metal centers. The dissimilarity of the spectra of the two forms is consistent with form I1 being six-coordinate. The most striking difference between the spectra of the two forms is in the $PO₂$ stretching region where form I has two strong bands and form I1 has three. The two bands observed for form I are quite readily assigned to the antisymmetric and symmetric $PO₂$ stretches.⁷ The intensity and position of the three bands in form II's spectra are characteristic of $PO₂$ stretching vibrations, but for C_{2v} or C_s symmetry of the phosphinate anion only two strong $PO₂$ stretching bands are expected. It may well be that form I1 contains more than one structural type of phosphinate.

The thermograms for both forms were recorded by dsc, and the results are given in Table I. The tempera-

tures of the transitions were found to vary with the thermal history of the sample as well as its size and heating rate, so that the values given are only approximate. The transition temperatures and energies of form I are in the range of those previously observed for Zn [OP(C₈H₁₇)₂O]₂⁸ which is isomorphous with form I. The assignment for the endothermic transformation at 54' for form I is based on that made for the isomorphous $Zn[OP(C_8H_{17})_2O]_2$ polymers, whereas the assignment of the transition at 111° was made after a careful study with a light polarizing microscope. The higher temperature transition for $Zn[OP(C_8H_{17})_2O]_2$ was previously assigned as a melting. However, even well above 170" birefringence was observed for both $Co[OP(C_8H_{17})_2O]_2$ and $Zn[OP(C_8H_{17})_2O]_2$. Thus the high temperature transition is most probably a transition from a crystalline form to. a paracrystalline form, similar to that described for other zinc and cobalt phosphinates.2 The thermogram of form I1 shows one transition which corresponds to its conversion to form I and which was found to be irreversible.

Although both forms are soluble in nonpolar solvents, infrared spectra indicate that form I1 converts irreversibly to form I upon dissolution and consequently no molecular weight data could be obtained for form 11. Form I1 is crystalline with a considerably different X-ray powder pattern than form I which can be either crystalline or amorphous. The *d* spacings for both forms are given in Table 11.

This appears to be the first example of a poly- (metal phosphinate) which exists both in four- and sixcoordinate forms. However the two isomers have some similarities to the well known α and β forms of Co(py)₂- $Cl₂$.⁹ In both systems the six-coordinate form converts

(7) R. **A.** Nyquist, *J. Mol. Slvuct.,* **2,** 111 (1968).

|--|--|

TABLE **I1** X-RAY POWDER PATTERNS OF $Co[OP(C_8H_{17})_2]O_2^a$,---Form **I-----** ,----Form **II--** *d*, \hat{A} *Ib Ib d*, \hat{A} *Ib Ib* $\begin{array}{cccc} 16.07 & & & 100 & & & 14.04 & & & 100 \\ 10.07 & & & & 40 & & & 10.32 & & & 30 \end{array}$ 10.07 40 10.32 30 9.10 30 8.81 80

^{*a*} Cu K_{α} radiation was used. ^{*b*} Relative intensities estimated visually.

to the four-coordinate form on heating or dissolution ; on the other hand, the conversion in the $Co(py)_{2}Cl_{2}$ system is reversible, whereas in the $Co(OP(C_8H_{17})_2O)_2$ system it is not.

Acknowledgment.-The author wishes to acknowledge B. B. Wayland for the use of his Gouy balance. This investigation was supported in part by the Office of Naval Research.

CONTRIBUTION FROM THE DEPARTMENT OF CHEMISTRY, HARVARD UNIVERSITY, CAMBRIDGE, MASSACHUSETTS 02138

Ab Initio **Self -Consistent Field Calculation of** the Energies of Formation of B_2H_6 and $B_2H_7^-$

BY JOHN H. HALL, JR., DENNIS S. MARYNICK, AND WILLIAM N. LIPSCOMB^{*}

Received May 1, **1972**

Accurate binding energies cannot usually be obtained from Hartree-Fock (HF) calculations because the changes in correlation energy are often of the same orders of magnitude as the corresponding changes in the HF energy.¹ The HF binding energies are still of interest, however, when some estimate of the sign and magnitude of the correlation energy changes can be made. We report in this note *ab initio* self-consistent field (SCF) calculations² of the energies of formation, ΔE_f , for $\rm B_2H_6$ and $\rm B_2H_7^-$.

