

Strong Trans Influence Methoxymethyl Ligand in B₁₂ Cobaloxime and Imine/Oxime Model Complexes: Structural, Spectroscopic, and Molecular Mechanics Investigations

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Received June 15, 1998

The CH₂OCH₃ ligand has a large trans influence comparable to that of bulky alkyl groups but lacks the complication of marked steric effects. However, crystals of model complexes with this ligand have proved difficult to obtain. The crystal structures of the cobaloximes reported here, Me₃BzmCo(DH)₂CH₂OCH₃·0.6CH₃OH (**1**) and 4-MepyCo(DH)₂CH₂OCH₃ (**2**) [DH = monoanion of dimethylglyoxime, Me₃Bzm = 1,5,6-trimethylbenzimidazole, and 4-Mepy = 4-methylpyridine], triple the number of cobaloxime structures with CH₂OCH₃. Also, these are the first structures in this class of models with an N-donor planar heterocyclic axial donor ligand, L. The planes of the Me₃Bzm and 4-Mepy ligands are almost perpendicular to the respective planes of the four equatorial DH nitrogen donors and bisect the (O⁻)N–Co–N(OH) angles. The Co–N_L bond distances average 2.096(2) Å, confirming the strong trans influence of CH₂OCH₃. Geometry optimization via molecular mechanics using MacroModel 5.0 and an AMBER-type force field was applied to both cobaloxime-type and imine/oxime-type B₁₂ models. In our initial work with some imine/oxime models, the calculated structures did not compare well to the solid-state structures. Therefore, adjustments to the force field were evaluated. The major adjustment that improved the fit of the computed and experimental structures was an ~10% reduction of the van der Waals (vdw) parameters for both the N donors and the C(sp²) atoms linked to the N donors; this adjustment may reflect the electron-withdrawing effect of the metal center. Analysis of ¹H–¹³C coupling constants of Me₃Bzm lends support to the concept that the atoms in the ligands were modified, but only slightly, by the metal center. To reproduce the dependence of geometric features on the trans influence, different force field parameters for L–Co bonds were used for compounds with weak (Cl) and those with strong (CH₂OCH₃) trans influence ligands. These small modifications allowed us to model the structural features of both classes of models well.

Introduction

Because organocobalt compounds are relevant to the chemistry and biochemistry of the vitamin B₁₂ coenzymes and because there is a large trans influence of the axial alkyl ligands, an extensive background exists on the solution and solid-state structural properties of B₁₂ model compounds.^{1–3} With the considerable range in axial bond lengths, ligand dissociation rates, etc., results from studies of organocobalt compounds are potentially useful in improving our fundamental knowledge of factors influencing the properties of metal complexes in general. Organocobalt compounds often serve as test cases for method development. One area of recent interest is molecular mechanics (MM) calculations, a tool of increasing importance for structural investigations of coordination and organometallic chemistry.^{4–19}

The widely used force fields (empirical potentials) are usually not well refined with respect to the atomic and molecular polarizability and the van der Waals (vdw) dimensions. MM suffers from incomplete information on the stretching, bending, and torsional potentials involving the bonds to the metal.

The relatively large database encompassing the diverse structural features of B₁₂ models offers considerable potential for developing force fields. An unfortunate aspect of these compounds is that the best trans influence alkyl ligands tend to be bulky (e.g., *i*-Pr), allowing steric effects to complicate the assessment of the trans influence. The CH₂OCH₃ ligand has a large trans influence comparable to *i*-Pr but lacks this steric

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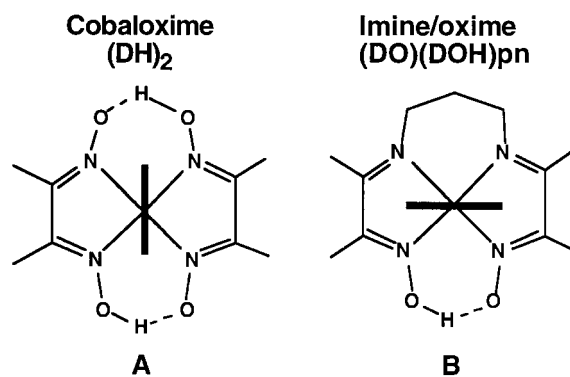
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Chart 1



complication.^{3,20–22} Although the powerful trans influence and trans effect of CH₂OCH₃ were discovered about a decade ago, most results are from the solution state since crystals of B₁₂ model systems with CH₂OCH₃ have proved difficult to obtain. In the cobaloxime class (Chart 1), the only reported structure does not have an N-donor heterocyclic axial ligand and therefore the usefulness of the data is limited.²⁰ The crystal structures of the cobaloximes reported here, Me₃BzmCo(DH)₂CH₂OCH₃·0.6CH₃OH (**1**) and 4-MepyCo(DH)₂CH₂OCH₃ (**2**) [DH = monoanion of dimethylglyoxime, Me₃Bzm = 1,5,6-trimethylbenzimidazole, and 4-Mepy = 4-methylpyridine], triple the number of structures available for this important class of models.

An important and dramatic difference between cobaloximes and the next most similar class of models, the imine/oxime class (I/O, Chart 1), is the orientation (different by 90°) of the axial N-donor heterocyclic ligand. We found in the initial phases of this work that the MM simulation with the standard original AMBER force field^{23–25} produced two minimized structures of almost equal energy for the I/O ion [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺. The structures differed by the orientation of L around the Co–N_L axis and by the orientation of the propanediyl (pn) bridge. From the experimental results (X-ray in the solid state; NMR in solution), the structures of [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺ and other I/O species^{3,26–28} have only the B orientation (Chart 1).

The initial assumption used in such calculations was that the ligand atom parameters were not altered by metal bonding. However, even for H-bonding, it has been reported that the vdw parameters (*R*^{*}, radius; *ε*^{*}, well depth) for a hydrogen with the potential for H-bonding must be smaller than those of aliphatic hydrogens.²⁹ This decrease is necessary since a significant part of the electron density of the hydrogen atom is shifted toward the heteroatom to which it is covalently linked. It is reasonable

that a similar appreciable decrease in the vdw dimensions of some ligand atoms should occur in complexes of transition metals in high oxidation states. In the MM simulations, orientation B was stabilized over orientation A by a reduction of the vdw dimensions of the donor atoms as well as the atoms directly linked to the donor atoms. This computational approach employing modified parameters has been tested by assessing how well the computed structures of other I/O complexes as well as several cobaloximes agree with the experimental structures.

