One-Pot Synthesis, Structure, and Unusual Luminescence of Novel One-Dimensional Lanthanide(III) Tetramethoxyborates

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Detailed studies of formation, structure, and luminescence characteristics of novel ionically linked sinusoidal chains of $[\{Na\}\{(sal)_3Ln(B(OCH_3)_4)Ln(sal)_3\}]_n$ [Ln = Nd, Sm, Tb-Lu, Y (4); Ln = Eu (5); Ln = Gd (6)] was undertaken to assess the versatility of tetramethoxyborate anions as consistent general synthons for one-dimensional materials. Crystal data for $[\{Na\}\{(sal)_3Nd(B(OCH_3)_4)Nd(sal)_3\}]_n$ [4(Ln = Nd)]: monoclinic; space group, C2/c (15), a = 20.884(2) Å, b = 12.454(1) Å, c = 37.663(5) Å, $\beta = 102.09(2)^\circ$, Z = 8; $[\{Na\}\{(sal)_3Eu(B(O CH_3)_4)-Eu(sal)_3\}]_n$ (5): monoclinic; space group, C2/c (15), a = 20.846(2) Å, b = 12.460(1) Å, c = 37.365(4) Å, $\beta = 102.17(1)^\circ$, Z = 8. Europium luminescence from $[\{Na\}\{(sal)_3Ln_{1-x}Eu_x(B(OCH_3)_4)Ln_{1-x}Eu_x(sal)_3\}]_n$ [Ln = Eu (5); Gd (7–11), Y (12), Dy (13–18), Dy/Gd (19)] reveals the presence of defect Eu³⁺ sites, possibly resulting from swapping regular Eu³⁺ and Na⁺ positions. Strong Eu³⁺-Eu³⁺ electronic interactions result in multidimensional energy migration and a dynamic excitation energy-transport regime. Efficient thermally activated regular Eu³⁺ to-defect Eu³⁺ back-energy transfer dominates at 297 K.

Introduction

Although simple tetra-alkoxyborate anions ([(RO)₄B]⁻ R = alkyl group) have been known for about a century,¹ relatively few reports on their coordination behavior exist. Most of the reported examples cover discrete complexes of groups one² and two³ elements for which the bidentate coordination behavior of the $[(RO)_4B]^-$ chelates is prominent. However, the $[(CH_3O)_4-$ B]⁻ anion displays remarkable versatility as a template for metallopolymers of one-dimensional structural motifs. In two separate serendipitous events, large lanthanide(III) ions (Ln =La-Nd), formed linear $[Ln(NO_3)_2((CH_3O_4B)(CH_3OH)_2]_n$ (1)⁴ polymers, whereas cadmium(II) formed helical chains of the salt $[Cd(tcm)B(OCH_3)_4] \cdot xCH_3OH$ (tcm = $[C(CN)_3]^{-1/5}$ (2). These broad similarities in structural motifs are observed despite vast differences in the chemical properties of hard lanthanide-(III) and soft cadmium(II) Lewis acids. The $[(CH_3O)_4B]^-$ ion and other tetra-alkoxyborates generally are thus attractive as potential synthons in the preparation, and rational crystal engineering, of materials in which strict control of the directionality of physical behavior such as magnetism and magnetic coupling,⁶ electron and electronic energy⁷ transport, or polariza-

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tion⁸ is desired. The problem is that the conditions under which the above-mentioned one-dimensional arrays are assembled by the tetramethoxyborate anions are enigmatic.

We have been interested in tetra-alkoxyborate and tetraphenoxyborate chelates as preparative templates for one-dimensional lanthanide(III) compounds with electronic energy-transport characteristics that could be readily manipulated. In the presence of efficient light-harvesting chromophores (light antennae),9 which are coupled to lanthanide(III) electronic states, and are therefore capable of sensitizing intense red europium(III) or green terbium(III) emission, one-dimensional arrays are of great interest in fluorescent sensors¹⁰ and diagnostic¹¹ applications. After numerous unsuccessful attempts, including one with 3, we discovered that one-pot methanolysis of boron species generated from Borax (Na₂[B₄O₅(OH)₄]•8H₂O) in the presence of salicylaldehyde (Hsal) and neodymium(III) nitrate produces a novel ionicaly linked polymeric sinusoidal one-dimensional $[{Na}{(sal)_3Nd(B(OCH_3)_4)Nd(sal)_3}]_n$ compound. We have successfully used this technique to prepare complexes of almost the entire lanthanide(III) series ([{Na}{(sal)_3Ln(B(OCH_3)_4)Ln- $(sal)_3$]_n [Ln = Nd, Sm, Tb-Lu, Y (4); Ln = Eu (5); Ln = Gd (6)]). To the best of our knowledge, this is the first consistent method for fabricating metallo-one-dimensional arrays using tetra-alkoxyborate synthons. With this breakthrough as well as series 1 and 2, it seems that design and preparation of

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Table 1. Numerical Assignments and Stoichiometries for $[{Na}{(sal)_3Ln'_{1-x}Ln_x(B(OMe)_4)Ln'_{1-x}Ln_x(sal)_3}]_n \cdot mH_2O$ Compounds

com-	$[\{Na\}\{(sal)_3Ln'_{1-x}Ln_x(B(OMe)_4)Ln'_{1-x}Ln_x(sal)_3\}]_n \cdot mH_2O$					
pound	Ln	Ln'	x			
4	Nd-Lu, Y	Nd-Lu, Y				
5	Eu	Eu				
6	Gd	Gd				
7	Eu	Gd	0.06			
8	Eu	Gd	0.16			
9	Eu	Gd	0.37			
10	Eu	Gd	0.51			
11	Eu	Gd	0.79			
12	Eu	Y	0.04			
13	Eu	Dy	0.03			
14	Eu	Dy	0.09			
15	Eu	Dy	0.16			
16	Eu	Dy	0.50			
17	Eu	Dy	0.58			
18	Eu	Dy	0.76			
19	Eu	Dy/Gd	≈ 0.29			

structurally diverse one-dimensional arrays using tetra-alkoxyborates is achievable and it is now possible for electronic properties of an extended series of compounds of low dimensionality, such as $[{Na}_{(sal)_3Ln(B(OCH_3)_4)Ln(sal)_3}]_n$, to be evaluated.

