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The Crystal and Molecular Structure of Di(tetra-*n*-butylammonium)cobalt(II) Bis(maleonitrile dithiolate) and the Geometry of the Divalent Cobalt(II) Bis(maleonitrile dithiolate) Ion¹

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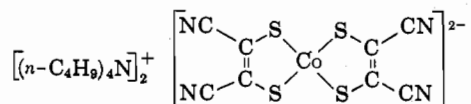
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The crystal and molecular structure of di(tetra-*n*-butylammonium)cobalt(II) bis(maleonitrile dithiolate) has been determined from an X-ray diffraction study of a single crystal specimen. A convenient triclinic cell, space group $P\bar{1}$, with $a = 10.77 \pm 0.01$, $b = 12.35 \pm 0.01$, $c = 9.81 \pm 0.01$ Å, $\alpha = 88.5 \pm 0.1$, $\beta = 114.8 \pm 0.1$, and $\gamma = 93.5 \pm 0.1^\circ$, contains one formula unit. Atomic parameters were refined by least-squares methods using full three-dimensional data out to $\sin \theta/\lambda = 0.482$. The anion is closely planar with the sulfur atoms in a square arrangement around the cobalt atom. The symmetry of the anion is mmm to a very close degree. Three of the *n*-butyl chains in the cation adopt the *trans* conformation while the other has the *gauche* conformation with a dihedral angle of 71.6° . The cobalt atoms are very well separated with the closest distance of approach as the shortest axial length of 9.81 Å. Positions of all 36 independent hydrogen atoms were determined from an electron density difference function.

Introduction

There has been considerable interest recently in the electronic energy levels of a series of compounds represented by the general formula $R'_z(M'S_4C_4R_4)^{-z}$, where $z = 0, 1, 2$; $M' = Co, Ni, Cu, Pd, Pt, Zn, Rh$; $R = CN, CF_3, C_6H_5$, etc.; and $R' = (n-C_4H_9)_4N^+$, $(CH_3)_4N^+$, $(C_2H_5)_4N^+$, etc. Interpretation of these results, however, has been somewhat hampered by the lack of any detailed knowledge of the molecular structure of these compounds and certain discrepancies have arisen in the electron spin resonance measurements.

$((n-C_4H_9)_4N)_2^+(Co(S_2C_4N_2)_2^{2-})$, abbreviated $((n-C_4H_9)_4N)_2^+(Co(MNT)_2^{2-})$, and shown schematically below, is a typical member of this series. The anion $(Co(MNT)_2)_2^{2-}$ was first prepared by Gray and co-workers,²



who found that the X-ray powder diffraction patterns for $((n-C_4H_9)_4N)_2^+(M'(MNT)_2^{2-})$ with $M' = Co, Pt, Ni$, and, later, Rh^3 were identical and differ considerably from that found with $M' = Zn$.

Electron spin resonance and magnetic susceptibility measurements on these compounds have been made by Davison, *et al.*,⁴⁻⁶ as well as by Gray, *et al.*,² and Billig, *et al.*⁷ Previous evidence^{8,9} indicated that four-co-

ordinated, low-spin complexes of Co were likely to be square-planar. This is based on magnetic measurements and apparently the only evidence obtained by X-ray diffraction techniques is the molecular structure of cobalt(II) phthalocyanine where Linstead and Robertson^{10,11} showed that the Co atom had a square-planar environment. Little further evidence seems to be available; *e.g.*, Wells.¹²

An X-ray structure determination was completed very recently by Eisenberg, *et al.*,¹³ on another member of this family, *viz.*, $((CH_3)_4N)_2^+(Ni(MNT)_2^{2-})$, and our work is in very close agreement with this.

Maki, *et al.*,¹⁴ have recently used the results of the present crystallographic investigation in a detailed report on the electronic structures of several members of this series of compounds; the principal axes of the *g*-value tensor and of the hyperfine tensor were associated with the symmetry directions of the anion.

