

CONTRIBUTION FROM THE LAWRENCE RADIATION LABORATORY AND DEPARTMENT OF CHEMISTRY,
UNIVERSITY OF CALIFORNIA, BERKELEY, CALIFORNIA

The Crystal Structure of Thorium Tetraiodide¹

BY ALLAN ZALKIN, J. D. FORRESTER, AND DAVID H. TEMPLETON

Received August 12, 1963

Thorium tetraiodide has a novel layer structure in the solid according to single-crystal X-ray diffraction data. The crystals are monoclinic, space group $P2_1/n$, with $a = 13.216 \pm 0.007$, $b = 8.068 \pm 0.006$, $c = 7.766 \pm 0.006$ Å, $\beta = 98.68 \pm 0.05^\circ$, $Z = 4$, $d_x = 6.00$ g./cc. Each thorium atom has eight iodine neighbors at an average distance of 3.20 Å. at the corners of an irregular polyhedron which is approximately a square antiprism. These polyhedra share edges and triangular faces to form layers which are only weakly bonded to each other.

Introduction

In 1950 we made and analyzed crystals of thorium tetraiodide. Weissenberg X-ray diffraction photographs of a crystal rotated about the b -axis showed that it is monoclinic and gave the unit cell dimensions. The space group was ambiguous because we did not have the $0k0$ data. Because of the complexity of the problem we suspended work on it. With the advent of high speed computers and improved techniques for intensity measurement we reactivated this problem and successfully determined the crystal structure.

In 1954 Jantsch, *et al.*,² reported on the basis of a powder diffraction pattern that ThI_4 might be tetragonal. D'Eye, *et al.*,³ concluded that the crystal is orthorhombic and reported cell dimensions. In the present work we show conclusively that ThI_4 is monoclinic in space group $P2_1/n$. An inspection of D'Eye's published $\sin^2 \theta$ values shows them to correspond to ours, but we cannot find any significant relation of the orthorhombic cell dimensions to the monoclinic lattice which we determine. The reproduction of the powder diagram in Jantsch's paper is so poor that we cannot comment concerning his data.

Experimental

Preparation of Crystals.— ThI_4 was prepared by heating thorium metal foil at about 500° in an iodine atmosphere in an evacuated Pyrex T-shaped tube for 1 week. The iodine was at the bottom of the tube stem at room temperature, and the thorium was in the side arm immersed in a furnace. As the ThI_4 was formed it distilled down the tube away from the hot reaction zone. At the higher temperatures near the reaction zone the material is orange colored, but it changes to yellow at a lower temperature. After reaction was completed the material was distilled to the other side arm and sealed off. Where the thorium foil had been there remained a gray powder which was presumably unreacted thorium and thorium oxide. As ThI_4 is very reactive to the atmosphere it was necessary to handle the material in a drybox; our drybox is charged with dry nitrogen.

The ampoule of ThI_4 was broken open, and small fragments of the yellow material were loaded into 0.3-mm. diameter vitreous silica capillaries. The capillaries were sealed and baked in a furnace at about 550° . There was enough ther-

mal gradient to allow crystals to grow on the cooler walls of the capillary.

X-Ray Diffraction.—A suitable small crystal plate of dimensions approximately $0.15 \times 0.10 \times 0.03$ mm. was attached to the wall of the capillary with its b -axis in the plate and along the axis of the tube. It was mounted in a goniostat on a General Electric XRD-5 apparatus equipped with a molybdenum X-ray tube, a scintillation counter, and a pulse-height discriminator. Cell dimensions and crystal settings were calculated with $\lambda (\text{K}\alpha_1) = 0.70926$ Å. Intensities were measured for the 1454 independent reflections permitted by the space group out to a limit of $\sin \theta/\lambda = 0.596$ ($2\theta = 50^\circ$). Of these intensities 106 were recorded as zero. The reflections were measured by counting the peak intensity for 10 sec. The settings were verified for selected reflections by step scanning of the angle θ .

