# Notes

# Possible Intermediate Formation of Cyclic Nitrous Oxide, N<sub>2</sub>O: Reaction of Oxygen Difluoride, OF<sub>2</sub>, with Cesium Azide, CsN<sub>3</sub><sup>†</sup>

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#### Introduction

Recently the chemistry and the structures of several covalently bound azides of the type  $X-N_3$  have been determined.<sup>1,2a,b</sup> These highly reactive compounds, often stable at low temperatures only, have been shown to decompose to form not X and N<sub>3</sub> through cleaving of the X-N bond (eq 1) but rather, the intermediate formation of XN and N2 has been shown to occur despite being a spin-forbidden process (eq 2). Despite the relatively strong N1-N2 bond in covalent azides of the type X-N1-N2-N3 for HN<sub>3</sub>, thermal fragmentation of this molecule is not induced by breaking the H-N3 bond but rather by dissociation into HN and  $N_2$ . Since the electronic ground state of XN (X = H, halogen) is a triplet state  $({}^{3}\Sigma^{-})$ , dissociation of XN<sub>3</sub> (<sup>1</sup>A') into XN (<sup>3</sup> $\Sigma^{-}$ ) and N<sub>2</sub> (<sup>1</sup> $\Sigma^{+}_{g}$ ) is spin-symmetry forbidden. Nevertheless, pyrolysis experiments have shown that HN3 decomposes into N2 and HN in the ground-state triplet  $(^{3}\Sigma^{-})$  state.^{2c} A theoretical study  $^{2d}$  on the energetics of the dissociation reaction of HN<sub>3</sub> (<sup>1</sup>A') yielding N<sub>2</sub> (<sup>1</sup> $\Sigma^+$ <sub>g</sub>) and HN  $(^{3}\Sigma^{-})$  at a very high level of theory using CASSCF and MCSCF-CI techniques has shown that the singlet-triplet coupling occurs because the HN<sub>3</sub> wave function in the region of the transition state can be considered an equal mixture of N<sub>2</sub> (X)·NH( $a^{1}\Delta$ ) and N<sub>2</sub> (X)·NH( $b^{1}\Sigma^{+}$ ). The calculated barrier for the dissociation (35.7 kcal mol<sup>-1</sup>) is in excellent agreement with the value of 36 kcal mol<sup>-1</sup> estimated from thermal dissociation studies.  $^{2e}$  This result suggests that the stabilities of  $XN_{3}$ molecules may be determined by the activation barrier for breaking the XN-N<sub>2</sub> bond, which in turn involves singlettriplet coupling along the reaction course. The accurate calculation of the dissociation barrier of XN3 yielding XN and N2 needs a very high level of theory that is not possible for larger atoms or substituent groups X at present time. However, Frenking et

al.<sup>2f</sup> have shown that for covalently bound main-group azides the reaction energy of the dissociation reaction can be calculated at the electron-correlated MP2 level of theory.

$$XN_3 \rightarrow X + N_3 \tag{1}$$

$$XN_3 \rightarrow XN + N_2$$
 (2)

In this context, as well as the preparation of new binary azides, the investigation of their decomposition products has also been regarded as of equal importance. The often low thermal stability of binary azides which results in characterization difficulties led naturally to the use of other modern techniques such as high level quantum chemical calculations (which have been shown to exhibit a high degree of accuracy for such systems) to investigate the possible structure and decomposition pathways.<sup>2</sup> Moreover, we have for some time been interested in the elucidation of the bond properties, by high-level quantum chemical ab initio computations combined with modern spectroscopic methods.

