closer to the truth because D is comparatively small.

In the simulations w_0 has been treated as empirical parameter. It is expected to be independent of temperature and field strength as long as the latter is high enough to justify the Boyle and Gabriel model despite of a finite zero-field splitting. Due to the zero-field splitting the level spacings are of course not really equidistant, but this may be compensated in part by an inhomogeneous broadening of the levels. For the spectra of the trinuclear complex recorded under 6.21 and 2.47 T it was indeed possible to get satisfactory simulations for the whole temperature range from 4.2 to 50 K with one single value for w_0. The zero-field splitting, rhombicity, and hyperfine parameters were taken as determined from the low-temperature spectra except that, for the sake of simplicity, the Euler angles α and β of the efg tensor have been disregarded. The solid lines in Figure 8 show the simulations obtained for $S = \frac{3}{2}$ and the parameters of Table IV with $w_0 = 2.6 \times 10^8$ s^{-1}. The estimated uncertainty of this value is about 20%.

The relaxation rate in the tetranuclear complex is obviously faster. The solid lines through the spectra in Figure 9 are simulations obtained for $S = \frac{3}{2}$ and the parameters of Table IV with $w_0 = 2.2 \times 10^9$ s^{-1}. Direct comparison with the trinuclear complex is possible, since the applied field of 6.21 T is so strong that the spin $\frac{1}{2}$ of the monomeric Cu(Mesalen) is coupled in a kind of Paschen–Back effect much stronger to the external field than to the spin $\frac{3}{2}$ of the trimer. The monomer may, however, mediate the cross-relaxation between two trimers, thus enhancing the relaxation rate by about 1 order of magnitude. This finding is not unexpected because it is well-known, e.g. from the work of Bhargava et al.,[24] that the presence of other paramagnetic ions shortens the spin–spin relaxation times substantially. Such an explanation is also consistent with the X-ray structure of the crystalline tetranuclear entity described in the preceding section.

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Multinuclear NMR Study of the Interaction of the Shift Reagent Lanthanide(III) Bis(triphosphate) with Alkali-Metal Ions in Aqueous Solution and in the Solid State[†]

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The paramagnetic ion-induced relaxation rate enhancements of ^{6}Li in adducts of Li$^+$ and Ln(PPP)$_2$$^{7-}$ complexes (Ln = Dy, Tm) in aqueous solution show that up to seven monovalent counterions can coordinate in the second coordination sphere of the Ln(III) ion to the outer oxygens of the triphosphate ligands. However, the pseudocontact ^{7}Li NMR shift data suggest that in the second coordination sphere some preference of the counterions for the axial region opposite the water ligand may exist. The estimated Ln^{3+}–Li$^+$ distances range from 5.1 to 5.9 Å for [Li$^+$]/[Ln(PPP)$_2$$^{7-}$] ratios of 0.2–7. This is supported by a two-dimensional nutation spectrum of a polycrystalline sample of Na$_7$La(PPP)$_2$, which indicates coordination of all Na$^+$ ions to phosphate oxygens.

Introduction

Several aqueous shift reagents (SRs) for metal cation NMR spectroscopy have been reported.[1-6] They have been shown to be useful in membrane transport biochemistry, for example in studies of alkali-metal ion transport across vesicles[7,8] and red blood cell membranes.[1,9-12] However, only two of them, Dy(TTHA)$^{3-}$ and Tm(DOTP)$^{5-}$, have found useful practical application for perfused heart studies[13,14] and in vivo rat brain ^{23}Na NMR spectroscopy.[15,16] Despite considerable discussion, there seems still to be some controversy regarding the mode of interaction of one of the SRs most widely used for NMR studies in cell systems, namely Dy(PPP)$_2$$^{7-}$, with alkali-metal ions in aqueous solution.[17-23] Moreover, it has been shown recently that Dy(PPP)$_2$$^{7-}$ can change ion distribution and transport across red blood cell membranes as well as membrane potential.[24,25] These effects may be related to the mode of binding of alkali metal cations to Dy(PPP)$_2$$^{7-}$. In this work we further analyze this problem by studying the ^{6}Li NMR spin–lattice relaxation times of Dy(PPP)$_2$$^{7-}$ and Tm(PPP)$_2$$^{7-}$

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† Abbreviations used in this paper are as follows: PPP^{5-}, tripolyphosphate; SRs, shift reagents; TTHA^{6-}, tetraethylenetetraminehexaacetate; DOTP^{8-}, tetraazacyclododecane-*N,N',N'',N'''*-tetrakis(methanephosphonate); MAS, magic angle spinning; NMR, nuclear magnetic resonance; LnIRE, lanthanide-induced relaxation rate enhancement.

