$$
^{1}/_{2}[PtH(\mu\text{-}NH_{2})(PPh_{3})]_{2} \xrightarrow{\text{-}NH_{3}} Pt(PPh_{3})_{2} \xrightarrow{CDCl_{3}} \text{cis-PtCl}_{2}(PPh_{3})_{2} (4)
$$

[PtH(
$$
\mu
$$
-NH<sub>2</sub>)(PEt<sub>3</sub>)]<sub>2</sub> + 2PEt<sub>3</sub>  $\rightarrow$   
<sup>4</sup>/<sub>3</sub>Pt(PEt<sub>3</sub>)<sub>3</sub> + <sup>2</sup>/<sub>3</sub>Pt + 2NH<sub>3</sub> (5)

however, the observed product is  $Pt(PEt<sub>3</sub>)<sub>3</sub>$  (eq 5), the reductive elimination product.<sup>10</sup> The elimination of  $NH_3$  from the anti isomer of  $[PH(\mu-NH_2)(PEt_3)]_2$  is found to be faster than from the **syn** isomer because the hydride and amide ligands are mutually cis in the anti isomer. In our initial premise summarized in eq **1,** we **proposed** that the first product of the reaction of the cationic complex trans- $[PtX(NH_3)L_2]ClO_4$  with strong bases was monomeric with a trans stereochemistry. We now find that such monomeric complexes trans-PtX(NH<sub>2</sub>)L<sub>2</sub> (X = H, Me) can be readily prepared with a bulky ligand L, such as  $PCy_3$  (eq 6). <br>trans-[PtX(NH<sub>3</sub>)(PCy<sub>3</sub>)<sub>2</sub>]ClO<sub>4</sub> + NaNH<sub>2</sub> ->

trans-
$$
\frac{\text{PrX(NH}_3)(PCy_3)_2}{\text{trans-PrX(NH}_2)(PCy_3)_2 + \text{NaClO}_4 + \text{NH}_3} \tag{6}
$$

 $X = H$ , Me

Indeed, solutions of these sterically stabilized compounds *trans-*PtX(NH<sub>2</sub>)(PCy<sub>3</sub>)<sub>2</sub> do not undergo substitution dimerization or reductive elimination at 25 °C.

These results show that sterically stabilized monomeric hydride and alkyl amide platinum(I1) complexes are thermodynamically

(10) '95Pt NMR shows a quartet ('J(PtP) = 4209 Hz) at **6** -4510 (see: Mann, **B.** E.; Musco, **A.** *J.* Chem. **SOC.,** Dalton Trans. **1980,** 776-785).

stable and that the trans amide hydrides undergo substitution dimerization prior to the reductive elimination of ammonia.

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Registry **No. syn-1,** 110719-63-6; anti-1, 110771-97-6; **syn-2,**  110719-64-7; anti-2, 110771-98-7; anti-3, 110719-65-8; syn-4, 110719- 67-0; anti-4, 110772-00-4; syn-5, 110771-99-8; anti-5, 110719-66-9; trans- [PtH(NH,)(PPh,),]CIO,, 32 109-33-4; *trans-* [PtH(NH,)- (PEt,),]C104, 110719-69-2; **trans-[PtMe(NH,)(PPh,),]CIO,,** 110719- 7 1-6; **trans-[PtMe(NH,)(PEt,),]ClO,,** 110719-73-8; trans-[PtMe-  $(NH_3)(PMePh_2)_2]CO_4$ , 110719-75-0; cis-PtMe $(NH_2)(PPh_3)_2$ , 110719-76-1; cis-PtMe(NH<sub>2</sub>)(PEt<sub>3</sub>)<sub>2</sub>, 110719-77-2; cis-PtMe(NH<sub>2</sub>)(PMePh<sub>2</sub>)<sub>2</sub>, 110719-78-3;  $cis-PtCl_2(PPh_3)_2$ , 15604-36-1;  $Pt(PEt_3)_2$ , 66916-63-0; trans-[PtH(NH<sub>3</sub>)(PCy<sub>3</sub>)<sub>2</sub>]CIO<sub>4</sub>, 110719-80-7; trans-PtH(NH<sub>2</sub>)(PCy<sub>3</sub>)<sub>2</sub>, 1 107 19-8 1-8; trans- [PtMe(NH3)(PCy,),] **C104,** 1 107 19-83-0; trans-PtMe(NH<sub>2</sub>)(PCy<sub>3</sub>)<sub>2</sub>, 110719-84-1.