An extensive calculation of the ΔE_f for the reaction
 $2BH_3(g) \longrightarrow B_2H_6(g)$ (1)

$$
2BH_3(g) \longrightarrow B_2H_6(g) \tag{1}
$$

has been reported by Gelus, Ahlrichs, Staemmler, and Kutzelnigg (GASK).³ This calculation employed a large Gaussian basis set and the correlation energy corrections were estimated using the independent electron pair approximation.^{3,4} The total ΔE_f was estimated to be -36.0 ± 5 kcal/mol. These authors attributed -11.5 kcal to the HF energy changes, and attributed -11.5 kcal to the HF energy changes, and -24.5 kcal to changes in the correlation energy; these values included estimates of -3.0 kcal/mol and -7.7

⁽⁸⁾ V. Giancotti, **A.** Ripamonti, and P. **4.** Temussi, *PoZyww,* **12,** 409 (1971).

^{(9) &}quot;Gmelins Handbuch der anorganischen Chemie," 8 Auflage, Syst.-Nr. 58, "Kobalt," Teil R, Erg.-Bd., Verlag Chemie, Weinheim, 1963, pp 65-70.

⁽¹⁾ R. Ahlrichs, *Theor. Chim.* Acta, **17,** 348 (1970).

⁽²⁾ C. C. J. Roothaan, *Rev. Mod. Phys.,* **23,** 69 (1951). (3) M. Gelus, R. Ahlrichs, V. Staemmler, and W. Kutzelnigg, *Chem. Phys. Lett.,* **7,** 503 (1970); a discussion of this and related **work** occurs in "The Electronic Structure of Atoms and Molecules" by H. F. Schaefer, **111,**

Addison-Wesley, Reading, Mass., 1972, pp 348-352. (4) **31.** Jungen and R. Ahlrishs, *Theor. Chim. Ada,* **17,** 339 (1970).

kcal/mol, respectively, for supposed inadequacies in the basis set.

Our calculation of the ΔE_f for B_2H_6 employed the experimental geometry for B_2H_6 ⁵ and an optimized geometry for BH_a . The B-H distance in BH_a , 1.19 \AA , is the same as that found previously⁶ from a minimum basis set. An extended basis was used (exponents in parentheses) for both $BH₃$ and $B₂H₆$, described by Laws, Stevens, and Lipscomb $(LSL):$ ⁷ H 1s (1.449) , H 2s (1.0205), H 2p (1.952), B 1s (7.338), B 1s (3.996), B 2s (1.724), B 2s (1.110), B 3s (4.796), B 2p (4.558), B 2p (1.753), B 2p (0.931), B 3d (1.490). The total energy for B_2H_6 is given by LSL as -52.8331 au, and our new total energy for BH3, calculated using the LSL basis, is -26.4014 au. From these energies we find that the near-Hartree-Fock limit of ΔE_i is -19.0 kcal/mol, in contrast to GASK's calculated value of - 8.5 kcal/mol and their estimate of the Hartree-Fock -8.5 kcal/mol and their estimate of the Hartree–Fock
limit of -11.5 kcal/mol.

Experimental values of ΔH_t for B₂H₆ range from -25 to -38 kcal/mol for kinetic measurements and from -39 to -59 kcal/mol for mass spectrometric results.⁸ A recent kinetic value, consistent with several reactions, is -35.5 kcal/mol,⁹ corresponding to a ΔE_f of -34.9 kcal/mol. -34.9 kcal/mol.
However, the estimated Hartree-Fock limit of -19.0

kcal/mol is not in the most favorable range of experimental values, probably largely because of change in the correlation energy in the formation of B_2H_6 from 2BH3. GASK's calculated value of this correlation energy is - 16.8 kcal/mol, but they add -7.7 kcal/mol as an estimate of the effect of inadequacy of their basis set on the correlation part. The result, -24.5 kcal/mol, is not necessarily the appropriate value to be added to our near-Hartree-Fock limit. Indeed, we suspect that GASK have overestimated the magnitude of the correlation correction, and we observe that their *calculated* correlation energy is about the difference between our near-Hartree-Fock limit and the most probable experimental value. However, we have not computed the correlation correction to our near-Hartree-Fock wave function, but such a procedure would give the most appropriate correlation correction to our value of -19.0 kcal/mol.¹⁰