Herein, we report the preparation, crystal structure, and intriguing luminescence decay dynamics of the $[{Na}_{(sal)_3Ln-(B(OCH_3)_4)Ln(sal)_3}]_n$ series.

Experimental Section

Materials. Salicylaldehyde (99%), standard pH 9.22 buffer tablets, and Borax ($Na_2[B_4O_5(OH)_4]\cdot 8H_2O$) were obtained from BDH. Lanthanide(III) nitrates were obtained by neutralizing the corresponding oxides (Aldrich, 99.99% purity) with nitric acid, followed by evaporation of water to near dryness. Methanol was of reagent grade.

Spectral and Quantitative Analyses. Carbon, hydrogen, and nitrogen analyses were obtained from MEDAC Ltd., Brunel Science Center, Surrey, UK; lanthanide concentrations were obtained by neutron activation analysis at the International Center for Environmental and Nuclear Sciences, University of the West Indies, Jamaica.¹² A Perkin-Elmer 1000 S FT-IR and HP 8451A diode array UV–Visible spectrometers were used to record IR and electronic absorption spectra, respectively. The powder X-ray diffraction analyses were obtained from A Siemens D5000 diffractometer.

Syntheses. Compounds **4**–**19** (Table 1) were prepared using standard pH 9.2 buffer tablets or Borax (Na₂[B₄O₅(OH)₄]•8H₂O), as shown in Scheme 1. Crystals were dried on tissue paper to maintain their crystalline integrity. Analyses¹³ showed the crystals to be anhydrous for **4** (Ln = Nd); monohydrates for **4** (Ln = Sm, Tm, Lu) and **5**; dihydrates for **4** (Ln = Tb, Dy, Ho, Yb) and **6**; and trihydrates for **4** (Ln = Y, Er). However, powder X-ray diffraction analyses (Figure 1) show (vide infra) that crystals of **4**–**6** are isomorphous irrespective of the hydration number. Further the luminescence decay kinetic values of terbium and europium samples prepared from CH₃OD/D₂O were similar to those of compounds prepared from CH₃OH, which confirms that the water found in the analyses is crystal rather than coordinated water.

Luminescence Studies. The luminescence and excitation spectra were recorded using an LS5B Perkin-Elmer Fluorescence spectrometer that is essentially similar to the one described earlier.¹⁴ The lumines-



Figure 1. A comparison of powder X-ray diffraction spectra of $[{Na}_{(sal)_3Ln(B(OMe)_4)Ln(sal)_3}]_n$ (4) (Ln = Nd, Ho, and Lu).

Scheme 1. Synthesis of $[{Na}{(sal)_3Ln(B(OMe)_4)Ln(sal)_3}]_n$ (4) (Ln = Nd-Lu, Y)



cence decay rates were measured using either our new Continuum Powerlite 8000 YAG laser and a matching ND6000 dye Laser, or a nitrogen dye laser, and an electronic and computational setup that was described previously.¹² The Powerlite 8000 YAG laser generates ca. 1200, 600, and 300 mJ pulses (5–8 ns) at 1064, 532, and 355 nm, respectively, whereas the ND6000 may generate up to 150-mJ pulses with the Rhodamine 6G dye. Indiscriminate excitation of samples was effected at 365 nm whereas selective excitation was achieved using a Rhodamine 6G dye tuned at 572 nm for Eu³⁺ (⁵D₀) and a coumarin 500 dye tuned at 515–535 nm, depending on temperature, for the Eu³⁺ (⁵D₁) state. Variable temperature measurements (8–320 K) were done using our new APD Cryogenics Inc. CSW-202 Displex Helium refrigerator system with the sample in contact with cryocon-conducting grease. Excitations using light with a wavelength of 337 nm (from a

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Table 2.	Crystallographic Data for
[{Na}{(sa	$l_3Ln(B(OMe)_4Ln(sal)_3)]_n^a$

data	5	4 [Ln = Nd]
formula	C46H42O16BNaEu2	C46H42O16BNaNd2
formula weight	1188.5	1173.1
color, habit	yellow prisms	pale olive prisms
crystal size, mm	$0.40 \times 0.33 \times 0.13$	$0.37 \times 0.20 \times 0.17$
<i>T</i> , K	293	293
diffractometer	Siemens P4/PC	Siemens P4/PC
wavelength, Å	0.71073	0.71073
lattice type	monoclinic	monoclinic
space group	C2/c, 15	C2/c, 15
cell dimensions		
<i>a</i> , Å	20.846(2)	20.884(2)
b, Å	12.460(1)	12.454(1)
<i>c</i> , Å	37.365(4)	37.663(5)
β , deg	102.17(1)	102.09(2)
<i>V</i> , Å ³	9487(2)	9579(2)
Ζ	8	8
$D_{\rm c}$, g cm ⁻³	1.664	1.627
F(000)	4704	4656
μ , cm ⁻¹	0.270	0.222
θ range, deg	1.9-25.0	1.9-25.0
no. of unique reflections	8364	8456
measured		
observed, $ F_{\rm o} > 4\sigma(F_{\rm o})$	6109	5873
absorption correction	semiempirical	semiempirical
-	$(\psi$ -scans)	$(\psi$ -scans)
maximum, minimum	0.21, 0.14	0.40, 0.38
transmission		
no. of variables	597	597
R_1^b	0.050	0.053
wR_2^c	0.099	0.101
weighting factors a, b^d	0.036, 9.127	0.026, 0.000
largest difference peak,	0.85, -0.61	0.50, -0.47
hole, eÅ ⁻³		