In an earlier study we reported on the preparation and characterization of nitrosyl azide, N<sub>4</sub>O, and its unimolecular fragmentation into N<sub>2</sub>O and N<sub>2</sub>. It was found through a combined quantum chemical study and experimentally observed gas-phase IR decomposition products that *trans-trans*-N<sub>4</sub>O decomposes via rotation into the *cis-cis* isomer and cyclic N<sub>4</sub>O to yield N<sub>2</sub>O ( $C_{\infty v}$ ) and N<sub>2</sub>.<sup>2-5</sup> Through this combination of theory and experiment the formation of a new binary isomer of N<sub>2</sub>O through this reaction was deemed highly unlikely. Therefore the ongoing quest for cyclic N<sub>2</sub>O continued due to the observation of trans-cis isomerization of the N<sub>4</sub>O moiety.

### **Results and Discussion**

In the present study we now report about the reaction of oxygen difluoride,  $OF_2$ , with cesium azide,  $CsN_3$ , with the aim of investigating the decomposition pathway with respect to the formation of cyclic  $N_2O$  as the observation and characterization of a binary oxygen diazide,  $O(N_3)_2$  would be expected only at extremely low temperatures. It was found that the reaction of  $OF_2$  with NaN<sub>3</sub> or CsN<sub>3</sub> did indeed yield dinitrogen and nitrous oxide according to eq 3 in the molar ratio of roughly 1:2. Both gaseous products were unambiguously identified by IR spectroscopy (N<sub>2</sub>O) and the typical gas discharge color of N<sub>2</sub> which is volatile at -196 °C.<sup>6</sup> Neither O<sub>2</sub> or NOF nor NO or NO<sub>2</sub> were detected as reaction products.

$$OF_2 + 2CsN_3 \xrightarrow{\Delta T} N_2O + 2N_2 + 2CsF$$
 (3)

However, the undertaking of this experiment was hoped to show the formation of  $N_2O$  and  $N_2$  which would then pose the interesting question of the decomposition pathway which was occurring. To explain the formation of nitrous oxide (NNO) from the reaction of  $OF_2$  with  $CsN_3$  we propose the intermediate

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<sup>&</sup>lt;sup>†</sup> Part of this work has already been presented as a lecture at the ACS 14th Winter Fluorine Conference, St. Petersburg Beach, FL, January 17–22, 1999. [1 kcal mol<sup>-1</sup> = 4.184 kJ mol<sup>-1</sup>, 1 Å =  $10^{-10}$  m (since these units are still used in computational chemistry we employed them throughout the paper).]

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Table 1. Structure, Energy, and Vibrational Frequencies of O(N<sub>3</sub>)<sub>2</sub> Computed ab Initio at Various Levels of Theory

