

# Synthesis and Crystal Structure of Zr<sub>2</sub>Te. Distinctions in Bonding to Isotypic Sc<sub>2</sub>Te and the Relationship to the Structures of Congeneric Hf<sub>2</sub>Te and Zr<sub>2</sub>Se

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Zr<sub>2</sub>Te is accessible by high-temperature synthesis. The structure of the zirconium-rich telluride was determined by means of powder X-ray diffraction to be orthorhombic, *Pnma* (No. 62), *Z* = 12, Pearson symbol *oP36*, *a* = 1995.0(2) pm, *b* = 382.36(2) pm, *c* = 1065.63(9) pm. Pairwise interpenetrating columns of *trans*-face-shared, centered Zr<sub>9</sub> cuboids, reminiscent of the bcc high-temperature form of zirconium can be recognized as the topologically characteristic structural feature. Tellurium atoms capping the remaining square faces complete the motif of a  ${}^1_{\infty}[\text{Zr}_8\text{Te}_4]$  double string running parallel [010]. The tellurium atoms are 7-, 8- and 9-fold coordinated by zirconium. The coordination figures represent mono-, bi- and tricapped distorted trigonal prisms, with zirconium atoms capping the square faces of the prisms. Extended Hückel calculations revealed distinctions in bonding in Zr<sub>2</sub>Te and the isotypic Sc<sub>2</sub>Te. According to Mulliken overlap populations, the heteronuclear interactions are similar in both tellurides. However, the lower valence electron concentration available for M–M bonding in Sc<sub>2</sub>Te is reflected in a considerable restriction of the attractive homonuclear interactions to one-dimensional metal cores, whereas in Zr<sub>2</sub>Te M–M bonding regions extend in space. The structure of Zr<sub>2</sub>Te is contrasted with two other types of bcc fragment structures adopted by the congeneric Hf<sub>2</sub>Te and Zr<sub>2</sub>Se. We show that the structural diversity observed for various dimetal chalcogenides is controlled by an intimate interplay of electronic and geometric factors.

## Introduction

In the past, a substantial number of binary metal-rich chalcogenides of the valence-electron-poor transition metals have been synthesized and characterized, e.g., Ti<sub>8</sub>S<sub>3</sub>,<sup>1</sup> Zr<sub>2</sub>S,<sup>2,3</sup> Hf<sub>2</sub>Se,<sup>2</sup>  $\alpha$ - and  $\beta$ -V<sub>3</sub>S,<sup>4</sup> Nb<sub>2</sub>Se,<sup>5</sup> and Ta<sub>2</sub>S.<sup>6</sup> The number of identified compounds in this category continues to grow. More recently discovered phases are, e.g., Sc<sub>8</sub>Te<sub>3</sub>,<sup>7</sup> Sc<sub>2</sub>Te,<sup>8</sup> Y<sub>8</sub>Te<sub>3</sub>,<sup>7</sup> Ti<sub>9</sub>Se<sub>2</sub>,<sup>9</sup> Ti<sub>11</sub>Se<sub>4</sub>,<sup>10</sup> Ti<sub>8</sub>Se<sub>3</sub>,<sup>11</sup>  $\beta$ -Ti<sub>2</sub>Se,<sup>12</sup> Zr<sub>3</sub>Te,<sup>13</sup> Hf<sub>2</sub>Te,<sup>14</sup> Hf<sub>3</sub>Te<sub>2</sub>,<sup>15</sup> Ta<sub>2</sub>Se,<sup>16</sup> Ta<sub>97</sub>Te<sub>60</sub>,<sup>17,18</sup> Ta<sub>181</sub>Te<sub>112</sub>,<sup>17,18</sup> and dodecagonal

(dd) Ta<sub>1.6</sub>Te.<sup>19,20</sup> Of these, the phases Sc<sub>2</sub>Te, Ti<sub>9</sub>Se<sub>2</sub>, Ti<sub>11</sub>Se<sub>4</sub>, Hf<sub>3</sub>Te<sub>2</sub>, Ta<sub>2</sub>Se, and the tantalum tellurides account for new, previously unknown structure types. Sc<sub>2</sub>Te, Ti<sub>9</sub>Se<sub>2</sub>, and Ti<sub>11</sub>Se<sub>4</sub> have as a common structural feature strings of M<sub>6</sub> octahedra connected in various ways according to the concept of condensed clusters.<sup>21</sup> Hf<sub>3</sub>Te<sub>2</sub> and Ta<sub>2</sub>Se represent layered materials with bcc-like metal stacks coated on both sides with chalcogen atoms. The Te-3Hf-Te and Se-4Ta-Se layers, respectively, are held together by weak interactions. dd-Ta<sub>1.6</sub>Te is the first identified quasicrystalline chalcogenide; Ta<sub>97</sub>Te<sub>60</sub> and Ta<sub>181</sub>Te<sub>112</sub> are its crystalline approximants. The M<sub>8</sub>Te<sub>3</sub> compounds of the third group elements Sc and Y<sup>7</sup> possess a particularly low number of valence electrons available for metal–metal bonding and are among the most metal-rich chalcogenides reported for the valence-electron-poor transition metals. Compared to the isostructural but valence-electron-richer Ti<sub>8</sub>Q<sub>3</sub> (Q = S, Se) phases,<sup>1,11</sup> distinctive differences in structure as well as chemical bonding were observed. The underlying reasons were discussed, emphasizing cooperative matrix and bonding effects based on differences in atomic sizes and valence electron concentration.<sup>7</sup> In the present paper we report the crystal and electronic structures of the new subtelluride Zr<sub>2</sub>Te,<sup>22</sup> which crystallizes isostructurally with Sc<sub>2</sub>Te. Since both compounds are tellurides and the atomic radii of zirconium and scandium are very close to one another so as to minimize matrix effects, the comparison of the bonding situation in Zr<sub>2</sub>Te and Sc<sub>2</sub>Te becomes especially instructive.

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Apart from the metal-rich phases  $Zr_2Te$  and  $Zr_3Te$ ,<sup>13</sup> the latest investigations in the Zr–Te system have brought to light  $Zr_{1.30}Te_2$ ,<sup>23</sup> and the recently discovered  $Zr_5Te_6$ ,<sup>24</sup> adopting a NiAs type-related structure.