Experimental Section

Materials and Methods. Me₃Bzm was prepared as previously reported.³⁰ All other reagents were purchased from Aldrich and used without further purification. H₂OC(DH)₂CH₂OCH₃ was prepared from pyCo(DH)₂CH₂OCH₃ by treatment with an acidic resin as previously described.^{20,31} Me₃BzmCo(DH)₂(*i*-Pr) was prepared as previously reported.²¹ ¹H NMR data were collected in CDCl₃ (Cambridge Isotope Laboratories, Inc.) on a GE QE-300 spectrometer and referenced to internal TMS. High-resolution one-bond proton–carbon coupling constants (¹*J*_{CH}) were obtained in CDCl₃ at 25 °C using *J*-coupled heteronuclear multiple quantum coherence spectroscopy (JHMQC)³² on a GE GN-600 Omega spectrometer.

Synthesis. Me₃BzmCo(DH)₂CH₂OCH₃ (**1**) was prepared by stirring H₂OC(DH)₂CH₂OCH₃ (0.17 g, 0.48 mmol) with Me₃Bzm (0.080 g, 0.50 mmol) overnight in MeOH (17 mL). The orange precipitate that formed was collected in 50% yield (0.16 g, 0.32 mmol), washed with H₂O and Et₂O, and recrystallized from MeOH/H₂O. ¹H NMR shifts (ppm): 18.64 (s, 2H, br), 7.98 (s, 1H), 7.96 (s, 1H), 7.08 (s, 1H), 4.35 (s, 2H), 3.75 (s, 3H), 3.19 (s, 3H), 2.36 (s, 3H), 2.35 (s, 3H), 2.09 (s, 12H). Anal. Calcd for C₂₀H₃₁N₆O₅Co·0.6 CH₃OH: C, 48.17; H, 6.55; N, 16.36. Found: C, 48.12; H, 6.38; N, 16.32.

4-MepyCo(DH)₂CH₂OCH₃ (**2**) was prepared by suspending H₂OC(DH)₂CH₂OCH₃ (0.39 g (1.11 mmol) in CH₂Cl₂ (25 mL) containing 0.12 mL (1.23 mmol) of 4-picoline. After stirring overnight, the mixture was filtered and the solvent evaporated to give 0.47 g of crude 4-MepyCo(DH)₂CH₂OCH₃. The product was recrystallized as orange prisms from MeOH/H₂O. ¹H NMR shifts (ppm): 18.36 (s, 2H, br), 8.46 (d, 2H), 7.12 (d, 2H), 4.34 (s, 2H), 3.16 (s, 3H), 2.33 (s, 3H), 2.11 (s, 12H). Anal. Calcd for C₁₆H₂₆N₅O₅Co: C, 44.96; H, 6.14; N, 16.39. Found: C, 44.73; H, 6.05; N, 16.25.

X-ray Crystallographic Data. A red crystal of **1** and a red-orange crystal of **2** were each mounted on a glass fiber and used for data collection on a Siemens P4 diffractometer. Cell parameters were calculated from the least-squares refinement of the angles of 30 (**1**) and 25 (**2**) randomly selected reflections (10 < 2θ < 30° (**1**) and 10 < 2θ < 35° (**2**)). Three standard reflections, measured every 2 h, had no significant decrease in intensity. Data were corrected for Lorentz–polarization effects and for absorption through the ψ scan technique, assuming the crystal had an ellipsoidal shape. Detailed crystallographic data are presented in Table 1.

Solution and Refinement of the Structures. The structures were solved through the Patterson and Fourier synthesis techniques and refined via full-matrix least-squares cycles. The scattering factors for neutral atoms were from SHELX 86³³ and SHELX 93.³⁴ All calculations were carried out on PC-Pentium Okidata computers using the SHELX 86,³³ SHELX 93,³⁴ and PARST 95³⁵ packages. The Co, N, C, and O

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Table 1. Selected Crystal Data and Structure Refinement^a for Me₃BzmCo(DH)₂CH₂OCH₃·0.6CH₃OH (**1**) and for 4-MepyCo(DH)₂CH₂OCH₃ (**2**)

	1	2
empirical formula	C _{20.6} H _{33.4} CoN ₆ O _{5.6}	C ₁₆ H ₂₆ CoN ₅ O ₅
fw	513.66	427.35
wavelength, Å	0.710 73	0.710 73
cryst syst	monoclinic	monoclinic
space group	<i>P</i> 2 ₁ / <i>n</i>	<i>P</i> 2 ₁ / <i>c</i>
unit cell dimens		
<i>a</i> , Å	8.1150(10)	8.398(1)
<i>b</i> , Å	15.8670(10)	29.235(3)
<i>c</i> , Å	19.266(2)	8.858(1)
β, deg	94.240(10)	115.62(1)
<i>V</i> , Å ³	2473.9(4)	1961.0(4)
ρ _{calc} , g cm ⁻³	1.379	1.448
absorption	0.738	0.912
coeff, mm ⁻¹		
<i>F</i> (000)	1076	896
cryst dimens, mm	0.40 × 0.30 × 0.40	0.30 × 0.40 × 0.15
θ range, deg	5.14 < 2θ < 50.00	5.38 < 2θ < 50.00
no. of indep reflns	4333 [<i>R</i> (int) = 0.0263]	3421 [<i>R</i> (int) = 0.0141]
final <i>R</i> 1, <i>wR</i> 2 ^b	0.0393, 0.1215	0.0434, 0.1153
[<i>I</i> > 2σ(<i>I</i>)]		
(all data)	0.0496, 0.1311	0.0555, 0.1324
largest diff. peak and hole	0.663 and -0.232 e Å ⁻³	0.610 and -0.470 e Å ⁻³

^a Details in common: 293 ± 2 K; *Z* = 4. ^b Weighting scheme as in ref 3.

atoms were refined anisotropically, whereas the H atoms were treated isotropically. The hydrogen atom of the oxime bridges (**1**, **2**) and those of the 4-Mepy ring (**2**) were located from the Fourier difference map. All the other H-atoms in **1** were set in calculated positions and allowed to ride on the atoms to which they are linked. The remaining H atoms of **2** were located in fixed positions via the HFIX option of SHELX 93. In **2** the thermal parameters of the H atoms were constrained to 1.2 times the values of the thermal parameters of the atoms to which they are bound.

