

Contribution from the Department of Chemistry and Materials Science Center, Cornell University, Ithaca, New York 14853-1301

## Bonding in Halocuprates

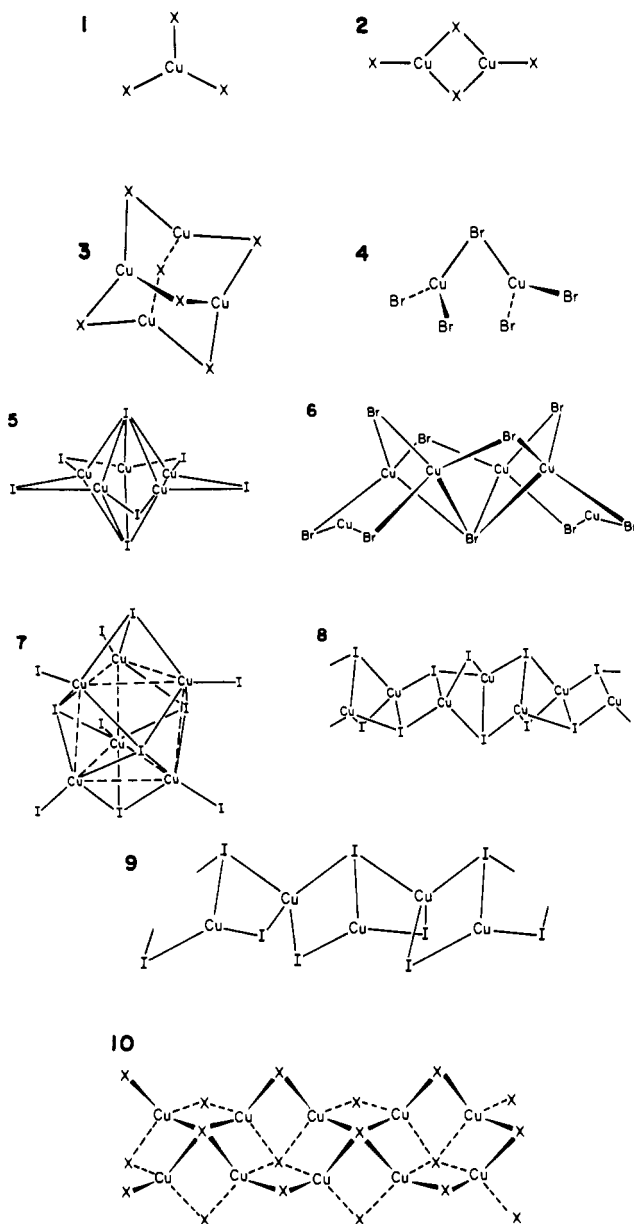
Lalitha Subramanian and Roald Hoffmann\*

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The unusual square pyramidal geometry around halogen atoms bridging four Cu atoms in halocuprates is studied within the extended Hückel framework. Calculations are performed on the discrete molecular anion  $[\text{Cu}_6\text{Br}_9]^{3-}$ , as well as on the one-dimensional polymeric  $[\text{Cu}_2\text{Br}_3]^-$  structure. We examine the bonding between the bridging Br and Cu atoms in planar and nonplanar model systems. Br should be capable of bridging as many as six Cu atoms. In the higher coordinations pyramidal geometry is favored; the bonding is perforce delocalized, an extension of three-center bonding. A few hypothetical structures of halocuprates with a tetrahedral and a square pyramidal  $\mu_4\text{-X}$  are predicted.

### Halocuprate Geometries

There is wide variety in the structures of copper(I) halocuprates,  $\text{Cu}_x\text{X}_y^{n-}$ . These molecules may form discrete geometries of varying nuclearity or polymeric extended systems. The halogens range from terminally bonded to bridging no less than five metals in the discrete structures, while in polymeric systems they bridge two, three, or four metal atoms. Structures 1-10 show some of the geometries observed.



Molecules 1-4 are discrete anions with halogen atoms in terminal and doubly bridging position, while in 5 the halides bridge

two and five coppers. 6 is a  $\text{Cu}_6\text{Br}_9^{3-}$  system where Br bridges two and four Cu atoms; in 7 there is an unusual capped prismatic structure. 8 is a polymeric iodocuprate where I bridges two and three Cu atoms, while the geometry around Cu is approximately tetrahedral. 9 is another polymer with Cu exhibiting distorted trigonal and tetrahedral geometries. In 10, we see an extended halocuprate where the halogens bridge two and four metal atoms with the geometry around the metal atom being tetrahedral. The  $\mu_4\text{-X}$  is at the apex of a square pyramid, as it is in the discrete molecule 6. Table I lists some structural parameters of a selection of halocuprates known.

This fantastic structural richness is exhibited almost without exception in the context of salts consisting of molecular or polymeric  $\text{Cu}_x\text{X}_y^{n-}$  anions and various, often organic counteranions. A series of systematic studies on halocuprates indicates that the anionic configuration, i.e., the coordination number of Cu(I) and the tendency of the anion to form extended structures, depends on the degree of dilution imposed on the ligands by the cation.<sup>1</sup> Bulky cations which have low well-screened charges favor the formation of discrete anions in which Cu(I) has a low coordination number. Small cations, in contrast, favor the formation of extended structures, where Cu(I) shows a higher coordination number of 4. A similar trend is observed for the halogentates as well.<sup>2-6</sup>

With an increase in the bulkiness of the metal or halogen used, the tendency of the anion toward polymerization increases, and the metal atom prefers four-coordination. Early work by Romming and Waerstad<sup>7</sup> showed that the intermediate in the Sandmeyer reaction consists of polymeric anion chains. Another system reported by Boehm et al.<sup>8</sup> is  $[(\text{CH}_3)_2\text{N}_2\text{CHN}_2(\text{CH}_3)_2][\text{Cu}_2\text{Br}_3]$

- (1) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, A 1989, 43, 39.
- (2) (a) Jagner, S.; Olson, S.; Stomberg, R. *Acta Chem. Scand.*, A 1986, 40, 230. (b) Helgesson, G.; Jagner, S. *J. Chem. Soc., Dalton Trans.* 1988, 2117. (c) Helgesson, G.; Jagner, S. *J. Chem. Soc., Dalton Trans.* 1990, 2413. (d) Andersson, S.; Helgesson, G.; Jagner, S.; Olson S. *Acta Chem. Scand.*, A 1989, 43, 946. (e) Jagner, S.; Helgesson, G. *Adv. Inorg. Chem.*, in press.
- (3) Gilmore, C. J.; Tucker, P. A.; Woodward, P. *J. Chem. Soc. A* 1971, 1337.
- (4) Meyer, H. J. *Acta Crystallogr.* 1963, 16, 788.
- (5) Peters, K.; von Schering, H. G.; Ott, W.; Seidenspinner, H.-M. *Acta Crystallogr., C* 1984, 40, 789.
- (6) Brink, C.; Stenfort Kroese, H. A. *Acta Crystallogr.* 1952, 5, 433.
- (7) Romming, Chr.; Waerstad, K. *Chem. Commun.* 1965, 299.
- (8) Boehm, J. R.; Balch, A. L.; Bizot, K. F.; Enemark, J. H. *J. Am. Chem. Soc.* 1975, 97, 501.
- (9) Asplund, M.; Nilsson, M.; Jagner, S. *Acta Chem. Scand.*, A 1983, 37, 57.
- (10) Asplund, M.; Nilsson, M.; Jagner, S. *Acta Chem. Scand.*, A 1984, 38, 57.
- (11) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, A 1986, 40, 52.
- (12) Andersson, S.; Jagner, S.; *Acta Chem. Scand.*, A 1985, 39, 799.
- (13) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, A 1985, 39, 515.
- (14) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, A 1985, 39, 577.
- (15) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, A 1985, 39, 297.
- (16) Asplund, M.; Jagner, S. *Acta Chem. Scand.*, A 1984, 38, 135.
- (17) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, A 1985, 39, 423.
- (18) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, A 1987, 41, 230.
- (19) Asplund, M.; Jagner, S. *Acta Chem. Scand.* A 1984, 38, 255.
- (20) Andersson, S.; Jagner, S. *Acta Crystallogr., C* 1987, 43, 1089.
- (21) Asplund, M.; Jagner, S. *Acta Chem. Scand.*, A 1984, 38, 411.
- (22) Asplund, M.; Jagner, S.; Nilsson, M. *Acta Chem. Scand.*, A 1982, 36, 751.