	$\mathrm{HF}^{a}$	$\mathrm{HF}^{a}$	$MP2^{a,c}$	B3LYP <sup>a</sup>	$CCD/D95V(d)^d$
symmetry	$C_1$	$C_2$	$C_2$	$C_2$	$C_2$
— <i>E/</i> au	401.353320	401.353320	402.594517	403.5792531	402.451087
<i>NIMAG</i>	0	0	0	0	0
<i>zpe/</i> kcal mol <sup>-1</sup>	18.8	18.8	17.0	16.5	17.2
$ \begin{array}{l} d(01-N2)/\text{\AA} \\ d(N2-N4)/\text{\AA} \\ d(N4-N6)/\text{\AA} \\ \angle (N2-01-N3)/\text{deg} \\ \angle (01-N2-N4)/\text{deg} \\ \angle (N2-N4-N6)/\text{deg} \\ \angle (N3-01-N2-N4)/\text{deg} \\ \angle (01-N2-N4-N6)/\text{deg} \end{array} $	1.393	1.393	1.447	1.452	1.446
	1.261	1.261	1.275	1.260	1.293
	1.088	1.088	1.145	1.128	1.134
	111.2	111.2	109.3	110.2	108.8
	106.5	106.5	106.8	107.3	105.0
	174.2	174.2	172.0	172.0	173.9
	88.7	88.7	82.5	89.8	86.8
	172.1	172.1	177.3	172.5	176.0
$v_1/cm^{-1}$ (rel. int.) <sup>b</sup> $v_2/cm^{-1}$ (rel. int.) <sup>b</sup> $v_3/cm^{-1}$ (rel. int.) <sup>b</sup> $v_4/cm^{-1}$ (rel. int.) <sup>b</sup> $v_5/cm^{-1}$ (rel. int.) <sup>b</sup> $v_6/cm^{-1}$ (rel. int.) <sup>b</sup> $v_7/cm^{-1}$ (rel. int.) <sup>b</sup> $v_8/cm^{-1}$ (rel. int.) <sup>b</sup> $v_9/cm^{-1}$ (rel. int.) <sup>b</sup> $v_{10}/cm^{-1}$ (rel. int.) <sup>b</sup> $v_{12}/cm^{-1}$ (rel. int.) <sup>b</sup> $v_{13}/cm^{-1}$ (rel. int.) <sup>b</sup> $v_{14}/cm^{-1}$ (rel. int.) <sup>b</sup> $v_{15}/cm^{-1}$ (rel. int.) <sup>b</sup>	$\begin{array}{c} 50 \ (2) \\ 89 \ (2) \\ 226 \ (0) \\ 333 \ (6) \\ 499 \ (2) \\ 598 \ (7) \\ 610 \ (26) \\ 767 \ (1) \\ 827 \ (3) \\ 1034 \ (9) \\ 1036 \ (355) \\ 1097 \ (133) \\ 1124 \ (2) \\ 2421 \ (1194) \\ 2446 \ (27) \end{array}$	$\begin{array}{c} 50 \ (2) \\ 89 \ (2) \\ 226 \ (0) \\ 333 \ (6) \\ 499 \ (2) \\ 598 \ (7) \\ 610 \ (26) \\ 767 \ (1) \\ 827 \ (3) \\ 1034 \ (9) \\ 1036 \ (355) \\ 1097 \ (133) \\ 1124 \ (2) \\ 2421 \ (1194) \\ 2446 \ (27) \end{array}$	$\begin{array}{c} 47 \ (0.5) \\ 76 \ (0.5) \\ 183 \ (0) \\ 264 \ (3) \\ 423 \ (2) \\ 483 \ (2) \\ 508 \ (8) \\ 679 \ (31) \\ 698 \ (0.5) \\ 768 \ (16) \\ 859 \ (0.5) \\ 1101 \ (114) \\ 1125 \ (6) \\ 2327 \ (623) \\ 2347 \ (34) \end{array}$	$\begin{array}{c} 46 \ (0.5) \\ 72 \ (14) \\ 190 \ (14) \\ 251 \ (5) \\ 416 \ (3) \\ 533 \ (3) \\ 547 \ (10) \\ 636 \ (54) \\ 666 \ (0.5) \\ 746 \ (2) \\ 862 \ (2) \\ 1097 \ (229) \\ 1125 \ (4) \\ 2168 \ (861) \\ 2190 \ (16) \end{array}$	$\begin{array}{c} 45 \ (1) \\ 81 \ (1) \\ 196 \ (0) \\ 291 \ (5) \\ 440 \ (2) \\ 540 \ (2) \\ 550 \ (11) \\ 688 \ (0) \\ 726 \ (1) \\ 863 \ (14) \\ 917 \ (1) \\ 1068 \ (242) \\ 1100 \ (8) \\ 2256 \ (943) \\ 2282 \ (31) \end{array}$

 ${}^{a}$  6-311++G(d) basis.  ${}^{b}$  Rel. IR int. in km mol<sup>-1</sup>.  ${}^{c}$  -*E*[MP4(SDTQ)//MP2]<sup>*a*</sup> = 402.669763 au.  ${}^{d}$  -*E*[CCSD(T)/D95V(d)//CCD/D95V(d)] = 402.522973 au.