Experimental

X-Ray Diffraction.—We are indebted to Dr. A. H. Maki of Harvard University for sending us some well-formed crystals of the complex $(n-C_4H_9)_4N)_2^+(CoS_4C_4(CN)_4)^{-2}$. The melting point, analysis, and conductivity of these crystals are described in a paper by Davison, *et al.*⁵

X-Ray photographs obtained by the Weissenberg technique and copper radiation established the diffraction symmetry of the crystal. A single crystal, in the form of a thin plate of approximate dimensions $0.24 \times 0.20 \times 0.10$ mm. and mounted about an a^* axis, was used for collecting the intensity data. Intensities were measured with a General Electric XRD-5 goniostat equipped with a scintillation counter and a pulse height discriminator. Mo $K\alpha$ radiation was used and the unit cell dimensions are based on $\lambda = 0.70926$ Å. for Mo $K\alpha_1$.

The 2213 independent reflections permitted by the space group in the sphere of reflection with $\sin \theta/\lambda$ less than 0.482 ($2\theta < 40^\circ$) were measured with counting times of 10 sec. each. Of these,

(1) Work done under the auspices of the U. S. Atomic Energy Commission.

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(3) E. Billig, S. I. Shupack, J. H. Waters, R. Williams, and H. B. Gray, *ibid.*, **86**, 926 (1964).

(4) A. Davison, N. Edelstein, R. H. Holm, and A. H. Maki, *ibid.*, **85**, 2029 (1963).

(5) A. Davison, N. Edelstein, R. H. Holm, and A. H. Maki, *ibid.*, **85**, 3049 (1963).

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(7) E. Billig, R. Williams, I. Bernal, J. H. Waters, and H. B. Gray, *ibid.*, **3**, 663 (1964).

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(9) B. N. Figgis and R. S. Nyholm, *ibid.*, 338 (1959).

(10) J. M. Robertson, *ibid.*, 615 (1935).

(11) R. P. Linstead and J. M. Robertson, *ibid.*, 1736 (1936).

(12) A. F. Wells, "Structural Inorganic Chemistry," Oxford University Press, 1962, p. 926.

(13) R. Eisenberg, J. A. Ibers, R. J. H. Clark, and H. B. Gray, *J. Am. Chem. Soc.*, **86**, 113 (1964).

(14) A. H. Maki, N. Edelstein, A. Davison, and R. H. Holm, *ibid.*, **86**, 4580 (1964).

TABLE I
FINAL COORDINATES AND ISOTROPIC THERMAL PARAMETERS, TOGETHER WITH THEIR STANDARD DEVIATIONS,
FOR ALL ATOMS EXCEPT HYDROGEN

Atom	<i>x</i>	$\sigma(x)$	<i>y</i>	$\sigma(y)$	<i>z</i>	$\sigma(z)$	<i>B</i> ^b	$\sigma(B)$
Co	0.0	...	0.0	...	0.0	...	<i>a</i>	
S(1)	0.0401	0.0002	0.1381	0.0002	0.1462	0.0002	<i>a</i>	
S(2)	-0.2156	0.0002	0.0303	0.0002	-0.1093	0.0002	<i>a</i>	
C(1)	-0.1165	0.0007	0.1930	0.0005	0.0876	0.0007	4.7	0.1
C(2)	-0.2283	0.0007	0.1463	0.0006	-0.0214	0.0008	5.0	0.2
C(3)	-0.1243	0.0008	0.2883	0.0007	0.1575	0.0009	6.0	0.2
C(4)	-0.3583	0.0010	0.1883	0.0007	-0.0717	0.0010	7.4	0.2
N(1)	-0.1273	0.0007	0.3662	0.0006	0.2172	0.0008	7.3	0.2
N(2)	-0.4643	0.0010	0.2247	0.0007	-0.1113	0.0010	10.5	0.2
N(3)	0.0867	0.0005	0.2783	0.0004	-0.2929	0.0006	4.4	0.1
C(5)	0.1928	0.0007	0.2589	0.0005	-0.1328	0.0007	4.6	0.1
C(6)	0.2939	0.0007	0.3510	0.0006	-0.0626	0.0008	5.3	0.2
C(7)	0.3797	0.0008	0.3242	0.0007	0.1021	0.0009	6.7	0.2
C(8)	0.4976	0.0010	0.4089	0.0008	0.1789	0.0011	9.2	0.3
C(9)	-0.0023	0.0007	0.1740	0.0006	-0.3418	0.0008	5.0	0.2
C(10)	-0.1167	0.0008	0.1779	0.0007	-0.5001	0.0009	6.9	0.2
C(11)	-0.2078	0.0009	0.0756	0.0008	0.4697	0.0010	8.0	0.2
C(12)	-0.3248	0.0010	0.0689	0.0008	0.3116	0.0011	8.6	0.2
C(13)	0.1538	0.0007	0.3116	0.0005	-0.3963	0.0007	5.0	0.2
C(14)	0.2348	0.0007	0.2286	0.0006	-0.4230	0.0008	5.4	0.2
C(15)	0.3224	0.0008	0.2751	0.0007	-0.5004	0.0009	6.6	0.2
C(16)	0.4075	0.0009	0.1952	0.0007	0.4737	0.0010	7.4	0.2
C(17)	0.0004	0.0007	0.3718	0.0006	-0.2930	0.0008	5.4	0.2
C(18)	-0.0671	0.0008	0.3599	0.0006	-0.1848	0.0009	6.5	0.2
C(19)	-0.1712	0.0009	0.4468	0.0007	-0.2126	0.0010	7.2	0.2
C(20)	-0.2991	0.0011	0.4290	0.0009	-0.3551	0.0012	10.4	0.3