### Experimental Section

**Preparation.** Due to their air-sensitivity, the tellurides were handled and stored under argon. The first indications of a new phase intermediate between  $Zr_5Te_4$ <sup>25–27</sup> and  $Zr_3Te$ <sup>13</sup> were obtained in arc-melting experiments: in as-cast samples of composition  $0.71 < x_{Zr} < 0.79$  an additional phase, now identified as  $Zr_2Te$ , was present along with Zr and  $Zr_3Te$ .<sup>13</sup>  $Zr_2Te$  could also be synthesized through the reduction of  $Zr_5Te_4$  with elemental Zr in the solid state (1170–1600 K), although not in a pure form. A better enrichment of the phase (>90%) was achieved by inductive heating, starting with the elements (Te, 99.999%, Fluka; Zr, 99.9%, Ventron,  $n_{Zr}:n_{Te} = 2:1$ , in total ca. 300 mg, 5 min), in sealed, argon-filled Ta tubes at 1970 K in a vacuum ( $P < 10^{-3}$  Pa), with subsequent quenching by radiative heat loss.

Energy dispersive analysis of emitted X-rays (CamScan CS 4DV, EDX system, Noran Instruments; 30 kV, Zr-L, Te-L, Ta-L) on two different samples (six measurements) gave the composition  $Zr_{2.0(1)}Te$ . In addition, small amounts of Ta were detected, with the mole fraction observed in separate measurements on different as well as identical probes strongly fluctuating ( $\langle x_{Ta} \rangle = 0.013(0.008)$ ). Furthermore, a reciprocal relationship was observed between the acceleration voltage and the detected Ta contents. Thus, we assume the detected Ta to be a surface phenomenon. Providing further support for this assumption are the lattice parameters of  $Zr_2Te$  from samples where Ta contamination can be ruled out, which do not systematically deviate from those of  $Zr_2Te$  prepared in Ta tubes.

**Powder X-ray Diffraction.** Guinier X-ray powder diffraction was used as a characterization method for phase identification. Guinier photographs were obtained with a Huber Guinier System 600 employing  $Cu\ K\alpha_1$  radiation. Silicon<sup>28</sup> was added to the samples as an internal standard. Lattice parameters were determined from Guinier diffraction data by least-squares refinement. X-ray powder diffractograms were recorded at ambient temperature on a computer-automated diffractometer (Siemens, D500), equipped with a diffracted-beam focusing graphitic monochromator, using  $Cu\ K\alpha_1$  and  $K\alpha_2$  radiation (50 kV, 30 mA). Intensities were detected by means of a scintillation counter. The data were processed with the local program DIFFRAKT.<sup>29</sup> Indexing of the diffraction pattern of  $Zr_2Te$  was accomplished with use of the program TREOR.<sup>30</sup> The lattice parameters and diffraction intensities suggested a phase isotopic with  $Sc_2Te$ .<sup>8</sup> As no single crystals of  $Zr_2Te$  could be obtained, the structural parameters were refined from a fit of the diffraction profile of a sample containing less than 2.5% weight fraction each of  $Zr_{1-x}Te_x$ , Ta, and Te, determined according to Hill and Howard.<sup>31</sup> The air-sensitive, reflection-rich  $Zr_5Te_4$ , initially present in the sample as a minor component, was allowed to decompose under atmospheric conditions into crystalline Te and an X-ray amorphous phase. The refinements were carried out using the modified Rietveld version LHPM,<sup>32–34</sup> applying a pseudo-Voigt profile fit. Atomic positions of the minority components were obtained from the literature (Zr and Ta,<sup>35</sup> Te<sup>36</sup>). The lattice parameters for  $Zr_2Te$ ,  $Zr_{1-x}Te_x$ , and Ta were obtained from Guinier data. The lattice parameters of  $Zr_2Te$  and those of the isostructural compounds  $Sc_2Te$ <sup>8</sup> and  $\beta$ - $Ti_2Se$ <sup>12</sup> are given

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**Table 1.** Comparison of Lattice Parameters (pm) and Cell Volumes ( $10^6\text{ pm}^3$ ) for Known Phases Crystallizing in the  $Sc_2Te$  Structure Type

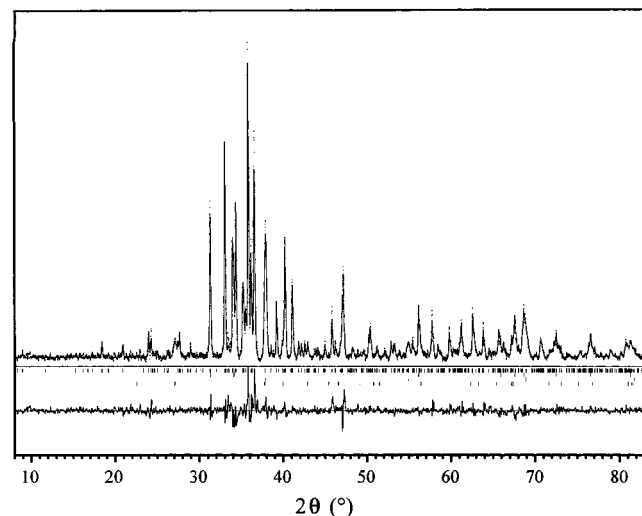
compd	a	b	c	V
$\beta$ - $Ti_2Se$ <sup>a</sup>	1793.4(20)	345.3(1)	952.6(10)	590(2)
$Zr_2Te$	1995.0(2)	382.36(2)	1065.63(9)	812.9(2)
$Sc_2Te$ <sup>b</sup>	2017.8(5)	391.86(7)	1067.5(2)	844.1(5)

<sup>a</sup> Reference 12. <sup>b</sup> Reference 8.

**Table 2.** Relevant Parameters of the Rietveld Analysis of  $Zr_2Te$ <sup>a</sup>

empirical formula	$Zr_2Te$
fw ( $\text{g mol}^{-1}$ )	310.048
space group (No.); Z	$Pnma$ (62); 12
$\rho_{\text{calc}}$ ( $\text{g cm}^{-3}$ )	7.600
angular range; step size ( $2\theta$ , deg)	8.5–85.0; 0.02
no. of phases	4
no. of variables (structural, $Zr_2Te$ )	54 (20)
no. of reflns	698
preferred orientation:	[010]; 1.271(6)
direction; parameter	
asymmetry parameter	0.8(1)
$R_B^b$ ; $R_p^c$ ; $R_{wp}^d$	0.056; 0.110; 0.143

<sup>a</sup> Lattice parameters in Table 1. <sup>b</sup>  $R_B = \sum |I_{\text{obs}} - I_{\text{calc}}| / \sum I_{\text{obs}}$ . <sup>c</sup>  $R_p = \sum |y_{\text{obs}} - y_{\text{calc}}| / \sum |y_{\text{obs}}|$ . <sup>d</sup>  $R_{wp} = [(\sum w_i (y_{\text{obs}} - y_{\text{calc}})^2) / (\sum w_i y_{\text{obs}}^2)]^{1/2}$ .



