

## Communications

The Structure of  $d^0$   $ML_6$  Complexes

In a study of the driving force for agostic bonding, one of us put forth<sup>1</sup> the idea that certain  $d^0$   $ML_6$  complexes were unstable at an octahedral geometry. Extended Hückel calculations on  $TiH_6^{2-}$  and other molecules found that distortions away from an octahedral geometry were in fact stabilizing. This was challenged in a recent communication<sup>2</sup> by Cameron, Fitzgerald, and Zerner on the basis of INDO and ab initio calculations for  $TiH_6^{2-}$ . The octahedral geometry was found to be more stable than a  $C_{2v}$  variant, one of the distortion modes originally considered;<sup>1</sup> however, the potential was soft when correlation was included at the second-order Møller-Plesset (MP2) level. In this communication we report ab initio calculations on  $TiH_6^{2-}$ ,  $CrH_6$ , and  $CrF_6$  at the MP2 and full MP4 levels.

The central thesis that certain  $d^0$   $ML_6$  complexes will distort from an octahedral geometry is derived from potential second-order Jahn-Teller forces that serve to mix the  $t_{1u}$  HOMO with the  $t_{2g}$  LUMO in these molecules. The normal modes available to do this are of  $t_{2u}$  and  $t_{1u}$  symmetry. Perhaps the most simple  $t_{2u}$  deformation is rotation about one 3-fold axis to yield a trigonal-prismatic,  $D_{3h}$ , geometry. As shown in Figure 1, this stabilizes two members of the filled  $t_{1u}$  set.<sup>3</sup> A further pyramidization from  $D_{3h}$  to  $C_{3v}$  geometry couples the filled  $a_2''$  and empty  $a_1'$  orbitals, stabilizing the former. This is, of course, but one of the multiple distortions available to an octahedral system. Many  $d^2$  complexes have structures along the  $O_h \rightarrow D_{3h}$  path.<sup>4</sup> A number of  $d^4$  molecules are distorted in a  $C_{2v}$  sense.<sup>5</sup> What is intriguing about the  $d^0$  case is that with a spherically symmetrical field at the metal, VSEPR and crystal field arguments clearly favor an octahedral arrangement of ligands. The  $C_{3v}$  and  $D_{3h}$  geometries are sterically much more demanding than  $O_h$ .<sup>6</sup>

We emphasize that some  $d^0$   $ML_6$  compounds will remain octahedral. The necessary conditions for distortion are 3-fold: First, for obvious steric reasons the ligands must not be bulky. Second, the  $t_{2g}-t_{1u}$  energy gap should be small since the stabilization afforded to the  $t_{1u}$  set is inversely dependent upon this energy difference.<sup>7</sup> In a one-electron sense, there are several lower lying orbitals that are destabilized along the distortion modes given in Figure 1, and these serve as a breaking force for the distortion.

A small  $t_{2g}-t_{1u}$  energy gap is favored when the ligands are strong  $\sigma$  donors, the ligands possess little or no  $\pi$ -donating capability, and the metal is not too electropositive.<sup>1</sup> Third, the metal-ligand bond should be covalent. An extreme, ionic case simply reduces to a point charges on a sphere problem. The mixing of  $t_{2g}$  into  $t_{1u}$  serves to increase the overlap (or covalency) in the M-L bond.<sup>1</sup> With these factors in mind we proceed to  $CrF_6$ ,  $TiH_6^{2-}$ , and  $CrH_6$ .

The ab initio calculations<sup>8</sup> were carried out with two basis sets. The first, referred to as basis I, started with Huzinaga's (4333/43/4) primitive set<sup>9</sup> for Ti and Cr. The d along with 4s and augmented 4p functions were replaced by ones optimized for molecular environments<sup>10,11</sup> to yield a basis of the form (4333/433/31). Standard 4-31G basis sets for H and F were used.<sup>12,13</sup> Basis II used a Wachters basis<sup>14</sup> for the metal, where the primitive d functions were replaced with an optimized (411) set<sup>10</sup> and the 4s/4p functions were replaced with an augmented set.<sup>11,15</sup> The H atoms used a 6-311 G\*\* basis.<sup>13,16</sup>

With use of basis I at the Hartree-Fock (HF) level, the optimized structure for  $TiH_6^{2-}$  with an  $O_h$  geometry was found to be in good agreement with that reported by Cameron, Fitzgerald, and Zerner.<sup>2</sup> A stationary point for a  $C_{3v}$  structure was not found; the geometry collapsed to  $D_{3h}$ . A summary of the geometries and energies for all systems in this study is reported in Table I. The  $O_h$  structure was found to be 10.7 kcal/mol more stable than the  $D_{3h}$  structure. Using Cameron, Fitzgerald, and Zerner's DZP basis,<sup>2</sup> we find essentially the same result with the  $O_h$  geometry to be 8.0 kcal/mol more stable than the optimized  $D_{3h}$  one.<sup>17</sup> With

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- (17) The optimal  $D_{3h}$  structure is one with  $r_1 = 1.83$  Å and  $\theta = 51.2^\circ$ . The vibrational frequencies calculated at the  $O_h$  geometry by using this basis are all real; however, the  $t_{2u}$  set is quite small ( $374$   $cm^{-1}$ ).