^a Treated anisotropically—see Table III. ^b In Å².

390 were assigned zero intensity and the maximum count was 12,350 counts/sec. for the 100 reflection. No corrections were made for either absorption or extinction. With $\mu = 5.9 \text{ cm.}^{-1}$ for Mo radiation, μR is 0.07 or less, making the absorption effects rather small.

Calculations were made with an IBM 7094 computer using a full matrix least-squares refinement program written by P. K. Gantzel, R. A. Sparks, and K. N. Trueblood, with minor modifications, and Fourier and distance programs written by Zalkin (all unpublished). We minimized the function $\sum w(|F_o| - |F_c|)^2 / \sum w F_o^2$, where F_o and F_c are the observed and calculated structure factors, respectively, and w is the weighting factor.

Atomic scattering factors were taken as the values given by Ibers¹⁵ for neutral Co, S, N, C, and H. Dispersion corrections¹⁶ of 0.4 and 0.1 electron were added to the Co and S scattering factors, respectively. The imaginary part of the dispersion correction is small and was ignored.

Results

Unit Cell and Space Group.—A convenient primitive unit cell contains one formula unit $((n\text{-C}_4\text{H}_9)_4\text{N})^+_2(\text{Co}(\text{S}_2\text{C}_4\text{N}_2)_2^{2-})$ and is triclinic with dimensions: $a = 10.77 \pm 0.01$, $b = 12.35 \pm 0.01$, $c = 9.81 \pm 0.01$ Å., $\alpha = 88.5 \pm 0.1$, $\beta = 114.8 \pm 0.1$, $\gamma = 93.5 \pm 0.1^\circ$, $V = 1182.0 \text{ Å.}^3$. With one formula unit per unit cell, the density calculated from the X-ray data is 1.16 g./cc., in comparison with the value of 1.14 g./cc. found by flotation methods using a mixture of benzene and carbon tetrachloride.

There are no systematic extinctions and the space group could be either P1 (C_1) or $P\bar{1}$ (C_i). Previous evidence indicated that the Co complex was likely to be planar and we favored the space group $P\bar{1}$ (C_i); the

success of our structure determination confirms this choice.

Determination of the Structure.—After about half of the data had been collected, a three-dimensional Patterson function and a three-dimensional Fourier with phases based on Co placed at the origin were calculated from the observed intensities after correction for Lorentz and polarization effects. From these two functions and a knowledge of the molecular geometry, it was possible to determine unambiguously the positions of all the independent atoms (2S, 4C, and 2N) in the anion. With the exception of the Co atom which is in a special position, 1 (a) 0, 0, 0, all the atoms are in general twofold positions, 2 (i) $\pm(x, y, z)$.

At this stage, all the data had been measured, and a second three-dimensional electron density function with phases based on the whole of the anion was calculated from 1627 observed reflections (whose intensity was greater than 3 counts/sec.) after correction for Lorentz and polarization effects. Using this function, together with information regarding bond distances and angles and rejecting any positions which did not lead to peaks in the Patterson function, all 17 independent atoms (1N and 16C) in the tetra-*n*-butylammonium ion were found in general positions.