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in sodium-free solutions containing increasing amounts of LiCl. We also studied by solid-state MAS ^{23}Na NMR spectroscopy the interaction of Na$^+$ ions with La(PPP)$_2$$^{7-}$.

Methods and Materials

DyCl$_3$, TmCl$_3$, D$_2$O, and sodium triphosphate were obtained from Aldrich Chemical Co. ^{6}LiCl (95.7% ^{6}Li) was from Oak Ridge National Laboratories, TN. All chemicals were used as obtained, except sodium triphosphate, which was purified by repetitive precipitation from 33% aqueous ethanol.[26] The purity of the resulting product was established by ^{31}P NMR spectroscopy.

Tetramethylammonium triphosphate was prepared from the sodium form by ion-exchange chromatography and recrystallized as described earlier.[21] The residual sodium content in crystallized tetramethyl-ammonium triphosphate was analyzed by atomic absorption and ^{23}Na NMR spectroscopy and was found to be 2%. ^{31}P NMR analysis of the above solution did not show any (phosphorous containing) impurities.

The shift reagents Dy(PPP)$_2$$^{7-}$ and Tm(PPP)$_2$$^{7-}$ were prepared as reported earlier.[1] The ratio of Dy^{3+} or Tm^{3+} to PPP^{5-} was maintained at 1:2. In all samples the concentration of the shift agent was held constant at 5 mM while ^{6}LiCl was varied between 1 and 100 mM. The pH of all samples was maintained at 7.5 and adjusted by using tetra-methylammonium hydroxide. No buffers were used. pH was checked before and after NMR measurements and was observed to be the same. pH values reported are not corrected for deuterium isotope effects. All the NMR samples had 15% (v/v) D$_2$O.

The Na$_7$La(PPP)$_2$ salts were prepared by mixing a solution of 1 mmol LaCl$_3$·6H$_2$O in 5 mL of water with a solution of 2 mmol of Na$_5$(PPP) in 5 mL of water. The solution was filtered and then poured into 100 mL of ethanol. The precipitate obtained was filtered off, washed with cold aqueous ethanol, and dried in the air.
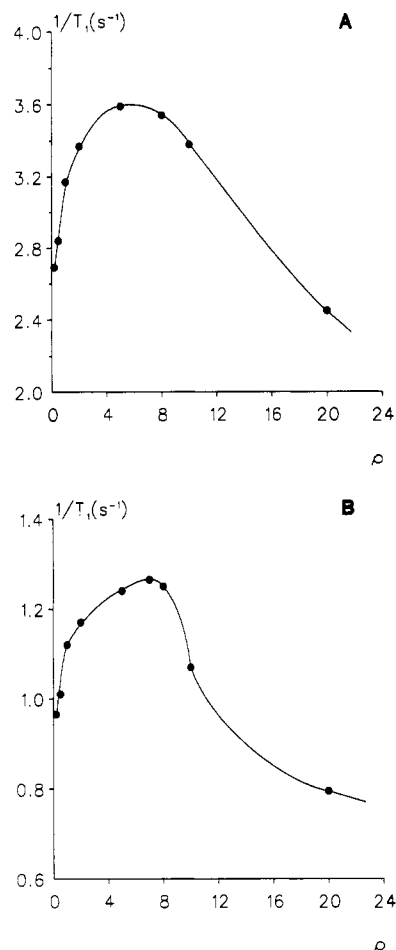
^{6}Li, ^{23}Na, and ^{31}P solution NMR measurements were made at 44.2, 79.4, and 121.4 MHz, respectively, on a Varian VXR-300 NMR spectrometer equipped with a 10-mm multinuclear probe and a variable-temperature unit. ^{6}Li T$_1$ measurements were made by using the inversion-recovery pulse sequence. The errors in the reported T$_1$ measurements were less than 5%. The spectral width was 8000 Hz, and the probe temperature was 25 °C. For ^{6}LiCl concentrations lower than 20 mM, 256 scans were used, while for higher concentrations 64 scans were used in T$_1$ measurements.