The molar energy of formation for the reaction
 $BH_4-(g) + BH_3(g) \longrightarrow B_2H_7(g)$ *(2)*

$$
BH_4^-(g) + BH_3(g) \longrightarrow B_2H_7^-(g) \tag{2}
$$

was calculated using optimized minimum basis sets of BH_4^{-11} BH₃,¹² and B₂H₇⁻. The geometry of B₂H₇⁻ has not been determined experimentally, but the proposed structure¹³ consists of two $BH₃$ groups connected

- **(5)** L. *S.* Bartelland B. L Carroll, *J* Chem Phys., **42, 1135 (1965)**
- (6) W. E. Palke and W. N. Lipscomb, $ibid.$, **45**, 3948 (1966).
- **(7)** E. Laws, R M Stevens, and W. N. Lipscomb, *J* Amer. Chem Soc., **94,4461 (1972)**

(8) See P. S. Ganguli and H. **A.** McGee, *J.* Chem. Phys , **BO, 4658 (1969),** and references therein.

(9) G. W. Mappes, S. **A.** Friedman, and T. P Fehlner, *J.* Phys. *Chem.,* **74,3307 (1970)**

(10) A heat of formation *of* **-35** 0 kcal/mol has been obtained by W **A** Goddard using G1 wave functions. This value includes an estimated correlation energy change correction, however, it cannot be rigorously separated into HF and correlation energy changes. W. A. Goddard III, public communication, Harvard-MIT colloquium, **1971**

(11) R. Hegstrom, W. Palke, and W. N. Lipscomb, *J. Chem. Phys.*, 46, **920 (1967)**

(12) E. Switkes, R. M. Stevens, and W. N. Lipscomb, $ibid.$, 51, 5299 (1969). **(13) (a) W N** Lipscomb, "Boron Hydrides," W **A** Benjamin, New York, N Y, **1963;** (b) H. Brown, P Stehle, and P Tierney, *J* Amer Chem

Soc., **79, 2020 (1957).**

by a linear B-H-B three-center bond. The nmr spectra¹⁴ indicate the presence of two $BH₃$ groups, but nothing about the nature of the B-H-B linkage could be determined. In particular, the possible asymmetry of the B-H-B bond could not be established or disproved.

We have calculated a geometry for staggered B_2H_7 from a minimum basis set of Slater-type orbitals, using a program previously described.¹⁵ The starting geometry and exponents were taken from the calculation on diborane.16 The geometry and exponents were optimized until a convergence of 10^{-3} au in the energy was achieved. All calculations were made at the selfconsistent field method, except for the optimization of the $H_t BH_t$ angle by the PRDDO method.¹⁷ However, all final results on B_2H_7 ⁻ were obtained by the SCF method. In Table I we present optimized bond

TABLE I

B₂H₇ Optimized Exponents, Bond Distances, AND ENERGY ANALYSIS

 a All quantities in these columns are in au.

distances, exponents, and the energy analysis. Although the B-H_t bond distance of 1.223 Å is longer than $B-H_t$ distances usually found in boron hydrides,¹⁸ this distance is exactly the average of the B-H distances in BH₄⁻ (1.255 Å)¹⁹ and BH₃ (1.19 Å).⁶ The B-H_b distance of 1.309 Å is shorter in B₂H₇⁻ than in B_2H_6 , but the smaller bridge hydrogen exponent indicates that the electron density associated with the B-H-B bond is more "spread out" in B_2H_7 than in B_2H_6 . The same trend is also seen in the H_t exponents, as one might expect for a negative ion. Our SCF calculations on other negative ions²⁰ also indicate that the bond distances are longer and the exponents smaller than in neutral molecules. The bridge hydrogen position was also optimized by moving the hydrogen atom relative to the two borons. To within the accuracy of our basis set, the B-H-B bond is symmetrical, and a nonlinear BH_bB structure is less stable.