Table I. Selected Halocuprate Structures

compound		coord no. of Cu(I)	Cu-X, $^{\circ}$ Å	Cu...Cu, Å	structural type	ref
cation	anion					
NBu <sub>4</sub> <sup>a</sup>	CuCl <sub>2</sub>	2	2.107		discrete linear	9
NBu <sub>4</sub>	CuBr <sub>2</sub>	2	2.227		discrete linear	9
NBu <sub>4</sub>	CuBrCl	2	2.195, 2.104		discrete linear	10
NPr <sub>4</sub> <sup>b</sup>	CuCl <sub>2</sub>	2	2.071		discrete linear	11
NPhMe <sub>3</sub> <sup>c</sup>	CuCl <sub>2</sub>	2	2.117		discrete linear	12
PPh <sub>3</sub> Et <sup>d</sup>	CuBr <sub>2</sub>	2	2.224		discrete linear	13
PPh <sub>3</sub> Pr	CuBr <sub>2</sub>	2	2.225		discrete linear	14
AsPh <sub>4</sub>	CuCl <sub>2</sub>	2	2.072		discrete linear	15
PPh <sub>4</sub>	CuCl <sub>2</sub>	2	2.080		discrete linear	16
PPh <sub>4</sub>	CuBr <sub>2</sub>	2	2.211		discrete linear	16
NEt <sub>4</sub>	Cu <sub>2</sub> Br <sub>4</sub>	3	2.319, 2.441	2.937	2 <sup>f</sup>	16
NPhMe <sub>3</sub>	Cu <sub>2</sub> Br <sub>4</sub>	3	2.310, 2.421	2.738	2	17
PMe <sub>4</sub>	CuBr <sub>3</sub>	3	2.365		1	18
PEt <sub>4</sub>	Cu <sub>2</sub> Br <sub>4</sub>	3	2.263, 2.423	2.870	2	19
PMePh <sub>3</sub>	Cu <sub>2</sub> Br <sub>4</sub>	3	2.337, 2.426	2.697	2	20
NPr <sub>4</sub>	Cu <sub>2</sub> I <sub>4</sub>	3	2.499, 2.571	2.698	2	21
NBu <sub>4</sub>	Cu <sub>2</sub> I <sub>4</sub>	3	2.514, 2.566	2.726	2	22
PMePh <sub>3</sub>	CuI <sub>3</sub>	3	2.537		1	23
AsPh <sub>4</sub>	Cu <sub>2</sub> I <sub>4</sub>	3	2.490, 2.578	2.663	2	24
NPr <sub>4</sub>	Cu <sub>4</sub> Br <sub>6</sub>	3	2.373	2.718	3	25
PBuPh <sub>3</sub>	Cu <sub>4</sub> Br <sub>6</sub>	3	2.410	2.719	3	26
PMePh <sub>3</sub>	Cu <sub>4</sub> I <sub>6</sub>	3	2.638	2.742	3	27
NMe <sub>4</sub>	Cu <sub>2</sub> Br <sub>5</sub>	3	2.381, 2.392	2.837	4	28
NMeBu <sub>3</sub>	Cu <sub>3</sub> Br <sub>7</sub>	4	2.345, 2.912	2.605	5 <sup>g</sup>	29
NPr <sub>4</sub>	Cu <sub>3</sub> I <sub>7</sub>	4	2.510, 2.640	2.580	5	30
NMeEt <sub>3</sub>	Cu <sub>6</sub> Br <sub>9</sub>	2, 4	2.285, 2.697	3.030	6	1
NEt <sub>4</sub>	[Cu <sub>6</sub> I <sub>11</sub> ]I	4	2.567, 2.731	2.961	7	31
NMe <sub>4</sub>	Cu <sub>2</sub> I <sub>3</sub>	4	2.517	2.452	polymeric chain 8	32
NMe <sub>4</sub>	Cu <sub>2</sub> Cl <sub>3</sub>	4	2.254, 2.494	2.869	polymeric chain 10	33
(Me <sub>2</sub> N <sub>2</sub> ) <sub>2</sub> CH	Cu <sub>2</sub> Br <sub>3</sub>	4	2.431, 2.600	3.03	polymeric chain 10	8
SMe <sub>3</sub>	Cu <sub>2</sub> I <sub>3</sub>	4	2.612, 2.724	2.954	polymeric chain 10	34
NEt <sub>4</sub>	Ag <sub>2</sub> Cl <sub>3</sub>	4	2.519, 2.679	3.348	polymeric chain 10	35
NMe <sub>4</sub>	Ag <sub>2</sub> Br <sub>3</sub>	4	2.612, 2.804	3.078	polymeric chain 10	35
S <sub>2</sub> C <sub>3</sub> (SCH <sub>3</sub> ) <sub>3</sub>	Cu <sub>2</sub> I <sub>3</sub>	3, 4	2.523, 2.730	2.551	polymeric chain 9	36
C <sub>23</sub> H <sub>17</sub> S	Cu <sub>2</sub> I <sub>3</sub>	4	2.651, 2.542	2.479	polymeric chain <sup>h</sup>	37

<sup>a</sup>Tetra-*n*-butyl. <sup>b</sup>Propyl. <sup>c</sup>Phenyl. <sup>d</sup>Ethyl. <sup>e</sup>Only the shortest distance is quoted. <sup>f</sup>Numbers refer to the text structures. <sup>g</sup>In the bromide structure the halides are not coplanar with the coppers. <sup>h</sup>This structure resembles chain 8.

wherein the anion has an extended structure. Since then several other polymeric systems have been made.

While we would love to understand in detail why a given cation leads to a specific counteranion aggregation mode of Cu<sup>+</sup> and X<sup>-</sup> entities, we cannot be optimistic about theory contributing decisively to an answer of that basic chemical question. The thermodynamics of aggregation is not easy to predict, especially in structures partly covalent and partly ionic, and the aggregation forces may be under kinetic control. The more restricted aim of this paper is to examine the peculiar bonding modes of the halides, in particular the strange  $\mu_4$  square pyramidal and  $\mu_5$  configurations. It is easy to understand, in the context of normal ideas about coordinate covalent bonding, the nature of terminal,  $\mu_2$ ,  $\mu_3$ , and tetrahedral  $\mu_4$  halides. But the nontetrahedral  $\mu_4$  and the  $\mu_5$  systems are at first sight strange.

A word about the copper-halogen bond metrics is in order. The Cu-X bond distances listed in Table I indicate that as the extent of bridging around the halogen atom increases, the bond distances also increase. The Cu-Cl distances vary from 2.107 to 2.254 and 2.494 Å as the Cl goes from terminal to bridging two and four Cu atoms. The Cu-Br distance increases from 2.227 Å to 2.421, 2.600, and 2.912 Å as the Br goes from terminal to bridging two, four, and five Cu atoms, respectively. The Cu-I distance changes from 2.499 to 2.571, 2.517, 2.724, and 2.640 Å as the I goes from terminal to bridging two, three, four, and five Cu atoms. Except for the  $\mu_3$ -I and the  $\mu_5$ -I cases, there seems to be a general tendency for the Cu-X distance to increase with bridging.

We begin our study with the bonding of nontetrahedral  $\mu_4$ -Br compounds, 6 and 10. In both these systems, the  $\mu_4$ -Br atom occupies the apex of a square pyramid formed by four Br-Cu bonds with a tetrahedral geometry around each Cu atom. We have performed the relevant calculation using extended Hückel method,<sup>38-43</sup> the parameters are listed in the Appendix.

In the polymeric Cu<sub>2</sub>Br<sub>3</sub><sup>-</sup> system,<sup>8</sup> the anions are well-separated from the organic cation. As can be seen in 10, the  $\mu_4$ -Br atoms are alternately above and below the plane of the Cu atoms. The  $\mu_2$ -Br atoms also alternate with the  $\mu_4$ -Br atoms. The  $\mu_4$ -Br-Cu distance varies from 2.58 to 2.609 Å while  $\mu_2$ -Br-Cu is 2.432 Å.

The  $\mu_4$ -Br-Cu distance in the extended structure 10 is approximately 0.1 Å shorter than that reported for the discrete

- (23) Bowmaker, G. A.; Clark, G. R.; Rogers, D. A.; Camus, A.; Marsich, N. *J. Chem. Soc., Dalton Trans.* **1984**, 37.  
 (24) Asplund, M.; Jagner, S. *Acta Chem. Scand.*, **A 1984**, 38, 297.  
 (25) Asplund, M.; Jagner, S. *Acta Chem. Scand.*, **A 1984**, 38, 725.  
 (26) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, **A 1986**, 40, 210.  
 (27) Bowmaker, G. A.; Clark, G. R.; Yuen, D. K. P. *J. Chem. Soc., Dalton Trans.* **1976**, 2329.  
 (28) Asplund, M.; Jagner, S. *Acta Chem. Scand.*, **A 1985**, 39, 47.  
 (29) Andersson, S.; Jagner, S. *J. Crystallogr. Spectrosc. Res.* **1988**, 18, 591.  
 (30) Hartl, H.; Mahdjour-Hassan-Abadi, F. *Angew. Chem., Int. Ed. Engl.* **1984**, 23, 378.  
 (31) Mahdjour-Hassan-Abadi, F.; Hartl, H.; Fuchs, J. *Angew. Chem., Int. Ed. Engl.* **1984**, 23, 514.  
 (32) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, **A 1985**, 39, 181.  
 (33) Andersson, S.; Jagner, S. *Acta Chem. Scand.*, **A 1986**, 40, 177.  
 (34) Andersson, S.; Jagner, S.; Nilsson, M. *Acta Chem. Scand.*, **A 1985**, 39, 447.  
 (35) Helgesson, G.; Jagner, S. *Acta Crystallogr.*, **C 1988**, 44, 2059.  
 (36) Asplund, M.; Jagner, S. *Acta Chem. Scand.*, **A 1984**, 38, 129.  
 (37) Batsanov, A. S.; Struchkov, Yu. T.; Ukhin, L. Yu.; Dolgoplova, N. A. *Inorg. Chim. Acta* **1982**, 63, 17.