The solid-state ^{23}Na NMR spectra were measured at 105.8 MHz on a Varian VXR-400 S spectrometer equipped with a Doty solid-state probe. Magic angle spinning with a rate of about 6 kHz was employed. The 2D-nutation spectra were recorded with 128 t$_1$ increments of 2 μs. The rf field strength was 60 kHz, which was determined by performing a nutation experiment with an aqueous solution of NaCl. Note that the solid-state experiments were conducted on the La compound and not Ln.

The speciations of species needed for the simulations of the curves of R$_1$ as function of ρ were computed with the use of a Lotus 1-2-3 program.[27]

Results and Discussion

^{6}Li Spin–Lattice Relaxation Studies in Solution.

For most alkali-metal ions the quadrupolar mechanism dominates the relaxation.[28] As a result the relaxation rates are usually so high that pertubations by paramagnetic ions are hard to measure accurately.[18] The ^{6}Li nucleus, however, has a very small quadrupole moment, and consequently its relaxation is extremely slow.[28,29] Therefore, ^{6}Li is a very suitable probe for studying the location

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Figure 1. Spin–lattice relaxation rates of ^{6}Li in the presence of the SRs Ln(PPP)$_2$$^{7-}$ (A, Ln = Dy; B, Ln = Tm), at 25 °C and pH 7.5. The SRs were used in the tetramethylammonium form. The Ln(PPP)$_2$ concentration was held constant at 5 mM, whereas the Li$^+$ concentration was varied between 1 and 100 mM by addition of ^{6}LiCl.

of counterions of negatively charged Ln(III) complexes via the lanthanide-induced relaxation rate enhancements (LnIRE).[18,30]

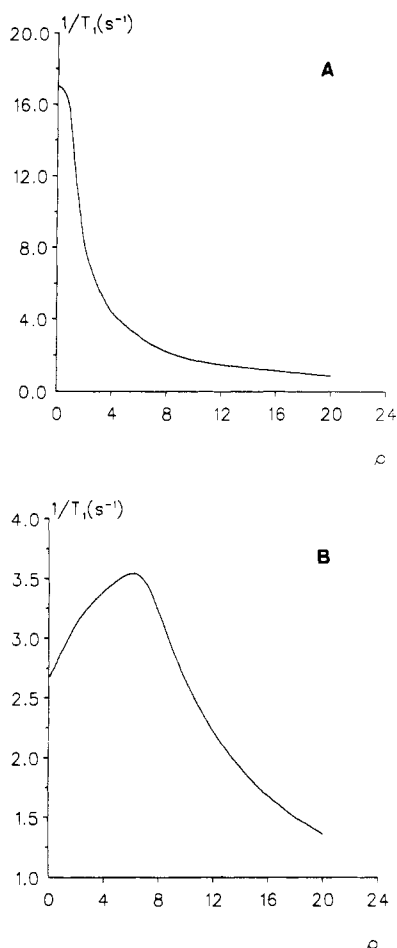Figure 1 shows the variation of the measured ^{6}Li spin–lattice relaxation rate of Na$^+$-free 5 mM solutions of Dy(PPP)$_2$$^{7-}$ or Tm(PPP)$_2$$^{7-}$ SRs at 25 °C and pH 7.5 when the concentration of LiCl present varied from 1.0 to 100 mM (the value of ρ = [Li$^+$]/[SR] changed from 0.2 to 20). The exchange rate between free and the various forms of SR-bound ^{6}Li was fast on the ^{6}Li NMR time scale; only a single ^{6}Li resonance was observed, the chemical shift of which was dependent on ρ.[21] The values of 1/T$_1$ (=R$_1$) increase by about 30% when ρ increases from 0.2 to 7. After a maximum at ρ = 7, R$_1$ decreases steadily, as a result of exchange of the fast relaxing bound ^{6}Li$^+$ nuclei with the slow relaxing ^{6}Li$^+$ nuclei in the bulk. The difference in magnitude of the relaxation rates for the Dy^{3+} and the Tm^{3+} complexes can be explained by the differences in electron spin relaxation times and effective magnetic moments of these lanthanides (see below).