The data were corrected for the Lorentz-polarization effects. These and other calculations were made with an IBM-7090 computer using the Zalkin data-processing and Fourier programs and our version of the Gantzel-Sparks-Trueblood least-squares program (all unpublished). The function minimized in least squares was $\sum w(|F_o| - |F_c|)^2/\sum w F_o^2$, where w is the weighting factor and F_o and F_c are the observed and calculated structure factors. Because of the fixed-time counting technique we took the weighting factors as unity. Atomic scattering factors were taken for neutral Th and neutral I (Ibers⁴) modified for dispersion by adding -6.0 and -0.5 electrons (Templeton⁵). The imaginary dispersion terms were neglected.

Correction for Absorption for Thin Flat Crystals.—The data were corrected for absorption by a method which is based on experimental measurement of the absorption effect for certain reflections. This method applies to thin flat crystals mounted with the goniometer axis parallel with the plane of the crystal. We have found this technique convenient for several substances which we have studied. In the present case it accomplished a considerable reduction in the discrepancies of observed and calculated structure factors.

If the crystal is in the shape of a thin circular disk, the absorption correction depends only on the angles P and Q between the incident and diffracted beams, respectively, and the normal to the disk. We make the approximation that the correction depends on the harmonic mean value m of $|\cos P|$ and $|\cos Q|$:

$$\frac{2}{m} = \frac{1}{|\cos P|} + \frac{1}{|\cos Q|}$$

The parameter m is inversely proportional to the distance through the crystal which is traversed by a photon scattered at the center of the crystal. This procedure treats alike the

(1) Work done under the auspices of the U. S. Atomic Energy Commission.

(2) G. Jantsch, J. Homayr, and F. Zemek, *Monatsh.*, **85**, 526 (1954).

(3) R. W. M. D'Eye, I. F. Ferguson, and E. J. McIver, *Congr. intern. chim. pure appl.*, 16^e Paris 1957, *Mem. sect. chim. minerale*, 341 (1958).

(4) J. A. Ibers in "International Tables for X-Ray Crystallography," Kynoch Press, Birmingham, England, 1962, Vol. III, p. 211.

(5) D. H. Templeton in "International Tables for X-Ray Crystallography," Kynoch Press, Birmingham, England, 1962, Vol. III, p. 216.

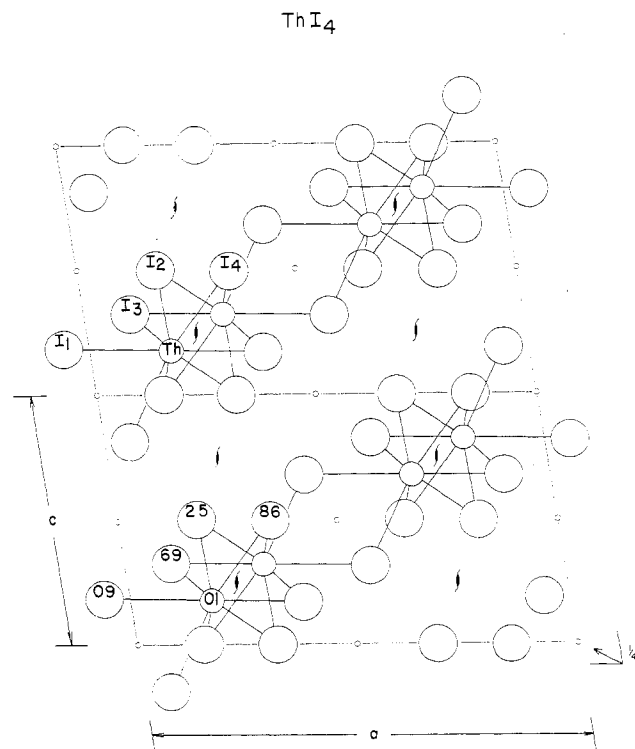


Fig. 1.—Projection of the ThI_4 structure onto the ac -plane. The numbers on atoms in the lower part of the figure are y -coordinates ($\times 100$).

cases in which the diffracted beam emerges from the same side or opposite side as the incident beam. The geometry of a measurement is determined by θ , the Bragg angle, and ϕ and χ , the settings of the spindle and vertical circle of the goniostat, respectively. By spherical trigonometry we have the relations

$$|\cos P| = |\cos \theta \cos \phi + \sin \theta \sin \phi \cos \chi|$$

$$|\cos Q| = |\cos \theta \cos \phi - \sin \theta \sin \phi \cos \chi|$$

if the crystal is perpendicular to the incident beam when $\theta = \phi = 0$. If the setting of ϕ is otherwise, a suitable constant is added to ϕ in the above equations.