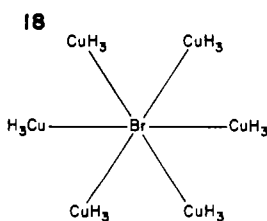
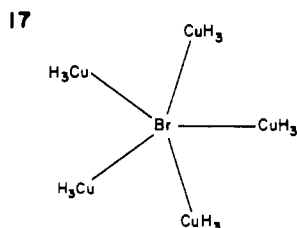
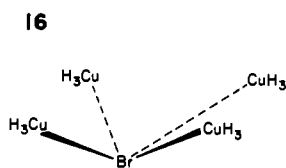
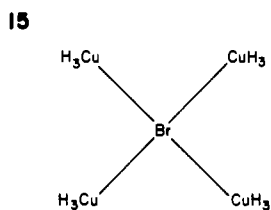
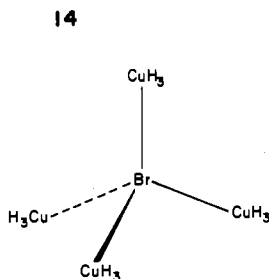
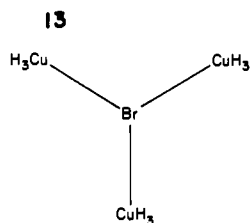
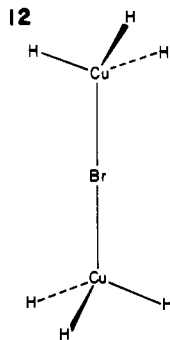
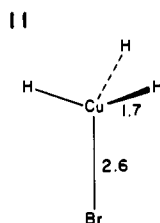
- (38) Hoffmann, R. *Solids and Surfaces: A Chemist's View of Bonding in Extended Structures*; VCH: New York, 1988.  
 (39) Hoffmann, R. *Angew. Chem., Int. Ed. Engl.* **1987**, 26, 846.  
 (40) Hughbanks, T.; Hoffmann, R. *J. Am. Chem. Soc.* **1983**, 105, 1150.  
 (41) Hoffmann, R.; Whangbo, M.-H. *J. Am. Chem. Soc.* **1978**, 100, 6093.  
 (42) Whangbo, M.-H.; Hoffmann, R.; Woodward, R. B. *Proc. R. Soc. London*, **A 1979**, 366, 23.  
 (43) Hoffmann, R. *Rev. Mod. Phys.* **1988**, 60, 601.

molecule  $[N(Me)(Et)_3]_3[Cu_6Br_9]$ . In the discrete molecule, the Cu...Cu distance along the  $a$ -axis is 3.199 Å and that along the  $c$ -axis is 2.862 Å. In both, the discrete molecule and the polymeric system, the Cu...Cu distance is large compared to that in the metal, which is 2.55 Å.

### Bonding in Model Bromocuprates

Our aim is 2-fold: (i) to understand the change in the Br-Cu bond lengths with variation in the number of atoms bridged by Br atom and (ii) to study the electronic structure and bonding in the unusual square pyramidal geometry around the  $\mu_4$ -Br.

To accomplish this, we have constructed model systems in which one Br atom is connected to varying number of Cu atoms, as shown in 11-18.



Structures 11-13, 15, 17, and 18 are models for  $\mu_1$ - $\mu_6$ -Br coordination, idealized for the moment to in-plane bonding for  $\mu_3$ - $\mu_6$ . 14 and 16 are alternative tetrahedral and square pyramidal  $\mu_4$  geometries, respectively. In each case the geometry around the Cu atoms is maintained as tetrahedral, three hydrogens completing the coordination sphere of each Cu. Consistent with a +1 oxidation state for Cu and -1 for H, each  $CuH_3$  group enters

Table II. Reduced Overlap Population of Br-Cu and Cu...Cu Bonds for the Various Models

type	overlap population			Cu...Cu dist, Å
	Br-Cu (one bond)	Br-Cu (sum over all bonds)	Cu...Cu	
$\mu_1$ -(11 <sup>a</sup> )	0.239	0.239		5.200
$\mu_2$ -(12)	0.217	0.434	0.000	4.503
$\mu_3$ -(13)	0.235	0.705	-0.001	4.246
tetr $\mu_4$ -(14)	0.243	0.972	-0.001	3.677
plan $\mu_4$ -(15)	0.218	0.872	-0.003	3.000
pyr $\mu_4$ -(16)	0.219	0.876	-0.004	3.056
$\mu_5$ -(17)	0.202	1.010	-0.031	2.600
$\mu_6$ -(18)	0.177	1.062	-0.015	

<sup>a</sup> Numbers refer to structures in the text.

the calculation with a 2- charge. In all the cases the Br-Cu distance has been maintained at 2.6 Å while the Cu-H distances were set at 1.7 Å. Note that we have kept the Br-Cu distance constant, knowing full well that in the real compounds it varies, increasing with the number of bonded Cu atoms. The philosophy of the approach is clear; i.e., we give the Cu-Br bond a chance to be equal in strength by making it of the same bond length in various geometries and then let the molecule tell us itself, through the overlap populations, what bond strength it desires. Table II shows the total and individual overlap population for the Cu-Br and the Cu...Cu bonds as well as the Cu...Cu distance for all the cases 11-18.

Thinking classically about these structures, we can consider the  $\mu_1$ ,  $\mu_2$ ,  $\mu_3$ , and tetrahedral  $\mu_4$  coordinations of Br as "normal" and the square pyramidal  $\mu_4$ ,  $\mu_5$ , and  $\mu_6$  as untypical. Indeed the overlap populations in the first normal group are larger, consistently so, than in the second group. The exception is the linear  $\mu_2$  geometry. We could try a further analysis, using our prejudices about hybridization and the notion that an  $sp^3$ - $sp^3$  bond should be weaker than an  $sp^3$ - $sp^2$  or  $sp^3$ - $sp$  bond. The hybridization at Br in 11 is ambiguous (is it  $sp^3$  or  $sp^2$ ?), but for 12, most everyone would pick  $sp$  at Br, and for 13  $sp^2$  and for 14  $sp^3$ . In fact the Cu-Br overlap population rises along the series 12  $\rightarrow$  13  $\rightarrow$  14, instead of falling. The classical hybridization argument for bond strength does not work for bridging Br. This is confirmed by the Cu-X distances in Table I.

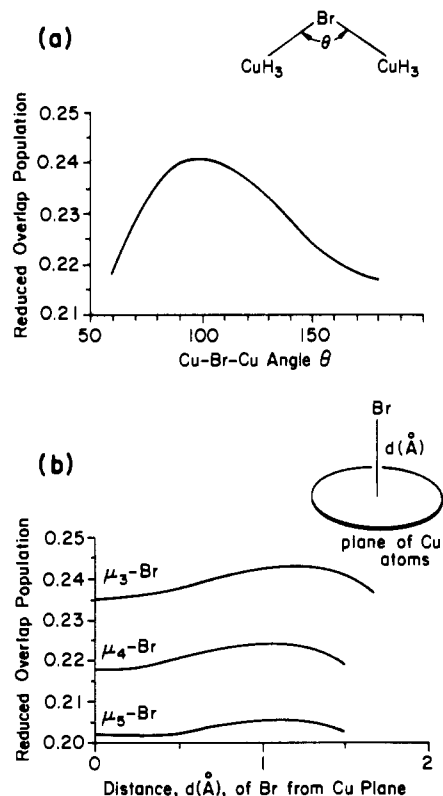
Hybridization is, of course, a qualitative valence bond notion, one of extreme heuristic value. The computed 4s and 4p orbital populations at Br in 12-14 do not match the expectations of  $sp$ ,  $sp^2$ , and  $sp^3$  taken seriously, (i.e., 50%, 33%, and 25% s character, respectively), but the 4p orbital population does increase along this series.