Supplementary Material Available: Atomic coordinates (Table lS), bond distances (Table 2S), bond angles (Table 3S), anisotropic thermal parameters (Table 4S), and hydrogen atom coordinates (Table 5S) (4 pages); observed and calculated structure factors (Table 6S) (37 pages). Ordering information is given on any current masthead page.



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## **Articles**

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## **Metallic versus Nonmetallic Properties of Ternary Chalcogenides Ta<sub>2</sub>MSe<sub>7</sub> (M = Ni, Pt)** and  $Ta_2NiX_5$   $(X = S, Se)$

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The electronic structures of layered chalcogenides  $Ta_2NiS_5$  and  $Ta_2NiS_6$  were examined by carrying out tight-binding band calculations. In agreement with experiment, Ta<sub>2</sub>NiS<sub>5</sub> and Ta<sub>2</sub>NiSe<sub>7</sub> are predicted to be semiconducting and metallic, respectively. The calculated band electronic structures of Ta<sub>2</sub>NiS<sub>5</sub> and Ta<sub>2</sub>NiSe<sub>7</sub> are best described by the oxidation formalisms (Ta<sup>5+</sup>)<sub>2</sub>- $(N^{10})(S^{2-})$ , and  $(Ta^{5+})_2(N^{12+})(Se^{2-})_5(Se^{2-})$ , respectively. In  $Ta_2NIS_5$ , the  $Ta^{5+}$  cations act as good  $\pi$  acceptors of each Ni atom, thereby stabilizing its zero oxidation state. The metallic properties of Ta<sub>2</sub>NiSe<sub>7</sub> originate from the fact that the bottom d-block band of the octahedral chains overlaps with the top Se p-block band of the trigonal-prismatic chains. The bottom d-block band of the trigonal-prismatic chains is empty, so that the Ta atoms of the trigonal-prismatic chains are not responsible for the metallic properties of  $Ta_2NiSe_7$ .

Recently, Sunshine and Ibers synthesized layered ternary chalcogenides  $Ta_2NiX_5$   $(X = S, Se)^2$  and  $Ta_2MSe_7$   $(M = Ni, Pt)^3$ In  $Ta_2NiX_5$  the Ta and Ni atoms have octahedral (OCT) and tetrahedral coordinations, respectively. According to the crystal structure of Ta<sub>2</sub>NiX<sub>5</sub> and the oxidation formalism  $(Ta^{4+})_2$ - $(Ni^{2+})(X^{2-})_5$ , it is expected<sup>2</sup> that Ta<sub>2</sub>NiX<sub>5</sub> is a one-dimensional metal and hence would exhibit a charge density wave (CDW) phenomenon. However, this  $Ta_2NiX_5$  compound is a diamagnetic semiconductor with the band gaps of  $0.13$  and  $0.36$  eV for  $X =$ S and Se, respectively,<sup>2</sup> and no evidence for a CDW phenomenon was found.<sup>4</sup> Thus, the oxidation formalism valid for  $Ta_2NiX_5$  is  $(Ta^{5+})_2(Ni^0)(X^2)$ , although this formalism was regarded unlikely2 because it implies the Occurrence of highly oxidized metal ions (i.e.,  $Ta^{5+}$ ) in the presence of unoxidized metal atoms (i.e.,  $Ni<sup>0</sup>$ ).

In Ta<sub>2</sub>MSe<sub>7</sub> ( $M = Ni$ , Pt) the M atoms are found in squarepyramidal coordination sites, while the Ta atoms are either in OCT or in bicapped trigonal-prismatic (BTP) coordination sites. The oxidation formalism  $(T_a^{5+})_2(M^{2+})(Se^{2-})_5(Se^{2-})$  suggests semiconducting properties for  $Ta_2MSe_7$ , but it is found to be metallic down at least to 100 K,<sup>3</sup> below which no measurements have been made. By analogy with the metallic properties of  $T_{a}Se_{3}$  and  $NbSe<sub>3</sub>$ <sup>5</sup> Sunshine and Ibers suggested<sup>3</sup> that the BTP chains are

<sup>(1) (</sup>a) Universite de Paris-Sud. (b) North Carolina State University. (2) Sunshine, S. **A,;** Ibers, J. **A.** Znorg. Chem. **1985,** *24,* 3611.