For reflections at $\chi = 90^\circ$, $m = |\cos P| = |\cos Q|$, and rotation of ϕ permits measurement of the intensities as a function of m over nearly the entire range 0 to 1. These data are fitted by a convenient empirical function of m which when multiplied by the measured intensities corrects them to a constant value for each reflection. This function is then used as a correction factor for the rest of the data.

In the present experiment, the function was chosen as

$$I(\text{corrected}) = \frac{I(\text{measured})}{1 + 5.25 \exp(-0.69/m)}$$

The values of this correction differed by as much as a factor of 3.6.

Results

Unit Cell and Space Group.—The crystals are monoclinic with unit cell dimensions: $a = 13.216 \pm 0.007$, $b = 8.068 \pm 0.006$, $c = 7.766 \pm 0.006$ Å, $\beta = 98.68 \pm 0.05^\circ$, $V = 819$ Å³. With four formula units (ThI_4) per unit cell, the density calculated from the X-ray data is 6.00 g./cc.

The systematic absence of $h0l$ reflections if $h + l$ is odd and of $0k0$ reflections if k is odd is characteristic of

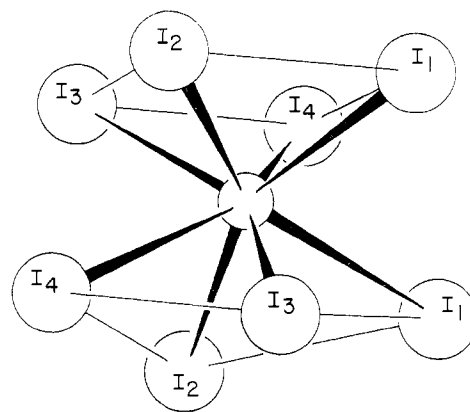


Fig. 2.—The approximately square antiprism arrangement of iodine about thorium in ThI_4 .

TABLE I
FINAL POSITIONAL PARAMETERS AND STANDARD DEVIATIONS FOR
 ThI_4

Atom	x	y	z	$\sigma(x)$	$\sigma(y)$	$\sigma(z)$
Th	0.1835	0.0149	0.1769	0.0001	0.0002	0.0002
I(1)	0.0587	0.9098	0.8094	0.0002	0.0003	0.0003
I(2)	0.1801	0.2535	0.4984	0.0002	0.0003	0.0003
I(3)	0.0972	0.6917	0.3251	0.0002	0.0003	0.0003
I(4)	0.1517	0.3638	0.0014	0.0002	0.0003	0.0003

space group $P2_1/n$ (C_{2h}^5). This space group is confirmed by the structure determination.

Determination of the Structure.—In order to obtain a trial structure we calculated a three-dimensional Patterson function. By inspection of this function with due regard to packing and interatomic distances we found a complete trial structure which accounted satisfactorily for the largest peaks. The thorium atom and all four iodine atoms in the asymmetric unit were placed in 4-fold general positions: $4(e)$, $\pm(x, y, z; 1/2 - x, 1/2 + y, 1/2 - z)$. The positional parameters obtained from the Patterson function were all within 0.03 of the corresponding final values given in Table I.

Least-squares refinement was started using this trial structure. Each atom was given an isotropic temperature factor of the form $\exp(-B\lambda^{-2} \sin^2 \theta)$, and only 622 reflections, each with unit weight, were used in the refinement. The conventional unreliability factor $R = \Sigma \|F_o\| - |F_c| / \Sigma \|F_o\|$ was 0.16 at the end of four cycles of refinement. Encouraged by this, we calculated four further cycles using all 1454 terms, each with unit weight. This calculation resulted in $R = 0.17$.