A Fourier difference synthesis at the isotropic refinement stage showed two new peaks near the carbon ligand from the methoxymethyl group in **1**; this was interpreted as statistical disorder of the methoxymethyl group. The peaks were assigned the scattering factors of oxygen and carbon and labeled as O(5B) and C(10B), respectively. As refined, the site occupation factors of the atoms relevant to the two orientations of the methoxymethyl ligand were restricted to sum up to 1. The final refinement gave occupancies of 0.457(9) and 0.543(9) for the A and B orientations, respectively.

Statistical disorder was also found for the methyl groups of the Me₃-Bzm ligand. Because residual electron density was found around the carbon atoms when just one set of three atoms was considered, a second set was allowed to refine through the AFIX 123 instruction of SHELX 93.³⁴ The site occupation factors of the two sets of H atoms were refined to 0.80(2) and 0.20(2), respectively. Since the thermal parameters of the methanol C(M) and O(M) atoms in **1** were relatively high at the end of the refinement procedure, the site occupation factor of both atoms was fixed at 0.6, in agreement with the analytical data. Selected geometrical parameters for the structures are reported in Tables 2 and 3.

Molecular Mechanics. The strain energies of the complexes have been computed as the sum $E_{\text{Tot}} = E_b + E_\theta + E_\phi + E_{\text{nb}} + (E_{\text{hb}}) + (E_\epsilon)$, where E_b , E_θ , E_ϕ , E_{nb} , E_{hb} , and E_ϵ are the bond length deformation, the valence angle deformation, the torsional angle deformation, the nonbonding interaction, the hydrogen-bond interaction, and the electrostatic contributions, respectively. The E_{hb} and E_ϵ contributions have been included only when specified. The algorithm and force field ("final quality") employed were those of AMBER²³⁻²⁵ contained in MacroModel³⁶ version 5.0 (MMOD) implemented on a Silicon Graphics

Table 2. Selected Bond Lengths (Å) for Me₃BzmCo(DH)₂CH₂OCH₃·0.6CH₃OH (**1**) and 4-MepyCo(DH)₂CH₂OCH₃ (**2**)

	1	2
Co-N(1)	1.882(2)	1.888(3)
Co-N(2)	1.882(2)	1.880(2)
Co-N(3)	1.888(2)	1.880(3)
Co-N(4)	1.886(2)	1.885(3)
Co-N(5)	2.093(2)	2.099(2)
Co-C(9)	2.033(3)	2.029(3)
O(1)-N(1)	1.359(3)	1.351(4)
O(2)-N(2)	1.335(3)	1.346(3)
O(3)-N(3)	1.368(3)	1.344(4)
O(4)-N(4)	1.336(3)	1.351(4)
O(5)-C(9)		1.336(5)
O(5)-C(10)		1.437(5)
O(5A)-C(9)	1.321(6)	
O(5A)-C(10A)	1.43(2)	
O(5B)-C(9)	1.292(7)	
O(5B)-C(10B)	1.36(2)	
N(1)-C(2)	1.290(4)	1.294(4)
N(2)-C(3)	1.305(4)	1.301(4)
N(3)-C(6)	1.294(4)	1.301(5)
N(4)-C(7)	1.296(4)	1.297(5)
N(5)-C(11)	1.401(3)	1.339(4)
N(5)-C(15)		1.343(4)
N(5)-C(20)	1.319(4)	
N(6)-C(18)	1.381(4)	
N(6)-C(19)	1.458(4)	
N(6)-C(20)	1.352(4)	
C(2)-C(3)	1.457(5)	1.458(5)
C(6)-C(7)	1.465(4)	1.460(6)

Table 3. Selected Bond Angles (deg) for Me₃BzmCo(DH)₂CH₂OCH₃·0.6CH₃OH (**1**) and 4-MepyCo(DH)₂CH₂OCH₃ (**2**)

	1	2
N(1)-Co-N(2)	81.56(11)	81.06(11)
N(1)-Co-N(3)	176.75(10)	178.04(11)
N(1)-Co-N(4)	97.69(11)	99.08(13)
N(1)-Co-N(5)	90.28(10)	91.85(10)
N(2)-Co-N(3)	99.63(11)	98.33(12)
N(2)-Co-N(4)	176.82(10)	177.60(11)
N(2)-Co-N(5)	93.84(10)	91.49(10)
N(3)-Co-N(4)	80.96(10)	81.44(13)
N(3)-Co-N(5)	92.66(9)	90.03(10)
N(4)-Co-N(5)	89.25(9)	90.91(10)
N(1)-Co-C(9)	90.87(13)	84.65(13)
N(2)-Co-C(9)	85.29(14)	89.85(14)
N(3)-Co-C(9)	86.22(13)	93.49(14)
N(4)-Co-C(9)	91.63(14)	87.78(14)
N(5)-Co-C(9)	178.45(13)	176.02(12)
Co-N(1)-O(1)	123.4(2)	123.0(2)
Co-N(1)-C(2)	117.1(2)	117.0(2)
Co-N(2)-O(2)	122.2(2)	123.0(2)
Co-N(2)-C(3)	115.8(2)	117.1(2)
Co-N(3)-O(3)	122.9(2)	123.1(2)
Co-N(3)-C(6)	117.3(2)	116.8(3)
Co-N(4)-O(4)	122.7(2)	122.3(2)
Co-N(4)-C(7)	116.9(2)	116.7(3)
Co-N(5)-C(11)	132.9(2)	121.4(2)
Co-N(5)-C(15)		121.9(2)
Co-N(5)-C(20)	122.0(2)	
Co-C(9)-O(5)		119.5(3)
Co-C(9)-O(5A)	120.6(3)	
Co-C(9)-O(5B)	123.0(4)	
C(9)-O(5)-C(10)		113.4(3)
C(9)-O(5A)-C(10A)	119.1(7)	
C(9)-O(5B)-C(10B)	118.2(8)	

Indigo workstation. Modification of the force field was necessary in order to account for the interactions between the metal center and the ligand atoms. The force field parameters for the stretching and bending vibrations were obtained through a "trial and error" method, which achieved the best agreement between calculated and observed structures.

