# Weak Ferromagnetism in a Three-Dimensional Manganese(II) Azido Complex, $[Mn(4,4'-bipy)(N_3)_2]_n$ (bipy = Bipyridine)

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## Introduction

Several families of molecular materials exhibiting spontaneous magnetization below a critical temperature ( $T_c$ ) have been discovered in the past several years.<sup>1</sup> The design of high- $T_c$ ferromagnetic materials has been realized by organizing transition metal ions into three-dimensional networks by use of organic bridging ligands such as cyanide and tetracyanoethenide [TCNE]<sup>•-</sup>.<sup>2,3</sup> Another intriguing synthetic approach to molecular magnets utilizes azide (N<sub>3</sub><sup>-</sup>) as a bridging ligand, due to its coordinative versatility in bridging two transition metal ions.<sup>4</sup> Magneto-structural correlations for azido-based magnetic materials have shown that the end-to-end (EE) bridging mode predominantly leads to antiferromagnetic coupling, and the endon (EO) mode to ferromagnetic coupling.<sup>4</sup> Examples include dinuclear,<sup>5</sup> triangular,<sup>6</sup> tetragonal,<sup>7</sup> 1-D,<sup>8</sup> and 2-D<sup>9</sup> nickel(II) compounds.

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In an effort to assemble high-dimensional compounds with a greater local spin to enhance the magnetic ordering temperatures, the Mn(II)-azido system has been examined extensively. Escuer and co-workers systematically studied several [MnL2- $(N_3)_2]_n$  compounds: 3-D for L = pyridine<sup>10</sup> and 2-D for L = 3-<sup>11</sup> and 4-acetylpyridine,<sup>12</sup> methyl-<sup>13</sup> and ethylisonicotinate,<sup>10</sup> and 4-cyanopyridine.<sup>11</sup> Others have also reported highdimensional complexes by use of ancillary multidentate ligands such as bipyrimidine (2-D, EO),<sup>14</sup> pyrazine (2-D, EO),<sup>15</sup>1,2bis(4-pyridyl)ethane (2-D, EE),<sup>16</sup> and 4,4'-bipyridine (3-D, EE)<sup>17</sup> as opposed to monodentate ligands, which promotes additional binding of the metal ions. Composed solely of azido 1,3-bridging units, [NMe<sub>4</sub>][Mn<sup>II</sup>(N<sub>3</sub>)<sub>3</sub>] was found to possess a distorted perovskite-like structure.<sup>18</sup> Recent examples, [Ni(5-methylpyrazole)<sub>4</sub>(N<sub>3</sub>)][ClO<sub>4</sub>]<sup>8</sup> and [Mn(4-acpy)<sub>2</sub>(N<sub>3</sub>)<sub>2</sub>]<sub>n</sub> (acpy = acetylpyridine),<sup>11</sup> are weak ferromagnets, whereas all other EE-bonded materials show only antiferromagnetic coupling without a spontaneous moment.

A recent communication by Shen et al. reported the synthesis, crystal structure, and preliminary magnetic properties of [Mn-(4,4'-bipyridine)(N<sub>3</sub>)<sub>2</sub>]<sub>n</sub>, **1**.<sup>17</sup> Their magnetic data suggested that **1** did not magnetically order down to 4 K and that only weak antiferromagnetic coupling was present. We found this particularly striking in that many of the known Mn(II)—azido complexes do show interesting magnetic behavior well above 4 K. As part of our studies on transition metal—azido complexes,<sup>15</sup> we reexamined the magnetic properties of 3-D Mn-(4,4'-bipyridine)(N<sub>3</sub>)<sub>2</sub>]<sub>n</sub>, **1**, using detailed magnetic susceptibility and magnetization measurements, and we clearly observe weak ferromagnetism below  $T_c = 42.5$  K, contrary to a previous report.<sup>17</sup> This work emphasizes the importance of taking low-field magnetic data to characterize the magnetic behavior of materials.