At this stage it became apparent that the shape of the crystal, combined with the high value of the absorption coefficient, was affecting the diffracted intensity in a systematic way. An absorption correction was applied

TABLE II
ANISOTROPIC THERMAL PARAMETERS ($\times 10^4$) AND AVERAGE ROOT MEAN SQUARE DISPLACEMENTS

Atom	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}	u , Å.
Th	22	42	31	-1	13	0	0.12
I(1)	25	80	45	2	15	-16	0.14
I(2)	30	67	37	-6	26	-3	0.14
I(3)	25	68	82	3	25	15	0.15
I(4)	34	59	51	-2	2	7	0.15

to the data as described earlier. After four more cycles of refinement with all the data R fell to 0.120. Correction of several blunders in the data measurements and in data card punching reduced R to 0.116 after four more cycles.

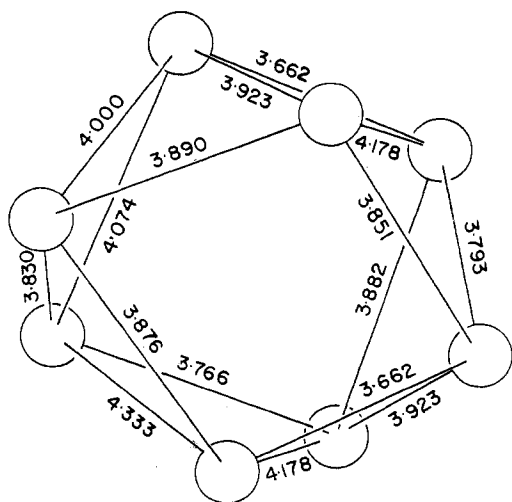


Fig. 3.—The iodine-iodine distances in the ligand polyhedron. Standard deviations of the distances are 0.005 Å.

Introduction of anisotropic thermal factors of the form $\exp(-\beta_{11}h^2 - \beta_{22}k^2 - \beta_{33}l^2 - 2\beta_{12}hk - 2\beta_{13}hl - 2\beta_{23}kl)$ for each of the five atoms resulted in only slight further improvement of the agreement. After four cycles with all the data, R was 0.09. In the last cycle no parameter changed by more than one unit in the sixth decimal place.

The final positional parameters and their standard deviations are listed in Table I. These standard deviations are estimated by the method of least squares assuming that the discrepancies represent random errors. The anisotropic thermal parameters and the average root mean square displacements of the atoms are reported in Table II. The observed structure factor magnitudes and the calculated structure factors are shown in Table III.

A comparison of the final atomic coordinates with those obtained in the first cycles with 622 reflections revealed that no atom moved more than about 0.01 Å.

Description of the Structure.—The ThI_4 consists of layers with the iodine atoms arranged in sheets parallel to $(\bar{1}01)$ and the thorium atoms located in alternate spaces between iodine sheets (Fig. 1). Each thorium atom has eight iodine neighbors arranged at the corners