^{*a*} Details in common: ω -scans, refinement based on F^2 . ^{*b*} $R_1 = \sum ||F_o| - |F_c|| \sum |F_o|$. ^{*c*} $wR_2 = \sqrt{\{\sum [w(F_o^2 - F_c^2)^2] / \sum [w(F_o^2)^2]\}}$. ^{*d*} $w^{-1} = \sigma^2(F_o^2) + (aP)^2 + bP$.

nitrogen laser) or 282 nm (from the frequency-doubled Rhodamine 6G dye laser) lead to degradation of the samples. Samples were thus routinely inspected for visible damage during the analyses, and duplicate runs were done for most of the luminescence decay curves to ensure reproducibility of decay behavior.

Structure Determination. Table 2 provides a summary of the crystal data, data collection, and refinement parameters for neodymium [4(Ln = Nd)] and europium (5) compounds. The structures were solved by direct methods and all the nonhydrogen atoms were refined anisotropically using full-matrix least-squares based on F^2 . The hydrogen atoms of the methoxy groups were located from ΔF maps and their positions optimized. The remaining hydrogen atoms were placed in calculated positions, and all the hydrogen atoms were assigned isotropic thermal parameters, $U(H) = 1.2 U_{eq}(C) [U(H) = 1.5 U_{eq}(C-Me)]$, and allowed to ride on their parent atoms. Computations were carried out using the SHELXTL PC program system.¹⁵

Results and Discussion

Syntheses and Structure. The successful preparation of members of series 4-6 is indicated by satisfactory elemental analyses¹³ and IR spectra, which show the presence of the carbonyl group (1634 cm⁻¹). The compounds have similar crystal morphologies. This feature together with the similarity in the powder X-ray diffraction patterns of the neodymium and holmium compounds (Figure 1) is indicative of their isomorphic nature. However, the spectrum of the lutetium compound (Figure 1), although it is qualitatively similar to those of the neodymium and holmium compounds, lacks the strong diffracting behavior

at low θ . This may indicate some subtle changes in the detailed crystal structure as a result of the lanthanide contraction. Thus single-crystal X-ray diffraction analyses of both the neody-mium—holmium and erbium—lutetium complexes are necessary to provide firm evidence for the nature of members of series **4**. Indeed single-crystal X-ray diffraction studies of the neody-mium(III) and europium(III) products confirm the formulation of the compounds as $[{Na}{(sal)_3Ln(B(OMe)_4)Ln(sal)_3}]_n$ as well as their one-dimensional structural motif (Figures 2 and 3).

The X-ray analyses of the neodymium(III) [4 (Ln = Nd)]and europium(III) (5) complexes shows that they are isomorphous (Tables 2 and 3), with only very minor differences in their comparative geometries reflecting the normal lanthanide-(III) contraction effects. The structures are ionically linked polymers and contain two independent lanthanide centers. Each of these centers is coordinated to three planar bidentate salchelates, the remaining two coordination sites in each case being occupied by a bridging tetramethoxyborate. The geometry at each lanthanide center is slightly distorted square antiprismatic; the principle distortion is an O(2)-Ln-O(3) angle of 58.3(2)° (Ln = Eu) [53.2(2)°, Ln = Nd] as a consequence of the bite of the tetramethoxyborate. The Ln-O coordination distances fall into two distinct groups. The coordination distances of the phenolate oxygen atoms [ranging between 2.286(7) and 2.340-(6) Å, Ln = Nd] are, as expected,¹⁶ significantly shorter than those of either the carbonyl [range, 2.397(6)-2.466(7), Ln = Eu; 2.437(6) - 2.522(8), Ln = Nd] or tetramethoxyborate oxygens [2.424(6)-2.457(5), Ln = Eu; 2.461(7)-2.502(6), Ln= Nd]. In all cases, the carbonyl groups substantially retain their double-bond character,¹⁶ the C=O distances being in the ranges 1.22(1)-1.24(1) Å (Ln = Eu) and 1.20(1)-1.24(1) Å (Ln = Nd). Likewise, the C-O (phenolate) bonds exhibit a characteristic slight shortening,¹⁶ the C–O distances being in the ranges 1.29(1)-1.30(1) Å (Ln = Eu) and 1.28(1)-1.34(1)Å (Ln = Nd). The bridging tetramethoxyborate ligands show marked departures from T_d geometry, the O-B-O angles ranging between 98.3(7) and $117.2(8)^{\circ}$ (Ln = Eu) and between 98.6(7) and 117.0(8) (Ln = Nd). These distortions are similar to those observed in the linear $[La(NO_3)_2(B(OCH_3)_4)(CH_3OH)_2]_n$ system.4

Each of the tetramethoxyborate-bridged ([(sal)₃Ln(B-(OMe)₄)Ln(sal)₃]⁻) 'dimers' has approximate C_2 symmetry, virtually down the molecular C_2 axis, as shown in Figure 2. The three phenolate oxygen atoms on each center are directed outward, away from the bridging tetramethoxyborates, thereby presenting a suitable charged binding site for the sodium cations. Indeed, pairs of [(sal)₃Ln(B(OMe)₄)Ln(sal)₃]⁻ dimers are linked via the sodium cation to form sinusoidal ionically linked polymeric chains (Figure 3) that extend in the crystallographic *c*-direction. The two linking sodium cations are nonequivalent, Na(1) positioned on a crystallographic inversion center and Na-(2) on a 2-fold axis.