**Figure 1.** X-ray powder diffractogram (rel intensities vs  $2\theta$ ) and a Rietveld profile fit of  $Zr_2Te$  and the minority components  $Zr_{1-x}Te_x$ , Ta, and Te. Measured (dots) and calculated (line) intensities with a difference plot (bottom). In the middle are shown the positions of the Bragg angles of  $Zr_2Te$ ,  $Zr_{1-x}Te_x$ , Ta, and Te (top to bottom).

in Table 1. Atomic positions of  $Sc_2Te$  were used as a starting set for  $Zr_2Te$  in the Rietveld refinement. Relevant parameters are summarized in Table 2. An X-ray powder diffractogram and the profile fit are shown in Figure 1. Positional parameters are listed in Table 3. Displacement parameters were separately refined for the Zr and Te atoms.

**Electronic Structure Calculations.** Band structure calculations on  $Zr_2Te$  and  $Sc_2Te$  were carried out within the tight-binding approximation, at 48 k points, using the extended Hückel method.<sup>37–39</sup> Valence-state ionization energies ( $H_{ii}$ 's) for Zr were obtained from a charge-iterative calculation on  $Zr_2Te$ . Other parameters have been cited elsewhere.<sup>8,40</sup> The parameters are listed in Table 4.

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**Table 3.** Atomic Coordinates and Isotropic Thermal Parameters ( $B_{\text{iso}}$  ( $10^4 \text{ pm}^2$ )) of Zr<sub>2</sub>Te<sup>a</sup>

atom	<i>x</i>	<i>z</i>	$B_{\text{iso}}^b$
Te1	0.0731(3)	0.8392(6)	1.5(2)
Te2	0.3696(3)	0.7190(6)	1.5(2)
Te3	0.2605(3)	0.0398(6)	1.5(2)
Zr1	0.3951(4)	0.1499(8)	0.56(14)
Zr2	0.3421(4)	0.4580(8)	0.56(14)
Zr3	0.2270(4)	0.7487(8)	0.56(14)
Zr4	0.0348(4)	0.1088(8)	0.56(14)
Zr5	0.1456(4)	0.4218(8)	0.56(14)
Zr6	0.4845(4)	0.8890(8)	0.56(14)

<sup>a</sup> All atoms in position 4c with  $y = 1/4$ . <sup>b</sup>  $B_{\text{iso}} = 8\pi^2 U^2$ .

**Table 4.** Parameters Used for Extended Hückel Calculations

	orbital	$H_{ii}$ (eV)	$\zeta_1^a$	$c_1^b$	$\zeta_2^a$	$c_2^b$
Zr	5s	-7.33	1.82			
	5p	-3.84	1.78			
	4d	-6.91	3.84	0.6213	1.505	0.5798
Sc	4s	-6.74	1.30			
	4p	-3.38	1.30			
	3d	-6.12	4.35	0.4228	1.700	0.7276
Te	5s	-21.20	2.51			
	5p	-12.00	2.16			

<sup>a</sup> Slater-type orbital exponents. <sup>b</sup> Coefficients used in double- $\zeta$  expansion.

## Results and Discussion

**Phase Relations.** The lattice parameters found for Zr<sub>2</sub>Te prepared with varying nominal composition of the reaction mixture gave no indication of a homogeneity range. According to DTA (Differential Thermal Analysis) measurements<sup>41</sup> there exists a eutectic at  $1616 \pm 3 \text{ K}$  in mixtures of Zr and Zr<sub>3</sub>Te at 21 at. % Te. De Boer and Cordfunke<sup>41</sup> found Zr<sub>3</sub>Te to melt incongruently at  $1627 \pm 6 \text{ K}$  to give Zr<sub>5</sub>Te<sub>4</sub> and liquid. Zr<sub>2</sub>Te was not observed, however, they reported difficulties in assigning the enthalpy effects near 35 at. % Te. One reason for this might be their method of synthesis used,<sup>41</sup> which most likely leads to minimal amounts of Zr<sub>2</sub>Te being produced. Thus, during the heating part ( $10 \text{ K min}^{-1}$ ) of the first DTA cycle one could not expect the enthalpy effect of Zr<sub>2</sub>Te to be observed. On cooling from 1925 K at a rate of  $10 \text{ K min}^{-1}$  we would however expect the peritectic formation of solid Zr<sub>2</sub>Te to be observed in a small temperature interval some degrees above the incongruent melting point of Zr<sub>3</sub>Te. This circumstance has to be elucidated in further experiments.

**Structural Features.** Zr<sub>2</sub>Te crystallizes in the Sc<sub>2</sub>Te type structure<sup>8</sup> in the orthorhombic space group *Pnma*. The structure is composed of nine crystallographically inequivalent atoms, Zr1–Zr6 and Te1–Te3, all of which lie in mirror planes located at heights  $y = 1/4, 3/4$  modulo 1. Interatomic distances are listed in Table 5. As emphasized in the projection of the Zr<sub>2</sub>Te structure along [010] in Figure 2a, four of the six crystallographically distinct Zr atoms (Zr1, Zr4, Zr5, Zr6) topologically represent a metal partial structure (Figure 3a), which is reminiscent of the bcc-Zr structure stable above 1139 K.<sup>42</sup> Pairwise interpenetrating columns of *trans*-face-shared, centered Zr<sub>9</sub> cuboids with additional Te atoms above the square faces extend along [010] presenting the motif of a  $1_{\infty}[\text{Zr}_8\text{Te}_4]$  double string. Accordingly, the *b* parameter (382.4 pm) is related to the cell parameter of bcc-Zr (360.9 pm at 1252 K).<sup>43</sup> In the metal core of the structure (Zr1, Zr5, Zr6), the Zr1–Zr6 and