It is interesting that the strongest Cu-Br bond is obtained for four Cu atoms tetrahedrally disposed around a Br. But there has been no report of a discrete bromocuprate where the geometry around a four-coordinate Br atom is tetrahedral! However, considering the solid-state cuprous halides,<sup>44</sup> CuCl, CuBr, and CuI crystallize in the zinc blende structure. At high temperatures, they transform to the wurtzite structure. In both these forms, the halogen atom has a tetrahedral environment. The argentous halides, AgF, AgCl, and AgBr, on the other hand, adopt the NaCl structure.<sup>44</sup> AgI prefers the wurtzite structure under normal conditions of temperature and pressure.<sup>44</sup> As for the aurous halides, structural information is available only for AuI, which consists of bent chains made up of alternating Au and I atoms.<sup>45</sup> The geometry around Au is linear while the Au-I-Au angle is 72°.

Returning to the halocuprate systems, the geometries 11-18 are idealized. In the real molecules where Br bridges two Cu atoms,  $\mu_2$  coordination is almost never linear and the  $\mu_4$  cases consistently pyramidal. This suggested a study of departures from linearity or planarity for the  $\mu_2$ - $\mu_6$  structures. Figure 1a shows

(44) Lawn, B. R. *Acta Crystallogr.* 1964, 17, 1341. Burley, G. J. *Phys. Chem.* 1964, 68, 1111.

(45) Jagodzinski, H. Z. *Kristallogr.* 1959, 112, 80.



**Figure 1.** (a) Reduced overlap population as a function of Cu-Br-Cu angle,  $\theta$ , for  $\text{Br}(\text{CuH}_3)_2$  system. (b) Reduced overlap population as a function of distance of Br atom from the plane of Cu atoms for the three cases of Br bridging three, four, and five Cu atoms in  $\text{Br}(\text{CuH}_3)_n$ .

a plot of the overlap population of the Cu-Br bond as a function of the bridging angle, Cu-Br-Cu, for the case of Br bridging two Cu atoms.

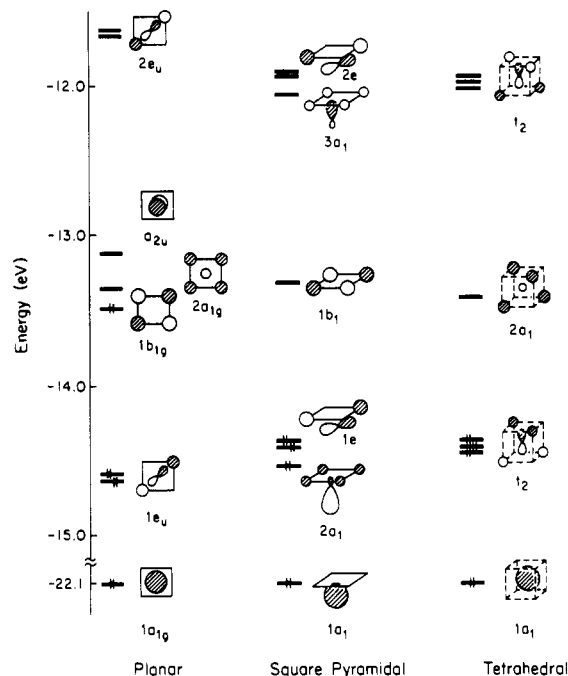
There is maximum overlap population at an angle of  $100^\circ$ . A maximum is also seen in Figure 1b, which shows overlap population as a function of distance of Br from the plane of Cu atoms for the cases of Br bridging three, four and five Cu atoms. The Cu-Br bond length is maintained constant while the Br atom is pulled out of the plane of the Cu atoms. The overlap population decreases as we go from  $\mu_3$  to  $\mu_4$  to  $\mu_5$ -Br. With Br pulled out of the plane of the Cu atoms, the reduced overlap population first increases till the distance of 1.15 Å and then decreases for all three cases. The total energy varies little with the angle (or distance) for the  $\mu_2$  (or  $\mu_5$ ) case. For Br bridging two Cu atoms the energy minimizes at a Cu-Br-Cu angle of  $100^\circ$ . For the cases  $\mu_3$ - $\mu_5$ , an energy minimum is obtained when the distance of Br atom from the plane of Cu atoms is 1.1 Å, which is also the distance at which the reduced overlap population is maximum.

Since it is possible that the total energy increases because the H atoms attached to the different Cu atoms come too close as the Cu-Br-Cu angle is decreased or as the distance of Br from the plane increases, we decided to repeat the calculations in the absence of H atoms. Similar results were obtained, with overlap population maxima and energy minima in some cases for a somewhat more out-of-plane Br-Cu<sub>n</sub> geometry.

On the basis of these calculations, we may predict structural preferences for Br bridging different number of Cu atoms. For  $\mu_2$ -Br, a bent geometry is preferred over a linear one. For the cases of Br bridging three, four, and five atoms, pyramidal structures are more likely than planar geometries.

Though we have used the same bond length for Cu-Br, we get different overlap populations when Br bridges different number of Cu atoms, depending on the orientation of the Cu atoms around Br.

The question remains: Why does one see a high coordination number such as 5 for Br, as well as the unusual pyramidal  $\mu_4$ -Br coordination mode? In Table II we list another bond indicator, the sum of all Cu-Br overlap populations. This increases steadily



**Figure 2.** Molecular orbitals for planar, square pyramidal, and tetrahedral  $\text{BrH}_4^{3+}$ .

from  $\mu_1$  to  $\mu_6$ . While it seems to saturate, that limit is reached only at  $\mu_6$ . Obviously, Br is capable of forming more than four effective bonds. We proceed to a more detailed examination of the bonding around Br in a still simpler model.

#### Insight from a Model $\text{BrH}_4^{3+}$ System

In order to study the geometry around the Br atom, a simple model,  $\text{BrH}_4^{3+}$  was considered. Calculations were performed on this system keeping the Br-H distance constant at 2.6 Å. Initially the Br atom was kept in the plane of the H atoms. The orbital diagram is shown in Figure 2. Many of the orbital features of these molecules have been discussed by Gimarc<sup>46</sup> and by Minkin and co-workers.<sup>47</sup>

$1a_{1g}$  is composed mainly of the Br s orbital while the  $e_u$  is a degenerate orbital having the  $p_x$  and  $p_y$  of Br mixing with appropriate linear combinations of H atomic orbitals. The highest occupied molecular orbital is the  $b_{1g}$ , a nonbonding orbital on Br. There is no Br orbital of  $b_{1g}$  symmetry. Similarly, the Br  $p_z$  is orthogonal to all H orbitals. Therefore, the eight electrons in the system occupy three bonding molecular orbitals and one nonbonding molecular orbital. So, in a localized picture, the molecule has only three pairs of bonding electrons to form four bonds.

When the Br atom is pulled 1.5 Å out of the plane of H atoms, the Br-H distance is still maintained at 2.6 Å. The  $a_{2u}$  containing the Br  $p_z$  orbital is now stabilized by bonding interactions and falls below the  $e_u$  set, thus making the nonbonding  $b_{1g}$  the lowest unoccupied molecular orbital. There is an effective level crossing along this seemingly minor geometrical distortion. The HOMO is stabilized by 0.9 eV compared to the planar system. Now there are four pairs of electrons occupying four bonding orbitals. The Br-H overlap population in the square planar geometry is 0.13, compared to 0.15 in the square pyramidal geometry. Note the great improvement in the HOMO-LUMO gap in the square pyramidal geometry. The square planar to square pyramidal distortion can be described as a typical second-order Jahn-Teller deformation.<sup>48</sup> A nonbonding molecular orbital in the square

(46) Gimarc, B. M. *Molecular Structure and Bonding: The Qualitative Molecular Orbital Approach*, Academic: New York, 1979.

(47) Minkin, V. I.; Simkin, B. Ya.; Minyaev, R. M. *Quantum Chemistry of Organic Compounds*; Springer-Verlag: Berlin, Heidelberg, Germany, 1990.

(48) Pearson, R. G. *Symmetry Rules for Chemical Reactions*, Wiley-Interscience: New York, 1976.

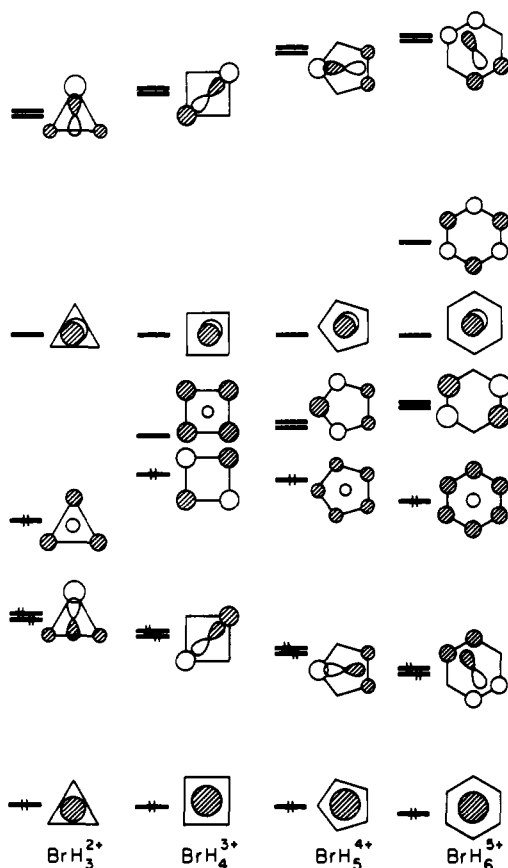


Figure 3. Molecular orbitals for planar  $\text{BrH}_n^{(n-1)+}$ ,  $n = 3, 4, 5$ , and  $6$ .