**<sup>(3)</sup>** Sunshine, **S. A.;** Ibers, J. **A.** *Znorg.* Chem. **1986, 25,** 4355.

<sup>(4)</sup> DiSalvo, F. J.; Chen, C. H.; Fleming, R. M.; Waszczak, J. V.; Dunn, R. *G.;* Sunshine, S. **A,;** Ibers, J. **A.** Less-Common *Met.* **1986,** *116,* **51.** 

**Table I.** Exponents and Parameters Used in the Calculations"

atom	orbital	$H_{ii}$ , eV	$\zeta_1$	$\zeta_2$	$c_1^b$	$c_2{}^b$	
Ta	6s	$-10.10$	2.28				
	6p	$-6.86$	2.24				
	5d	$-12.10$	4.76	1.94	0.6597	0.5589	
Ni	4s	$-9.70$	2.10				
	4p	$-5.15$	2.10				
	3d	$-13.49$	5.75	2.30	0.5798	0.5782	
Se	4s	$-20.50$	2.44				
	4p	$-13.20$	2.07				
S	3s	$-20.00$	1.817				
	3p	$-13.30$	1.817				

**A** modified Wolfsberg-Helmholz formula'0 was used to calculate the off-diagonal  $H_{ij}$  values.  $b$  Contraction coefficients used in the double- $\zeta$  expansion.

responsible for the metallic properties of  $Ta<sub>2</sub>MSe<sub>7</sub>$ , which was proposed to occur via charge transfer from the Se atoms to the Ta atoms of BTP chains, thereby creating Ta<sup>4+</sup> cations.

In the present study, we perform tight-binding band calculations<sup>6</sup> on Ta<sub>2</sub>NiS<sub>5</sub> and Ta<sub>2</sub>NiSe<sub>7</sub> to gain some insight into the origin of their apparently puzzling electrical properties. The atomic parameters employed in our band calculations based upon the extended Hückel method<sup>7</sup> are summarized in Table I.

## **Results and Discussion**

**A. Ta<sub>2</sub>NiX<sub>5</sub> (X = S, Se).** The Ta<sub>2</sub>NiX<sub>5</sub> phase contains layers of composition  $Ta_2NiX_5$ . It is convenient to consider each  $Ta_2NiX_5$  layer in terms of the double OCT chain,  $Ta_2X_6$  depicted in **1.** These  $Ta_2X_6$  chains can be condensed to form a  $Ta_2X_5$  net



by sharing their apical X atoms  $(X_a)$ . When the Ta-X<sub>a</sub>-Ta linkages are bent, every two adjacent apical X atoms form a tetrahedral hole with every two adjacent basal  $X$  atoms  $(X_b)$ . Then the  $Ta_2NiX_5$  layer is obtained by filling all the tetrahedral holes with Ni atoms (see 2). In the  $Ta_2X_6$  double OCT chains 1, the



Ta-Ta distances  $r_1$  and  $r_2$  are fairly long (i.e., greater than 3.40 and 3.49 Å for  $X = S$  and Se, respectively<sup>2</sup>) so that there is no significant Ta-Ta bonding. Around each Ni atom, however, there are four Ta atoms in a square-planar arrangement with a very

- *(5)* (a) Meerschaut, A,; Rouxel, J. **In** *Crystal Chemistry and Properties of Materials with Quasi-One-Dimensional Structures;* Rouxel, J., Ed.; Reidel: Dordrecht, The Netherlands, **1986;** p 205. (b) Whangbo,
- (6) Whangbo, M.-H.; Hoffmann, R. *J. Am. Chem. SOC.* **1978,100,6093. M.-H.** *Ibid.,* **p** 27.
- **(7)** Hoffmann, R. *J. Chem. Phys.* **1963,** *39,* **1399.**



**Figure 1.** Dispersion relations of the bands calculated for a  $Ta_2NiS_5$ layer, where *E,* refers to a band gap. **In** units of the reciprocal vectors  $a^*$  and  $c^*$ , the vector points  $\Gamma$ , *X*, and *Z* are defined as  $\Gamma = (0, 0)$ ,  $X =$  $(a^*/2, 0)$ , and  $Z = (0, c^*/2)$ .



**Figure 2.** Calculated band gap as a function of the valence-shell ionization potentials of the Ni 3d and Ta 5d orbitals.

short Ta-Ni distance (i.e., 2.747 and 2.8 **13** *8,* for **X** = *S* and Se, respectively<sup>2</sup>).