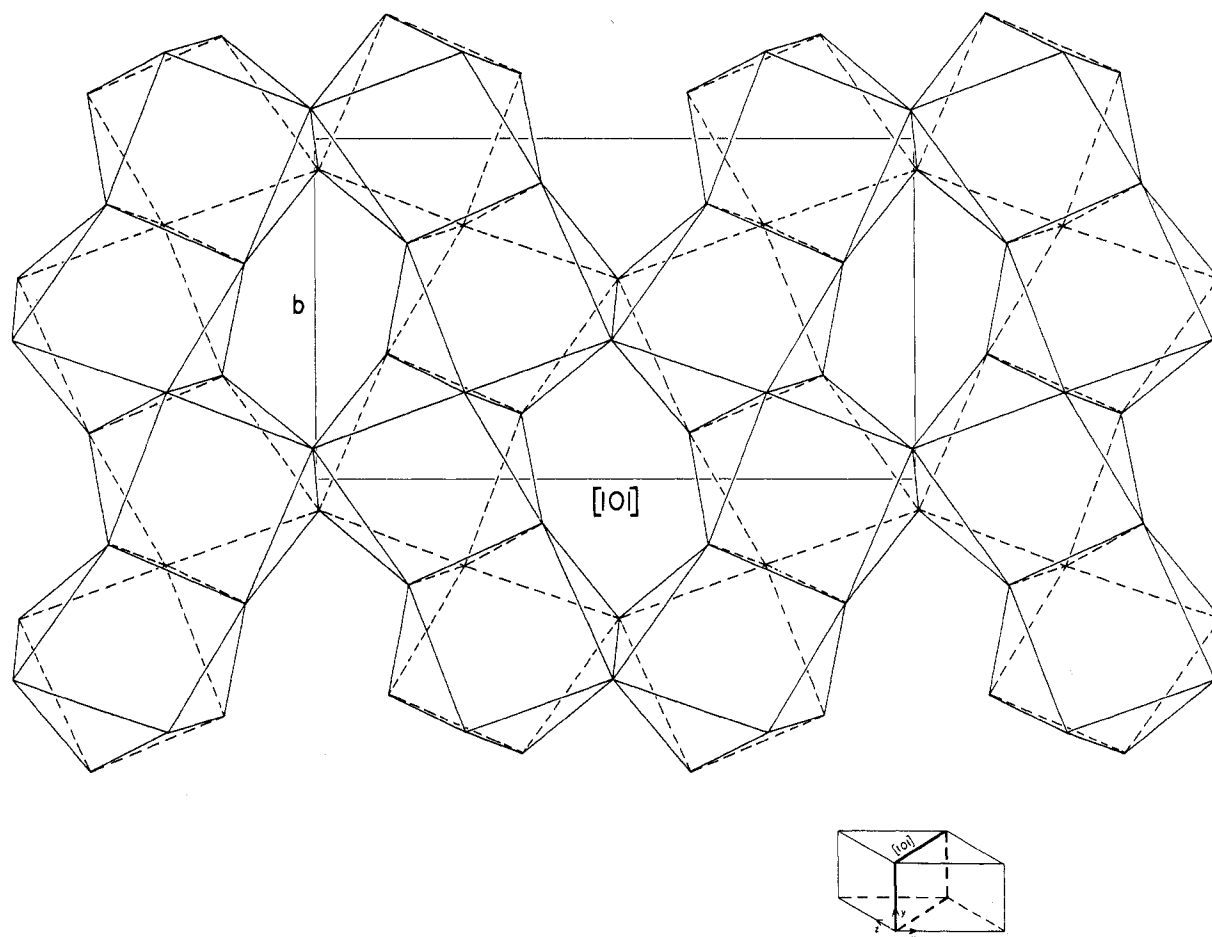


Fig. 4.—The structure of one layer of ThI_4 , indicated by linked polyhedra. A thorium atom is at the center of each polyhedron, and an iodine atom is at each corner.

TABLE VI
IODINE-IODINE DISTANCES BETWEEN LAYERS

Atoms	Distance, ^a Å.
I(1)-I(2)	4.155
I(1)-I(3)	4.252
I(2)-I(3)	4.125
I(2)-I(4)	4.079

^a Standard deviations are 0.006 Å. or less.

oxygen sheets in the BaTi₄O₉ structure,⁶ but the stacking of these sheets is quite different. We know of no other substance with a structure which resembles the entire arrangement.

The layers are bonded together only by weak forces. The I-I contacts between layers (Table VI) all exceed 4.07 Å. The orientation of these layers conforms to the platy {10 $\bar{1}$ } habit of the crystals.

(6) D. H. Templeton and C. H. Dauben, *J. Chem. Phys.*, **32**, 1515 (1960).

CONTRIBUTION FROM THE DEPARTMENT OF CHEMISTRY,
LAFAYETTE COLLEGE, EASTON, PENNSYLVANIA

The Interaction between Ammonium Fluoride and Metal Fluorides as Compressed Powders

BY DAVID S. CROCKET, III, AND ROBERT A. GROSSMAN

Received September 9, 1963

Finely powdered mixtures of metal fluorides interact with NH₄F to form double salts of the type NH₄MF₃ (where M = Co⁺², Cu⁺², Mn⁺², Cd⁺², Ni⁺², or Zn⁺²) at pressures of 100,000 p.s.i. at 25°. In the case of Co, Cu, Ni, and Zn compounds heating to 110° caused the formation of (NH₄)₂MF₄ type compounds which are isostructural with K₂NiF₄.

Introduction

The interchange of anions in powdered mixtures subjected to pressure has been shown by Meloche and Kalbus,¹ Ketelaar,² and others. Nonnenmacher and Merke³ have also shown that complex ions can be formed under pressure in the formation of KBr disks used in infrared spectral determinations.