Table III. Overlap Population of Br-H Bond in  $\text{BrH}_n^{(n-1)+}$

type	overlap population of Br-H (one bond)
$\mu_3\text{-BrH}_3^{2+}$	0.136
plan $\mu_4\text{-BrH}_4^{3+}$	0.130
$\mu_5\text{-BrH}_5^{4+}$	0.100
$\mu_6\text{-BrH}_6^{5+}$	0.087
tetr $\mu_4\text{-BrH}_4^{3+}$	0.161
pyr $\mu_4\text{-BrH}_4^{3+}$	0.150

planar geometry (the  $p_z$  of Br) is transformed into a bonding one upon pyramidalization.

When the geometry around the Br atom is taken as tetrahedral (right side of Figure 2), maintaining the Br-H distance at 2.6 Å, the HOMO is the  $t_2$  set and is very close in energy to the HOMO of the square pyramidal system.

At this juncture, it is interesting to study the molecular orbitals for the planar  $\mu_3$ ,  $\mu_4$ ,  $\mu_5$ , and  $\mu_6\text{-BrH}_n^{(n-1)+}$  model systems. These are shown in Figure 3.

In all the planar models considered, there are a total of eight electrons. Except for  $\mu_4\text{-BrH}_4^{3+}$ , the highest occupied molecular orbital is slightly antibonding with respect to Br and H. In all the cases there are three bonding molecular orbitals. The overlap population for these systems as well as for the tetrahedral and square pyramidal  $\mu_4$  systems is shown in Table III.

The overlap population is a maximum for the tetrahedral  $\mu_4$  case. Similar to observations for the model systems 13, 15, 17, and 18, the overlap population decreases with an increase in the number of atoms bonded to Br. So there is saturation, but it sets in slowly. The overlap population of pyramidal  $\mu_4$  lies between that of the planar and tetrahedral models.

The increased Br-H overlap population of the HOMO is an important factor for Br atom preferring a nonplanar geometry. However, it gives us no clue as to why a square pyramidal geometry is preferred over a tetrahedral one.

A more detailed analysis of the nature of orbitals indicated that, upon pyramidalization of each of the systems, one  $a_1$  orbital is

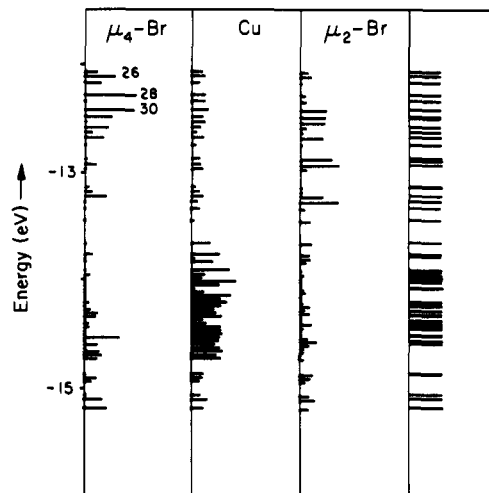
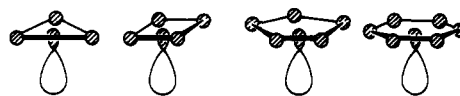


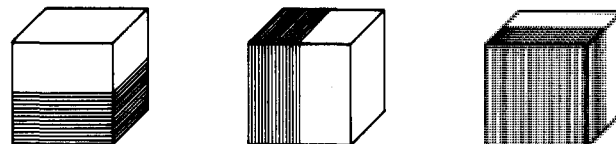
Figure 4. Projection of orbital contribution of  $\mu_4\text{-Br}$ , Cu, and  $\mu_2\text{-Br}$  atoms to the molecular orbitals of the discrete system  $\text{Cu}_6\text{Br}_3^{3-}$ .

strongly stabilized and the degenerate  $e$  set is destabilized. The final shape of the critical  $a_1$  orbital is shown in 19. Mixing in



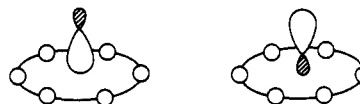
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of  $s$  character to form a lone pair pointing away from the ligands is clearly a stabilizing feature for this orbital as is the  $p_z$  bonding with the hydrogens, which is turned on when the Br moves out of the plane. Another way to describe this orbital is it is the third  $p$ -type "united atom" orbital of an  $\text{XH}_n$  system, 20, the other ones being the  $e$  set.



20

Still another way to describe the presence of four bonding or nonbonding orbitals in  $\text{XH}_n$  systems is to invoke multicenter bonding, an extrapolation of the three-center bonding concept exploited so fruitfully by Lipscomb for boron hydrides.<sup>49</sup> If one forms  $sp$  hybrids along the  $C_n$  axis at the Br, one can form one bonding combination (21) and one lone pair (22), as discussed



21

22

above. The two unhybridized  $p_{xy}$  orbitals then form multicenter bonds with an appropriate symmetry-adapted ligand combination (23) of  $e$ -type symmetry. There is a total of four bonding or nonbonding orbitals.

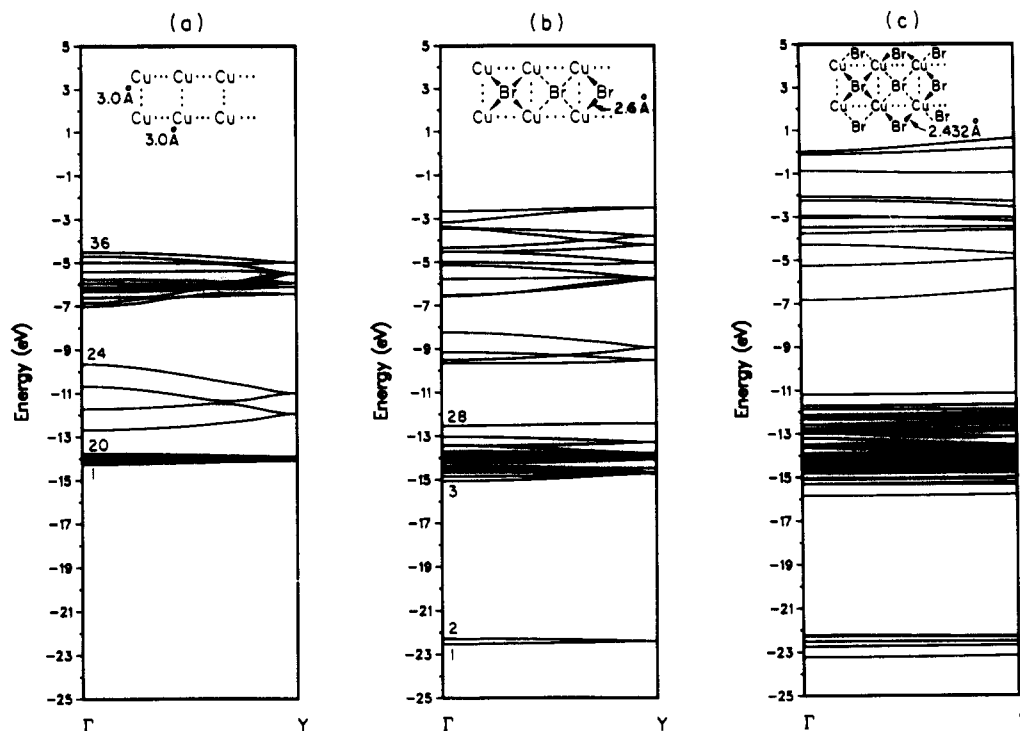


23

#### The Square Pyramidal Molecular System: $\text{Cu}_6\text{Br}_3^{3-}$

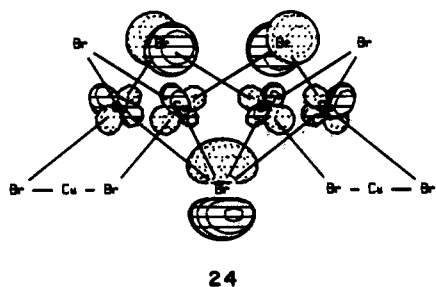
A central feature of the pyramidal  $\mu_n\text{-BrH}_n$  model systems of the previous section is the existence of a well-defined lone pair. It becomes interesting to see if this orbital feature is to be found

(49) Lipscomb, W. N. *Boron Hydrides*, Benjamin; New York, 1963.