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Starting values were obtained by applying Badger's rule,^{37,38} using the parameters from the Herschbach and Laurie³⁹ (stretching) and Halgren's (bending) equations.^{40,41} The initial values for the parameters of the latter equations were based on those found for some sulfoxide compounds of ruthenium.⁴²

The new force field parameters are reported in Table 4. The total strain energy, E_{Tot} (Table 5), was minimized through the Polak-Ribiere conjugate gradient minimization method until the root-mean-square value (rms) of the first derivative vector was <0.0025 kcal/mol Å. The starting structures were those found for the solid state via single-crystal X-ray diffraction^{3,26,27,43} for [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺, [py-Co((DO)(DOH)pn)Cl]⁺, [pyCo((DO)(DOH)pn)CH₃]⁺, [pyCo((DO)(DOH)pn)CH₂OCH₃]⁺, pyCo(DH)₂Cl, Me₃BzmCo(DH)₂CH₂OCH₃, and 4-MepyCo(DH)₂CH₂OCH₃. The C(sp²)-H and C(sp³)-H bond distances for this MM investigation were all set to 1.101 and 1.113 Å, respectively, as found in the original AMBER.²³⁻²⁵

The E_e contribution was included in simulations of [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺. Differences between the energy values of the minima computed with and without the electrostatic contribution were <0.2 kcal/mol. Therefore, all the calculations were carried out in the absence of the E_e term and solvation effects. Atomic charges were evaluated by using GAUSSIAN-92/DFT.⁴⁴ Energetics, charges, and population analyses reported have been calculated using the hybrid density functional theory B3LYP method^{45,46} with the LANL2DZ basis set.^{44,47} The geometrical parameters were those found in the solid-state structure as determined by X-ray diffraction. No geometry optimization was applied. Solvent (chloroform) effects were treated via the GB/SA model⁴⁸ of MMOD for all the molecules, at the first stage of the calculations.

Results

Structural Results. The ORTEP representations⁴⁹ for Me₃-BzmCo(DH)₂CH₂OCH₃ (**1**) and 4-MepyCo(DH)₂CH₂OCH₃ (**2**) are presented in Figures 1 and 2, respectively. The cobalt atom is at the center of a distorted octahedron in both cases.

(a) Equatorial Ligand. The two DH moieties are nearly planar in both structures with a dihedral angle (α) between the planes of 2.79(7) and 3.88(9)° for the Me₃Bzm and 4-Mepy complexes, respectively. The metal atom deviates slightly (0.0502(4) and 0.0353(4) Å, respectively) from the plane of the equatorial nitrogen donors, toward N(5). The hydroxyl group of the methanol molecule in the crystal structure of **1** is linked to the O(2) atom (O(M)(0.5 + x, 0.5 - y, -0.5 + z)···O(2), 2.83(8) Å; O(M)-H···O(2), 115(5)°).

(b) Me₃Bzm and 4-Mepy Ligands. The atoms of the axial N-donor ligand in both complexes are strictly coplanar, the largest deviation being 0.029(4) Å for C(14) in **2**. The orientation of the L plane with respect to the equatorial plane can be described by the torsion angle " ϕ " defined in the systems with

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Table 4. Selected Force Field Parameters (Atom and Bond Symbols as in MMOD³⁶ and AMBER²³⁻²⁵)

stretching ^a	r_0 (Å)	k_r (kcal/(mol Å ²))	
Co-N2	1.90	150	
Co-N2 <i>trans</i> to Cl	2.00	120	
Co-N2 <i>trans</i> to CH ₂ OCH ₃	2.10	110	
Co-N2 <i>trans</i> to CH ₂ Cl	2.10	110	
Co-C3	2.00	127	
Co-Cl	2.20	105	
Cl-C3	1.77	240	
O3-C3	1.47	350	
N2-OM/O3	1.34	350	
N2=C2	1.30	450	
N2-C2	1.40	350	
C3-C3	1.53	310	
N2-C3	1.45	355	
C2-C3	1.50	335	
C2*C2 conjugate systems	1.39	469	
C2=C2	1.34	480	
H-C3	1.11	331	
H-C2	1.10	340	
H-O3	0.96	553	
O3-Lp (lone pair)	0.70	400	
bending ^a	θ_0	k_θ (kcal/(mol rad ²))	
Co-N2-OM/O3	120	35	
Co-N2*C2	120	40	
Co-N2-C3	120	30	
Co-C3-Cl	109.5	50	
Lp-O3-00	112	40	
00*Co*00 <i>trans</i>	180	10	
00*Co*00 <i>cis</i>	90	30	
torsional parameters ^c (kcal/mol)	V ₁	V ₂	V ₃
Co-N2=C2-C2	0	5.3	0
Co-N2=C2-C3	0	5.3	0
Co-N2=C2-H1	0	5.3	0
Co-N2-O3-H2	0	1.0	0
Co-N3-C3-C3	0	0	0.8
Co-N3-C3*C2	0	0.6	0.4
Co-N3-C3-H1	0	0	0.4
Lp-O3-00-00	0	0	0.07
van der Waals parameters ^d	R*	ϵ^*	
C2	1.85	0.12	
C2 ^e (linked to N donors)	1.70	0.11	
C3	1.80	0.06	
C3 ^e (linked to Co)	1.65	0.05	
N2	1.75	0.16	
N2 ^e (linked to Co)	1.55	0.12	
Co	2.35	0.35	

^a Substructure (coordination sphere and I/O or cobaloxime system). Co(-C3)(-N2)^b-N2=C2-C2=N2(-1)-C3-C3-N2(-1)=C2-C2=N2(-1)-O3-H2 Co(-C3)(-N2)^b(-N2(-OM)=C2-C2=N2(-1)-O3-H2)-N2(-OM)=C2-C2=N2(-1)-O3-H2. ^b Nitrogen donor from the axial L ligand. ^c All the potentials around the metal-donor bonds were taken equal to zero. ^d In general, the values were those of MMOD for all the atoms. Here we report the new values used in this work as well as the values from MMOD or the values used for all the other atoms of the same type. ^e New values.