Assuming that isotropic relaxation occurs, that the contact contribution to the LnIRE is negligible,[31] the value of R$_1$ for ^{6}Li$^+$ bound to a Ln(PPP)$^{7-}$ complex can be related to the distance between the Ln(III) ion and the ^{6}Li ion by using the reduced Solomon–Bloembergen equation:[32,33]

$$R_1 = \frac{4}{3}(\mu_0/4\pi)^2(\mu^2\gamma^2\beta^2 T_{1e}/r^6) \tag{1}$$

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Figure 2. Simulated curves of ^6Li spin–lattice relaxation rates in the presence of a $Ln(PPP)_2^{7-}$ complex. (A) shows the model in which only one Li^+ binds strongly. Each triphosphate ligand is assumed to bind both Ln^{3+} and Li^+ ions.[17,19,22,23] Parameters used are as follows: concentration SR = 5 mM, stability constant adduct K_1 = 10 000 M^{-1}, $R_{1\ bound}$ = 17.4 s^{-1}, $R_{1\ free}$ = 0.05 s^{-1}. (B) shows the model in which seven Li^+ ions are bound in the second coordination sphere of Ln^{3+} in the $Ln(PPP)_2^{7-}$ complex. Parameters used are as follows: concentration $Ln(PPP)_2^{7-}$ = 5 mM, stepwise stability constants K_1 = 4000, K_2 = 2000, K_3 = 1800, K_4 = 1600, K_5 = 1400, K_6 = 1200, K_7 = 1000 M^{-1}. The corresponding Li^+ complexes are assumed to have the following R_1 values: 2.8, 3.0, 3.2, 3.6, 3.7, and 3.8 s^{-1}. For free Li^+ an R_1 value of 0.05 was assumed.

Here $\mu_0/4\pi$ is the magnetic permeability under vacuum, μ is the effective magnetic moment, γ is the magnetogyric ratio, β is the Bohr magneton, and T_{1e} is the electron spin relaxation time. This latter parameter has been found to be rather independent of the ligation of the Ln(III) cation.[32-34] If the adduct would have the previously suggested binuclear structure, in which each triphosphate is bound to both Ln^{3+} and M^+,[17,19,22,23] the distance $Ln^{3+}–Li^+$ would be about 4 Å.[23] Then for Ln = Tm, an R_1 value of 11.3 s^{-1} can be calculated with eq 1 and a T_{1e} value of 4.6 × 10^{-13} s.[32] Analogously it can be calculated that in the corresponding $LiDy(PPP)_2^{6-}$ system, the R_1 value of Li^+ is 17.3 s^{-1}. These calculated relaxation rates are considerably larger than the experimental ones. The curves of R_1 as a function of ρ were simulated with the use of the above-mentioned bound values of R_1, a value of 0.05 s^{-1} for free Li^+,[29] and eq 2, where f_{bound} and

$$R_1 = f_{bound}R_{1\ bound} + f_{free}R_{1\ free} \qquad (2)$$

f_{free} are the mole fractions of bound and free Li^+, respectively, which were calculated by using a stability constant of the complex of 10 000 M^{-1}. The calculated curves have shapes that differ dramatically from the experimental ones (see e.g. Figure 2A). The simulated curves show a steep decrease of R_1 at $\rho > 1$, whereas

the experimental ones have a maximum at ρ about 7. On the basis of these shapes and the lower R_1 values, the binuclear structure of the $Ln(PPP)_2^{7-}–Li^+$ adduct can be rejected. The experimental data strongly suggest that $Ln(PPP)_2^{7-}$ complex binds rather strongly up to seven Li^+ ions in the second coordination sphere of the Ln^{3+} ion, which is in agreement with a previous proposal.[18] The increase of R_1 between $\rho \approx 0.2$ and 7 indicates that successive binding of the seven Li^+ ions to $Ln(PPP)_2^{7-}$ causes the $SR–Li_n$ structure to become more compact, probably as a result of the electrostatic repulsion between the ligands upon charge neutralization by the Li^+ ions. Simulation of the dependence of R_1 of ρ for the $Dy(PPP)_2^{7-}–Li^+_n$ adduct with a procedure similar to that described above, by using bound values for R_1 of 2.8–3.8 for $n = 1–7$ and a somewhat arbitrary chosen set of stepwise stability constants for the respective adducts (K_n = 4000–1000 M^{-1}), gives a curve that is in good agreement with the experimental one (see Figure 2B). A fitting of the experimental data with stability constants and relaxation rates was not attempted, since the model with seven bound alkali-metal ions requires 15 variables (7 stepwise stability constants and 8 R_1 values). The good agreement between the experimental and simulated curve indicates, however, that the estimates of the bound R_1 values were correct. Then with eq 1, it can be calculated that for the $Dy(PPP)_2^{7-}–Li^+_n$ system the $Dy^{3+}–Li^+$ distances vary between 5.4 and 5.1 Å, going from $n = 1$ to 7. Analogously it can be calculated that in the $Tm(PPP)_2^{7-}$ system the $Tm^{3+}–Li^+$ distance is between 5.9 and 5.7 Å. These distances show that the seven monovalent cations are coordinated to the outer oxygens of the triphosphate ligands.