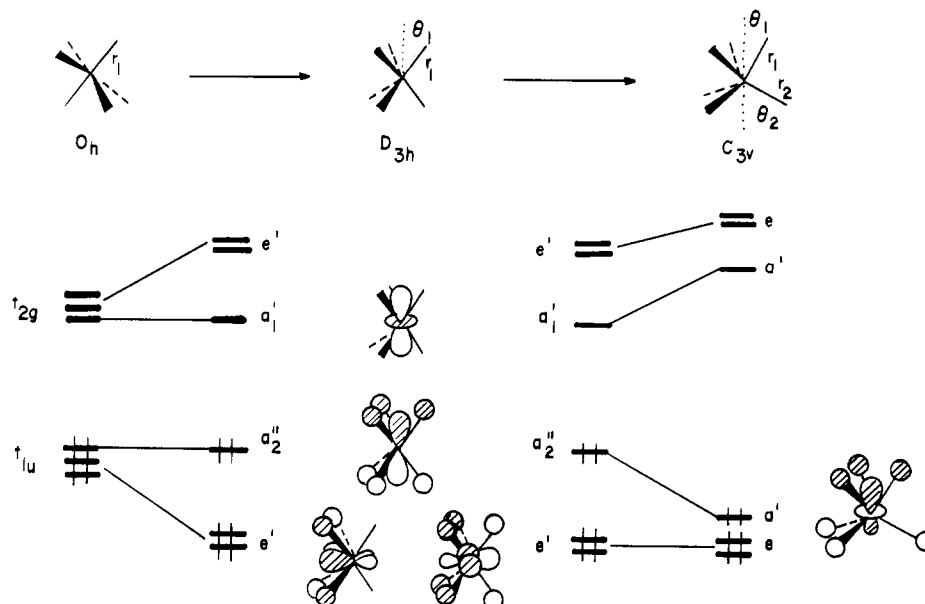


Figure 1. Orbital correlation diagram for a trigonal twist from an  $O_h$  to a  $D_{3h}$  geometry and pyramidization to a  $C_{3v}$  structure.

Table I. Summary of the Geometries and Energies with Basis I

molecule	geometry <sup>a</sup>				method	total energy <sup>b</sup>	rel energy <sup>c</sup>
	$r_1$	$r_2$	$\theta_1$	$\theta_2$			
TiH <sub>6</sub> <sup>2-</sup>	$O_h$	1.85			HF	-850.743 68	0
	$D_{3h}$	1.84			HF	-850.726 64	10.7
	$O_h$	1.86			MP2	-851.033 75	1.3
	$D_{3h}$	1.84			MP2	-851.035 79	0.0
CrH <sub>6</sub>	$O_h$	1.69			MP2	-1045.463 51	155.3
	$D_{3h}$	1.58			MP2	-1045.679 57	19.7
	$C_{3v}$	1.52	1.55	36.8	MP2	-1045.711 03	0.0
CrF <sub>6</sub>	$O_h$	1.71			MP2	-1639.087 79	0.0
	$D_{3h}$	1.73			MP2	-1639.058 31	18.5

<sup>a</sup>The geometrical parameters are defined in Figure 1;  $r$  values are in angstroms and  $\theta$  values in degrees. <sup>b</sup>Total energies are listed in hartrees. <sup>c</sup>Relative energies are listed in kcal/mol.

use of basis II, this energy difference is 9.1 kcal/mol (see Table II). Including correlation reverses this prediction. The optimized structure for the  $D_{3h}$  geometry is 1.3 kcal/mol more stable than the  $O_h$  structure at the MP2 level with basis I and 3.2 kcal/mol with basis II. This difference is increased further to 9.2 kcal/mol at the full fourth-order Møller–Plesset (MP4) level with basis II. Thus, the relative energies appear to be sensitive to the level of correlation.