L = Me₃Bzm as the torsion angle B2-N(5)-Co-N*, where N* is the midpoint between N(1) and N(4). The " ϕ " values for **1** and **2** are 3.9(2) and 4.8(3)°, respectively; these values are very close to that found for Me₃BzmCo(DH)₂(CH(CN)CH₂-CN).⁵⁰

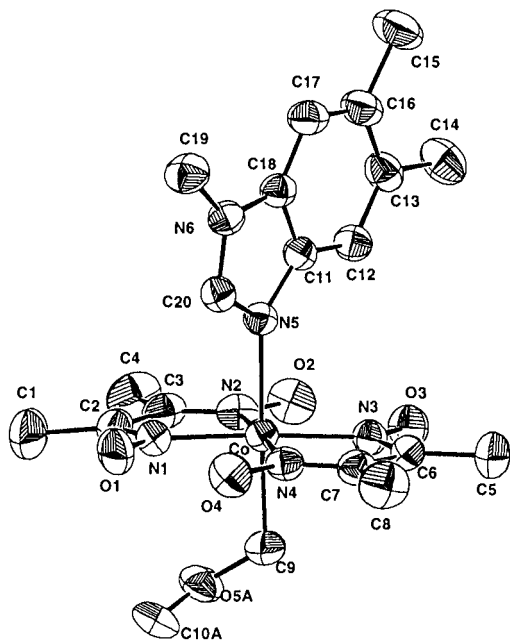
(c) Axial Fragment and Me₃Bzm or 4-Mepy Coordination. The Co-N(5) and Co-C(9) distances are 2.093(2) and 2.033-

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Table 5. Total Strain Energy (kcal/mol) for Selected Optimized Structures

molecule	structure type	
	A ^a	B ^a
[Me ₃ BzmCo((DO)(DOH)pn)CH ₂ Cl] ⁺	50.49	50.29
[pyCo((DO)(DOH)pn)CH ₂ OCH ₃] ⁺	24.22	18.44
[pyCo((DO)(DOH)pn)CH ₃] ⁺	19.51	18.61
[pyCo((DO)(DOH)pn)Cl] ⁺	21.27	19.51
Me ₃ BzmCo(DH) ₂ CH ₂ OCH ₃	43.53	46.01
4-MepyCo(DH) ₂ CH ₂ OCH ₃	13.57	15.21
pyCo(DH) ₂ Cl	11.97	14.19

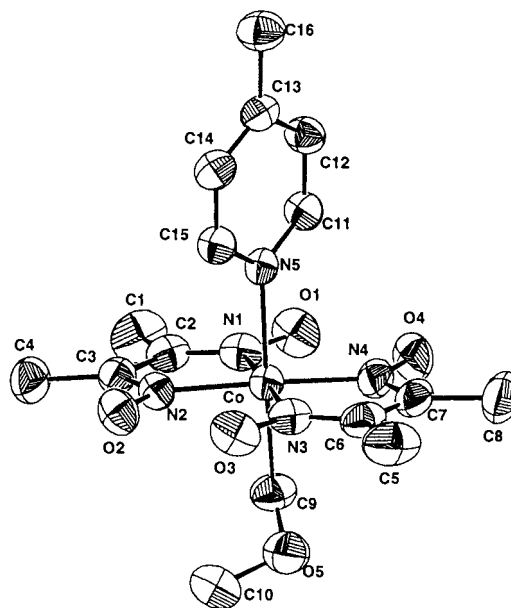
^a For I/O models, all A and B structures are *anti*-CH₂ and *syn*-CH₂, respectively.

**Figure 1.** ORTEP drawing (thermal ellipsoid; 50% probability) and labeling scheme for non-hydrogen atoms of **1**.

(3) Å for **1**, and 2.099(2) and 2.029(3) Å for **2**. The Co–C bond is significantly shorter for **1** and **2**, than for the R = *i*-Pr analogue. In contrast, the Co–N values are similar to those in the corresponding R = *i*-Pr complexes.^{21,22} NMR coupling constants obtained from the JHMQC spectra of Me₃Bzm cobaloximes have been shown to correlate well with Co–N(5) distances and are unrelated to Co–C bond distances.³² The ¹J_{CH} values for the *i*-Pr and CH₂OCH₃ complexes in solution (Table 6) are similar, especially for B2, B7, and B12, Chart 2; thus, the ¹J_{CH} values and the crystallographic results agree well.

The angles Co–N(5)–C(20) and Co–N(5)–C(11) for Me₃-BzmCo(DH)₂CH₂OCH₃ are 122.0(2) and 132.9(2)°, respectively. The angle C(20)–N(5)–C(11) is 105.1(2)°. Values previously reported for these angles in other complexes are in the ranges 120.8(4)–122.7(2), 132.2(3)–134.4(3), and 104.7(2)–105.3(3)°.²¹ The dihedral angles (γ) of the Me₃Bzm and the 4-Mepy planes with the plane of the equatorial nitrogen atoms are 87.77(6) and 87.2(1)°, respectively, showing that the metal atom lies almost in the plane of the axial ligand (deviation = 0.0457(4) and 0.0780(5) Å, respectively).