Zr6–Zr6 contacts (310 pm) are shorter than the closest contact in the metal itself ( $d_{\text{bcc}} = 313 \text{ pm}$ ;  $d_{\text{hcp}} = 318 \text{ pm}$ ).<sup>42</sup> Other contacts in the core unit range from 320 to 379 pm. Alternatively, in compliance with Simon's concept of condensed clusters,<sup>21</sup> the columnar fragments also represent chains of distorted octahedral Zr<sub>6</sub>Te<sub>8</sub> clusters with Zr atoms replacing four of the eight Te atoms. Each octahedron is condensed via *cis*-standing edges to two others, resulting in double chains of the columnar fragments. The double chains are further linked along [001] via homonuclear Zr4–Zr4 contacts (331 pm), as shown in Figure 6b. The metal atoms on the remaining two sites, Zr2 and Zr3, can be viewed as forming a zigzag chain along [010] with  $d_{\text{Zr2-Zr3}} = 325 \text{ pm}$ . Note, that this chain can also be seen as a part of another type of a columnar bcc fragment built up by Zr2 and Zr3 together with Zr1 and Zr5 of two separate double chains. Two further bcc fragments can be detected in the structure: Zr3 connecting to Zr1 and Zr5, and thereby forming a part of the bcc structure in the double chains, and, a fragment built up by Zr2 and two Zr4. These fragments, along with the double chain, can in part be recognized in Figure 2b, where Zr atom pairs with Mulliken overlap population (MOP) values  $> 0.090$  (vide infra) are connected.

The Te atoms are 7-, 8-, and 9-fold (Te2, Te3, and Te1, respectively) coordinated by Zr, the coordination figures representing mono-, bi-, and tricapped distorted trigonal prisms, with Zr atoms capping the square faces of the prisms (Figure 3b). The bi- and tricapped trigonal prisms are fused along [010] by sharing of triangular Zr<sub>3</sub> faces. These fragments are condensed via common edges, building twin columns of capped trigonal prisms running along [010] (see Figure 6b). The monocapped prisms lie perpendicular to the others, sharing edges, and thereby building columns along [010] of oppositely oriented prisms with every other prism centered with a Te atom. An examination of the average Zr–Te distances in the capped, Te centered trigonal prisms, amounting to 293.2 pm (Te1), 290.2 pm (Te2), 297.2 pm (Te3), and 293.5 pm in the average, reveals a close resemblance to the Te coordination figures in the contiguous phases Zr<sub>3</sub>Te and Zr<sub>5</sub>Te<sub>4</sub>. In Zr<sub>3</sub>Te,<sup>13</sup> where solely tricapped trigonal prisms are observed, the corresponding average distance is 294.5 pm, and, in Zr<sub>5</sub>Te<sub>4</sub>,<sup>27</sup> where the Te coordination figure constitutes a *hypho*-Zr<sub>6</sub>Te fragment of a tricapped trigonal prism, this distance is 295.1 pm. The shortest Te–Te contact in Zr<sub>2</sub>Te is 375 pm (Te2–Te3), well above the second nearest neighbor distance in elemental Te (349.5 pm).<sup>36</sup>

**Electronic Structure: Comparison with Sc<sub>2</sub>Te.** According to the total DOS diagrams for Zr<sub>2</sub>Te and the isotypic Sc<sub>2</sub>Te shown in Figure 4, the Fermi level lies in both cases in regions that are dominated by Sc and Zr d orbital contributions. Hence, both materials should have metallic properties. The Fermi level in Sc<sub>2</sub>Te lies near the top of a peak in the DOS whereas in Zr<sub>2</sub>Te it lies near a local minimum. The greater dispersion of the projected Zr 4d states around the Fermi level compared to the Sc 3d states can be related to a greater interaction integral<sup>44</sup> of the Zr atoms in Zr<sub>2</sub>Te. Using the cohesive energy of the metal as an indicator for the interaction integral, the dispersions of the projected 3d and 4d states of Sc and Zr, respectively, reflect the difference in the enthalpies of atomization:  $\Delta_a H_{298}^0(\text{Zr}) = 608.8 \text{ kJ/mol}$ ;  $\Delta_a H_{298}^0(\text{Sc}) = 377.8 \text{ kJ/mol}$ .<sup>45</sup> This difference leads to slightly shorter homonuclear metal–metal mean distances in Zr<sub>2</sub>Te, expressing itself mainly in the shorter *b* axis.

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**Table 5.** Relevant Interatomic Distances and Mulliken Overlap Populations (MOP) for Zr<sub>2</sub>Te (MOP Values in Square Brackets Refer to Sc<sub>2</sub>Te)

atom no.			d (pm)	MOP	atom no.			d (pm)	MOP
M6	M6	2×	310.3(9)	0.406 [0.362]	M2	M5	1×	394(1)	0.040 [–]
M1	M6	2×	309.9(9)	0.354 [0.228]	M2	M3	1×	386(1)	0.039 [0.013]
M1	M6	1×	330(1)	0.291 [0.167]	M5	M5	2×	382.36(2)	0.032 [–]
M1	M5	2×	319.7(9)	0.290 [0.174]	M4	M5	1×	400(1)	0.029 [0.008]
M5	M6	2×	324.2(9)	0.276 [0.146]					
M2	M3	2×	324.5(1)	0.274 [0.086]	M3	Te2	1×	286(1)	0.357 [0.296]
M4	M4	2×	331.1(1)	0.223 [0.028]	M2	Te2	1×	284(1)	0.356 [0.320]
M1	M3	2×	327.1(9)	0.212 [0.121]	M2	Te1	2×	285.0(8)	0.323 [0.274]
M2	M4	2×	350.2(1)	0.134 [0.015]	M5	Te2	2×	290.1(8)	0.315 [0.254]
M1	M2	1×	345(1)	0.131 [0.072]	M2	Te3	2×	293.4(8)	0.307 [0.283]
M5	M6	1×	379(1)	0.117 [0.174]	M4	Te2	2×	294.5(8)	0.301 [0.280]
M4	M6	2×	356.7(1)	0.108 [0.036]	M1	Te1	2×	285.1(8)	0.297 [0.234]
M6	M6	2×	382.36(2)	0.101 [0.045]	M3	Te3	2×	294.5(8)	0.290 [0.225]
M3	M3	2×	382.36(2)	0.100 [0.036]	M5	Te3	2×	295.7(8)	0.280 [0.236]
M3	M5	2×	367.6(1)	0.095 [0.009]	M4	Te1	2×	293.2(8)	0.279 [0.273]
M1	M1	2×	382.36(2)	0.064 [0.018]	M6	Te2	1×	292(1)	0.266 [0.239]
M1	M4	1×	379(1)	0.060 [0.026]	M1	Te3	1×	293(1)	0.249 [0.246]
M4	M4	2×	382.36(2)	0.056 [–]	M4	Te1	1×	297(1)	0.249 [0.267]
M2	M2	2×	382.36(2)	0.052 [–]	M6	Te1	1×	301(1)	0.194 [0.127]
M3	M5	1×	384(1)	0.045 [0.010]	M3	Te3	1×	317(1)	0.161 [0.173]
M2	M4	1×	391(1)	0.042 [–]	M3	Te1	1×	322(1)	0.112 [0.073]