A least-squares refinement with all 26 of these atoms, each having an isotropic temperature factor of the form $\exp(-B\lambda^{-2} \sin^2 \theta)$, and using the 1627 terms each with unit weight resulted in a conventional unreliability factor $R = \sum ||F_o| - |F_c|| / \sum |F_o|$ of 0.18 after four cycles of refinement. Interatomic distances and angles appeared at this stage to be very reasonable, but there

(15) J. A. Ibers in "International Tables for X-Ray Crystallography," Kynoch Press, Birmingham, England, 1962, Vol. III, p. 202.

(16) D. H. Templeton, *ibid.*, p. 215.

TABLE II
 FINAL COORDINATES AND ISOTROPIC THERMAL PARAMETERS FOR THE HYDROGEN ATOMS

Atom	<i>x</i>	<i>y</i>	<i>z</i>	<i>B</i> , Å. ²	Atom	<i>x</i>	<i>y</i>	<i>z</i>	<i>B</i> , Å. ²
H(1) C(5)	0.13	0.23	-0.08	1.1	H(1) C(13)	0.22	0.39	-0.35	4.6
H(2)	0.24	0.20	-0.15	1.9	H(2)	0.08	0.33	-0.47	5.4
H(1) C(6)	0.35	0.37	-0.12	1.9	H(1) C(14)	0.29	0.20	-0.34	0.5
H(2)	0.25	0.42	-0.06	4.0	H(2)	0.18	0.17	-0.49	4.9
H(1) C(7)	0.32	0.32	0.16	4.9	H(1) C(15)	0.38	0.33	-0.44	3.4
H(2)	0.41	0.26	0.11	0.3	H(2)	0.26	0.31	-0.58	3.3
H(1) C(8)	-0.43	0.42	0.12	4.6	H(1) C(16)	0.47	0.16	0.56	5.2
H(2)	-0.47	0.35	0.21	1.6	H(2)	0.37	0.13	0.43	4.2
H(3)	-0.53	0.48	0.16	3.6	H(3)	0.46	0.24	0.43	9.9
H(1) C(9)	-0.04	0.15	-0.25	3.4	H(1) C(17)	-0.06	0.37	-0.40	3.1
H(2)	0.06	0.12	-0.34	0.9	H(2)	0.06	0.44	-0.28	2.9
H(1) C(10)	-0.18	0.21	-0.48	1.0	H(1) C(18)	0.01	0.36	-0.07	7.7
H(2)	-0.07	0.18	-0.57	7.9	H(2)	-0.12	0.29	-0.20	2.7
H(1) C(11)	-0.25	0.08	0.55	7.0	H(1) C(19)	-0.21	0.44	-0.12	7.0
H(2)	-0.14	0.05	0.45	2.9	H(2)	-0.14	0.52	-0.23	3.7
H(1) C(12)	-0.38	0.11	0.33	3.5	H(1) C(20)	-0.28	0.44	-0.44	4.4
H(2)	-0.27	0.04	0.26	5.3	H(2)	-0.33	0.37	-0.32	5.5
H(3)	-0.33	0.14	0.25	8.8	H(3)	-0.36	0.49	-0.33	22.8

were several large discrepancies in the observed and calculated structure factors, some of the temperature factors were somewhat abnormal, and we felt that the agreement should be much better. However, it was then realized that we had neglected to apply a special position factor to the Co form factor which was thus twice as large as it ought to be. Correction of this blunder and three cycles of refinement speedily reduced *R* to 0.11. Several errors in the data taking and in the card punching were corrected and the Co atom given an anisotropic temperature factor of the form $\exp(-\beta_{11}h^2 - \beta_{22}k^2 - \beta_{33}l^2 - 2\beta_{12}hk - 2\beta_{13}hl - 2\beta_{23}kl)$, with $4\beta_{ij} = a_i^*a_j^*B_{ij}$, a_i^* being the length of the *i*th reciprocal axis. With this notation the anisotropic thermal parameters B_{ij} are in the units (Å.²) which are used for isotropic thermal parameters *B* in the temperature factors of the form $\exp(-B\lambda^{-2} \sin^2 \theta)$. Four cycles of refinement with all 2213 terms included resulted in *R* = 0.142. The terms were weighted so that the 1823 nonzero terms were given unit weight and the 390 zero terms were given 1/4 weight.