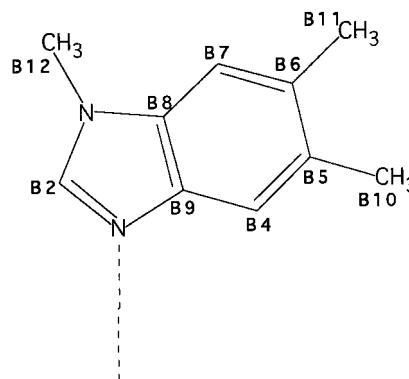
The CH₂OCH₃ moiety in Me₃BzmCo(DH)₂CH₂OCH₃ (**1**) is affected by statistical disorder, with CH₂–O(5A) almost eclipsed on Co–N(1) and CH₂–O(5B) nearly eclipsed on Co–N(4) [torsion angles O(5A)–C(9)–Co–N* and O(5B)–C(9)–Co–N* are 61.0(4) and –29.0(5)°, respectively]. The torsion angle O(5)–C(9)–Co–N* is 134.6(3)° for 4-MepyCo(DH)₂CH₂OCH₃

**Figure 2.** ORTEP drawing (thermal ellipsoid; 50% probability) and labeling scheme for non-hydrogen atoms of **2**.**Table 6.** ¹J_{CH} (Hz) for Me₃BzmCo(DH)₂R Complexes

position	Me ₃ Bzm ^a	CH ₂ OCH ₃	<i>i</i> -Pr	CH ₃ ^a	Cl ^a	Me ₃ BzmH ⁺ ^a
B2	204.3	207.8	207.9	208.2	209.9	218.6
B4	159.6	163.5	163.3	163.3	164.8	166.5
B7	158.5	159.0	159.0	159.2	160.4	163.2
B10	126.0	126.1	126.1	126.2	126.6	overlap
B11	126.1	125.8	125.7	126.0	126.5	overlap
B12	139.6	140.1	140.1	140.4	141.2	143.4

^a Previously reported values, ref 32.

Chart 2



(**2**), such that C(9)–O(5) and Co–N(3) are eclipsed. The two CH₂–O bond distances for **1** average 1.306(6) Å, whereas the O–CH₃ bond lengths average 1.40(2) Å. The CH₂–O and O–CH₃ lengths are 1.336(5) and 1.437(5) Å, respectively, for **2**. The CH₂–O bond distance is much shorter than that of a single C(sp³)–O bond (average, 1.418(16) Å⁵¹); this suggests that some shift of the oxygen lone pairs into the C–O bond occurs. A similar effect was found for some imino ether complexes of platinum(II).⁵²

Molecular Mechanics. The total strain energies obtained for various geometries are discussed below. The optimized geometries of seven compounds computed with the force field

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parameters reported in Table 4 are in good overall agreement with the solid-state structures determined by X-ray diffraction (Supporting Information). In fact, except for a few metric parameters for the often disordered methoxymethyl group, the computed and experimental bond lengths and angles differ by <0.06 Å and 5°, respectively. The quality of these fits and the force field parameters relevant to the bond stretching, bending, and torsion interactions generally agree well with those in MM studies previously reported for other metal complexes.^{4–19,53} Inclusion of the electrostatic term did not play an important role in predicting the conformation of the energy-minimized structures. In contrast, the nonbonding parameters were crucial in accounting for the stability of the B conformation with respect to the A conformation (Chart 1) as discussed below.

With the caveat that the conformation and metric parameters for the methoxymethyl group are not well defined in solution or solid states, we note, as an aside, two observations concerning this group and arising from the MM analysis of **1** and **2**. [The experimental bond distance for O–CH₃ (1.359(11) Å) found in the solid state for [pyCo((DO)(DOH)pn)CH₂OCH₃]⁺ is not comparable to the other found structures and computed structures containing the CH₂OCH₃ group because of the relatively high thermal parameter of the methyl carbon atom ($U = 0.1538 \text{ \AA}^2$) due to thermal or statistical disorder (Figure 2).³ Thus, it is not too surprising that the largest discrepancies between computed and experimental values were found for the H₂C–O(CH₃) distance (0.16 Å) of the alkyl ligand, and the H₂C–O–CH₃ and Co–C–O bond angles (7.7 and 7.3°, respectively).] First, in order for MM geometry optimization to reproduce nicely the orientation of CH₂OCH₃ in the solid state, the two lone pairs of the oxygen atom must be included in the calculation. Without the lone pairs, the minimum energy structure has the CH₂–O bond exactly bisecting the (O)N–Co–N(O) angles, and the O–CH₃ bond points away from the equatorial plane. Second, the geometry of the CH₂–O(CH₃) moiety is reproduced poorly by using the classical force field parameters for sp³-hybridized carbon and oxygen atoms, suggesting the hypothesis that the CH₂–O linkage possesses some double-bond character.

Discussion

Our structural results confirm that the cobaloximes with pyridine and benzimidazole ligands have the A-type conformation, regardless of the strength of the trans influence of the ligand opposite to the N-donor heterocyclic ligand. Consequently, these cobaloxime complexes are not so useful as the I/O complexes studied here for evaluating the values of the force field parameters. We therefore begin this discussion with the I/O computational findings. Recall that for the I/O complexes the experimental geometry found in the solid state is invariably B-type. We further note that the propanediyl (pn) bridge is flexible and that the middle CH₂ group can be either *syn* or *anti* to the L ligand with respect to the equatorial plane. These conformations will be specified as *syn*-CH₂ and *anti*-CH₂, respectively.

Our initial computations employed the vdw parameters for C, H, and N in the original AMBER force field. The total strain energy for the B-type optimized structure of [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺ is 55.07 kcal/mol. In this case, the starting geometry for the minimization was that of the solid-state structure, *syn*-CH₂-B-type. When the starting geometry was *anti*-

CH₂, the absolute minimum was represented by an *anti*-CH₂-A-type structure with a total strain energy of 52.42 kcal/mol. Simulations on other I/O derivatives carried out on the basis of the vdw parameters of original AMBER or MM2 force fields always produced results similar to those reported above for [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺: the *anti*-CH₂-A-type structure was more stable than any B-type structure.