The Te p orbital contributions lie in both cases almost exclusively in the region between  $-14.5$  and  $-11$  eV. There are also small Te p orbital contributions well above the Fermi level.

The crystal orbital overlap population (COOP) curves (Figure 4) for heteronuclear metal–tellurium interactions show comparable bonding behavior for both compounds. On the other hand, it can be seen that homonuclear metal–metal d orbital interactions are much more pronounced, and that fewer metal–metal bonding states remain unoccupied in Zr<sub>2</sub>Te than in Sc<sub>2</sub>Te. As the 4d orbitals of Zr are, compared to the 3d orbitals of Sc, more diffuse and more exposed, and considering the difference in valence electron concentration *vec* available for metal–metal bonding in the two materials (Zr<sub>2</sub>Te, *vec* = 3; Sc<sub>2</sub>Te, *vec* = 2), this can be anticipated. In the region between  $-14.5$  and  $-11$  eV there are minor Zr–Zr bonding interactions, whereas in this region bonding and antibonding Sc–Sc interactions cancel out.

The Mulliken overlap population (MOP) for a given pair of atoms can be seen as an expression for the strength of bonding interactions. MOP values for pairs of atoms in both structures are listed in Table 5, for Zr<sub>2</sub>Te according to decreasing MOP, with a parallel listing of the corresponding values for Sc<sub>2</sub>Te. In plotting MOP vs  $d_{M-M}$  and MOP vs  $d_{M-Te}$  for metal–metal and metal–tellurium pairs respectively, the distinctions in homonuclear and heteronuclear bonding are demonstrated (Figure 5). The MOP values for metal–tellurium interactions scatter around a common straight line, with bond length being inversely proportional to bond strength. In contrast, the MOP values for homonuclear metal–metal interactions scatter around two separate curves. Considering the near coincidence of the MOP vs  $d_{M-Te}$  curves for both structures, we conclude that the distinct MOP vs  $d_{M-M}$  curves reflect differing MOP values at a given distance rather than differing interatomic distances at a given MOP value.

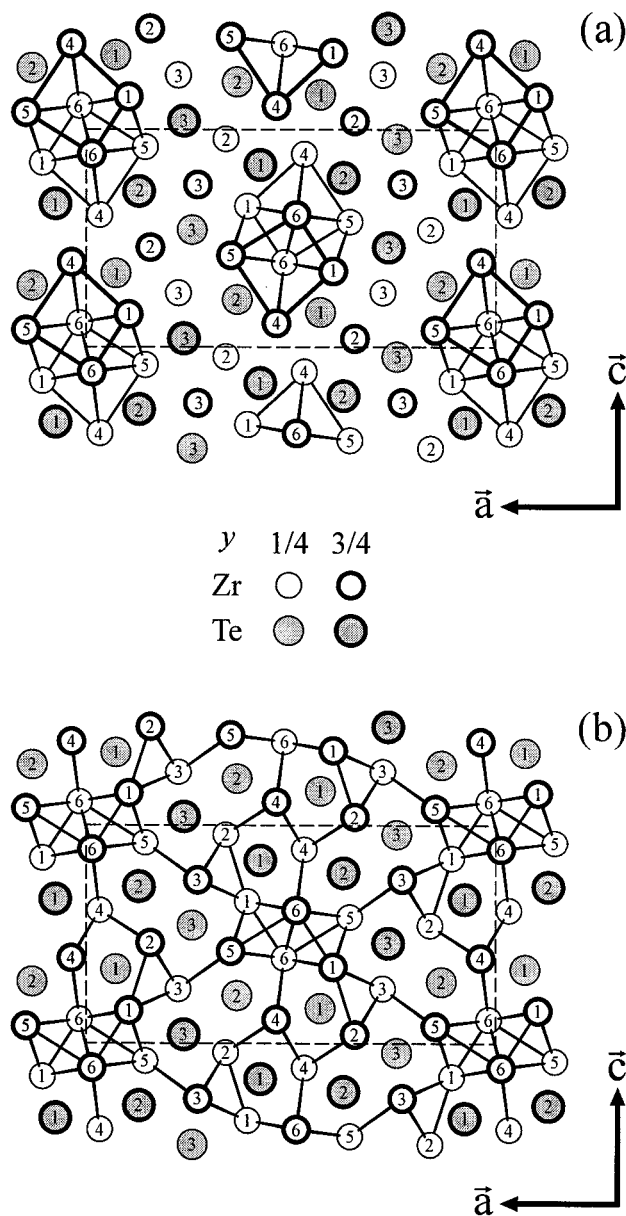
A qualitative comparison of the two diagrams in Figure 5 gives an idea of the greater importance of metal–tellurium interactions for Sc<sub>2</sub>Te over Zr<sub>2</sub>Te, or, conversely, the greater importance of metal–metal interactions in Zr<sub>2</sub>Te over Sc<sub>2</sub>Te. Exploring this point further, it can be seen from Table 5 that the Sc–Sc overlap populations around Sc2, Sc3, and Sc4, the Sc atoms with the most Te neighbors, namely five, are small, much smaller than the corresponding values in Zr<sub>2</sub>Te. Compare for example the metal pairs ( $d_{Zr-Zr}$ , pm) 2–3 (324.5), 3–5

(367.6), 4–4 (331.1), 2–4 (350.2), and 4–6 (356.7). Sc2, Sc3, and Sc4 are strongly oxidized through Te and, as opposed to the equivalent Zr atoms, lack valence electron density in order to enter into extensive bonding interactions with the surrounding Sc atoms. They will therefore, compared to the equivalent Zr atoms, be to a much greater extent fixed in place by heteronuclear rather than homonuclear covalent interactions. The bonding electrons of Sc1, Sc5, and Sc6, on the other hand, seem to be more or less delocalized inside the core unit built up by these very atoms, leading to extended metal–metal bonding interactions inside this unit.