**Figure 5.** Calculated band structures for (a) the Cu sublattice, (b) the sublattice containing the Cu atoms and the  $\mu_4$ -Br atoms, and (c) the complete lattice.

from a calculation on an actual pyramidal  $\mu_4$ -Br containing molecule such as  $\text{Cu}_6\text{Br}_9^{3-}$  shown in 6. Calculations on this molecule gave a band of Br 4s levels at  $-22.0$  eV and Cu 3d levels around  $-14.0$  eV. The bonding and nonbonding Br 4p levels mix with Cu 3d and 4s in the region between  $-12.0$  and  $-15.0$  eV. The orbital energies are shown in Figure 4. This is analogous to the projection of density of states for extended systems.<sup>38</sup> The horizontal length of a "stick" measures the contribution of the specified atom to the composition of each molecular orbital. The first column shows the projection of molecular orbitals having a  $\mu_4$ -Br contribution. The second column shows the contribution of one of the Cu atoms, and the third corresponds to  $\mu_2$ -Br contribution. Though many levels are very close in energy and may even be degenerate, they have been shown as individual "sticks" in the first three columns of the figure. In order to get a feeling for the energy level distribution in the energy range  $-16.0$  to  $-12.0$  eV, the levels are marked on the extreme right of the figure at appropriate energy values. Atomic orbitals of  $\mu_4$ -Br atoms are distributed over numerous molecular orbitals. The biggest  $\mu_4$ -Br concentrations close to the HOMO are in the orbitals marked 26, 28, and 30. Further analysis of these three molecular orbitals showed that while molecular orbitals 26 and 28 have  $p_x$  and  $p_y$  contributions, molecular orbital 30 contains  $p_z$  of  $\mu_4$ -Br. This orbital is of interest to us since our intention is to find a molecular orbital with the lone pair character on  $\mu_4$ -Br. A representation<sup>50</sup> of orbital 30 is shown in 24. While the Cu contributions in this molecular orbital are small, it has substantial  $\mu_2$ -Br character mixed in.



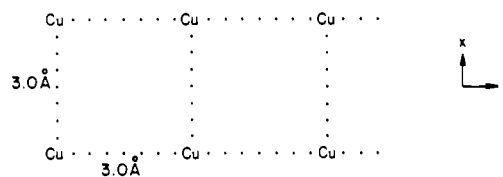
24

We proceed next to the band structure of the extended polymer 10.

#### The One-Dimensional Polymer: $\text{Cu}_2\text{Br}_3^-$

In order to understand the bonding in the extended structure 10, we have averaged the bond lengths of all the chemically equivalent bonds. In the real system the Cu...Cu distances are  $2.86$  Å along the chain direction and  $3.09$  Å perpendicular to the chain direction. We have considered an idealized system with all the Cu...Cu distances equal to  $3.0$  Å. The  $\mu_4$ -Br-Cu distance is  $2.6$  Å and the  $\mu_2$ -Br-Cu distance is  $2.432$  Å. The basic unit cell is made up of  $\text{Cu}_2\text{Br}_3^-$  units. Initially we look at the Cu sublattice and then we study the interaction of the  $\mu_4$ -Br with Cu by introducing the Br atoms in the center. Finally, we include the  $\mu_2$ -Br atoms in the lattice.

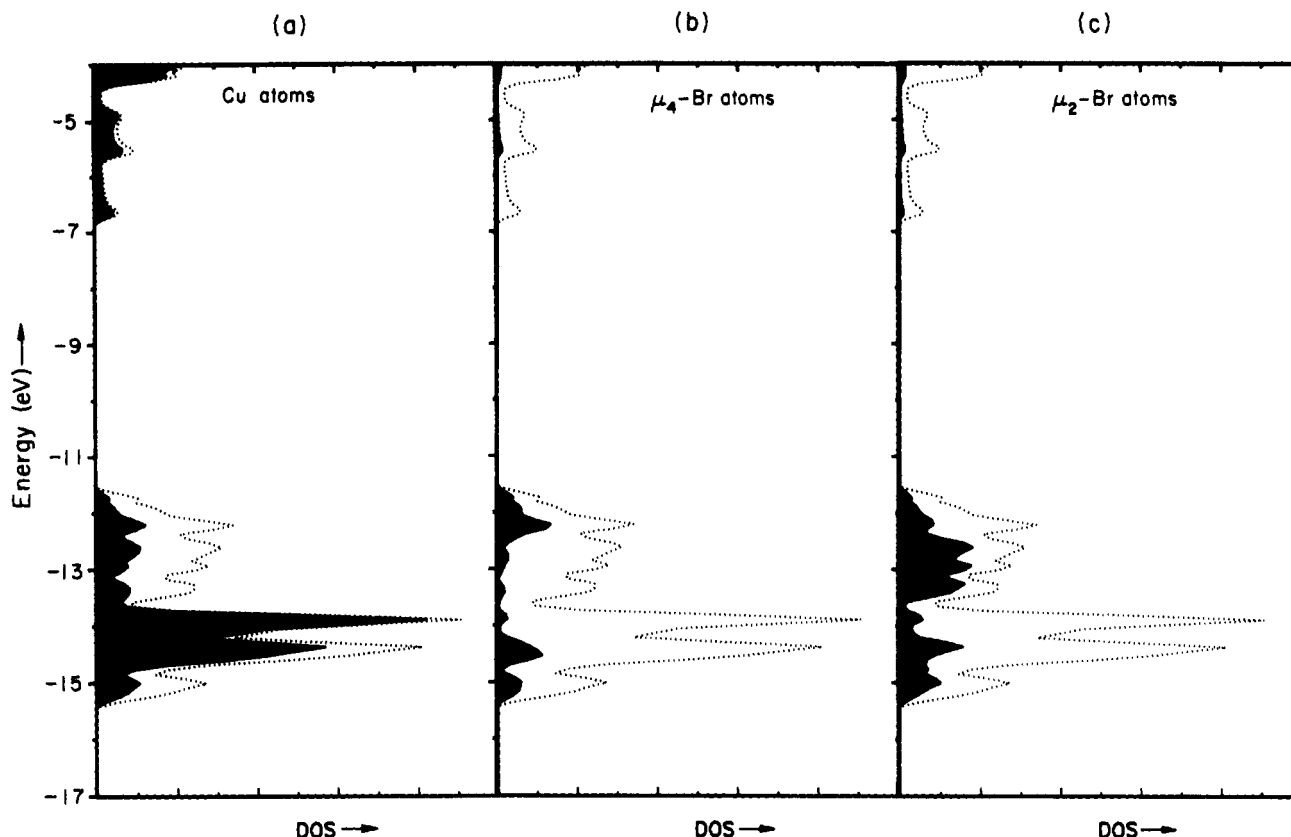
To facilitate the subsequent comparison with the complete system, which has a glide plane of symmetry, we have considered four Cu atoms in a unit cell as shown below. The chain extends along the  $y$ -direction, as indicated in 25. The band structure of



25

this chain is seen in Figure 5a. The group of 20 bands located at  $-14.0$  eV consists of Cu d-bands, which show little dispersion. Bands 21–24 are the in-phase and out-of-phase linear combinations of Cu s-orbitals. Bands 25–36 are due to the p-orbitals. The total overlap population is 0.02 for Cu...Cu perpendicular to the chain and is the same along the chain.

Proceeding to the one-dimensional chain with the inclusion of the  $\mu_4$ -Br, we have maintained all the  $\mu_4$ -Br-Cu distances at  $2.6$  Å. Since the Br p-orbitals are close in energy to the Cu d-orbitals, there is some interaction between them. The band structure is shown in Figure 5b for  $\text{Cu}_4\text{Br}_2^{2+}$  chain. Here, only the  $\mu_4$ -Br atoms have been introduced in the lattice. The bands in the region  $-7.0$  to  $-2.0$  eV are combinations of Cu p-orbitals and Br p-orbitals.



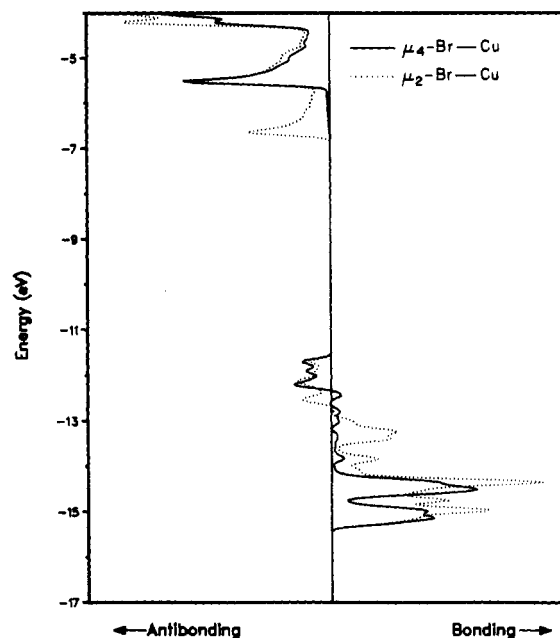
**Figure 6.** Calculated DOS for the  $(\text{Cu}_2\text{Br}_3)_n$  polymeric system. The projections (shaded areas) indicate the contributions to this DOS from (a) Cu atoms, (b)  $\mu_4$ -Br atoms, and (c)  $\mu_2$ -Br atoms.