An electron density difference function, with all the atoms except hydrogen subtracted out, was calculated using the results of this refinement for the terms with $\sin \theta/\lambda < 0.4$. With the exception of several peaks due to an incorrect description of the thermal motion of the sulfur atoms, the highest peaks in this function all corresponded to reasonable positions for the 36 independent hydrogen atoms, all in twofold general positions. Two cycles of least-squares refinement were run with the two sulfur atoms, as well as the cobalt, having anisotropic thermal parameters, and *R* fell to 0.126.

The number of parameters now exceeded the capacity of our program, so we refined parts of the structure separately with the remainder held fixed. First, the H atoms were included but not refined, while all the other parameters were refined, and *R* fell to 0.103. Then all atoms except H were held fixed and the 36 H atoms, each with an isotropic temperature factor, were refined and resulted in *R* = 0.097. This procedure was repeated twice to reduce *R* to 0.091.

TABLE III

FINAL ANISOTROPIC THERMAL PARAMETERS OF THE COBALT AND THE TWO SULFUR ATOMS^a

Atom	<i>B</i> ₁₁	<i>B</i> ₂₂	<i>B</i> ₃₃	<i>B</i> ₁₂	<i>B</i> ₁₃	<i>B</i> ₂₃ ^b
Co	3.8	5.4	4.1	0.2	1.2	0.0
S(1)	4.3	6.0	5.0	0.1	0.9	-0.6
S(2)	4.2	7.6	5.6	0.3	0.9	-1.8

^a Estimated standard deviations are all approximately 0.1 Å.². ^b In Å.².

We are indebted to a referee who noticed that we had by error used the wrong sign for the dispersion corrections to the atomic scattering factors of cobalt and sulfur. Two additional cycles of least-squares refinement of the heavy atoms after correction of this mistake improved the agreement slightly, with a final *R* of 0.090. No atom moved more than 0.002 Å. The thermal parameters (other than cross terms) of cobalt and sulfur increased 0.3 and 0.1 Å.², respectively, while those of carbon and nitrogen decreased about 0.1 Å.². In the final cycle, no coordinate shifted more than 0.00007 and no thermal parameter by more than 0.01 Å.².

The parameters resulting from the last cycle are listed in Tables I, II, and III. The observed and calculated structure factors before correction of the dispersion error are compared in Table IV.

Standard deviations of the parameters of the heavy atoms were calculated assuming that the discrepancies in the structure factors represent random errors and with neglect of the effect of the H parameters. Limitations of our computer program did not permit us to estimate the standard deviations of the hydrogen atom parameters by the method of least squares. From the resulting bond distances we estimate the standard deviations of the H atom coordinates to be 0.1 to 0.2 Å.

Description of the Structure.—A projection down the *b* axis of the over-all structure is shown in Fig. 1. The planar Co(MNT)₂²⁻ ion has a center of symmetry at the cobalt atom and the two tetra-*n*-butylammonium ions are related to each other by a center of symmetry. The ions are interlaced to form layers in the *bc* plane.