The intrachain Ln...Ln separations are (given to 2 decimal points to highlight the magnitude of their differences) 6.23 Å, Ln = Eu (6.31 Å, Ln = Nd) within each dimer, and 6.57, 6.84 Å, Ln = Eu (6.61, 6.88 Å, Ln = Nd), respectively, across the linking Na⁺ centers. These distances are all shorter than any of the interchain Ln...Ln distances, the shortest of which is 7.22 Å, Ln = Eu (7.28 Å, Ln = Nd), for the Ln(1)....Ln(1) separation. The polymer chains are closely packed with the crests of one

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Figure 2. ORTEP diagram of the [(sal)₃Eu(B(OMe)₄)Eu(sal)₃]⁻ anion showing 50% probability thermal ellipsoids for all nonhydrogen atoms.



Figure 3. The one-dimensional sinusoidal motif of the $[{Na}_{(sal)_3Ln}(B(OMe)_4)Ln(sal)_3]_n$ ionically linked polymeric chain. Hydrogen and carbon atoms have been omitted for clarity.

chain lying in the troughs of the next. There is no solvent included in the structures of 4(Ln = Nd) and 5.

Although the common synthetic route to complexes of $[(RO)_4B]^-$ with groups 1 and 2 metal ions has been via the reaction of $(RO)_3B$ with the RO⁻ anion under strict exclusion of moisture and air,^{2,3} methanolysis of commonly used reagents such as sodium tetrahydroborate and trihydrocyanoborate (leading to 1),⁴ tetraphenylborate (resulting in 2),⁵ and in this study, Borax, seems to proceed well without those limitations.

Series 1 was obtained in a mixture with amino-imino products but attempts to obtain pure products of the series or complexes of smaller lanthanides (Ln = Sm-Lu, Y) failed.⁴ In sharp contrast, the formation of 4 with Ln = large lanthanides(La and Pr) seems unfavorable, whereas compounds of the smaller $Nd^{3+}-Lu^{3+}$ are readily obtained. This preference for smaller Ln^{3+} ions by system 4 was revealed conclusively by the mixed lanthanide(III) complexes, $([{Na}_{(sal)_3Ln_{1-x}Eu_x} (B(OMe)_4)Ln_{1-x}Eu_x(sal)_3]_n [x = 1 (5); Ln = Gd (6-10), Y$ (12) or Dy (13-18)], the metal contents of which were established by nuclear activation analyses.12 The higher selectivity for incorporation of the smaller dysprosium(III) ions (1.03 Å for 8-fold coordination¹⁷) in crystalline compounds [{Na}-{ $(sal)_3Dy_{1-x}Eu_x(B(OMe)_4)Dy_{1-x}Eu_x(sal)_3$ }]_n over the significantly larger europium(III) (1.07 Å) is evident from the cation discrimination index¹⁸ of $I_{\rm D} \approx 1.3$ (Figure 4). For system



Ln = Dv

Figure 4. Cation discrimination indices for crystalline ([{Na}-{ $(sal)_3Ln_{1-x}Eu_x(B(OMe)_4)Ln_{1-x}Eu_x(sal)_3$]_n complexes [Ln = Gd (6–10) or Dy (13–18)].

 $[{Na}{(sa)_3Gd_{1-x}Eu_x(B(OMe)_4)Gd_{1-x}Eu_x(sa)_3}]_n$, in which Gd^{3+} and Eu^{3+} have more similar ionic radii (1.06 Å for Gd^{3+} in 8-fold coordination¹⁷) the cation discrimination index accordingly is $I_D \approx 1$ (Figure 4). Further, lanthanide(III) ions have been complexed with tetrahydroborate anion $[H_4B]^-$ numerous times in the presence of CH_3O^- anions. Although some limited methanolysis of the tetrahydroborate anion was reported, no lanthanide(III) compounds of the $[(CH_3O)_4B]^-$ anion or its

$$x_{\rm A} = (X_{\rm A}P_{\rm A})/(X_{\rm A}P_{\rm A} + X_{\rm B}P_{\rm B}); \quad x_{\rm B} = (X_{\rm B}P_{\rm B})/(X_{\rm A}P_{\rm A} + X_{\rm B}P_{\rm B})$$
(1)

where X_A and X_B are the initial mole fractions of cations A and B, respectively, in the mother-reaction mixture; x_A and x_B are mole fractions of cations A and B in the crystalline product, respectively; P_A and P_B are probabilities for the incorporation of cations A and B, respectively, in the crystalline product. For comparative purposes, we set cation $A = Eu^{3+}$ and cation $B = Gd^{3+}$ or Dy^{3+} and rearranged eq 1 to eq 2:

$$I_{\rm D}(X_{\rm B}/X_{\rm A}) + 1 = 1/x_{\rm A} \tag{2}$$

- where $I_{\rm D} = (P_{\rm B}/P_{\rm A})$ = discrimination index of cation B over A for incorporation into the crystalline product.