An anomaly regarding the Sc1–Sc6 (348.6 pm) and Sc5–Sc6 (351.9 pm) interactions becomes evident in the MOP vs  $d_{M-M}$  diagram of Figure 5. Compared to the Zr1–Zr6 coordinate, the Sc1–Sc6 coordinate is shifted toward a lower overlap population and a longer distance. Opposedly, the values for Sc5–Sc6 are shifted to a higher overlap population and a shorter interatomic distance. A tentative explanation would be the increased metal coordination of M1 through M2 and M3, as compared to the metal coordination of M5. In the case of electron-poor Sc this could lead to some weakening of the interaction of Sc1 with Sc6. As a consequence, the long Sc1–Sc6 distance then would allow Sc5 to interact more closely with Sc6.

In summary, on the basis of the extended Hückel calculations, it can be said that in Zr<sub>2</sub>Te the metal–metal bonding interactions spread throughout the whole body of the structure, as shown in Figure 2b, building an extensive network, whereas in Sc<sub>2</sub>Te the metal–metal interactions are concentrated on the core unit leaving other parts of the structure dominated by Sc–Te interactions. From this viewpoint one gains the impression that there exist two, to a certain extent uncoupled, electronic systems in this bcc fragment structure, one that provides the electrons for the one-dimensionally extended M–M bonding regions in the metallic core and a second one that may serve as an electron reservoir for additional M–M bonding. This reservoir seems to be nearly empty for Sc<sub>2</sub>Te, and almost filled for Zr<sub>2</sub>Te.

**Structural Comparison.** The Sc<sub>2</sub>Te type structure seems to allow a considerable amount of flexibility in the number of valence electrons available for homonuclear M–M bonding interactions. This can be seen from the way in which the bonding interactions in Sc<sub>2</sub>Te and Zr<sub>2</sub>Te adjust to the structure, leaving it topologically intact, instead of forming two different structure



**Figure 2.** Projection of the  $Zr_2Te$  structure along [010]. The unit cell is indicated by a broken line. (a) Pairwise interpenetrating columns of *trans*-face-shared centered  $Zr_9$  cuboids are emphasized. (b)  $Zr$ - $Zr$  contacts with a Mulliken overlap population (MOP)  $> 0.090$  are emphasized, revealing the three-dimensional character of the bonding interactions in the metal partial structure of  $Zr_2Te$ .

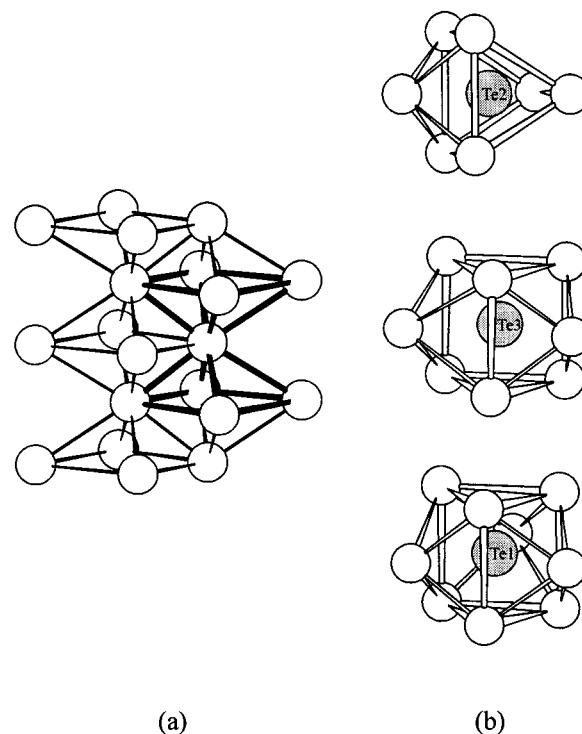
types. Therefore, it is interesting to investigate why the homologous telluride  $Hf_2Te^{14}$  and the congeneric zirconium selenide  $Zr_2Se^{46}$  adopt structures which differ from that of  $Zr_2Te$  (Figure 6). Obviously there are additional factors conducting the structure type formed in each case.

Let us first examine the distinctions in the structures of  $Zr_2Te$  and  $Hf_2Te$ , beginning with some remarks concerning the metal-rich chemistry of Zr and Hf. Despite the extensive similarities in the general chemical behavior of the two homologous metals Zr and Hf,<sup>47</sup> stoichiometrically equivalent, metal-rich chalcogenides and pnictides of both metals have, apart from  $Zr_7P_4^{48}$  and  $Hf_7P_4^{49}$  been found to form different structure

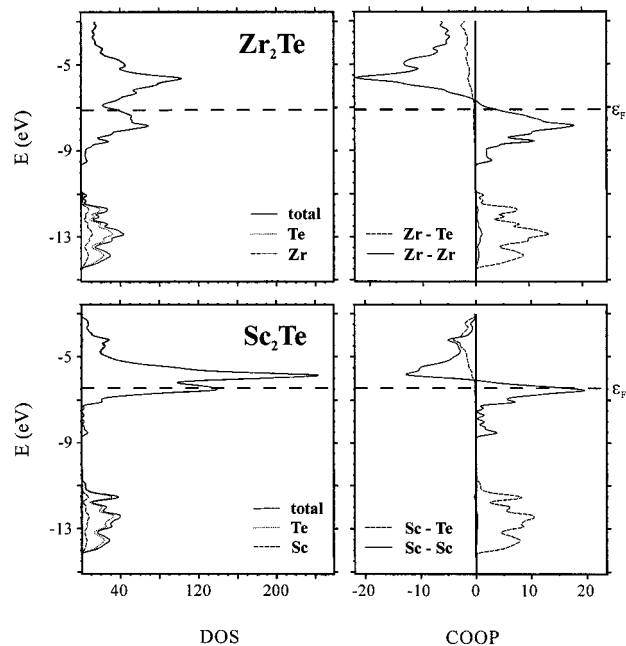
(46) Franzen, H. F.; Norrby, L. J. *Acta Crystallogr.* **1968**, B24, 601.

