

Er<sub>7</sub>Ni<sub>2</sub>Te<sub>2</sub>: The Most Rare-Earth Metal-Rich Ternary Chalcogenide

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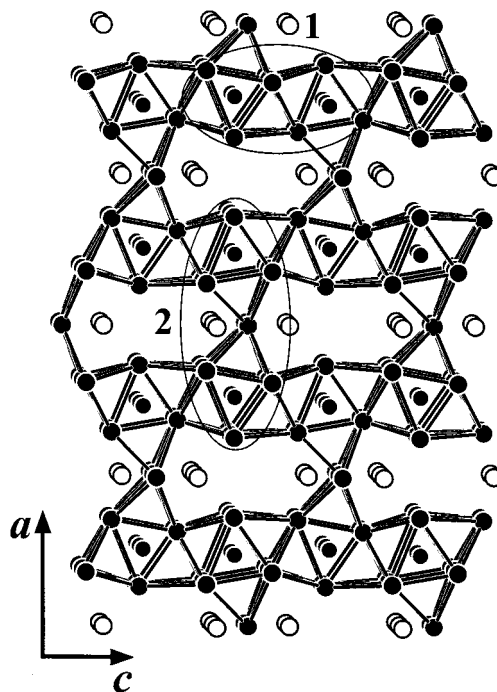
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Synthetic exploration of metal-rich chalcogenides of group IV and V early transition metals incorporating the Brewer–Wengert concept of polar intermetallic bonding<sup>1</sup> has provided a great variety of new chemistry. Recent extension of this research to the electron-poorer rare-earth metal systems has produced Sc<sub>5</sub>Ni<sub>2</sub>Te<sub>2</sub>,<sup>2</sup> Sc<sub>6</sub>MTe<sub>2</sub> (M = Mn, Fe, Co, Ni, Pd),<sup>3</sup> Dy<sub>6</sub>MTe<sub>2</sub> (M = Fe, Co, Ni),<sup>4</sup> Y<sub>5</sub>M<sub>2</sub>Te<sub>2</sub> (M = Fe, Co, Ni),<sup>5</sup> R<sub>5</sub>M<sub>2</sub>Te<sub>2</sub> (R = Gd, Dy, Er, M = Co, Ni),<sup>6</sup> and R<sub>6</sub>MTe<sub>2</sub> (R = Gd, Er, M = Co, Ni, Ru).<sup>7</sup> The new compound Er<sub>7</sub>Ni<sub>2</sub>Te<sub>2</sub> reported here represents the most metal-rich of the ternary rare-earth chalcogenides. Rare-earth elements and many of their compounds exhibit a range of fascinating physical properties originating from localized 4f electrons. The magnetic properties of rare-earth systems have received extensive attention due to their technological importance and scientific challenges. Magnetic properties of the title compound were measured and analyzed in this report.

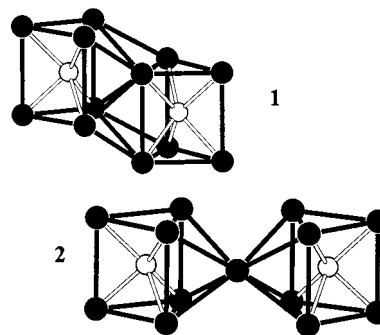
Er<sub>3</sub>Ni and NiTe<sub>2</sub> were prepared by direct reaction of stoichiometric elements (Er foil, Ames Lab 99.99%; Te powder, Alfa 99.99%; Ni powder, Alfa 99.95%) via arc melting and solid-state reactions in sealed silica tubes, respectively. Appropriate ratios of Er, Er<sub>3</sub>Ni, and NiTe<sub>2</sub> were wrapped in molybdenum foil, which was sealed in flame-baked silica tube under vacuum. This vessel was heated between 800 and 1100 °C for 7 days; the product contained Er<sub>7</sub>Ni<sub>2</sub>Te<sub>2</sub> in 80–90% yield, with ErTe as the only other observable phase in a Guinier powder diffraction pattern. The target compound is also accessible in quantitative yield by heating the above vessel at 850 °C for 4 weeks or by using molybdenum foil within a niobium tube as the reaction container. Black rods suitable for single-crystal X-ray studies were obtained, and the compound was found to adopt a new structure type with space group *Imn2* (No. 44). Synthesis of the cobalt analogue afforded an isostructural compound in high yield.