⁽¹⁸⁾ Competitive incorporation of two lanthanide(III) metal ions (A and B) in a mixed crystalline product can be expressed in terms of their respective selectivities:¹²

Table 3. Temperature Evolution of Decay Rates^{*a*} for $[\{Na\}\{(sal)_3 \in u_x Ln'_{1-x}(B(OMe)_4 \in u_x Ln'_{1-x}(sal)_3\}]_n$

<i>T</i> (K)	$Ln = k_1/s \times 10^4$	Eu $k_2/s \times 10^3$	$Ln = k_1/s \times 10^4$	$\frac{\mathrm{Gd}^{b}}{k_{2}/\mathrm{s}\times10^{3}}$	$Ln = k_1/s \times 10^4$	$rac{\mathrm{Y}^c}{k_2/\mathrm{s} imes10^3}$	$Ln = k_1/s \times 10^4$	${ m Dy}^d$ $k_2/{ m s} imes 10^3$	$Ln = k_1/s \times 10^4$	${ m Gd/Dy^e}\ k_2/{ m s} imes 10^3$
8-50		2.2		2.2		2.2	1.0			
55		2.2								
60				2.2		2.2	1.0	3.1	1.1	3.4
65		2.2					1.0	3.2		
70						2.2				
75							f	3.3		
80				2.2			-		1.0	4.0
85		2.2				2.3				
100				2.2			1.1	3.6	1.1	4.0
105		2.2				2.3				
110				2.2						
120							1.1	3.7		
125		2.2				2.3				
140				2.3		2.3	1.2	4.0	1.1	4.1
145		2.2								
155						2.3				
160										
165				2.2		2.3	1.4	4.0		
180							f	4.1	1.2	5.5
185		2.2						1.0		
200		2.2		2.2		2.4	1.4	4.2		
205		2.2				2.4				
210							15	12		
213		2.2		2.4		2.4	1.5	4.5	1.0	6.2
220		2.5		2.4		2.4			1.0	0.2
225		2.4				2.5				
240		2.5		27		2.7	17	4.4		
245		3.0		3.0		2.0	1.7	7.7		
250	1.8	3.4		3.5		3.3	1.7	5.1		
255	f	3.8		3.7		3.5	117	011		
260	2.0	4.4		4.1	1.5	3.6	1.9	5.9	1.4	7.5
265	2.0	5.3		4.5	f	4.1				
270	f	6.3	1.7	5.1	Ĭ.8	4.2	2.0	6.5	1.6	7.9
275	2.3	7.4			2.2	5.5	2.5	7.7	2.5	8.8
280	2.5	9.1	2.4	5.8	2.5	5.7	f	8.7	2.4	7.9
285	2.4	11.1			3.3	7.4	3.0	10.4	3.0	8.8
290	3.4	13.2	2.9	10.4	3.4	7.2	3.6	12.0	3.4	9.9
295	4.5	17.5			5.0	11.8	5.9	15.2	3.7	10.5
297	6.7	16.7	4.3	10.0						
300	8.3	20.4					8.3	19.6	4.0	11.8
310	9.1	28.6								
320	14.3	40.0								

^{*a*} Tails of decay curves were fit as single exponentials; the fast components were derived from double-exponential approximations. ^{*b*} x = 0.06. ^{*c*} x = 0.04. ^{*d*} x = 0.58. ^{*e*} Gd_{0.29}Dy_{0.42}Eu_{0.29}. ^{*f*} Nonlinear Marquardt fit not satisfactory.

derivatives were isolated.¹⁹ The contrasting synthetic results of series 1 and 4 are intriguing. Experiments designed to determine the influence of thermodynamic and kinetic stabilities of corresponding Ln^{3+} coordination polyhedra as well as extended supramolecular interactions are in progress. This dearth of information notwithstanding, the consistent formation of complexes of 4 in high yield and over much of the lanthanide(III) series is very encouraging. Better understanding of the coordination behavior of $[(RO)_4B]^-$ anions and exploitation of their versatility as general synthons for one-dimensional materials should be of great future interest.

Luminescence Spectroscopy

Structurally, compounds of the $[{Na}{(sal)_3Ln(B(OMe)_4)}-Ln(sal)_3}]_n$ series are unequivocally one-dimensional. But the tight trough-into-crest crystal packing of the sinusoidal chains results in some interchain distances (shortest, 7.22 Å for 5) that are comparable with the intrachain ones (longest, 6.84 Å for 5). It was thus interesting to determine whether intrachain Ln³⁺-Ln³⁺ electronic coupling resulting from through-bond interac-

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tions provides a better conduit for energy transfer than interchain electronic coupling facilitated by crystal packing forces and through-space interactions. Stronger intrachain coupling would constrain excitation energy transport in [{Na}{(sal)_3Ln(B(OMe)_4)-Ln(sal)_3}]_n to one dimension, otherwise multidimensional energy transport would dominate.²⁰ We thus studied in detail the luminescence spectral and decay dynamical properties of the emitting [{Na}{(sal)_3Eu(B(OMe)_4)Eu(sal)_3}]_n (5) chains and their mixed metal analogues 7–19 (Table 1) to establish the nature and dynamics of electronic interactions in this structural motif.

The gadolinium(III) compound (6) shows broad ligand (sal⁻) absorptions that lead to complex emission (350-600 nm) (Figure 5) at both 297 and 77 K. This ligand emission overlaps

$$k = C/R^n \tag{3}$$

where C = multipolar coupling constant responsible for energy transfer; n = the order of the multipolar interaction (6, 8, or 10 for dipole–dipole, dipole–quadrupole, and quadrupole–quadrupole interactions, respectively), and R = metal–metal separation.

⁽²⁰⁾ When energy transfer from a donor to an acceptor is predominantly multipolar,²¹ the transfer rate (*k*) is expressed as:



Figure 5. Electronic spectra for compounds 5 and 6.