Figure 1. Ab initio computed  $C_2$  structure of  $O(N_3)_2$  (cf. Table 1).

formation of  $O(N_3)_2$  which decomposes into  $N_2O$  and two equivalents of  $N_2$ . The structure, energy and vibrational frequencies of the hitherto unknown  $O(N_3)_2$  were computed ab initio at various levels of theory (Table 1).<sup>7</sup> Initially calculations were performed without symmetry constraints but in all cases the structure was optimized to  $C_2$  symmetry (Figure 1).<sup>7</sup> It can be stated that the molecule  $O(N_3)_2$  in  $C_2$  symmetry possesses a true minimum (*NIMAG* = 0) on its potential energy hypersurface and is vibrationally stable. Therefore  $O(N_3)_2$  is an intrinsically stable but elusive molecule since the overall decomposition reaction according to eq 4 is thermodynamically highly favorable. The calculated total energies of  $O(N_3)_2$ ,  $N_2O(C_{\infty v})$ , and  $N_2$  can be used to predict theoretically the energy of reaction 4. After correction for zero point energies (*zpe*,  $\Delta U^{\text{vib}} = -4.2$ kcal mol<sup>-1</sup>), rotational ( $\Delta U^{\text{rot}} = +3/2RT$ ), translational correction ( $\Delta U^{\text{tr}} = +3RT$ ) and the work term ( $p\Delta V = +2RT$ ), the reaction

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<sup>(6)</sup> Caution: Covalent azides are usually highly explosive and often are very toxic. ClF<sub>3</sub>, F<sub>2</sub>, and OF<sub>2</sub> are strong oxidizers. Although no explosion was observed in our studies, all covalent azides and high oxidizers should be handled with respect and fully protective clothing should be worn at all times. A stainless steel (T316SS) reaction vessel was pre-passivated with ClF3 (1 bar, 100 °C, 12 h) and loaded in the drybox with 1.05 g (6.0 mmol) of dry CsN<sub>3</sub> (Dynamite Nobel). The vessel was then cooled to -196 °C, all N<sub>2</sub> was pumped off and 0.11 g (2.0 mmol) OF2 were condensed onto the CsN3. The reaction vessel was then warmed to room temperature and heated for 15 min. to 250 °C. After this the reaction vessel was cooled to -196 °C and all products volatile at this temperature were pumped off and identified by their typical gas-discharge color. N2 was identified as the only volatile reaction product at this temperature. Then the reaction vessel was allowed to warm to room-temperature again and a gas-phase IR spectrum was recorded [gas IR (2 Torr, 25 °C, 10 cm cell,  $\nu$  in cm<sup>-1</sup>): 3480 m (NNO), 2565 m ( $2 \times \nu_3$ , NNO), 2460 w, 2225 vs ( $\nu_1$ , NNO), 1285 s (v<sub>3</sub>, NNO), 825 m (v<sub>3</sub>, FOF), 590 w (v<sub>2</sub>, NNO)]. Finally all volatile products were pumped off and the steel vessel was taken into the drybox and the remaining solid material identified as a mixture of unreacted CsN3 (excess, see above) and CsF (F- in CsF was detected after the residue was dissolved in water using a fluoride sensitive electrode). Qualitatively the same results were obtained when NaN3 or Me<sub>3</sub>SiN<sub>3</sub> were used as the azide compound.

<sup>(7)</sup> The structures of all compounds considered were "pre"-optimized using semiempirical calculations. All semiempirical calculations were carried out with the program package HyperChem8 at the semiempirical PM39 level of theory using a VSTO-3G\* basis set. Finally, the structures, energies, vibrational frequencies and zero point energies were computed ab initio with the program package Gaussian 94.10 All structures were fully optimized at the HF and electron correlated MP2 and CCD (using double substitutions from the Hartree-Fock determinant)<sup>11</sup> as well as at the DFT hybrid B3LYP level of theory.<sup>12</sup> In addition, the energy of O(N<sub>3</sub>)<sub>2</sub> at the MP2 structure was computed at the MP4(SDTQ) level of theory and in addition at the CCD structure at the CCSD(T) level. For all HF, MPn and B3LYP calculations a 6-311++G(d) basis set was used.<sup>13</sup> The coupled cluster calculations were performed with a Dunning/Huzinaga valence double- $\zeta$  basis with the addition of a set of d-type polarization functions to each atom (D95V(d))14 whereas for computing the potential energy hypersurface at the MP2 level of theory a 6-31+G(d) basis set was used.