The situation for CrH<sub>6</sub> is more striking. Stationary points for  $O_h$ ,  $D_{3h}$ , and  $C_{3v}$  geometries were located (Table I), and the relative stabilities increase in the order given. The  $C_{3v}$  structure is computed to be 174.6 kcal/mol more stable than the  $O_h$  one at the MP4 level with basis II (Table II)! In fact, both the  $O_h$  and  $D_{3h}$  geometries are higher order saddle points, while the  $C_{3v}$  structure is a local minimum on the potential energy surface for CrH<sub>6</sub>.<sup>18</sup>

To show that trigonal-prismatic or  $C_{3v}$  structures are not an artifact of the basis sets or methods used, we have also carried out optimizations of CrF<sub>6</sub> with basis I at the MP2 level. Again, starting from  $C_3$  or  $C_{3v}$  geometries resulted in collapse to a  $D_{3h}$  structure that was 18.5 kcal/mol higher in energy than the  $O_h$  structure. While CrF<sub>6</sub> is a known molecule, its instability has frustrated a precise determination of the structure (an octahedral geometry was found in an Ar matrix).<sup>19</sup> However, a large number

Table II. Summary of the Energies with Basis II<sup>a</sup>

molecule	method	total energy <sup>b</sup>	rel energy <sup>c</sup>
TiH <sub>6</sub> <sup>2-</sup>	$O_h$	-851.585 22	0.0
	$D_{3h}$	-851.570 65	9.1
	$O_h$	-851.789 07	3.2
	$D_{3h}$	-851.794 16	0.0
CrH <sub>6</sub>	$O_h$	-851.837 82	9.2
	$D_{3h}$	-851.852 42	0.0
CrF <sub>6</sub>	$O_h$	-1046.446 46	158.9
	$D_{3h}$	-1046.658 06	26.1
	$C_{3v}$	-1046.699 75	0.0
	$O_h$	-1046.496 68	174.6
	$D_{3h}$	-1046.747 08	17.5
	$C_{3v}$	-1046.774 89	0.0

<sup>a</sup>The geometries used were taken from the optimized structures at the basis I/MP2 level. <sup>b</sup>Total energies are listed in hartrees. <sup>c</sup>Relative energies are listed in kcal/mol.

of d<sup>0</sup> MX<sub>6</sub> molecules, and this includes MoF<sub>6</sub>,<sup>20</sup> WF<sub>6</sub> in the gas and solid states,<sup>21</sup> and the isoelectronic TiF<sub>6</sub><sup>2-</sup> in three different crystalline environments,<sup>22</sup> are octahedral.<sup>23,24</sup>

(18) The calculated frequencies at the MP2 level with basis I were computed to be all real for the  $C_{3v}$  geometry, the  $D_{3h}$  geometry had  $e'$  and  $a_2''$  imaginary frequencies, and the  $t_{1u}$  and  $t_{2g}$  frequencies for the octahedral geometry were imaginary.

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The diversity of behavior offered by these three  $d^0$   $ML_6$  complexes is easily accommodated by the ideas presented previously. In  $CrH_6$  the ligands are strong  $\sigma$  donors and the bonding is reasonably covalent. The fluorine ligands in  $CrF_6$  are much weaker  $\sigma$  donors (and have some  $\pi$ -donating capability) hence, a return to an octahedral structure is favored. Notice in Table I that the optimized Cr-H bond lengths are much shorter at the  $C_{3v}$  (or  $D_{3h}$ ) geometry than at  $O_h$ , in agreement with our idea that a distortion from  $O_h$  symmetry will lead to larger M-L overlap. Exactly the reverse is seen for  $CrF_6$ , where the bonding is much more ionic. The situation for  $TiH_6^{2-}$  is intermediate. Certainly Ti is more electropositive than Cr and the bonding in  $TiH_6^{2-}$  is more ionic than in  $CrH_6$ .<sup>25</sup> Thus, deformation all the way to  $C_{3v}$  is not favored. However, our computations at the highest level point to a  $D_{3h}$  geometry being more stable than an  $O_h$  one, contrary to a previous assertion.<sup>2</sup> We have recently been informed<sup>26</sup> that the X-ray structure of the related  $ZrMe_6^{2-}$  molecule has been completed and that it is indeed trigonal prismatic, in accord with our computations on  $TiH_6^{2-}$ . A study of the full potential energy surfaces for these and other  $d^0$   $ML_6$  molecules will be forthcoming.

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### Binuclear Polyhydrides of Tantalum. Reductive Desulfurization of a Bridging Thioether Ligand and Formation of a $Ta_2(\mu-S)$ Unit. X-ray Crystal Structure of $[(Pr^i)_2P(CH_2)_3PPr^i]_2TaHCl_2(\mu-S)(\mu-H)_2$

Pentahydride complexes of tantalum of the type  $TaH_5P_4$  ( $P_2 = dmpc^{1a}$  and  $P = PMe_3^{1b}$ ) are members of a large family of phosphine-stabilized polyhydride complexes having the formula  $MH_5P_y$ . Other members of this family include  $WH_6P_3$ ,<sup>2a</sup>