These observations prompted us to consider the appropriateness of using the original nonbonded AMBER parameters for the donors and eventually for the atoms directly bound to the donors. We reasoned that some drift of electron density toward the metal center due to σ and possibly π donation can change the vdw dimensions (R^*) and the well depth (ϵ^*) of some ligand atoms. This concept led us to investigate modifying these force field parameters, which, in fact, are usually the most problematic in MM analysis. Both the function involved (namely Exp-6 (MM3, MM2),^{54–56} Lennard-Jones 6-9 (VFF),^{57,58} 6-12 (VFF, AMBER,^{23–25} CHARMM⁵⁹) and the R^* and ϵ^* parameters^{6,60–62} are often a matter of discussion and debate. The AMBER minimizing option of MMOD works through the 6-12 function and uses the parameters reported in the literature.^{23–25} As evidenced by Hagler et al.²⁹ and by Kollman et al.²³ the R^* value (1.00 Å) for potentially H-bonding hydrogens (N–H, O–H, S–H) is significantly smaller (~0.3 Å, 27%) than the aliphatic H atoms (1.37 Å), owing to the shift of their electron density to the heteroatom to which they are attached. A significant reduction was also applied to the values of ϵ^* (0.038 kcal/mol, aliphatic; 0.020 kcal/mol, H(N,O,S)).

When the vdw parameters (R^* and ϵ^*) of the C and N ligating atoms and the C atoms linked to these donor atoms were decreased by ~10% from the values normally used for the free ligands in the original AMBER (Table 4), E_{Tot} of the B form for I/O models decreased with respect to that of the A form. This amount of reduction in ϵ^* has been employed elsewhere.¹⁷ For example, calculations with the reduced vdw parameters gave total strain energies of 50.49 and 50.29 kcal/mol for *anti*-CH₂-A-type and *syn*-CH₂-B-type structures of [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺, respectively (Table 5). Since the A form has never been observed experimentally for I/O models, reductions of the values of R^* even greater than ~10% might be reasonable.

The relative stability of A- and B-type structures was not greatly influenced by reducing the R^* value for the metal atom below 2.35 Å. When the value was significantly increased, the geometry of the full molecule was not reproduced well. Thus the value $R^*(\text{Co})$ appears to be appropriate. Furthermore, values reported in the literature for Co and other metal complexes compare well with those used in this work.¹³

The force field developed with [Me₃BzmCo((DO)(DOH)pn)CH₂Cl]⁺ was successfully applied to [pyCo((DO)(DOH)pn)CH₂OCH₃]⁺, and the analysis showed interesting features. The X-ray structure of [pyCo((DO)(DOH)pn)CH₂OCH₃]⁺ is the *syn*-CH₂-B-type, with the CH₂–O vector projection almost bisecting the N–Co–N angle of the pn bridge. Optimization does not change significantly this structure, which converges to 18.44

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kcal/mol. Another optimized *syn*-CH₂-B-type structure (18.14 kcal/mol) has the CH₂-O vector almost bisecting the (O)N-Co-N(O) angle, the O-CH₃ vector pointing away from the equatorial system, and the pn bridge *syn* to *py*. Finally, optimizing an *anti*-CH₂-A-type structure and a *syn*-CH₂-B-type structure with a similar CH₂OCH₃ conformation [CH₂-O eclipsed with Co-N(O-) and with O-CH₃ almost parallel with the equatorial donors plane] leads to a total strain energy of 24.22 and 21.32 kcal/mol, respectively. The *anti*-CH₂-A-type and *syn*-CH₂-B-type structures have minima at 24.22 and 21.32 kcal/mol, respectively. Thus, the B-type structures are preferred for [pyCo((DO)(DOH)pn)CH₂OCH₃]⁺ also.

Solid [pyCo((DO)(DOH)pn)Cl]⁺ and [pyCo((DO)(DOH)pn)CH₃]⁺ both have an *anti*-CH₂-B-type structure.^{26,27} On optimization, the structures maintain this conformation and converge to 21.27 and 19.98 kcal/mol, respectively. Although the corresponding *anti*-CH₂-A-type forms converge to lower energies, the *syn*-CH₂-B-type structures were again most favored (Table 5). Thus, the order *anti*-CH₂-B-type > *anti*-CH₂-A-type > *syn*-CH₂-B-type appears to be quite independent of the trans influence.

The structures of **1** and **2** optimized such that the common A-type conformation of L in cobaloximes was maintained (Table 5), thus showing that our force field makes the correct predictions. Furthermore, the root-mean-square of the all-atom superimposition of the X-ray and computed structures is smaller than 0.015 Å. Finally, the structure of pyCo(DH)₂Cl converged to 11.97 and 14.19 kcal/mol for the A- and B-type structures, respectively (Table 5).

The equatorial bond lengths of the coordination sphere of cobaloximes do not differ appreciably.^{1,2} As a consequence, differences in the electronic nature and bulkiness of R and L do not influence the Co-N_{eq} bond lengths (comparisons of bond angles also show a close similarity). For example, the experimental Co-N_{eq} ranges are 1.888(3)–1.906(3) Å, for pyCo(DH)₂Cl; 1.880(2)–1.888(3) Å, for 4-MepyCo(DH)₂CH₂OCH₃; and 1.882(2)–1.888(2) Å, for Me₃BzmCo(DH)₂CH₂OCH₃. The computed Co-N_{eq} ranges were 1.913–1.935, 1.920–1.924, and 1.917–1.922 Å for the three compounds, respectively. The computed values were slightly larger (average 0.03 Å) than the experimental values. When the same analysis was carried out for the I/O structures, the average difference between computed and experimental values was even smaller (0.02 Å), an improvement we attribute to the larger experimental distances for Co-N(pn) (average, 1.914 Å) than for Co-N(O) (average 1.886 Å). Our force field uses the same values for the Co-N(sp²) bonds with no distinction between N(O) and N(pn). The equilibrium bond distance (*r*₀) of 1.90 Å (Table 4) found by using the procedure reported above therefore seems acceptable as a compromise for N(O) and N(pn)-sp² hybridized donors. For simulating cobaloximes only, better agreement was obtained when a smaller value for *r*₀ (e.g., 1.886 Å) was applied.