It should be noted that the previously reported Dy(III)-induced ^7Li shifts of the $Li^+–Dy(PPP)_2^{7-}$ system as a function of ρ also showed a break at ρ = 7.[21] The induced shifts, however, decreased steeply between ρ = 0.2 and 7. If it is assumed that the induced ^7Li shifts are of a pseudocontact origin and that internal reorientations in the complexes result in effective axial symmetry, the induced shifts are determined by
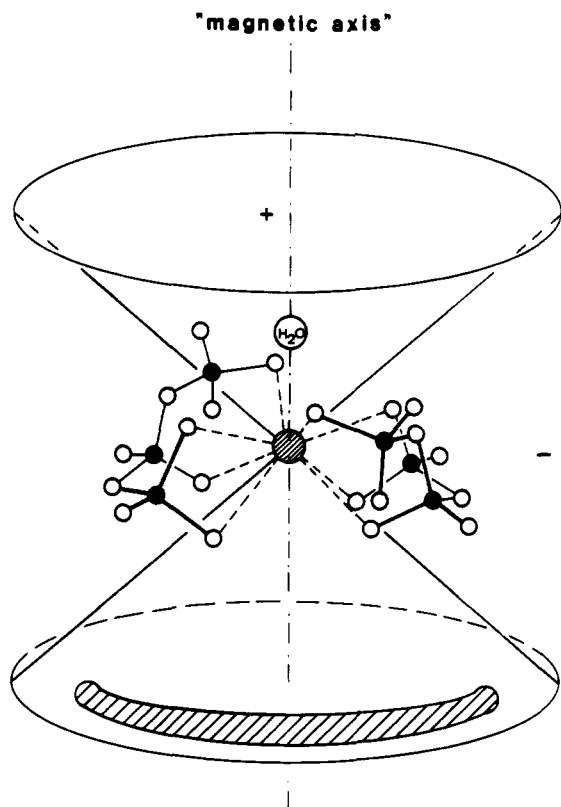
$$\Delta = k(3\cos^2\theta - 1)/r^3 \qquad (3)$$

where Δ is the induced shift and θ is the angle between the vector r and the effective magnetic axis. The variation of the chemical shifts as function of ρ indicates that at low ρ values binding sites with a high Δ value are occupied by Li^+ preferentially. As outlined previously,[18,20,21] the signs of the pseudocontact portion of the Ln(III)-induced ^{31}P and ^{17}O shifts indicate that the triphosphate ligands are in the equatorial region ($55 < \theta < 125°$), whereas the water ligand is in the axial region ($0 < \theta < 55°$, $125 < \theta < 180°$) (see Figure 3). Similarly, the sign of the induced shift for the counterion nucleus indicates that the counterion also is located in the axial region. The coordination geometry of nine-coordinated Ln^{3+} complexes can often best be described by a tricapped trigonal prism. If it is assumed that this geometry also applies here, and if the binding fashion of the triphosphate ligands established previously[18] is taken into account, a model can be constructed in which the water molecule and oxygens of the central phosphates of the triphosphate ligands occupy the capping positions and oxygens of the terminal phosphates the remaining positions (see Figure 3). Then the effective magnetic axis is most likely in the direction of the $Ln^{3+}–water$ bond, and the ^{31}P nuclei are indeed in the equatorial region. The crowding in the axial region opposite to the water ligand is relatively small, which may explain some preference of the counterions for this region.

Further addition of Li^+ leads to exchange of the Li^+ ions on this preferred site with those at other locations. Since each site has a different θ value, this will result in a decrease of the ^7Li shifts due to averaging of the 3 $\cos^2\theta - 1$ term in eq 3. In this way only one of the bound Li^+ ions contributes mostly toward the observed ^7Li shifts. The ^6Li relaxation measurements, however, unambiguously show that about seven monovalent counterions can be bound in the second coordination sphere of Ln(III) in the Ln-$(PPP)_2^{7-}$ complexes.