(47) Greenwood, N. N.; Earnshaw, A. *Chemistry of the Elements*, 2nd edition; Butterworth-Heinemann: Oxford, 1997.

(48) Ahlén, P.-J.; Rundqvist, S. Z. *Kristallogr.* **1989**, 189, 149.



**Figure 3.** (a) Mutually interpenetrating bcc columns or, alternatively, *cis*-edge-condensed strings of compressed octahedra. A centered  $Zr_9$  cuboid is emphasized. (b) Coordination environments of the Te atoms consist of mono-, bi-, and tricapped distorted trigonal prisms of Zr atoms.

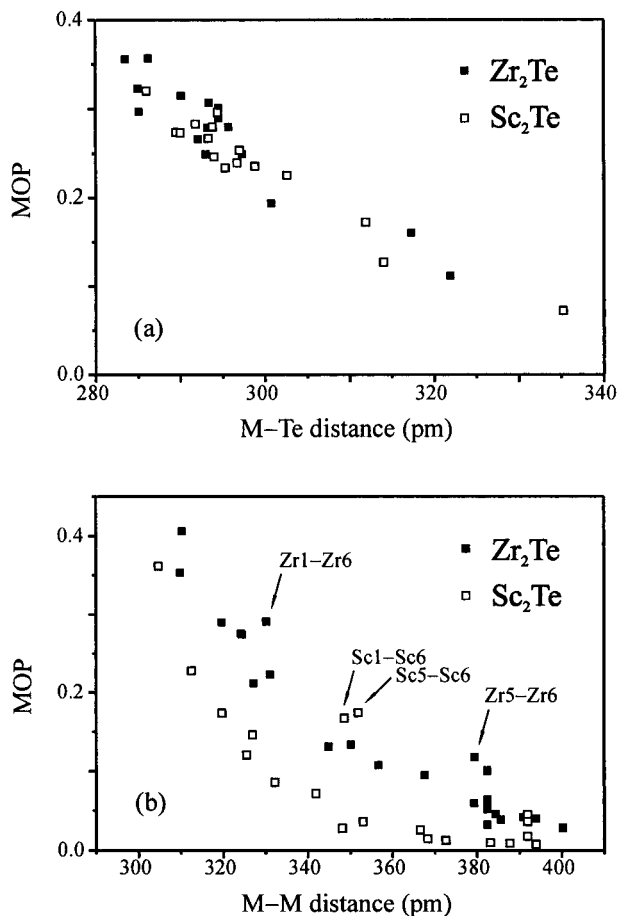


**Figure 4.** Densities of states (DOS) and crystal orbital overlap population (COOP) curves for  $Zr_2Te$  and the isotopic, valence-electron-poorer  $Sc_2Te$ . COOP curves for  $Zr$ - $Te$ ,  $Zr$ - $Zr$ ,  $Sc$ - $Te$ , and  $Sc$ - $Sc$  contacts up to 322, 400, 335.2, and 404.9 pm, respectively, are shown. Levels up to the Fermi level (broken, horizontal line) are filled; levels to the right of the vertical line are bonding, to the left are antibonding. The Fermi level for  $Zr_2Te$  lies at  $-7.12$  eV, for  $Sc_2Te$  at  $-6.45$  eV.

types, e.g.,  $Zr_2Se^{46}$  vs  $Hf_2Se$ ,<sup>2</sup> and  $Zr_2P^{50}$  vs  $Hf_2P$ .<sup>51</sup> Other metal-rich pnictides and chalcogenides, e.g.,  $Zr_3Te$ ,<sup>13</sup>  $Hf_3Te_2$ ,<sup>15</sup>

(49) Kleinke, H.; Franzen, H. F. *Angew. Chem., Int. Ed. Engl.* **1996**, 35, 1934.

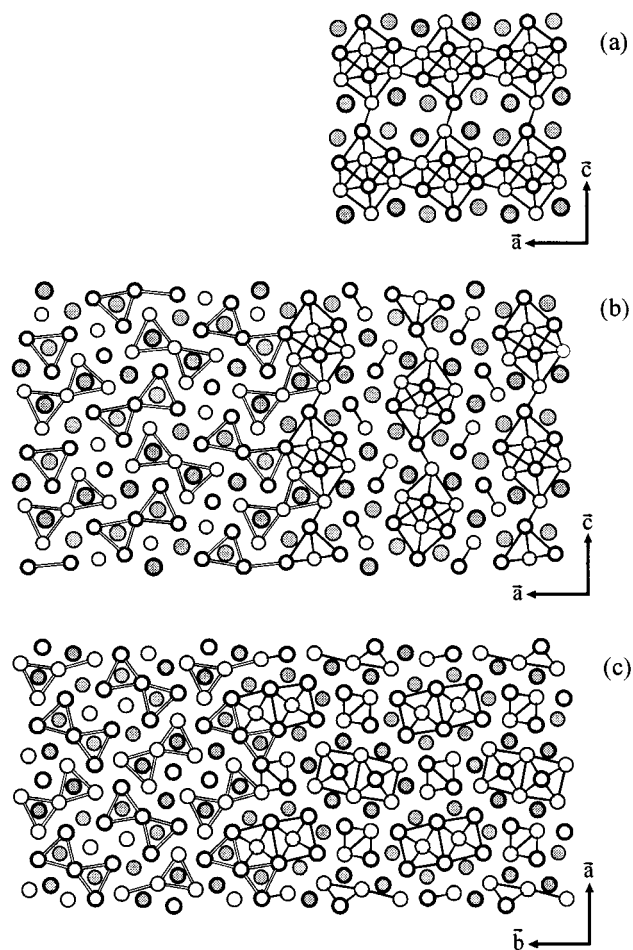
(50) Ahlén, P.-J.; Rundqvist, S. Z. *Kristallogr.* **1989**, 189, 117.



**Figure 5.** Mulliken Overlap Populations (MOP) in Zr<sub>2</sub>Te and Sc<sub>2</sub>Te plotted versus (a) M-Te and (b) M-M distance.