Figure 1 shows the dispersion relations for some of the bands calculated for a  $Ta_2NiS_5$  layer. In general, these bands are dispersive along the chain axis (i.e., the  $a$  axis) but relatively flat along the direction perpendicular to the chain axis. Our calculations show that  $Ta_2NiS_5$  has a direct band gap at  $\Gamma$ . As summarized in Figure 2,  $Ta_2NiS_5$  is calculated to have a direct band gap at  $\Gamma$  for any reasonable choice of the valence-shell ionization potentials (VSIP's) of the Ni 3d and the Ta 5d orbitals. It should be noted that the direct band gap at  $\Gamma$  persists even when the VSIP value of the Ni 3d orbital is raised to as high as that of the Ta 5d orbital. This phenomenon basically occurs since the symmetry of the  $Ta_2NiX_5$  layer is low. For any linear combination of Ta d orbitals, therefore, there exists a corresponding linear combination of Se p orbitals with the same symmetry. The interaction between such levels would widen their energy difference, thereby leading to a band gap. Consequently,  $Ta_2NiX_5$  is predicted to be a semiconductor, and the oxidation formalism  $(Ta^{5+})_2$ - $(Ni<sup>0</sup>)(X<sup>2</sup>)<sub>5</sub>$  is appropriate for Ta<sub>2</sub>NiX<sub>5</sub>. In organometallic compounds, the Ni<sup>o</sup> oxidation state is found with strong  $\pi$  acceptors.<sup>8</sup> By analogy, it is reasonable to suggest that the empty  $t_{2g}$  levels of each Ta<sup>5+</sup> cation act as good  $\pi$ -acceptor levels for the Ni<sup>0</sup> atoms. This interpretation **is** consistent with the observation that the Ta-Ni distances are quite short and also with our computational result that the overlap population for the Ta--Ni contact is considerably larger than that for the  $Ta \cdot Ta$  contact (e.g., 0.09 vs  $0.02$  in Ta<sub>2</sub>NiS<sub>5</sub>). Our calculations show that the Ta d-block bands

<sup>(8) (</sup>a) Malatesta, L.; Cenini, **S.** *Zeroualent Compounds of Metals;* Academic: New York, **1974.** (b) Cotton, **F.** A,; Wilkinson, *G. Advanced Inorganic Chemistry,* 4th ed.; Wiley: New York, **1980; p** 797.



**Figure 3.** Dispersion relations of the bands calculated for a  $Ta_2NiSe_7$ layer, where the dashed line refers to the Fermi level. In units of the reciprocal vectors  $b^*$  and  $c^*$ , the vector points  $\Gamma$ , *Y*, and *Z* are defined as  $\Gamma = (0, 0), Y = (b^*/2, 0), Z = (0, c^*/2),$  and  $M = (b^*/2, c^*/2).$ 

of Ta<sub>2</sub>NiX<sub>5</sub> (i.e., the upper part of Figure 1) along the  $\Gamma \rightarrow X$ line are practically the same as those of the isolated double OCT chain Ta<sub>2</sub>X<sub>6</sub> (1)<sup>9</sup> and the latter lie slightly lower in energy ( $\sim$  0.1) eV). This supports our suggestion that the  $t_{2g}$  levels of  $Ta^{5+}$  act as  $\pi$ -acceptor levels for  $\widetilde{Ni}^{\widetilde{0}}$  in Ta<sub>2</sub>NiX<sub>5</sub>.

**B.** Ta<sub>2</sub>MSe<sub>7</sub> (M = Ni, Pt). The Ta<sub>2</sub>MSe<sub>7</sub> phase consists of layers of composition  $Ta_2MSe_7$ . Convenient building blocks for a Ta<sub>2</sub>MSe<sub>7</sub> layer are the Ta<sub>2</sub>Se<sub>6</sub> double OCT chains (1) and the double BTP chains **3.** Chains **1** and **3** can be brought together



to form a layer of composition  $Ta_2Se_7$  in such a way that every three Se atoms of 1 (i.e., two apical and one basal Se atoms) form a square pyramid with every two adjacent capping Se atoms of **3.** When such square-pyramidal sites are occupied by the metal atoms M (see 4), we obtain the  $Ta_2MSe_7$  layer.