with electronic absorptions of luminescent Tb³⁺ and Eu³⁺ ions thereby meeting the principal Dexter-Förster condition²¹ for ligand-to-metal energy transfer which may culminate in the 'antenna effect'. When excited in the ligand absorption region $(\lambda_{\text{exc}} = 360 \text{ nm})$, the europium(III) complex (5) yields red Eu³⁺ (⁵D₀) emission, albeit rather weak (Figure 5). The excitation spectra of Eu³⁺ (⁵D₀) emission ($\lambda_{em} = 612$ nm) are dominated by ligand absorptions but very strange differences exist which depend on both temperature and europium(III) ion concentration (Figure 6). For example, at 297 K, the excitation spectrum of Eu^{3+} (⁵D₀) emission ($\lambda_{em} = 612$ nm) from the pure europium-(III) compound **5** features prominent direct Eu³⁺ ${}^{5}D_{1} \leftarrow {}^{7}F_{i}$ (*j* = 0, 1) and weaker ${}^{5}D_{2} \leftarrow {}^{7}F_{i}$ (*i* = 0, 1) transitions superimposed on a broad sal⁻ absorption tail which extends into the green region ($\lambda > 480$ nm) (Figure 6, compound 5). In sharp contrast to this behavior, the prominence of these direct europium(III) ${}^{5}D_{1,2} \leftarrow {}^{7}F_{0,1}$ transitions decreases systematically with decreasing concentration of europium(III) ions in crystals of the mixed $[{Na}{(sal)_{3}Gd_{1-x}Eu_{x}(B(OMe)_{4})Gd_{1-x}Eu_{x}(sal)_{3}}]_{n}$ (7–9) (Figure 6). Concomitantly, the low-energy sal⁻ green absorption tail takes on a systematic blue shift with decreasing europium-(III) concentration and becomes insignificant at $\lambda \approx 480$ nm for compound 7 (Figure 6, compound 7) where the Eu^{3+} concentration is the least. The 77 and 297 K excitation spectra of 7 and 12 (Figure 6) and that of 5 at 77 K (Figure 5) are similar. Further, the ${}^{5}D_{0} \rightarrow {}^{7}F_{i}$ emission spectrum (Figure 5) is



Figure 6. The 297 K excitation spectra for compounds 5, 7–9, and 12.

dominated by the so-called hypersensitive²² ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$ peak and, quite unusually, the forbidden ${}^{5}D_{0} \rightarrow {}^{7}F_{0}$, ${}^{7}F_{3}$ transitions²³ are more intense than the allowed ${}^{5}D_{0} \rightarrow {}^{7}F_{4}$ electric dipole. It seems that the geometrical disposition of the various oxygen atoms around the 8-fold coordinated Eu³⁺ ion produces an effective crystal electric field of very low symmetry; hence, the observed Eu³⁺ (${}^{5}D_{0}$) emission spectral profile in the 580–630 nm region (Figure 5). However, symmetry and group theory arguments that allow for weak ${}^{5}D_{0} \rightarrow {}^{7}F_{3}$ intensity do simultaneously predict greater probabilities for electric dipole ${}^{5}D_{0} \rightarrow {}^{7}F_{j}$ (j = 2, 4, 6) transitions.²²

Thus, arguments based on symmetry of the coordination polyhedron around Eu^{3+} and impurity emission cannot explain satisfactorily the unusual intensities of the ${}^{5}D_{0} \rightarrow {}^{7}F_{3}$ transition.

To shed more light on the nature of the emitting europium-(III) species and the energy-transport dynamics in the sinusoidal $[{Na}{(sal)_3Ln(B(OMe)_4)Ln(sal)_3}]_n$ chains, we studied the temperature evolution of the decay behavior of the resulting emission in detail and found that the luminescence decay kinetics fall into two regimes (Table 3): low (8 < *T* < 220 K) and high (*T* > 220 K) temperature.

Luminescence Decay Dynamics of Eu^{3+} (⁵D₀) at Low **Temperature (**T < 220 K). Direct excitation of the Eu³⁺ (⁵D₀) level at 572 nm in 7 and 12, in which Eu^{3+} ions are well separated by the effectively inert Gd³⁺ and Y³⁺ ions, and in the pure europium compound 5 produces red emission that decays exponentially for 8 < T < 220 K (Table 3). The spontaneous decay rate is ca. 2.2×10^3 s⁻¹ whereas a weak temperature dependence corresponding to a phonon-type thermal barrier of ca. 100 cm⁻¹ is observed. The 77 K decay curves of Eu^{3+} (⁵D₀) emission from other mixed gadolinium(III)/europium(III) compounds 8-11 are also exponential with similar decay rates $(2.2 \times 10^3 \text{ s}^{-1})$. Indiscriminate (at 365 nm) and direct (at 510–535 nm) excitation of the Eu³⁺ (⁵D₁) level in 5 produces an excitation buildup of ca. 7.7×10^5 s⁻¹ which is of the same order of magnitude as the decay rate of the Eu^{3+} (⁵D₁) emission (ca. 7.7 \times 10⁵ s⁻¹). This means that ligand-to-Eu³⁺ $({}^{5}D_{1})$ energy transfer at low temperature is faster than the Eu³⁺ $({}^{5}D_{1})$ -to-Eu $^{3+}$ $({}^{5}D_{0})$ energy transfer, and confirms the involvement of Eu^{3+} (⁵D₁) in the sensitization of Eu^{3+} (⁵D₀) emission that is inferable from the excitation spectra (Figures 5 and 6).

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Figure 7. The 77 and 297 K (insert) decay curves ($\lambda_{em} = 612$ nm; $\lambda_{exc} = 572$ nm) of Eu³⁺ (⁵D₀) in mixed metal europium/dysprosium compounds **5** and **13–18**.



Figure 8. Stern–Volmer plots [ratio of decay rates vs mole fraction of Ln acceptor (δ) for the major Eu³⁺ (⁵D₀) emission in mixed metal europium/dysprosium compounds **13–18**]. ${}^{a}k_{Dy}/k_{Eu} = 1 + 4.54 \delta_{Dy} [k_{Eu} = decay rate of Eu³⁺ (⁵D₀) in$ **5** $at 77 K; <math>k_{Dy} = decay rate of Eu³⁺ (⁵D₀) in$ **7** $at 297 K; <math>\delta_{type2Eu} = \delta_{Eu}$ for comparative purposes; $k_{type2Eu} = decay$ rate of long-lived component of Eu³⁺ (⁵D₀) in **5** at 297 K; $k_{type2Eu} = 1 - 0.59 \delta_{Dy} [k_{Eu} = decay rate of Eu³⁺ (⁵D₀) in$ **5** $at 297 K; <math>k_{type2Eu} = decay$ rate of long-lived component of Eu³⁺ (⁵D₀) in **5** at 297 K; $k_{type2Eu} = decay$ rate of long-lived component of Eu³⁺ (⁵D₀) of compounds **13–18** at 297 K].