Bands 26–28 are localized largely on Br p-orbitals with a small contribution from Cu d-orbitals. Bands 1 and 2 correspond to Br s-orbitals. Bands 3–8 are combinations of Cu s and Br p-orbitals while bands 9–22 are mainly Cu d-orbitals. At  $\Gamma$ , band 29 of Figure 5b is similar to band 24 of Figure 5a, since there is no Br contribution to this band. Bands 30–32 have all increased in energy at  $\Gamma$  compared to the respective bands 21–23 in Figure 5a because of antibonding interactions of Br p-orbitals with Cu s-orbitals. The total overlap population between Cu and Br is 0.26, for the Cu–Cu interactions perpendicular to the chain it is 0.01, and that along the chain is 0.03. The observed bond length in the real system is 3.09 Å perpendicular to the chain and 2.86 Å along the chain direction. This trend seems to be reflected in the total overlap populations, though these are small.

The  $\mu_2$ -Br atoms were introduced next, with the band structure shown in Figure 5c, and the bonding in the system was studied with the help of decompositions of densities of states (DOS) and crystal orbital overlap populations (COOP). Figure 6 shows the DOS of Cu,  $\mu_4$ -Br, and  $\mu_2$ -Br. It is rather difficult to identify any crystal orbital or band as the  $\mu_4$ -Br lone pair, though there is a peak in the  $\mu_4$ -Br contribution around -12.0 eV.

The COOP curves for the various Br–Cu interactions are shown in Figure 7. Just below the Fermi level (which is at -11.7 eV), both of the Br–Cu interactions are antibonding. However, the total overlap population is bonding in all cases. Around the region of -12.0 eV, the  $\mu_4$ -Br–Cu interaction is slightly antibonding and might be considered to be similar to the molecular orbital 30 shown in structure 24 for the discrete molecule. From the DOS projections (for single atoms), the orbitals around this energy value were found to have a greater contribution from each  $\mu_4$ -Br atom (about 20%) than from each Cu atom (about 4%). Further breakdown of the projection of DOS in terms of orbitals indicated that, around this region, there is greater contribution from the  $p_z$  orbital of  $\mu_4$ -Br than from other orbitals. Magnified projections of  $p_x$ ,  $p_y$ , and  $p_z$  orbitals of  $\mu_4$ -Br along with the total density of states is shown in Figure 8.

The crystal orbital overlap population is 0.20 for  $\mu_4$ -Br–Cu and 0.35 for  $\mu_2$ -Br–Cu. The total overlap for  $\mu_2$ -Br–Cu is larger than



**Figure 7.** Crystal orbital overlap population (COOP) curves for  $\mu_4$ -Br–Cu (solid line) and  $\mu_2$ -Br–Cu (dotted line).

that for  $\mu_4$ -Br–Cu, and in order to compare their bond strengths, the calculation was repeated on the polymeric chain while keeping all the Br–Cu distances equal to 2.6 Å. The crystal orbital overlap population is 0.22 for  $\mu_4$ -Br–Cu and 0.25 for  $\mu_2$ -Br–Cu. Here again the overlap population is larger for  $\mu_2$ -Br–Cu than for  $\mu_4$ -Br–Cu, though the difference is not as large as in the previous calculation.

#### Tetrahedral versus Square Pyramidal $\mu_4$ -Br

One possible alternative structure for a polymeric chain with the same stoichiometry ( $\text{Cu}_2\text{Br}_3^-$ ) where the  $\mu_4$ -Br is tetrahedral

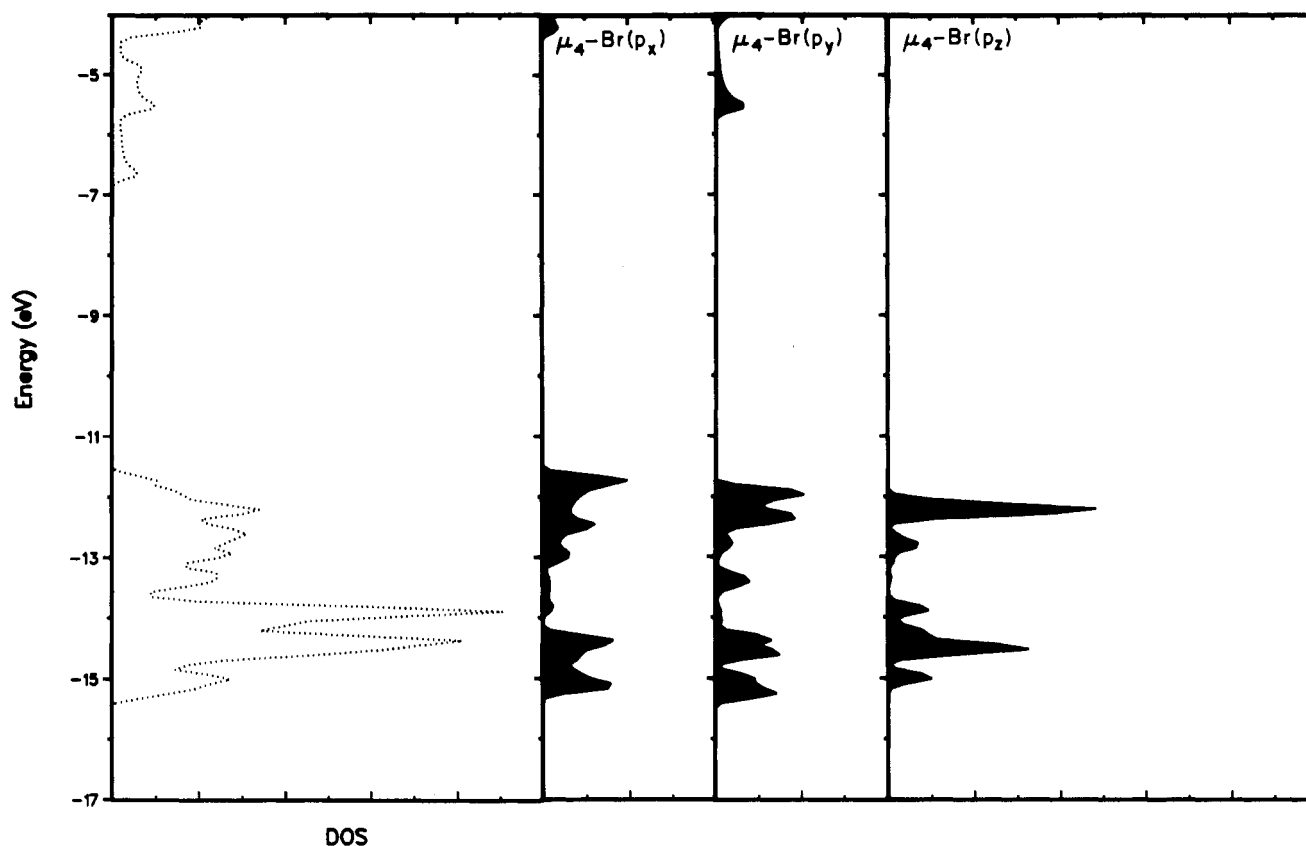
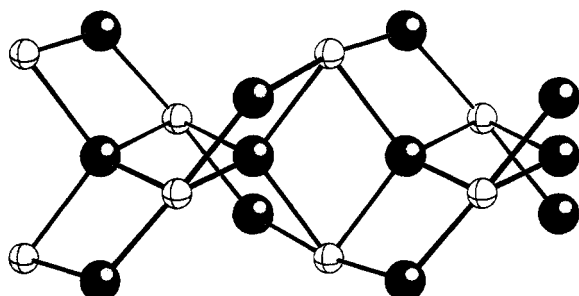


Figure 8. Calculated DOS for the  $(\text{Cu}_2\text{Br}_3)_n$  polymeric system. The magnified projections (shaded areas) indicate the contributions to this DOS from  $p_x$ ,  $p_y$ , and  $p_z$  atomic orbitals of  $\mu_4$ -Br atoms.

is **26**. Here the large dark circles denote Br atom while the smaller circles represent Cu atoms.



26

In an attempt to construct a geometry similar to the one-dimensional chain **10**, here the  $\mu_4$ -Br-Cu distance is maintained at 2.6 Å while the  $\mu_2$ -Br-Cu distance is 2.43 Å. Calculations revealed that this structure was destabilized by about 10.0 eV compared to the corresponding one with a pyramidal  $\mu_4$ -Br!