The structure<sup>8</sup> of Er<sub>7</sub>Ni<sub>2</sub>Te<sub>2</sub>, projected along the [010] direction, is shown in Figure 1. The basic structural unit is a distorted, Ni-centered, tricapped trigonal prism (TTP) of Er that is fused with like prisms by sharing the Er<sub>3</sub> triangular bases to form an infinite [(Er<sub>3</sub>Er<sub>3/2</sub>Er<sub>3/2</sub>)Ni]<sub>∞</sub> chain that propagates along the *b* axis. These infinite chains are condensed to form corrugated layers across the *bc* plane by sharing Er zigzag chains, which are composed of capping Er atoms on one TTP and inner Er on the adjacent TTP



**Figure 1.** Approximate [010] projection of the Er<sub>7</sub>Ni<sub>2</sub>Te<sub>2</sub> structure. Er and Ni atoms are shown as dark circles with and without bonds, respectively. Te are open circles. Er–Te and Er–Ni bonds are not shown for clarity. Circled regions **1** and **2** are depicted below.

(**1**). Finally, vertex condensation forms links between individual layers that result in the overall 3-D structure (**2**). Tellurium atoms reside between the metal–metal bonded layers, on *bc* planes.

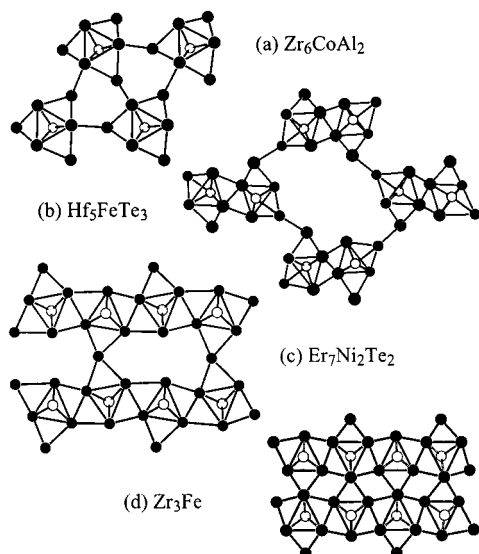


This work is placed in context by examining the structural relationship between Er<sub>7</sub>Ni<sub>2</sub>Te<sub>2</sub> and three other structures: Zr<sub>6</sub>CoAl<sub>2</sub>-type (e.g., Zr<sub>6</sub>MTe<sub>2</sub>,<sup>9</sup> Sc<sub>6</sub>MTe<sub>2</sub>,<sup>3</sup> Gd<sub>6</sub>MTe<sub>2</sub>),<sup>7</sup> Hf<sub>5</sub>FeTe<sub>3</sub>,<sup>10</sup> and Zr<sub>3</sub>Fe.<sup>11</sup> All four structures can be constructed using centered TTPs as fundamental building blocks, but following different

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- (8) Single-crystal (0.27 × 0.03 × 0.02 mm) data of Er<sub>7</sub>Ni<sub>2</sub>Te<sub>2</sub> were collected at 110 K on a Bruker SMART CCD diffractometer, which indicated a body-centered orthorhombic cell. A total of 1823 reflections (5° < 2θ ≤ 58°) were collected, of which 770 were unique (*R*<sub>int</sub> = 0.0995). The structure was refined in the orthorhombic space group *Imn2* (No. 44, *Z* = 2), lattice parameters *a* = 15.345(3) Å, *b* = 3.8377(8) Å, *c* = 9.438-(2) Å. Direct methods were used to locate all the atoms, and anisotropic refinement converged to *R*<sub>1</sub> = 0.0424, *wR*<sub>2</sub> = 0.0994.

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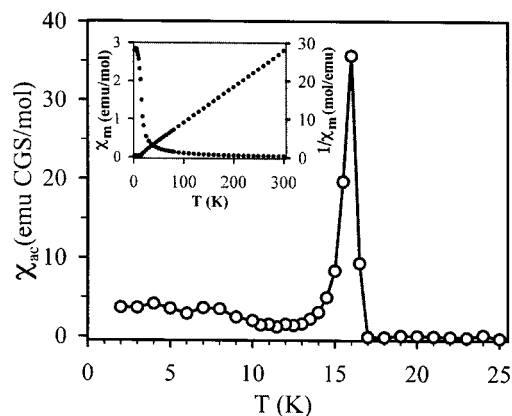
**Figure 2.** Arrangements of TTPs into (a) single chains in the  $Zr_6CoAl_2$  structure, (b) double chains in the  $Hf_5FeTe_3$  structure, (c) layers in the  $Er_7Ni_2Te_2$  structure, and (d) a 3-D network in the  $Zr_3Fe$  structure. Early transition metals are shown as larger circles, and late transition metals are smaller. Some metal–metal bonds are omitted for clarity.