#### **Experimental Section**

 $Mn(ClO_4)_2 \cdot 6H_2O$ , 4,4'-bipyridine, and NaN<sub>3</sub> were uses as purchased (Aldrich). Infrared spectra (400–4000 cm<sup>-1</sup>) were recorded on a Perkin-Elmer Spectrum 1000 FT-IR spectrometer with samples prepared as KBr pellets. C, H, and N elemental analysis were performed on an Elemental Anaylsen systeme GmbH Vario EL. Magnetic measurements were taken on a Quantum Design MPMS-5T magnetometer as previously described.<sup>19</sup>

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Figure 1. (a) Molecular structure of 1 showing only the  $Mn^{II}$  coordination sphere and atom-labeling scheme. (b) Perspective view of the crystal structure of  $Mn(4,4'-bipy)(N_3)_2$ , 1. Shaded, open, and filled spheres represent Mn, N, and C atoms, respectively.

[**Mn(4,4'-bipyridine)(N<sub>3</sub>)<sub>2</sub>]**<sub>n</sub>, **1**, was prepared in a way similar to that described in ref 17, by mixing solutions of Mn(ClO<sub>4</sub>)<sub>2</sub>·6H<sub>2</sub>O, 4,4'-bipyridine, and NaN<sub>3</sub> in methanol (61% yield). Single crystals suitable for X-ray diffraction study were obtained from slow addition of sodium azide in methanol with a syringe pump during 3 days into the methanol solution of manganese perchlorate and 4,4-bipyridine. Anal. Calcd for C<sub>10</sub>H<sub>8</sub>MnN<sub>8</sub>: C, 40.69; H, 2.73; N, 37.96. Found: C, 40.32; H, 2.71; N, 37.8. IR (KBr, cm<sup>-1</sup>):  $\nu$ (C–H), 3059 (w);  $\nu_{as}$ (N<sub>3</sub>), 2076 (vs);  $\nu_{s-1}$ (N<sub>3</sub>), 1323 (vw);  $\nu$ (4,4'-bipy), 1601 (s), 1411 (s), 815 (s). To verify that **1** was identical to that reported earlier,<sup>17</sup> we redetermined the structure; space group =  $P4_{1}2_{1}2$  (No. 92), M = 295.18, a = 8.236(1) Å, c = 16.937(3) Å, V = 1148.8(3) Å<sup>3</sup>, Z = 4,  $\rho$ (calcd) = 1.707 g cm<sup>-3</sup>, R1= 0.0309, wR2 = 0.0687.

### **Results and Discussion**

In light of our findings (vide infra), we verified the crystal structure<sup>17</sup> of  $\mathbf{1}$  as it is relevant to the interpretation of the observed magnetic behavior. The molecular structure of 1 showing the coordination environment of the Mn<sup>2+</sup> ion is given in Figure 1a, which has slightly elongated octahedral Mn<sup>II</sup> centers composed of four EE-bonded N3- ligands, and two trans-coordinated 4,4'-bipyridine molecules. When extended in 3-D, an unprecedented network structure is formed, Figure 1b. The IR spectrum shows a very strong band at 2076 cm<sup>-1</sup> for the  $v_{asym}(N_3)$  mode, which is consistent with a structure containing an end-to-end bridging ligand. The azide symmetric stretch,  $v_{sym}(N_3)$ , is not active for the symmetrical EE coordination mode. However, **1** has a very weak  $v_{sym}(N_3)$  band at 1323 cm<sup>-1</sup> due to asymmetry of the EE bridges in the complex. A similar result was observed in the  $[Ni(2,2'-bipy)_2(N_3)]_n(ClO_4)_n$ chain complex.<sup>21</sup>

We verified the crystal structure of 1,<sup>17</sup> which has slightly elongated octahedral Mn<sup>II</sup> centers composed of four EE-bonded N<sub>3</sub><sup>-</sup> ligands, and two *trans*-coordinated 4,4'-bipyridine molecules, which, when extended in 3-D, afford an unprecedented network structure, Figure 1. The tetragonal symmetry (*P*4<sub>1</sub>2<sub>1</sub>2)