The exponential nature of decay curves of quenched Eu³⁺ (⁵D₀) emission is indicative of efficient energy migration on the Eu³⁺ sublattice. To establish the dimensionality of this regime in crystals of [{Na}{(sal)_3Eu(B(OMe)_4)Eu(sal)_3}]_n, we studied the luminescence decay behavior of mixed Eu³⁺/Dy³⁺ compounds **13–18** (Figure 7). Quenching of 77 K Eu³⁺ (⁵D₀) emission by Dy³⁺ ions (compounds **13–18**) follows Stern–Volmer kinetics²⁴ (Figure 8; 77 K), the decay curves deviating marginally from single-exponential behavior throughout the entire concentration range of Dy³⁺ studied. This confirms that the energy-transport regime in **5** and **13–18** is dominated by efficient energy migration that culminates in less efficient quenching by traps (e.g., Dy³⁺). Eu³⁺–Eu³⁺ interactions are therefore much stronger than the Eu³⁺–Dy³⁺ interactions, and



Figure 9. Decay curves of Eu³⁺ (${}^{5}D_{0}$) emission ($\lambda_{em} = 612$ nm; $\lambda_{exc} = 572$ nm) in **5** and mixed metal compounds **9**, **17**, and **19**.

the energy-transport regime is nearly dynamic.^{25,26} It is well known that introduction of energy scatterers into donor– acceptor arrays dominated by one-dimensional energy-migration regimes produces a dramatic effect on decay-curve profiles.^{27,28} However, neither the decay curves of **17** and **19** nor those of **5** and **9** show the expected dramatic differences at 77 K (Figure 9 insert). This means that energy migration in [{Na}{(sal)_3Eu-(B(OMe)_4)Eu(sal)_3}]_n is multidimensional, and through-space multipolar Eu³⁺–Eu³⁺ interactions are more efficient conduits of excitation energy than through-bond interactions of the sinusoidal chains. The temperature evolution of luminescence decay rates of compounds **17** and **19** were very similar (Table 3); both exhibited weak temperature dependence for 50 < *T* < 220 K with a phonon-type thermal barrier of roughly 40 cm⁻¹.

From the slope of the 77 K Stern–Volmer plot (Figure 8; 77 K), the limiting Eu³⁺ (⁵D₀)-to-Dy³⁺ energy-transfer rate is ca. $9 \times 10^3 \text{ s}^{-1}$. This parameter is comparable with the value of ca. $6 \times 10^3 \text{ s}^{-1}$ obtained from perfectly exponential decay curves of Eu³⁺ (⁵D₀) in compound **13**. Assuming that Eu³⁺ (⁵D₀)-to-Dy³⁺ energy transfer is predominantly dipolar, the dipole–dipole coupling constant responsible for the transfer is ca. $4 \times 10^{-52} \text{ m}^6 \text{ s}^{-1}$ which is reasonable.²⁶

Luminescence Decay Dynamics of Eu³⁺ (⁵D₀) at High Temperature (T > 220 K). After direct excitation at 572 nm, luminescence from isolated Eu³⁺ (⁵D₀) states of compounds 7 and 12 decays exponentially for 220 < T < 260 K. However, the decay behavior exhibits a strong temperature dependence throughout the 220 < T < 300 K region and is not exponential for T > 260 K. For both 7 and 12, the thermal barriers for quenching processes of the long- and short-lived red Eu³⁺ (⁵D₀) emission (monitored at 612 nm) are ca. 2.1 × 10³ (Figure 10) and 3.1 × 10³ cm⁻¹, respectively. Eu³⁺ (⁵D₀) emission decay behavior of crystalline 5 is similar to that of 7 and 12 up to 250 K. Thereafter, a stronger temperature dependence is observed with decay curves deviating very marginally from single exponential behavior (Figure 11) because of the appearance of a short-lived component (Table 3). The long-lived

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Figure 10. The temperature evolution of the trapping rates of the major (slowly decaying) Eu³⁺ (⁵D₀) emission component for **5**, **7**, and **12** in the high-temperature regime (210 < T < 300 K). k_s = spontaneous decay rate = 2.2×10^3 s⁻¹; k_t = trapping rate = $k_o - k_s$ (k_o = observed rate).



emission ($\lambda_{em} = 612 \text{ nm}$; $\lambda_{exc} = 572 \text{ nm}$) from **5** (8–320 K). component decay is single exponential over six lifetimes. Thermally activated quenching of both the fast- and slowdecaying components require similar activation energies, ca. 3.1 \times 10³ and 3.3 \times 10³ cm⁻¹ (Figure 10), respectively. To shed more light on the nature of the thermally activated quenching process and its dependence on the nature of Ln in [{Na}- $\{(sal)_3Ln(B(OMe)_4)Ln(sal)_3\}]_n$ we studied the mixed compounds 13-18. In sharp contrast to the behavior at 77 K, the 297 K decay curves of Eu³⁺ (⁵D₀) emission from mixed dysprosium/europium compounds 13-18 (Figure 7, insert) are not single exponential and have initial slopes similar to that of 5; but those compounds with the most dysprosium traps have the most intense Eu^{3+} (⁵D₀) emission (Figure 7, insert). As a result, the long-lived component of Eu^{3+} (⁵D₀) emission from compounds 5 and 13-18 at 297 K exhibits an extraordinary anti-Stern-Volmer dependence on the concentration of Dy³⁺ (Figure 8). This anomaly is corroborated by the comparative luminescence decay behavior of 5, 9, 17, and 19 (Figure 9) in which the initial decay rates are of similar magnitude, and the tail decay rates are in the unexpected order of 9 < 19 < 17 < 175. The activation energy for thermally activated quenching of the short- and long-lived Eu³⁺ (${}^{5}D_{0}$) emission from 17 (220 <

T < 300 K) are very similar to those of **5** (ca. 3×10^3 cm⁻¹). That of **19** (ca. 1.6×10^3 cm⁻¹) is more similar to those of **7** and **12** (ca. 2.1×10^3 cm⁻¹).