**Table 2.** Comparison of Different Isomers of  $N_2O$ , ab Initio Calculated at the MP2/6-311++G(d) Level of Theory

	linear NNO $(C_{\infty v})$	linear NNO $(C_{\infty v})$ , exptl	$\begin{array}{c} \text{cyclic NON} \\ (C_{2v}) \end{array}$	linear NON $(D_{\infty h})$
-E/au	184.290030		184.169357	184.183851
$E_{\rm rel}/\rm kcal\ mol^{-1}$	0.0		+75.7	+66.6
d (NN)/Å	1.163	1.12615	1.214	
d (NO)/Å	1.181	$1.186^{15}$	1.511	1.342
∠(NON)/deg			47.4	
∠(ONN)/deg			66.4	
NIMAG <sup>a</sup>	0	0	0	0
<i>zpe</i> /kcal mol <sup>-1</sup>	6.58	6.70	4.13	5.26
$\hat{\nu}_1/\text{cm}^{-1}$ (IR int.) <sup>b</sup>	2243 (366)	2225 s <sup>16</sup>	966 (16)	2398 (2647)
$\nu_2/\text{cm}^{-1}$ (IR int.) <sup>b</sup>	538 (3)	$590 \text{ w}^{16}$	287 (3)	368 (11)
$\nu_3/cm^{-1}$ (IR int.) <sup>b</sup>	1286 (11)	1285 s <sup>15,16</sup>	1636 (3)	546 (0)

<sup>*a*</sup> NIMAG = number of imaginary frequencies. <sup>*b*</sup> IR intensities are given in km mol<sup>-1</sup>.



Figure 2. Two-dimensional potential energy hypersurface representing the unimolecular decay of  $O(N_{3})_2$  to yield cyclic  $N_2O$  and  $N_2$  computed at MP2/6-31+G(d) level of theory.

enthalpy of the gas-phase reaction according to eq 4 was calculated at the optimized MP2/6-311++G(d) geometries to be  $\Delta H^{\circ}_{298}(2) = -187.7$  kcal mol<sup>-1</sup>.

$$O(N_3)_2(C_2) \rightarrow N_2 O(C_{\infty v}) + 2N_2(D_{\infty h})$$
(4)

However, according to the gas-phase IR data<sup>6</sup> the experimentally observed nitrous oxide has undoubtedly a N–N–O ( $C_{\infty\nu}$ ) (and not a N–O–N ( $D_{\infty h}$ )) connectivity. We therefore propose for the unimolecular decomposition in the gas phase a dissociation of O(N<sub>3</sub>)<sub>2</sub> into 2 equiv of N<sub>2</sub> and initially formed cyclic N<sub>2</sub>O ( $C_{2\nu}$ ) which then rearranges to give the more stable linear isomer

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**Figure 3.** Relative stabilities of  $O(N_3)_2$  with respect to  $N_2$  and various oxides of nitrogen computed at MP2/6-311++G(d) level of theory. For comparison, the experimentally measured relative stabilities are  $E^{rel}(NNO + 2N_2) = 0.00$  kcal mol<sup>-1</sup>;  $E^{rel}(NO + \frac{5}{2}N_2) = 1.97$  kcal mol<sup>-1</sup>; and  $E^{rel}(NO + N_2 + N_3) = 106.97$  kcal mol<sup>-1.17</sup>