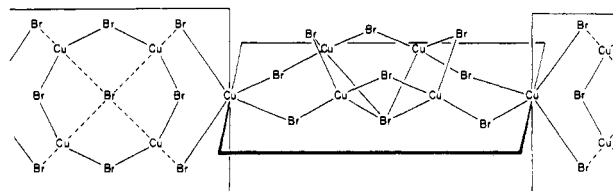
Why is this system so unstable? The total DOS, magnified projection of the orbitals of tetrahedral  $\mu_4$ -Br, and the COOP curves are shown in Figure 9 for this hypothetical polymeric chain. The Fermi level falls at -8.46 eV for this system as compared to -11.71 eV for system **10**. Comparing the COOP curves in Figures 7 and 9, it can be seen that the  $\mu_4$ -Br-Cu interaction is antibonding in the region -12.5 to -14.0 eV in the latter while it is bonding in the former. The COOP for  $\mu_2$ -Br-Cu is 0.36 while that of the tetrahedral  $\mu_4$ -Br-Cu is 0.10. These values are to be compared with the COOP of 0.20 obtained for the pyramidal  $\mu_4$ -Br-Cu and 0.35 for the  $\mu_2$ -Br-Cu for the real system, **10**. In the tetrahedral system, **26**, there is a great decrease in  $\mu_4$ -Br-Cu bonding interactions and a slight increase in Br...Br and Cu...Cu antibonding interactions.

It is clear that this hypothetical structure, constructed in an attempt to find a polymer with a tetrahedral  $\mu_4$ -Br, is under much strain. The coordination at each Cu is highly distorted, and the

set of six angles around Cu is 64.7, 106.5, 70.5, 170.0, 106.5, and 64.7, clearly very far away from either tetrahedral or square planar. We think the problem is at the copper, not the bromide. But it is hard to think of an unstrained tetrahedral Br chain with the same stoichiometry.

#### Hypothetical Structures

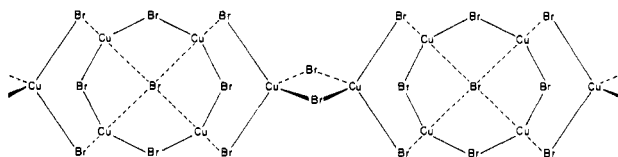
It would be nice to keep ahead of the richness that nature offers to us even within this limited realm of copper halide polymers. So we have begun to think of some hypothetical structures based on the heretofore observed geometrical features in halocuprates. Our first idea was to try to extend the structure shown for the discrete molecule **6**. Keeping the  $\mu_4$ -Br at the apex of a square pyramid and maintaining tetrahedral coordination around all the Cu atoms, one comes to **27**. Since all Cu atoms are assumed to



27

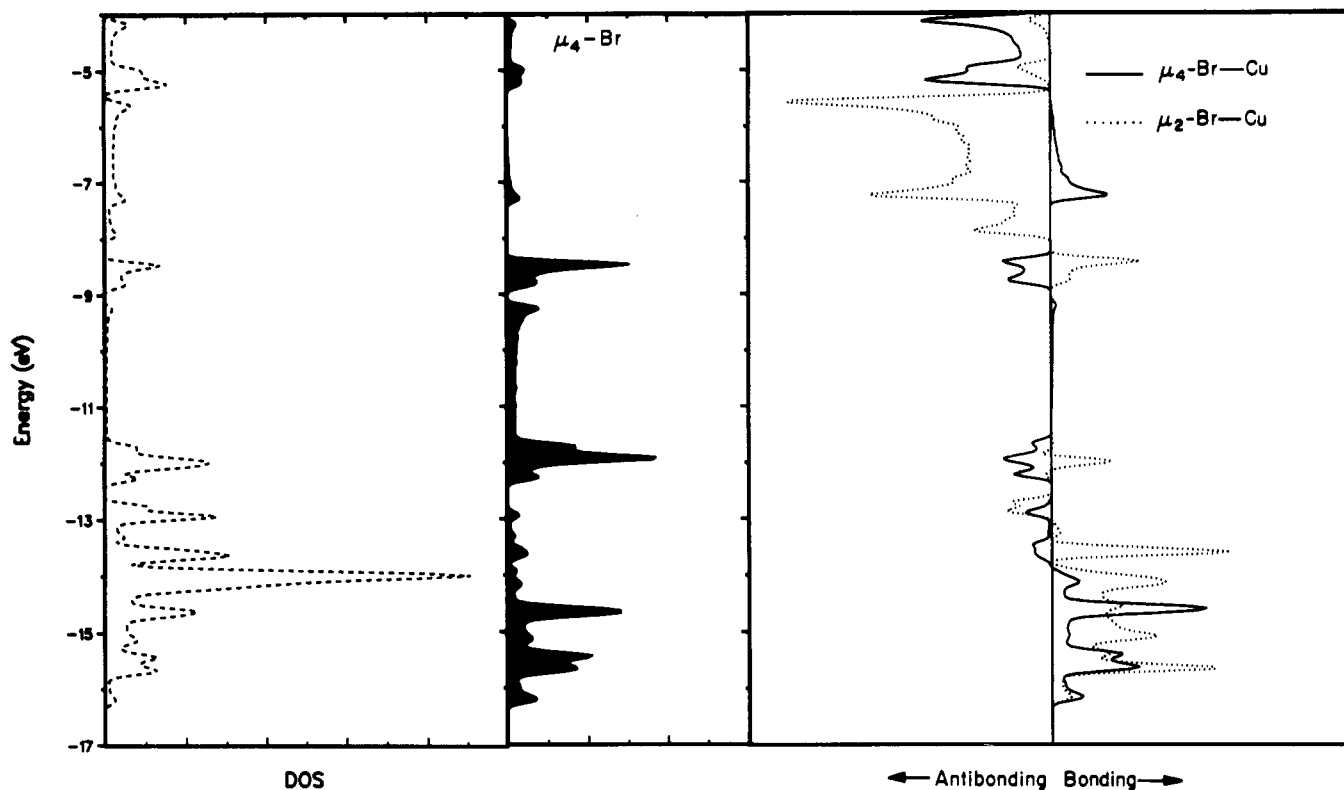
be tetrahedral, if the first unit of  $\text{Cu}_5\text{Br}_9^{4-}$  is in one plane the next unit will be in a plane perpendicular to the first.

Alternately, we could bridge two monomer units with two Br atoms as shown in **28**. The stoichiometry is  $\text{Cu}_6\text{Br}_{11}^{5-}$ .



28





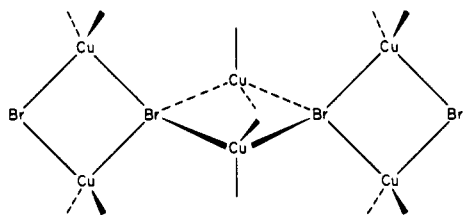
**Figure 9.** First panel: calculated total DOS for the hypothetical polymer with tetrahedral geometry around  $\mu_4$ -Br. Second panel: magnified projection of  $\mu_4$ -Br atoms. Third panel: COOP plot for the Br-Cu interactions. The  $\mu_4$ -Br-Cu interaction is represented by a solid line, and  $\mu_2$ -Br-Cu, by a dotted line.

**Table IV.** Extended Hückel Parameters

atom	orbital	$H_{ii}$ , eV	slater exponent $\zeta^a$	ref
Cu	3d	-14.00	5.95 (0.5935), 2.30 (0.5742)	51
	4s	-11.40	2.20	
	4p	-6.06	2.20	
Br	4s	-22.07	2.59	52
	4p	-13.10	2.13	

<sup>a</sup> Exponents and coefficients (in parentheses) in a double- $\zeta$  expansion of the metal d-orbitals.

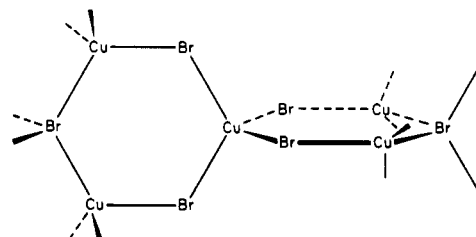
Further possibilities can be imagined for a one-dimensional structure involving a tetrahedral  $\mu_4$ -Br. One simple way this may occur is shown in **29**. However, now there are terminal Br (or other ligand) atoms attached to Cu atoms.



**29**

Another possibility is a linear polymer with six-membered rings,

**30.** This is a one-dimensional piece of the zinc blende structure of extended copper halides.



**30**

It is highly probable that the next copper halide polymeric system made will not be one of these, but a still different one. This is the fun of chemistry.

**Acknowledgment.** We thank Jane Jorgensen for her excellent drawings. Our research was supported by the National Science Foundation through Research Grant CHE-8912070.

#### Appendix

The extended Hückel parameters used in this work are shown in Table IV.

Registry No. 4, 54854-76-1; 6, 121500-82-1.

(51) Hay, P. J.; Thibeault, J. C.; Hoffmann, R. *J. Am. Chem. Soc.* **1975**, *97*, 4884.

(52) Alvarez, S.; Mota, F.; Novoa, J. *J. Am. Chem. Soc.* **1987**, *109*, 6586.