Solid-State ^{23}Na NMR Studies of $Na_7La(PPP)_2$. In the solid-state MAS ^{23}Na NMR spectrum of $Na_7La(PPP)_2$ only a single broad line ($\Delta\nu_{1/2}$ = 1.3 kHz) was observed. The NMR spectra
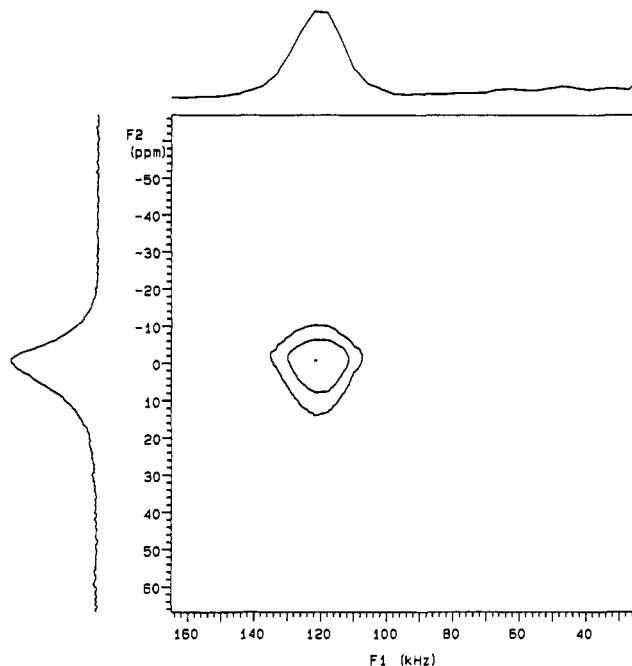
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"magnetic axis"

Figure 3. Schematic representation of the structure of the Ln-(PPP)$_2$$^{7-}$–alkali-metal ion adduct. A tricapped trigonal-prismatic coordination of the Ln^{3+} ion is assumed. ● = P, and O = O.

Figure 4. 2D nutation spectrum of Na$_7$La(PPP)$_2$ (contour plot), together with its projections on the F_1 and F_2 axis. The 2D spectrum contains information about chemical shift and quadrupole interactions along the F_2 axis and information exclusively about the quadrupole interactions along the F_1 axis. The rf field strength used was 60 kHz.

of quadrupolar spin nuclei such as ^{23}Na ($I = {}^3/_2$) contain information on both the chemical shift and the quadrupole interaction, which depends on the local symmetry around the ^{23}Na nucleus, giving in principle direct information on the nature of the Na$^+$ sites. The two-dimensional solid-state nutation NMR spectrum allows one to separate the quadrupolar interaction from the chemical shift interaction.[35-37] These spectra have been calculated theoretically for various half-integer quadrupolar spins, including $I = {}^3/_2$.[37]

Figure 4 shows the two-dimensional MAS ^{23}Na nutation spectrum obtained for a polycrystalline sample of Na$_7$La(PPP)$_2$. No separation is observed in the F_1 direction; only a single peak at [2ω_{rf},0] is present. A comparison with simulated spectra[37] shows that the absence of peaks at ω_{rf} and between ω_{rf} and 2ω_{rf} indicates that all ^{23}Na nuclei have a nuclear quadrupolar coupling constant (QCC) of at least 1.4 MHz. This implies that in the solid state the seven ^{23}Na counterions are all in environments with high electric field gradients. These data, therefore, suggest that all

Na$^+$ ions are coordinated to phosphates; for Na$^+$ ions with an octahedral coordination with water a lower QCC should be expected. It has, for instance, been reported that in a single crystal of Na$_4$P$_2$O$_7$·10H$_2$O the Na$^+$ ions that are fully coordinated by water have a QCC of 0.467 MHz, whereas those coordinated to both waters and pyrophosphate have a QCC of 1.907 MHz.[38]

The proposed coordination of seven Na$^+$ ions to the outer oxygens of the triphosphate ligand is supported by the crystal structure of Na$_5$(PPP)·14H$_2$O, in which four of the five different Na$^+$ ions are coordinated to triphosphate oxygens.[39]

In conclusion, in both liquid and solid phases up to seven alkali-metal cations bind to the second coordination sphere of Ln-(PPP)$_2$$^{7-}$ complexes via the outer oxygens of the triphosphate ligands (Figure 3) rather than having one preferential binding site in the proximity of the lanthanide ion.[17,19,22,23] However, in the second coordination sphere some preference of the counterions for the axial region opposite the water ligand may exist.

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