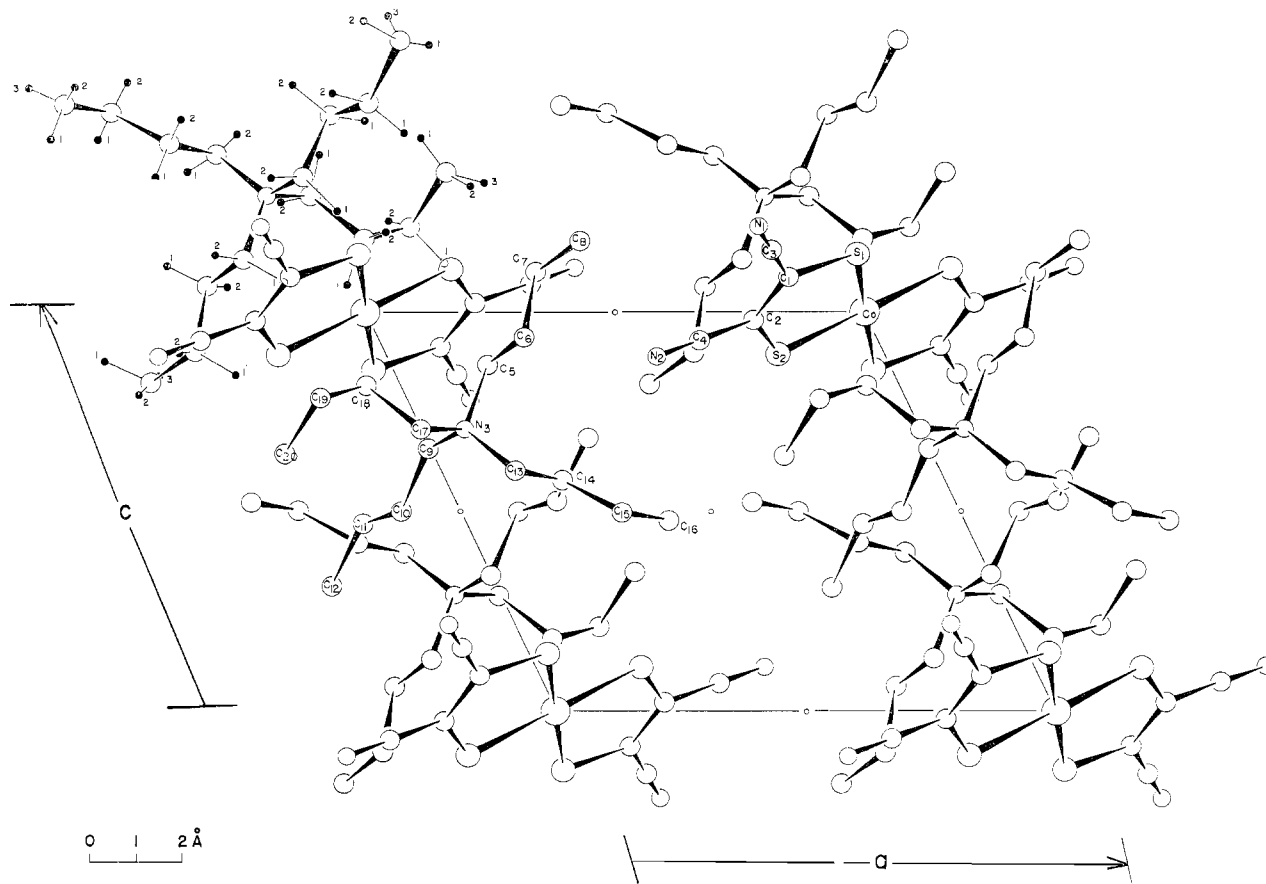


Fig. 1.—Projection of the structure of $((n\text{-C}_4\text{H}_9)_4\text{N})^+_2(\text{Co}(\text{S}_2\text{C}_4\text{N}_2)_2)^{2-}$ down the b axis. The small black circles represent examples of each type of hydrogen atom.

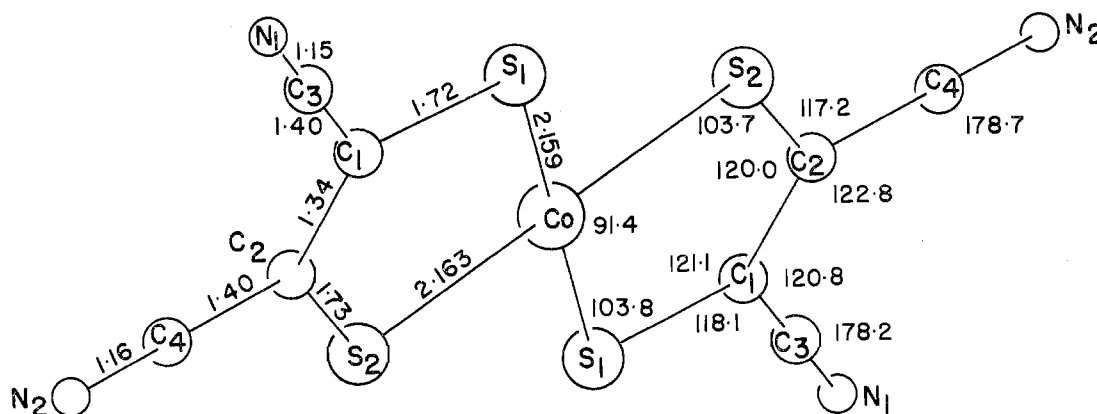


Fig. 2.—Dimensions of the $\text{Co}(\text{MNT})_2^{2-}$ ion. See Table V for the estimated standard deviations.