For $B_2H_7^-$, the calculated ΔE_f is -25.1 kcal/mol. This value is in good agreement with the experimental ΔH_f of -31.0 ± 8 kcal/mol.²¹ We can guess the sign and magnitude of the correlation energy by noting that the correlation energy should increase with the number of three-center bonds or, equivalently, that the correlation energy should increase with the number of next neighbor bond interactions. Since the number of three-center bonds (and next neighbor interactions)

(14) D F Gaines, Inorg Chem **,2, 523 11963)**

- **(15)** R M Stevens, *J.* Chem Phys **,62, 1937 (1970)**
- (16) E. Switkes, R. M. Stevens, and W. N. Lipscomb, *ibid.*, 51, 2085 **(1969)**
- (17) T. A. Halgren and W. N. Lipscomb, Proc. Nat. Acad. Sci. U. S., **68, 1853 (1972).**

(18) See ref 11 and **A** Tippe and **W** Hamilton, *lnovg* Chem , **8, 469 (1969)**

(19) P T Ford and R E Richards, *Dzscuss* Faiaday *SOL* , **19,230 (1955)** (20) D. S. Marynick, J. H. Hall, and W. N. Lipscomb, unpublished results

(21) W. G. Evans, C. E. Halloway, K. Sukumarabandhu, and D. Mc-Danie1,Inorg. Chem *,I,* **1746 (1968)**

increases in both reactions, the calculated energies of formation should be higher (less negative) than the experimental values. Also, the correlation correction to the HF ΔE_f should be smaller for $\rm B_2H_7$ ⁻ than for $\rm B_2H_6$.

Our results are summarized in Table 11.

TABLE I1 SUMMARY OF *AHf* CALCULATIONS

	Total energies, au	Hf , keal/mol
BH,	-26.4014	-19.0°
B_2H_6	$-52.8331c$	
BH ₃	-26.3533^{d}	
$\rm BH_{4}$ -	-26.9232 ^e	$-25.4b$
$B_2H_7^-$	-53.3170	

^a For reaction 1. ^b For reaction 2. ^c Reference 7. ^d Reference 12. **6** Reference 11.

We believe that the basis set for B_2H_0 is of such quality that our value for the energy of formation is essentially a Hartree-Fock value. In view of the reasonable agreement with experiment for the calculated energy of formation of $B_2H_7^-$, we predict that the B-H-B bond is symmetric in the gas phase. However, in the crystal phase, forces could conceivably be sufficient to produce some asymmetry.22

Acknowledgments.—We wish to thank the Office of Naval Research and the Advanced Research Projects Agency for support and the National Science Foundation for a fellowship for D. S. **&I.** We also thank I. M. Pepperberg and D. A. Kleier for assistance in these calculations.

(22) NOTE ADDED IN Proof.-Inclusion of all 14 single and 210 double excitations from the valence shells of minimum Slater basis SCF wave iunctions for BHa, B2H6, and possible transition states has yielded the following results. This correlation correction is -10.6 kcal for reaction 1. Further studies with expanded basis sets will give a more appropriate correlation correction to be added to the HF value of -19.0 kcal for reaction 1. The best transition state discovered so far is centrosymmetric, of C_{2h} symmetry, having two very unsymmetrical hydrogen bridges and a $B \cdots B$ distance of 3.0 **A,** and is less stable than 2BH3 by **2.6** kcal (study in progress by I. M. Pepperberg, D. **A.** Dixon, and W. N. Lipscomb).