With this in mind it was decided to investigate a series of known fluoride double salts as to the interaction of powdered mixtures of their components under various conditions of temperature and pressure.

Experimental

Reagents.—The NH₄F used throughout was J. T. Baker or Mallinckrodt analytical grade reagent. It was further purified by dissolving in methyl alcohol and reprecipitating by the addition of ether. Cobalt(II) fluoride and manganese(II) fluoride were prepared by the reaction of the carbonates with HF followed by oven heating to remove water. The infrared spectra of these compounds showed them to be anhydrous. Copper(II) fluoride and nickel fluoride were formed by reaction of the metal bromide with ClF₃ at 250–300°. X-Ray diffraction patterns gave lines attributable only to the fluorides.

Complex Fluorides.—The known fluoride complexes NH₄-ZnF₃, NH₄CoF₃, NH₄CuF₃, NH₄CdF₃, and NH₄MnF₃ for comparison purposes were prepared by reaction of the metal bromides with ammonium fluoride in methyl alcohol as fully described by Haendler, *et al.*⁵

Analytical.—Copper, zinc, nickel, and cobalt were determined using standard EDTA titrations.

Ammonia was determined by distillation from strongly alkaline solutions into cold boric acid solution, followed by titration with hydrochloric acid, using brom cresol green indicator.

Fluorine was determined by distillation from sulfuric acid solution⁶ followed by titration of aliquots of distillate with cerium(III) chloride, using murexide indicator.⁷

Infrared spectra were determined using a Beckman IR-5 spectrophotometer with sodium chloride optics. Samples were run as Halocarbon mulls⁸ on calcium fluoride plates.

X-Ray diffraction patterns were taken using a G.E. RDX-5 diffractometer with filtered copper K α and iron K α radiation.

Procedure for Determination of Complex Formation in the Mixtures of NH₄F and Metal Fluorides.—The particular metal fluoride involved was well ground with a large excess of ammonium fluoride in an agate mortar and pestle. The mixture was then placed in a die constructed of a 0.5-in. i.d. steel bushing⁹ and two hardened steel pins, one 0.5 in. long, the other 2.5 in. long; the diameter was such as just to allow passage through the bushing. The mixture was then pressed into a pellet by subjecting it to a pressure of about 100,000 p.s.i. in a Carver 10-ton bench press. Depending upon the circumstances the pellet was held at this pressure for a certain length of time or removed immediately. Again depending upon the purpose of the particular run the pellet might or might not be heated. The pellet was then reground, leached with methyl alcohol to remove the unchanged ammonium fluoride, and dried with anhydrous ether. Since the complex fluorides are very insoluble in methyl alcohol, and the metal fluorides are for the most part insoluble in methyl alcohol also, these remained as a residue. The formation of any complex was then shown from a comparison of the infrared spectrum or X-ray powder pattern of the residue with that of an authentic sample, and the amount of complex formed could be determined from an ammonia analysis, since the only ammonium ion remaining in the residue was tied up in the complex.

(1) V. W. Meloche and G. E. Kalbus, *J. Inorg. Nucl. Chem.*, **6**, 104 (1958).

(2) J. A. A. Ketelaar, C. Haas, and J. van der Elshen, *J. Chem. Phys.*, **24**, 624 (1956).

(3) W. Nonnenmacher and R. Merke, *Chem. Ber.*, **93**, 1279 (1960).

(4) J. Parcher, B.S. Thesis, Lafayette College, Easton, Pa., 1962.

(5) (a) H. M. Haendler, F. A. Johnson, and D. S. Crocket, *J. Am. Chem. Soc.*, **80**, 2662 (1958); (b) D. S. Crocket and H. M. Haendler, *ibid.*, **82**, 4158 (1960); (c) A. E. Baker and H. M. Haendler, *Inorg. Chem.*, **1**, 127 (1962).

(6) W. B. Huckabay, E. T. Welch, and A. V. Metter, *Ind. Eng. Chem., Anal. Ed.*, **19**, 154 (1947).

(7) C. Brunisholz and J. Michod, *Helv. Chim. Acta*, **37**, 598, 1546 (1954).

(8) D. S. Crocket and H. M. Haendler, *Anal. Chem.*, **31**, 626 (1959).

(9) Accurate Bushing Co., Garwood, N. J.