The $\text{Co}(\text{MNT})_2^{2-}$ Ion.—The dimensions of the anion are illustrated in Fig. 2 and compared in Table V with those found in the complex $((\text{CH}_3)_4\text{N})^+_2(\text{Ni}(\text{MNT})_2^{2-})$ by Eisenberg, *et al.*¹³ The agreement is very close, with the Co-S and Ni-S distances in the two complexes equal to well within the experimental accuracy.

The anion is very closely planar and deviations from planarity are small enough to be chemically insignificant. The three atoms Co, C(1), and C(2) define a plane of equation (triclinic coordinate system)

$$-4.996x - 6.634y + 7.975z = 0$$

The distances of the various atoms from this plane are

listed in Table VI. S(1) is somewhat out of the plane of Co, S(2), C(1), and C(2) and makes the five-membered ring slightly puckered, and the cyanide groups are bent a small amount out of the plane. None of these deviations, however, is as much as 0.1 Å.

The chemically equivalent, but crystallographically nonequivalent bonds in the anion, *viz.*, S(1)-C(1) and S(2)-C(2), C(1)-C(3) and C(2)-C(4), C(3)-N(1) and C(4)-N(2), differ by no more than two standard deviations. Although the ion is required to have only a center of symmetry (C_i), the geometry is such that it does not differ significantly from mmm (D_{2h}).

The arrangement of the sulfur atoms around the co-

TABLE V

A COMPARISON OF INTERATOMIC DISTANCES (Å.) AND ANGLES IN THE TWO IONS $\text{Co}(\text{MNT})_2^{2-}$ AND $\text{Ni}(\text{MNT})_2^{2-}$

Atoms	$\text{Co}(\text{MNT})_2^{2-}$		$\text{Ni}(\text{MNT})_2^{2-}$	
	Distance	E.s.d.	Distance	E.s.d.
Ni/Co-S(1)	2.159	0.003	2.16	0.008
Ni/Co-S(2)	2.163	0.003	2.16	0.008
S(1)-C(1)	1.715	0.007	1.75	<i>a</i>
S(2)-C(2)	1.731	0.007	1.75	<i>a</i>
C(1)-C(3)	1.40	0.010	1.44	<i>a</i>
C(2)-C(4)	1.40	0.010	1.42	<i>a</i>
C(1)-C(2)	1.34	0.010	1.30	0.02
C(3)-N(1)	1.15	0.010	1.13	0.02
C(4)-N(2)	1.16	0.010	1.13	0.02

Atoms	Angles	E.s.d.	Angles ^b
S(1)-Co/Ni-S(2)	91.4°	0.1°	91.7°
Co/Ni-S(1)-C(1)	103.8°	0.3°	103°
S(1)-C(1)-C(2)	121.1°	0.5°	123°
C(1)-C(2)-S(2)	120.0°	0.6°	120°
Co/Ni-S(2)-C(2)	103.7°	0.3°	103°
S(2)-C(2)-C(4)	117.2°	0.6°	...
S(1)-C(1)-C(3)	118.1°	0.5°	...
C(2)-C(1)-C(3)	120.8°	0.7°	120°
C(1)-C(3)-N(1)	178.2°	0.8°	179°
C(1)-C(2)-C(4)	122.8°	0.7°	124°
C(2)-C(4)-N(2)	178.7°	1.0°	177°

^a Intermediate between ± 0.008 and ± 0.02 Å. ^b E.s.d. all approximately $\pm 1.0^\circ$.

TABLE VI

DISTANCES OF THE VARIOUS ATOMS (Å.) IN THE $\text{Co}(\text{MNT})_2^{2-}$ ION FROM THE PLANE THROUGH ATOMS Co, C(1), AND C(2)

Atom	Distance from plane	E.s.d. of the atomic position	E.s.d. of the plane at a particular atom
Co	0	0	0
S(1)	+0.049	0.002	0.004
S(2)	+0.004	0.002	0.004
C(1)	0	0.007	0.007
C(2)	0	0.007	0.007
C(3)	-0.035	0.008	0.010
C(4)	-0.031	0.009	0.010
N(1)	-0.062	0.007	0.012
N(2)	-0.058	0.009	0.012

atom (C-N-C) average 109.5° with a spread of $\pm 3^\circ$ and are very close to the tetrahedral angle.