$WH_4P_4$ ,<sup>2b</sup>  $ReH_7P_2$ ,<sup>2c</sup>  $ReH_5P_3$ ,<sup>2d</sup>  $OsH_4P_3$ ,<sup>2e</sup>  $IrH_3P_3$ ,<sup>2f</sup> and  $IrH_5P_2$ ,<sup>2g</sup> the stoichiometry of these complexes is dependent on the method of preparation, the size of the phosphine ligand, and the amount of the phosphine present. A characteristic feature of all of these systems is that they are coordinatively saturated, and thus any reactivity by these complexes must be preceded by either phosphine dissociation<sup>3</sup> or loss of  $H_2$ .<sup>4</sup> In an effort to induce coordinative unsaturation, we have examined the effect of bulky, chelating bidentate phosphines as a possible strategy to make these complexes phosphine deficient. In this paper we describe an approach we have used in an attempt to generate polyhydrides of tantalum stabilized by only two phosphine donors per tantalum. However, the success of the reaction depends critically on the choice of starting material.

As a starting point, the ability to oxidatively add  $H_2$  to low-valent binuclear Ta(III) complexes<sup>5</sup> appeared attractive. Since the binuclear Ta(III) complex  $Ta_2Cl_6(SMe_2)_3$  (**1**) has been reported<sup>6</sup> to be a precursor to low-valent, bioctahedral complexes containing chelating diphosphine ligands, we examined simple displacement reactions of **1**. However, reaction of **1** with the bulky, chelating ligand 1,3-bis(diisopropylphosphino)propane (dipp) does not go to completion even at high temperatures and results in impure complexes having the approximate formula  $[(dipp)_2TaCl_2]_2(\mu-SMe_2)(\mu-Cl)_2$  with dangling phosphines<sup>7</sup> (by  $^{31}P\{^1H\}$  NMR spectroscopy). An alternative synthesis was attempted. The Na/Hg reduction of  $Ta_2Cl_6(SMe_2)_3$  in the presence of 2 equiv of dipp under 4 atm of  $H_2$  provided air-sensitive, black crystals in 70% yield. One recrystallization from toluene at  $-30^\circ C$  provided an analytically pure sample<sup>8</sup> of a material having the formula  $(dipp)_2Ta_2H_2Cl_2S$  (**2**). The  $^{31}P\{^1H\}$  NMR spectrum<sup>9</sup> at  $25^\circ C$  shows two doublets at 18.8 and 15.6 ppm with some fine structure, indicating that two inequivalent, coordinated phosphines are present. The  $^1H$  NMR spectrum<sup>10</sup> of **2** shows a broad triplet at 12.4 ppm and a broad multiplet at 7.7 ppm, both integrating for two protons, which we attribute to two bridging and two terminal hydrides respectively; both of these resonances are absent in the deuterium analogue  $(dipp)_2Ta_2D_4Cl_2S$  (**2-d**). The IR spectrum of **2** (KBr disk) shows two strong broad bands at 1759 and 1553  $cm^{-1}$  that shift to 1263 and 1114  $cm^{-1}$ , respectively, for **2-d**.

The solid-state structure<sup>11</sup> of **2** (Figure 1a) shows that there

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- (9)  $^{31}P\{^1H\}$  NMR measurements were recorded at 121.4 MHz on a Varian 300 XL spectrometer. Chemical shifts ( $\delta$ ) are in ppm from external  $P(OMe)_3$  set at 141.0 ppm.
- (10)  $^1H$  NMR for **2** ( $C_6D_6$ , ppm):  $H_b$ , 12.39 (br t, 2,  $J_{PH} = 48$  Hz);  $H_t$ , 7.70 (br m, 2); ligand resonances,  $PCHMe_2$ , 2.82 (overlapping d of sept, 2);  $PCHMe_2$ , 2.55 (overlapping d of sept, 2);  $PCHMe_2$ , 2.06 (overlapping d of sept, 4); 1.64-1.03 (m, 60).
- (11) Black crystals of **2** were grown from concentrated THF/hexanes solutions at  $-30^\circ C$ .  $[(dipp)_2TaHCl]_2(\mu-S)(\mu-H)_2$  crystallizes in the monoclinic space group  $P2_1/c$ ,  $a = 10.730$  (2) Å,  $b = 15.653$  (2) Å,  $c = 13.402$  (3) Å,  $\beta = 110.78$  (2)°,  $\rho_c = 1.61$  g  $cm^{-3}$  and  $Z = 2$ . The structure was solved by conventional heavy-atom methods and was refined by successive difference Fourier map calculations to  $R = 0.0614$  and  $R_w = 0.0665$  for 2365 absorption-corrected reflections with  $F_o^2 > 3\sigma F_o^2$ . Hydrogen atoms were placed in idealized positions during the final refinement. All non-hydrogen atoms were refined anisotropically. Data were collected at  $24^\circ C$  on a four-cycle Syntex P2<sub>1</sub> diffractometer with graphite-monochromatized Mo  $K\alpha$  radiation.

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