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**Figure 2.**  $\chi T(T)$  ( $\bigcirc$ ) and reciprocal molar magnetic susceptibility,  $1/\chi$ -(T) ( $\bullet$ ), for **1** at 500 Oe.

of the complex precludes a center of inversion. A variety of channel sizes and shapes are produced including triangular, rectangular, and hexagonal. Interestingly, solvent molecules do not intercalate into the cavities nor does a second interpenetrating network form. The 4,4'-bipyridine ligands are not coplanar but are twisted about the C–C bond by 31.9(2)°, and the Mn– $N_{bipy}$  bond distances are inequivalent [2.256(4) and 2.304(4) Å]. The two kinds of Mn– $N_{azide}$  bond distances are also different from each other [2.188(4) and 2.208(4) Å], which is consistent with the observation of the symmetric stretching mode in IR. The intranetwork Mn•••Mn separations are 5.941 (via azide) and 11.647 Å (via 4,4'-bipy).

The 2–300 K magnetic susceptibility,  $\chi$ , of a polycrystalline sample of **1** could be fit to a Curie–Weiss law above 60 K with g = 2 and  $\theta = -100$  K, indicating strong nearest-neighbor antiferromagnetic interactions, Figure 2. At 300 K,  $\chi T(T)$  has a value of 3.86 emuK/mol and decreases upon cooling due to antiferromagnetic coupling. A sharp increase in  $\chi T(T)$  occurs near 45 K whereby a sudden increase is observed, reaching a maximum value of 6.91 emu K/mol at ~30 K. This behavior is characteristic of a spontaneous moment due to spin canting. Below ~30 K,  $\chi T(T)$  decreases abruptly down to 2 K owing to increasing antiferromagnetic interactions and/or saturation effects. Moreover, we and others have observed analogous

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 Table 1. Summary of the Structure and Magnetic Properties of Representative Extended Azido Structures<sup>a</sup>

compound <sup>a</sup>	bridging mode	magnetic behavior	<i>J</i> , K	<i>T</i> <sub>c</sub> , K	ref
Two-Dimensional (2-D)					
$[Mn_2(bpm)(N_3)_4]_n$	EO	ferro-antiferro <sup>d</sup>	g	g	14
$[Mn(pyz)(N_3)_2]_n$	EO	ferro-antiferro <sup>d</sup>	0.61	<2	15
$[Mn(etinc)_2(N_3)_2]$	EO-EE	ferro-antiferro <sup>b</sup>	g	g	10
		weak ferro <sup>c</sup>	g	16	11
$[Mn(4-cnpy)_2(N_3)_2]_n$	EO-EE	weak ferro <sup>c</sup>	g	18	11
$[Mn(3-acpy)_2(N_3)_2]_n$	EO-EE	weak ferro <sup>c</sup>	g	16	11
$[Mn(minc)_2(N_3)_2]_n$	EO-EE	antiferro <sup>c</sup>	-3.23	g	13
$[Mn(4-acpy)_2(N_3)_2]_n$	EE	antiferrof	-5.52	g	12
		weak ferro <sup>c</sup>	g	28	11
$[\mathrm{Mn}(\mathrm{bpe})_2(\mathrm{N}_3)_2]_n$	EE	antiferrof	-18	g	16
Three-Dimensional (3-D)					
$[Mn(py)_2(N_3)_2]_n$	EE	antiferro <sup>b</sup>	-1.94	g	10
		weak ferro <sup>c</sup>	g	40	11
$[Mn(4,4'-bipy)_2(N_3)_2]_n$	EE	antiferrof	g	g	17
- · · · · · · · · · · · · · · · · · · ·		weak ferro <sup>e</sup>	g	41	h
$[N(CH_3)_4][Mn(N_3)_3]_n$	EE	antiferrof	-2.5	g	18