condensation schemes (Figure 2). In the  $Zr_6CoAl_2$ -type structure, bonding between the capping and apical atoms of adjacent TTPs interconnects single TTP chains. In the  $Hf_5FeTe_3$  structure, double chains are formed by edge condensation of single chains, and the so-formed double chains are stitched together into a 3-D network by forming longer bonds between the capping atoms not involved in condensation. As described previously, formation of condensed intermetallic layers becomes characteristic in the  $Er_7Ni_2Te_2$  structure. Successive layers are shifted relative to one another along the  $c$  direction and are linked together by sharing capping atoms such that Te atoms are accommodated at two chemically inequivalent sites between the layers. A similar but more symmetrical layer architecture is found in the structure of  $Zr_3Fe$ , which in turn is closely related to hexagonal-close-packing metals. Adjacent layers in the  $Zr_3Fe$  structure stack vertically in Figure 2d, sharing one Zr atom in every TTP unit. This 3-D network is even more condensed due to the formation of extensive Zr–Zr bonds between the layers.

Electronic band structure calculations for  $Er_7Ni_2Te_2$  were performed with the extended Hückel method<sup>12</sup> using the YAeHMOP package.<sup>13</sup> The Fermi level intersects a prominent conduction band having mainly Er 5d and 6s character, indicating that this material is a metallic conductor, as expected. The majority of the Ni 3d states lie at lower energies, consistent with significant polarity of the Er–Ni bonds. COOP (crystal orbital overlap population) calculations indicate Er–Ni that bonding is optimized in that the Fermi level is at the crossover of the Er–Ni bonding and antibonding states. Er–Er bonding appears to impose weaker constraints on the structure because shorter Er–Er bonds do not always possess higher overlap populations. This treatment agrees well with the previous structural description. The persistence of the late-transition-metal-centered TTP unit in various intermetallic compounds is a result of the strong heterometallic bonding. With increasing late-transition-metal contents, TTP units link together to form single chains, then double chains, then layers or 3-D networks by angular distortion of the bonds between the early transition or rare-earth metals.

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**Figure 3.** The ac susceptibility of polycrystalline  $Er_7Ni_2Te_2$ . Inset: temperature dependence of  $\chi_m$  and  $\chi_m^{-1}$ .

Magnetic measurements on a polycrystalline sample (12.74 mg) were carried out with the use of a Quantum Design (model MPMS-5) SQUID magnetometer. Temperature-dependent susceptibility data were collected from 2 to 300 K at a field of 1000 G. The magnetic susceptibility data ( $\chi$ ) were corrected for the paramagnetic impurity (ErTe) contribution,<sup>12</sup> for sample holder contribution, and for the intrinsic diamagnetic contribution;<sup>15</sup> temperature independent paramagnetism was found to be negligible within the experimental error.

Between 100 and 300 K,  $\chi$  is fit with the Curie–Weiss expression,  $\chi = C/(T - \theta)$ ; a  $\chi^{-1}$  versus  $T$  plot (Figure 3, inset) gave  $C = 10.6 \pm 0.2$  and  $\theta = 4.1(4)$  K. The effective moment,  $\mu_{\text{eff}} \approx \sqrt{8C}$ , was  $9.21 \pm 0.09 \mu_B$ , to be compared with  $9.58 \mu_B$  calculated for free  $Er^{3+}$  ( $^4I_{15/2}$ ). Reported  $\mu_{\text{eff}}$  values for  $Er_2Te_3$ ,<sup>14</sup>  $ErTe$ ,<sup>14</sup> and  $Er$  metal<sup>16</sup> are  $9.50 \pm 0.14$ ,  $9.35 \pm 0.14$ , and  $8.98 \pm 0.18 \mu_B$ , respectively. The more numerous the Er–Er bonds, the lower the observed  $\mu_{\text{eff}}$ . The effective moment for  $Er_7Ni_2Te_2$  is at the expected place in this series.

An ac susceptibility measurement demonstrates an ordering transition at  $T_C = 16.5$  K (Figure 3). Ordering in  $Er_7Ni_2Te_2$  and in  $Er$  metal ( $T_C = 18$  K) can be attributed to the RKKY exchange interaction between the Er moments, which is indirect in nature and is mediated through itinerant conduction electrons.<sup>17–19</sup> This is the first metal-rich rare-earth chalcogenide for which a magnetic transition has been observed.

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**Supporting Information Available:** X-ray crystallographic files for  $Er_7Ni_2Te_2$  in CIF format. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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