N=N=O ( $C_{\infty v}$ ) (Table 2). The energies of activation separating cyclic NON ( $C_{2v}$ ) from (i) linear NON ( $D_{\infty h}$ ) and (ii) linear NNO ( $C_{\infty v}$ ) have also been calculated. At MP2/6-31+G(d) level the barrier for the reaction from cyclic NON ( $C_{2v}$ ) to form linear NON ( $D_{\infty h}$ ) is 53 kcal mol<sup>-1</sup> whereas it is only 26 kcal mol<sup>-1</sup> for the reaction from cyclic NON ( $C_{2v}$ ) to form linear NNO ( $C_{\infty v}$ ). However, it has to be stressed that these energy potentials have been calculated for the singlet state only and that the triplet state lies energetically very close, especially for the high energy regions on the reaction pathway for the interconversion of cyclic NON into linear NNO or NON, respectively.

This proposition is an interesting and reasonable one when it is remembered that as was stated above that azides decompose initially into XN and N<sub>2</sub>. Therefore, the proposed intermediately formed oxygen diazide would be expected to decompose according to eq 5. The formation of linear N<sub>2</sub>O it was thought most likely to occur through a cyclization forming cyclic N<sub>2</sub>O which would then from the more stable linear N<sub>2</sub>O which was observed as the main reaction product along with N<sub>2</sub>.

$$O(N_3)_2(C_2) \rightarrow NON(C_{2\nu}) + 2N_2$$
(5)

To account for the interconversion of  $O(N_3)_2$  into  $N_2$  and cyclic N<sub>2</sub>O we calculated the corresponding two-dimensional energy hypersurface at MP2/6-31+G(d) level (Figure 2). Figure 2 clearly shows that the molecule  $O(N_3)_2$  represents a local minimum at this hypersurface and that cyclic  $N_2O(C_{2\nu})$  and two equivalents of  $N_2$  are already 111.6 kcal mol<sup>-1</sup> lower in energy (MP2/6-311++G(d), Figure 3). The experimentally observed products  $2N_2$  and linear  $N_2O(C_{\infty v})$ , however, lie again 75.7 kcal mol<sup>-1</sup> below cyclic N<sub>2</sub>O ( $C_{2v}$ ) and two equivalents of N<sub>2</sub> and therefore 187.3 kcal mol<sup>-1</sup> below the initial compound  $O(N_3)_2$ . The relative stabilities of  $O(N_3)_2$  with respect to  $N_2$  and various oxides of nitrogen are summarized in Figure 3 (MP2/ 6-311++G(d)). It is noteworthy that the linear symmetric isomer of N<sub>2</sub>O, i.e. NON  $(D_{\infty h})$  has almost the same total energy as cyclic  $C_{2\nu}$  N<sub>2</sub>O [ $\Delta E$ (MP2/6-311++G(d)) = 9.1 kcal mol<sup>-1</sup> in favor of linear NON in  $D_{\infty h}$  symmetry] (Figure 3, Table 2). The intermediate formation of this isomer, however, would not explain the final formation of the stable NNO  $(C_{\infty v})$ .

From the present study the following conclusions can be drawn:

(i) OF<sub>2</sub> reacts with CsN<sub>3</sub> to yield the experimentally observed products dinitrogen (N<sub>2</sub>) and nitrous oxide (linear N<sub>2</sub>O,  $C_{\infty\nu}$  symmetry).

(ii) The formation of linear N<sub>2</sub>O as the product of the reaction of OF<sub>2</sub> with CsN<sub>3</sub> can best be explained by the intermediate formation of O(N<sub>3</sub>)<sub>2</sub> which then decomposes in an unimolecular reaction to yield N<sub>2</sub> and cyclic N<sub>2</sub>O ( $C_{2v}$ ) the latter of which then rearranges to form the more stable linear isomer of N<sub>2</sub>O ( $C_{\infty v}$ ).

(iii) The still elusive molecule  $O(N_3)_2$  represents a true minimum in  $C_2$  symmetry on its potential energy hypersurface at all levels of theory applied.

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