Zr<sub>14</sub>P<sub>9</sub>,<sup>52</sup> and Hf<sub>3</sub>P<sub>2</sub>,<sup>53</sup> possess no counterpart of the respective homologous transition metal. In the quasi-binary systems Hf-Zr-P<sup>54</sup> and Ta-Nb-S<sup>55</sup> it has been shown that the 5d metals generally prefer the metal sites which allow more metal-metal bonding. The differences in bonding in metal-rich Zr and Hf phosphides have been extensively discussed before,<sup>49</sup> referring among others to the effects of the greater expansion of the 5d orbitals of Hf compared to the 4d orbitals of Zr, leading to stronger metal-metal bonds and eventually to more metal-metal contacts in Hf compounds. This behavior can be verified with help of the aforementioned M<sub>2</sub>X compounds (M = Zr, Hf; X = Se, P) as well as Ta<sub>2</sub>Se<sup>16</sup> and Nb<sub>2</sub>Se.<sup>5</sup> Accompanying more metal-metal contacts and stronger metal-metal bonds are larger uninterrupted regions consisting solely of metal atoms. As pointed out earlier,<sup>14</sup> this concomitantly leads to a reduction in the average coordination number of the chalcogen or pnictogen atoms involved.

The homologous telluride Hf<sub>2</sub>Te<sup>14</sup> (Nb<sub>2</sub>Se type)<sup>5</sup> adopts a structure which in its basic assemblage is related to that of Zr<sub>2</sub>Te (Figure 6a and b) but nonetheless expressing important differences. The motif of double strings linked in the second dimension over corners appears in both structures. As can be seen in Figure 6 every second one of these "sheets" in Zr<sub>2</sub>Te has to be rotated 180° about the *c* axis, shifted by *c*/2 and *b*/2



**Figure 6.** Comparison of the crystal structures of (a) Hf<sub>2</sub>Te, (b) Zr<sub>2</sub>Te, and (c) Zr<sub>2</sub>Se. On the left, twin columns of bi- and tricapped trigonal Zr prisms are emphasized, whereas on the right the bcc partial structure is stressed. Shaded circles represent chalcogen atoms, open circles represent metal atoms. Thin-bordered circles are located in the plane, thick-bordered at 1/2 lattice parameter above the plane.

and connected in the third dimension, leaving out one Te atom (Te3) and the zigzag Zr chain (Zr2 and Zr3), to arrive at the Hf<sub>2</sub>Te structure. As opposed to Zr<sub>2</sub>Te, Hf<sub>2</sub>Te contains no capped trigonal prisms about Te although a related motif can be recognized, the Te coordination representing *hypho*-Hf<sub>6</sub>Te fragments of tricapped trigonal prisms, Hf<sub>9</sub>Te. Inspecting the metal substructure of Hf<sub>2</sub>Te reveals two-dimensional puckered slabs of bcc arrays spreading parallel (001), the Te atoms lining channels extending throughout the structure along [010]. This gives way to a more highly metal-metal bonded structure than in Zr<sub>2</sub>Te and is in accordance with the general conclusion that in metal-rich chalcogenides and pnictides of the valence-electron-poor transition metals the 5d metals are more efficient metal-metal bonding elements than the 4d metals.

Comparing Zr<sub>2</sub>Te with the congeneric zirconium selenide Zr<sub>2</sub>Se<sup>46</sup> (Ta<sub>2</sub>P type)<sup>56</sup> is also instructive (Figure 6b and c). Zr atoms on four of the six distinct sites form face-shared double strings of bcc columns, as opposed to the mutually interpenetrating bcc columns found in Zr<sub>2</sub>Te. Atoms on the two additional sites form a chain of compressed octahedra sharing edges. The chalcogen atoms are mono-, bi-, and tricapped trigonal prismatic coordinated with Zr. The bi- and tricapped trigonal prisms are also fused over trigonal faces to columns,

(51) Lundström, T.; Ersson, N.-O. *Acta Chem. Scand.* **1968**, *22*, 1801.

(52) Tergenius, L.-R.; Nörläng, B.; Lundström, T. *Acta Chem. Scand.* **1981**, *A35*, 693.

(53) Kleinke, H.; Franzen, H. F. *Acta Crystallogr.* **1996**, *C52*, 2127.

(54) Kleinke, H.; Franzen, H. F. *J. Solid State Chem.* **1998**, *136*, 221.

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(56) Nylund, A. *Acta Chem. Scand.* **1966**, *20*, 2393.

the columns being edge condensed to double strings as found in  $Zr_2Te$ . The monocapped trigonal prisms lie perpendicular to these strings. This arrangement thus maintains the chalcogen coordination found in  $Zr_2Te$ , whereas shorter average Zr–Zr contacts are realized. Compared to Te, the smaller atomic radius of Se probably plays an important role regarding the type of structure formed. Interestingly,  $Ti_2Se$  forms both structure types,  $\alpha$ - $Ti_2Se$ <sup>57</sup> ( $Ta_2P$  type) and  $\beta$ - $Ti_2Se$ <sup>12</sup> ( $Sc_2Te$  type), indicating relatively small stability differences between these two structure types.  $Ti_2Te$  has not been described in the literature.

### Conclusions

The new binary zirconium-rich telluride  $Zr_2Te$  has been uncovered and its structure determined by means of X-ray powder diffractometry.  $Zr_2Te$  is isostructural with  $\beta$ - $Ti_2Se$  as well as  $Sc_2Te$ , which has less valence electrons available for metal–metal bonding. As both Zr and Sc have almost the same atomic radii, the contribution of matrix effects to bonding differences in  $Zr_2Te$  and  $Sc_2Te$  is minimized and the influence

of other factors, such as the number of valence electrons, can be assessed. Therefore, a study of the bonding situation in  $Zr_2Te$  and  $Sc_2Te$  was undertaken. According to extended Hückel calculations, metal–tellurium bonding interactions are comparable for  $Sc_2Te$  and  $Zr_2Te$ , whereas metal–metal bonding is anisotropically reduced from a three-dimensional network in  $Zr_2Te$  to quasi-one-dimensional units in  $Sc_2Te$ . Despite the robustness of the  $Sc_2Te$  type structure with respect to variations in the valence electron concentration available for M–M bond formation, the congeneric phases of  $Zr_2Te$ ,  $Zr_2Se$  and  $Hf_2Te$ , adopt distinct types of structures.

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**Supporting Information Available:** A table of relevant parameters of the Rietveld analysis and a table of  $2\theta$ -values, indexation, calculated and observed relative X-ray diffraction intensities of  $Zr_2Te$ . This material is available free of charge via the Internet at <http://pubs.acs.org>.

(57) Weirich, T. E.; Pöttgen, R.; Simon, A. *Z. Kristallogr. – NCS* **1996**, *211*, 928.