<sup>*a*</sup> etinc = ethyl isonicotinate, 4-cnpy = 4-cyanopyridine, 3-acpy = 3-acetylpyridine, 4-acpy = 4-acetylpyridine, minc = methyl isonicotinate, bpm =2,2'-bipyrimidine, pyz = pyrazine, bpe = 1,2-bis(4pyridyl)ethane, py = pyridine, bipy = bipyridine. <sup>*b*</sup> 15 kOe. <sup>*c*</sup> 100 Oe. <sup>*d*</sup> 1000 Oe. <sup>*e*</sup> 500 Oe. <sup>*f*</sup> Magnetic field was not reported. <sup>*k*</sup> Not reported. <sup>*h*</sup> This work.

phenomena in the canted antiferromagnet  $Mn^{II}[N(CN)_2]_2$ .<sup>22</sup> Compound **1** was previously reported as an antiferromagnet.<sup>17</sup> However, for a weak ferromagnet due to spin canting, the magnetic behavior is quite field-dependent. As summarized in Table 1, some Mn–azide compounds showing characteristic antiferromagnetic behavior at higher magnetic field exhibit weak ferromagnetism at lower magnetic field.

Long-range magnetic order is clearly evident from low-field magnetization and ac susceptibility measurements. The divergence of the zero-field (ZFC) and field-cooled (FC) M(T) data below  $T_c = 42.5$  K indicates irreversible behavior arising from the formation of a magnetically ordered state. The presence of a nonzero  $\chi_{ac}''(T)$  component in the  $\chi_{ac}(T)$  data suggests hysteretic effects and an uncompensated moment, Figure 3.  $\chi_{ac}'-(T)$  exhibits behavior analogous to that of the ZFC/FC data as expected in the low-field dc limit. Small magnetization values at 2 K indicate the presence of a weakly spin canted (noncollinear) moment. The strong magnetic exchange observed for **1** is attributed to the Mn–NNN–Mn interaction as opposed to exchange via the 4,4'-bipy ligand, which would lead to a more 2-D magnetic behavior.

The 2 K isothermal magnetization displays an unusual hysteresis curve, Figure 4. Initially, the magnetization rises rapidly to a value of ~200 emu Oe/mol ( $H_{dc} < 500$  Oe). Beyond this M(H) rises quasi-linearly to ~1.2 T, whereby an inflection point is observed with an increased slope which is likely attributed to a spin flop transition. Further evidence for a spin reorientation comes from a clear discontinuity in the dM(H)/dH data, Figure 5, although single-crystal magnetization measurements are needed to confirm this. At the highest field measured (50 kOe), a low magnetization value of 1750 emu Oe/mol is achieved and saturation does not occur. A coercive

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**Figure 3.** Temperature dependencies of the in-phase (open symbols),  $\chi'$ , and out-of-phase (filled symbols),  $\chi''$ , ac susceptibility of **1** measured in an oscillating field of 10 ( $\bigcirc$ ,  $\bullet$ ), 100 ( $\times$ , +), and 1000 ( $\triangle$ ,  $\blacktriangle$ ) Hz with an amplitude of 1 Oe.



Figure 4. Hysteresis loop of 1 taken at 2 K.



Figure 5. Derivative of the isothermal magnetization with respect to dc field of 1 at 2 K.

field of 40 Oe and remanent magnetization of 200 emu Oe/mol were also observed, Figure 4 (inset). Extrapolating the reversible branch of the hysteresis curve to H = 0 gives a saturated moment of 256 emu Oe/mol, and from this, we estimate the spin canting angle,  $\gamma$ , to be ~0.5°.<sup>23</sup> At fields greater than ~1000

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Oe the spontaneous moment is destroyed, rendering a typical antiferromagnet, which is likely how Shen et al.<sup>17</sup> acquired their results.

Spin canting may arise provided a mechanism for antisymmetric exchange and/or single-ion anisotropy is present.<sup>24</sup> As this system lacks a center of symmetry and possesses distorted Mn<sup>II</sup> sites, it is reasonable to expect that the spontaneous moment arises from a Dzyaloshinski–Moriya interaction.<sup>25</sup>

However, to better understand the detailed aspects of the magnetic properties of **1**, neutron diffraction experiments are needed. Additionally, we are studying the remainder of the first-row transition metal series of  $M(4,4'-bipy)(N_3)_2$  (M = Fe, Co, Ni, Cu).

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