6</sub><sup>-</sup>. Characterizing data for the complex are given.

### Introduction

In the preceding papers,<sup>1–3</sup> we reported the preparation and properties of the new chlorine oxyfluoride, ClF<sub>3</sub>O. Since ClF<sub>3</sub> is a powerful oxidative fluorinating agent, it was anticipated that ClF<sub>3</sub>O would be similarly reactive as a fluorinating, as well as an oxygenating agent. This paper describes some of the reaction chemistry of ClF<sub>3</sub>O.

### Experimental Section

**Apparatus.**—Experimental techniques used in these studies were essentially the same as those described elsewhere.<sup>1–4</sup> Debye-Scherrer powder patterns were taken using a Philips Norelco instrument, Type No. 12046, with copper K $\alpha$  radiation and a nickel filter. Samples were sealed in quartz capillaries (~0.5 mm o.d.). Gas chromatographic analyses were performed using the method of Dayan and Neale.<sup>5</sup> A Hanovia utility lamp (catalog no. 30620) was used for the photolyses.

**Materials.**—The preparation and purification of ClF<sub>3</sub>O, Cl<sub>2</sub>O,

ClOSO<sub>2</sub>F, HNF<sub>2</sub>, and NF<sub>2</sub>CFO are described elsewhere.<sup>1,6–9</sup> Chlorine (Matheson Co.), tetrafluorohydrazine (Allied Chemical Corp.), and platinum hexafluoride (Ozark Mahoning Co.) were purchased and purified by fractional condensations.

**Reactions of Chlorine Trifluoride Oxide. With Chlorine.**—At room temperature, chlorine and ClF<sub>3</sub>O do not interact. An equimolar mixture (100 cm<sup>3</sup>, 4.46 mmol each) of the two when heated at 200° for 16 hr in a 30-ml stainless steel cylinder did result in complete consumption of the ClF<sub>3</sub>O. The products were separated by fractional condensation at -142 and -196° and analyzed by infrared spectroscopy and gas chromatography. Chlorine monofluoride was the principal product (262 cm<sup>3</sup>, 11.7 mmol) with only a small amount of ClO<sub>2</sub>F (6.7 cm<sup>3</sup>, 0.3 mmol) and unreacted Cl<sub>2</sub> (15.7 cm<sup>3</sup>, 0.7 mmol) being observed. Non-condensables (O<sub>2</sub>) were also produced. When kept at 100° for 2 days, ClF<sub>3</sub>O reacted with Cl<sub>2</sub> only partially (~30%) to give ClF<sub>3</sub>, ClF, and ClO<sub>2</sub>F as the major reaction products.

**With Dichlorine Monoxide.**—Chlorine trifluoride oxide (76.5 cm<sup>3</sup>, 3.42 mmol) and an equal quantity of Cl<sub>2</sub>O were separately condensed at -196° into a Kel-F reactor fitted with a Teflon valve. (Previous experiments in stainless steel cylinders resulted in extensive decomposition of the Cl<sub>2</sub>O, apparently due to

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