CONTRIBUTION No. 1930 FROM THE CENTRAL RESEARCH DEPARTMENT, EXPERIMENTAL STATION, E. I. DU PONT DE NEMOURS AND COMPANY, WILMINGTON, DELAWARE 19898

Electronic Effects of Phosphorus Ligands on the Protonation of NiL4 Complexes

BY CHADWICK **A.** TOLMAN

Received May 16, 1972

Novel cationic nickel hydride complexes formed by protonation of $Ni[POEt)_3]_4^1$ and $Ni[Ph_2PCH_2CH_2$ - $PPh₂$]₂² by strong acids have been reported recently. Solutions of $HNi[POEt)₃1₄$ ⁺ are catalytically active in the coupling of butadiene and ethylene to form hexadienes³ and in olefin isomerization.⁴ We now wish to report that a wide range of similar complexes of the type $HNiL₄$ ⁺ can be formed by these protonation reactions.

(4) C. A. Tolman, *ibid.,* **94,** 2994 (1972)

Equilibrium constants for reaction 1 using H_2SO_4 in
 $H^+ + Nil_4 \stackrel{K}{\Longleftarrow} HNil_4^+$ (1)

$$
H^{+} + Nil_{4} \stackrel{K}{\Longleftarrow} HNil_{4}^{+} \tag{1}
$$

CH30H are shown to depend on the electron donoracceptor character of the phosphorus ligands L.

Addition of strong acids like H_2SO_4 to solutions of zerovalent nickel complexes of the type NiL4 forms $HNiL₄$ ⁺, as shown by the appearance of a symmetrical quintet^{1,5} at high field in the proton nmr spectra (Table I). Though the chemical shift shows no systematic

*^a*Formed *in situ* by the addition of HzS04 to solutions of NiLa. The latter were prepared as described by C. A. Tolman, *J. Amer. Chem. Soc.,* **92,** 2956 (1970). $\frac{b}{c}$ Measured at -20° with respect to internal tetramethylsilane using a Varian HA-100 spectrometer. ^c Uncertainty ± 0.5 Hz, unless noted otherwise. ^d Hydride first reported in ref 1. *e* Hydride first reported in ref 2.

variation, the value of J_{PH} increases quite regularly as the electron-withdrawing ability of the phosphorus ligands increases, judging from the carbonyl stretching frequencies measured for $Ni(CO)_3L$ complexes.⁶

Equilibrium constants for hydride formation (eq 1) with selected compounds were measured spectrophotometrically using H_2SO_4 in CH₃OH at 0° , as described earlier for $Ni[P(OEt)₃]$ ⁷ The results are shown in Table II and Figure 1. The value of K for $Ni[Ph_2-$

HYDRIDE FORMATION CONSTANTS WITH H_2SO_4 IN CH₃OH at 0°

Infrared data from ref 6. b Determined at ambient temperature. \cdot 0.1 *M* L added. See ref 7. *d* Reported earlier in ref 7. **^e**No change in **uv** spectrum on adding 1 *M* HzS04. *f* Value for PPhzEt.

 $PCH_2CH_2PPh_2|_2$ is most uncertain because of the extremely low solubility of the complex in CH₃OH $(\sim 10^{-4}$ *M*). The hydride formed was, however, exceedingly stable.

As expected, the equilibrium constant decreases as the phosphorus ligands remove more electron density from the metal. With $Ni[P(OCH₂CCl₃)₃]$ ₄, the complex with the most strongly electron-withdrawing ligands of

⁽¹⁾ W. C. Drinkard, D. R. Eaton, J. P. Jesson, and R. V. Lindsey, Jr., *Inoug. Chem.,* **9,** 392 (1970).

⁽²⁾ R. A. Schunn, *ibid.,* **9,** 394 (1970).

⁽³⁾ C. **A.** Tolman, *J. Amev. Chem. SOL.,* **92,** 6777 (1970).

⁽⁵⁾ The appearance of the NIH resonance as a quintet is attributed to rapid intramolecular ligand exchange in these labile presumably trigonalbipyramidal five-coordinate complexes. The phenomenon has recently been studied in detail in HM(PF₈)& complexes by P. Meakin, J. P. Jesson, F. N. Tebbe, andE. L. Muetterties, *ibid.,* **93,** 1797 (1971).

⁽⁶⁾ C. A. Tolman, *ibid.,* **92,2953** (1970).

⁽⁷⁾ C. A. Tolman, *ibid.,* **92,** 4217 (1970).