For the high energy B-type computed structures of the cobaloxime molecules, the butterfly distortion of the equatorial planes is, as expected, larger than that for the A-type (computed and experimental) structures; this bending results primarily from the repulsions of the H(L) atoms and the atoms of the equatorial moieties. Dihedral angles, α, and displacement from the plane of the four N_{eq} for the computed B-type molecules follow: 19.6°, 0.065 Å for Me₃BzmCo(DH)₂CH₂OCH₃; 14.5°, 0.039 Å for 4-MepyCo(DH)₂CH₂OCH₃; and 20.9°, 0.108 Å for pyCo(DH)₂Cl. These values indicate that the bending would be greatest for weak trans influence ligands and suggest that Me₃-Bzm has a greater steric size than *py*-type ligands. However,

this B-type computed structure of Me₃BzmCo(DH)₂CH₂OCH₃ has Co-N(5)-C(20) and Co-N(5)-C(11) bond angles of 121.3 and 132.3°, respectively, in excellent agreement with the values computed for the more stable A-type structure (122.1 and 131.0°). These values are similar to those found for the X-ray A-type structure (122.0 and 132.9°, respectively). One might have expected that a lopsided ligand like Me₃Bzm would display a larger dissymmetry in the Co-N_L-C(N) and Co-N_L-C(C) angles for the B- than the A-type structures for cobaloximes. For example, the increased steric interaction between H12 (i.e., B4H, Chart 2) and the equatorial plane atoms might be released by an opening of the Co-N_L-C(C) angle. This opening does not occur and it seems likely that the dissymmetry in these angles in the A-type cobaloximes already reflects a balance of steric size between the small and large parts of the lopsided Me₃Bzm ligand. In support of this concept, the Co-N_L-C(N) and Co-N_L-C(C) bond angles reported for I/O (B-type) X-ray structures are ca. 122 and 133°.^{3,63}

The ¹J_{CH} values for Me₃Bzm bound in a series of cobaloximes increase as the trans R group becomes less electron-donating, a result indicating slightly more s character contributed to the hybrid orbitals involved in the CH bond. This rehybridization is consistent with increased electron donation from Me₃Bzm to the Co. More specifically, the ¹J_{CH} changes may be interpreted in terms of an increase in the electronegativity of the C substituent (i.e., N bound to Co) favoring bonding with C hybrid orbitals with more p character.⁶⁴ The percentage of s character (ρ) in the C hybrid orbitals in CH bonds has been examined for other systems in terms of ¹J_{CH} values according to the equation ¹J_{CH} (Hz) = 500ρ.⁶⁵ For the position showing the strongest response to the trans influence of R (i.e., B2, Chart 2), we observed an increase in ρ from 40.86% for the free ligand to an average of 41.69% for the complexed ligand to 43.72% for the protonated ligand. Such changes in hybridization of the ring carbons should also be reflected as a small decrease in the vdw parameters. Reductions in R* and ε* for the B2 carbon atom in particular were found necessary in this study for the MM reproduction of experimental structures. However, although we are using very conservative reduction in parameters (the minimum needed to get the I/O B-type conformation lower in energy than the *anti*-CH₂-A-type conformation), the relatively modest changes in ¹J_{CH} suggest that this conservative approach is justified. Furthermore, we must not forget that MM methods are still simply tools and one should not expect to fine-tune parameters from one complex to the next. Indeed, the variation in the ¹J_{CH} values for complexed ligand is small, consistent with the standard assumption that this fine-tuning is not needed.

However, the use of the simplest MM model makes simulating trans influences in coordination and organometallic compounds a difficult, if not impossible, task. To account for the different trans influence of various ligands, different force field parameters for the Co-N_L bond were used for Cl, 2.00 (*r*₀, Å), 120 (*k*_r, kcal/mol), and for the R ligands, 2.10 (*r*₀, Å), 110 (*k*_r, kcal/mol) (Table 4). With this modification, the computed and X-ray Co-N_L bond distances agreed within 0.06 Å.

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In conclusion, MM analysis of I/O B₁₂ model compounds has led to the discovery of the importance of modifying the vdw parameters of the ligand atoms close to the metal. By applying a ~10% reduction of the parameters, we found that the calculations can model the experimentally observed B-type structure favored for I/O models. This discovery would not have been made if the MM calculations had been restricted to cobaloximes, for which the preferred conformation is clearly of the A-type. When implemented in MM analysis, this concept allowed us to reproduce cobaloxime structures nicely. Although we have increased substantially the number of structures with the methoxymethyl ligand, the diversity of the conformations and the disorder problem preclude a thorough analysis of bonding and the development of force field parameters to treat this ligand. Consequently, the poorest fit in our MM analysis involves this group. Nevertheless, even this poor fit could not be made without including the lone pairs on the ether oxygen atoms. The apparent insensitivity of the dissymmetry of the Co–N–C bond angles to the A- or B-type structure seems to be accommodated well by the force field. The computations further predict that there would be substantial butterfly bending in the cobaloximes if the planar L N-donor heterocyclic ligands were able to adopt the B-type conformation and that this bending would increase as the trans influence of the other axial ligand

decreased. There is no direct experimental test for this prediction. However, the prediction seems reasonable since, for I/O models, butterfly bending does seem to reflect the trans influence, and butterfly bending is small in an I/O model in which a lariat maintains the axial N-donor ligand in the A-type conformation.⁶⁶

Acknowledgment. R.C. thanks Consiglio Nazionale delle Ricerche (CNR, Roma) and Università di Siena for financial support; Centro di Calcolo at Università di Siena for (CUCES) for the use of the Silicon Graphics Indigo, and Mr. Francesco Berrettini, Centro Analisi e Determinazioni Strutturali (CIADS) at Università di Siena, for X-ray data collection. We thank Suzette M. Polson for the synthesis of the Me₃BzmCo(DH)₂CH₂OCH₃. This work was supported by NIH Grant GM 29225 (to L.G.M.). Instrument purchases were funded by the NIH and the NSF.

Supporting Information Available: Tables of complete bond lengths and angles for X-ray and computed structures, hydrogen atom coordinates, and anisotropic thermal parameters (28 pages).

IC980666I

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