In the butyl chains, the mean C-C distance is 1.515 Å. Applying a thermal correction, assuming that each carbon atom "rides" on its neighbor nearer to the central nitrogen atom, increases this distance by 0.010 Å. This assumption is reasonable as the temperature factors of the carbon atoms increase progressively along each carbon chain (see Table I). The mean value of the twelve N-C-C and C-C-C angles is 113.1° , somewhat greater than the tetrahedral angle. The four N-C-C angles all have higher than average values.

Most of the hydrogen atoms are suitably placed with respect to bond distances and angles. The C-H bond lengths range from 0.8 to 1.1 Å., with the average of the 36 values being 0.98 Å. Several of the bond angles deviate substantially from tetrahedral, but the majority

TABLE VII

DISTANCES AND ANGLES INVOLVING THE CARBON AND NITROGEN ATOMS IN THE TETRA-*n*-BUTYLAMMONIUM GROUP

Atoms	Distance, ^a		Atoms	Angles, ^b
	Å.			
N(3)-C(5)	1.53		C(5)-N(3)-C(9)	106.3
N(3)-C(9)	1.52		C(5)-N(3)-C(13)	111.7
N(3)-C(13)	1.51		C(5)-N(3)-C(17)	109.4
N(3)-C(17)	1.53		C(9)-N(3)-C(13)	112.9
			C(9)-N(3)-C(17)	110.3
			C(13)-N(3)-C(17)	106.3
Mean C-N dist.	1.52		Mean C-N-C angle	109.5

Atoms	Dis-	Dis-	Atoms	Angles, E.s.d.,	
	tance, ^c	tance, ^d		deg.	deg.
Å.		Å.			
C(5)-C(6)	1.49	1.49	N(3)-C(5)-C(6)	115.6	0.5
C(6)-C(7)	1.53	1.54	C(5)-C(6)-C(7)	108.9	0.6
C(7)-C(8)	1.53	1.56	C(6)-C(7)-C(8)	112.5	0.7
C(9)-C(10)	1.53	1.54	N(3)-C(9)-C(10)	114.0	0.6
C(10)-C(11)	1.51	1.52	C(9)-C(10)-C(11)	109.5	0.7
C(11)-C(12)	1.54	1.54	C(10)-C(11)-C(12)	114.3	0.8
C(13)-C(14)	1.48	1.49	N(3)-C(13)-C(14)	115.7	0.6
C(14)-C(15)	1.52	1.53	C(13)-C(14)-C(15)	112.4	0.6
C(15)-C(16)	1.48	1.49	C(14)-C(15)-C(16)	114.2	0.7
C(17)-C(18)	1.52	1.53	N(3)-C(17)-C(18)	115.3	0.6
C(18)-C(19)	1.54	1.54	C(17)-C(18)-C(19)	111.4	0.7
C(19)-C(20)	1.51	1.53	C(18)-C(19)-C(20)	113.3	0.8
Mean C-C	1.515	1.525	Mean angle	113.1	
dist.					

^a E.s.d. all ± 0.01 Å. ^b E.s.d. all $\pm 0.5^\circ$. ^c E.s.d. all ± 0.01 Å. ^d Corrected for thermal vibration assuming that atom C(*n* + 1) "rides" on atom C(*n*).

TABLE VIII

THE DIHEDRAL ANGLES FOR THE FOUR BUTYL CHAINS IN THE TETRA-*n*-BUTYLAMMONIUM ION

Plane (1) defined by atoms			Plane (2) defined by atoms			Dihedral angle, deg.
C(5)	C(6)	C(7)	C(6)	C(7)	C(8)	
C(9)	C(10)	C(11)	C(10)	C(11)	C(12)	1.6
C(13)	C(14)	C(15)	C(14)	C(15)	C(16)	1.3
C(17)	C(18)	C(19)	C(18)	C(19)	C(20)	71.6

are very close. In particular, except for the hydrogen atoms on the terminal carbon atoms of the four *n*-butyl chains, all the hydrogen bond angles are within 15° of the expected value; consistent with the estimated standard deviations of the hydrogen atom positions. It is noteworthy that all the hydrogen atoms except one have calculated temperature factors less than those of the carbon atoms with which they are associated.

X-Ray powder diffraction photographs of material from the sample studied in this investigation and of material used in the e.s.r. experiments (kindly supplied by Dr. N. Edelstein) are consistent but do not agree with the *d*-spacings reported by Davison, *et al.*⁵

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