Figure **3** shows the dispersion relations for some of the bands calculated for a  $Ta_2NiSe_7$  layer. In general, these bands are dispersive along the chain axis (Le., the *b* axis) but relatively flat along the direction perpendicular to the chain axis. There are two bands, a and b, that are partially filled. Bands a and b are primarily represented by the Se p orbitals of the BTP chains and by the Ta d orbitals of the OCT chains, respectively. The two



**Figure 4.** Fermi surfaces calculated for a  $Ta_2NiSe_7$  layer: (a) hole pockets of the top selenium p-block band; (b) the electron pocket of the bottom d-block band.

bands switch their orbital character near  $\Gamma$  along the  $\Gamma \rightarrow Y$  line. In short, the top of the highest lying Se p-orbital band overlaps slightly with the bottom of the lowest lying Ta d-block band, so that  $Ta_2NiSe_7$  is a semimetal. In  $Ta_2NiSe_7$ , the bottom d-block and the top Se p-block bands overlap essentially because these bands arise from the orbitals of different chains. In  $Ta_2NiX_5$ , however, the corresponding two bands arise from the orbitals of the same chain, which leads to a strong mixing between the two and hence to a band gap.

Bands b and c are derived primarily from the d orbitals of the double OCT chains and are lower lying in energy than bands d and e derived mainly from the d orbitals of the double BTP chains. Consequently, this study shows that the d-block bands of the BTP chains are not responsible for the metallic properties of  $Ta_2NiSe_7$ , unlike the case of  $TaSe_3$  and  $NbSe_3$ . It is the bottom d-block band of the double OCT chains and the top p-block band of the BTP chain Se atoms that are largely responsible for electrical conduction in Ta<sub>2</sub>NiSe<sub>7</sub>.

The Fermi surfaces associated with bands a and b are shown in parts a and b, respectively, of Figure 4. Band a provides two hole pockets of a cigar shape centered at  $(\pm 0.055b^*, 0)$ . Each hole pocket, elongated along the  $c^*$  axis, has the dimension of roughly  $0.01b^* \times 0.7c^*$ . Note that the two hole pockets are nested by the wave vector  $q \approx (0.1b^*, 0)$ . Band b leads to an electron pocket of an ellipse shape, elongated along the *b\** axis, which has the dimension of roughly  $0.08b^* \times 0.15c^*$ . Therefore, the  $Ta<sub>2</sub>MSe<sub>7</sub>$  phase is expected to be a two-dimensional metal, although it is made up of the OCT and BTP chains.

## **Concluding Remarks**

This band electronic structure study predicts the layered chalcogenides  $Ta_2NiX_5$  and  $Ta_2MSe_7$  to be semiconducting and metallic, respectively, which is in agreement with experiment. The oxidation formalism appropriate for  $Ta_2NiX_5$  is  $(Ta^{5+})_2$ - $(Ni^0)(X^2)$ <sub>5</sub>, and the Ta<sup>5+</sup> cations act as good  $\pi$  acceptors, thereby stabilizing the  $Ni<sup>0</sup>$  atoms. The oxidation formalism appropriate for Ta<sub>2</sub>MSe<sub>7</sub> is  $(Ta^{5+})_2(Ni^{2+}) (Se^{2-})_5 (Se^{2-})$ , and its metallic properties originate from the fact that the lowest lying tantalum d-block band overlaps slightly with the highest lying selenium p-block band. Therefore,  $Ta_2MSe_7$  is a semimetal. The partially filled d-block band does not originate from the BTP but from the OCT chains. The two partially filled bands give rise to closed Fermi surfaces, so that  $Ta_2MSe_7$  is a two-dimensional metal. Nevertheless, the two hole pockets arising from the selenium p-block band are nested by the vector  $q \approx (0.1b^*, 0)$ , and therefore an interesting CDW phenomenon is expected.

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**Registry No. NiTa<sub>2</sub>S<sub>5</sub>**, 98104-46-2; NiTa<sub>2</sub>Se<sub>7</sub>, 104548-68-7.

<sup>(9)</sup> For an analysis of the band structure of a double **OCT** chain, see: Canadell, E.; Whangbo. M.-H. *Inorg. Chem.* **1986**, 25, 1488. **(10) Ammeter, J. H.; Bürgi**, H.-B.; Thibeault, J.; Hoffmann, R. *J. Am.* 

*Chem. SOC.* **1978,** *100,* **3686.**