To explain the above-mentioned results, we propose that there are two types of Eu³⁺ ions in [{Na}{(sal)₃Eu(B(OMe)₄)Eu- $(sal)_3$]_n (5) chains. Type 1 Eu³⁺ occupies the regular europium-(III) sites in chains of 5 as shown in Figure 3 and is responsible for the Eu^{3+} (⁵D₀) excitation profile exhibited by compounds **5** at 77 K (Figure 5) as well as 7 and 12 at 297 K (Figure 6). **Type 2** Eu^{3+} sites are crystal defects, the population of which increases with the concentration of europium in compounds 5-12 and 13-19. These defect sites would in this case be responsible for the dominant ${}^5D_{1,2} \leftarrow {}^7F_{0,1}$ transitions and the green absorption tail ($\lambda > 480$ nm) seen in Figure 6 (compounds **5** and 7-9). There is good spectral overlap between the broad 500-600 nm emission tail (Figure 5) and the Eu^{3+ 5}D_{0.1.2} \leftarrow $^{7}F_{0.1}$ absorptions (Figure 6); hence, the strong absorption tail at 480 nm in the excitation spectra (Figure 6) is most likely associated with the defect site. Most probably Type 2 Eu^{3+} arises from swapping Eu³⁺ and Na⁺ sites, which is reasonable, in view of similarities in their cationic radii (ca. 1 Å for 6-fold coordination).¹⁷ It is unlikely that the defect **Type 2** Eu^{3+} site is a hydrate, because samples prepared from deuterated solvents (CH₃OD and D₂O) did not exhibit the dramatic decrease in luminescence decay rates usually seen when coordinated H₂O is replaced by D₂O.^{9b,29} Instead, differences in **Type 1** and defect **Type 2** sites are expected to arise from the differential in site symmetry (at most D_{4d} and ca. O_h , respectively) and the nature of coordinated atoms (vide supra). The main Eu^{3+} (⁵D₀) emission in 5 is caused by Type 1 Eu^{3+} and is sensitized as well as quenched via back energy transfer by thermally activated processes involving both the ligand absorption (at 480 nm) and **Type 2** Eu^{3+} . This explains well the excitation spectral features (Figure 6, compounds 5 and 7-9) and the thermally activated quenching processes (Figure 11 and Table 3) as well as their dependence on the concentration of europium(III). For comparative purposes, when a simple proportionality between the population of Type 2 Eu^{3+} and the europium concentration in the crystals is assumed and the Stern-Volmer plot revised accordingly, the normal profile is obtained (Figure 8).

Within this scheme, the short-lived component may be attributed to direct **Type 1 Eu³⁺**-to-**Type 2 Eu³⁺** energy transfer, the moderate efficiency of which precludes complete homogenization of the europium ion population by energy migration. The thermal barriers of about 2×10^3 and 3×10^3 cm⁻¹ correspond to the energies required to bridge the Eu^{3+} (⁵D₀) with the upper Eu^{3+} (⁵D₁) and the ligand state responsible for the absorption at 480 nm, respectively. Most likely, back energy transfer to the ligand states associated with the Type 2 Eu³⁺ site predominantly quenches Eu³⁺ (⁵D₀) emission from compounds rich in Eu^{3+} . Transiting the Eu^{3+} (⁵D₁) state may be essential for quenching of Eu^{3+} (⁵D₀) emission for compounds with too few Type 2 Eu³⁺sites such as 7, 12, and 19. It is remarkable that the green tail overlaps well with the Tb^{3+} (⁵D₄) state; this may be the reason for efficient quenching of terbium-(III) emission from compound $[{Na}_{(sal)_3}Tb(B(OMe)_4)Tb (sal)_3$]_n.³⁰

General Remarks

We have shown conclusively how hitherto elusive onedimensional Ln^{3+} compounds of the $[(CH_3O)_4B]^-$ chelate can

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be prepared consistently and throughout much of the lanthanide series. The formation of both linear (1) and sinusoidal (4–19) systems suggests that a variety of one-dimensional motifs may be accessible through use of the $[(CH_3O)_4B]^-$ chelate to bridge Ln^{3+} ions. The challenge seems to be finding suitable complementary counteranions to stabilize one-dimensional compounds of desired Ln^{3+} ions. The luminescence spectra and energy-transport dynamics exhibited by compounds 4–19, especially the dominance of through-space against through-bond Eu^{3+} – Eu^{3+} electronic interactions, are intriguing and require better understanding. For these reasons and the extended structural elegance of 1, 2, and 4–19, further synthetic and luminescence work is in progress.

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Supporting Information Available: Tables S1–S11 lists elemental analyses for 4–19, atomic positions, bond angles and distances, anisotropic displacement parameters, and hydrogen coordinates for 4 (Ln = Nd) and 5; Figure S1 shows the stereoscopic view of 5; Figure S2 shows the ORTEP diagram of 4 (Ln = Nd). This material is available free of charge via the Internet at http://pubs.acs.org.

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