

Synthesis, Crystal Structure, and Magnetic Properties of Two Cyano-Bridged Bimetallic 4f–3d Arrays with One-Dimensional Chain and Two-Dimensional Brick Wall Molecular Structures

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The design, synthesis, and structural and magnetic properties of the two new bimetallic complexes $[\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_2\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ and $[\text{Sm}(\text{DMF})_2(\text{H}_2\text{O})_3\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ (DMF = *N,N*-dimethylformamide) are presented. $[\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_2\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ was prepared by the reaction between Sm^{3+} and $[\text{Cr}(\text{CN})_6]^{3-}$ in a methanol–DMF solution, while $[\text{Sm}(\text{DMF})_2(\text{H}_2\text{O})_3\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ was prepared by the reaction between Sm^{3+} , DMF, and $[\text{Cr}(\text{CN})_6]^{3-}$ in the molar ratio of 1:2:1 in H_2O . $[\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_2\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ crystallizes in the monoclinic space group of $P2(1)/c$ with $a = 13.161(3) \text{ \AA}$, $b = 12.928(3) \text{ \AA}$, $c = 19.174(4) \text{ \AA}$, $\beta = 109.82(3)^\circ$, and $Z = 4$, while $[\text{Sm}(\text{DMF})_2(\text{H}_2\text{O})_3\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ is in the triclinic space group $P\bar{1}$ with $a = 7.7535(1) \text{ \AA}$, $b = 9.4307(3) \text{ \AA}$, $c = 16.2905(5) \text{ \AA}$, $\alpha = 94.1590(14)^\circ$, $\beta = 100.0597(18)^\circ$, $\gamma = 100.9154(18)^\circ$, and $Z = 2$. The structure of $[\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_2\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ possesses a cyano-bridged one-dimensional (1D) zigzag chain structure with alternating $\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_2$ and $\text{Cr}(\text{CN})_6$ moieties. $[\text{Sm}(\text{DMF})_2(\text{H}_2\text{O})_3\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ consists of cyano-bridged neutral brick wall-like layers in which each $[\text{Cr}(\text{CN})_6]^{3-}$ unit uses three cyanide groups in the meridional arrangement to connect with three $[\text{Sm}(\text{DMF})_2(\text{H}_2\text{O})_3]^{3+}$ units. Each $[\text{Sm}(\text{DMF})_2(\text{H}_2\text{O})_3]^{3+}$ unit, in turn, links three $[\text{Cr}(\text{CN})_6]^{3-}$, generating a flat brick wall-like structure. Magnetic studies on $[\text{Sm}(\text{DMF})_2(\text{H}_2\text{O})_3\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$ showed a magnetic-phase transition temperature below 4.2 K and a coercive field of 100 Oe at 1.85 K, while no long-range magnetic ordering was observed in the 1D complex $[\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_2\text{Cr}(\text{CN})_6] \cdot \text{H}_2\text{O}$.

Introduction

Recently, there has been considerable interest in lanthanide(III) hexacyanoferrates and the analogous hexacyanocobaltate as well as hexacyanochromium complexes because of their potential as catalytic and semiconductive materials.¹ In 1976, a series of cyano-bridged three-dimensional (3D) rare-earth hexacyanometalates(III) $\text{LnM}(\text{CN})_6 \cdot n\text{H}_2\text{O}$ ($M = \text{Fe}$ or Cr , $n = 4$ or 5) were crystallized and subjected to single-crystal X-ray and magnetic susceptibility measurements.² Ferrimagnetic ordering was observed for $\text{LnCr}^{\text{III}}(\text{CN})_6 \cdot n\text{H}_2\text{O}$ ($\text{Ln} = \text{Nd}$, Sm , Gd , Tb , Dy , Ho , Er , and Tm), and $\text{TbCr}(\text{CN})_6 \cdot 4\text{H}_2\text{O}$ has the highest magnetic-phase transition temperature of 11.7 K. Incorporation of organic ligands into the 3D solids would produce 4f–3d species with rich and interesting molecular structures.^{3–7} This strategy has also been successfully employed for the preparation of hybrid 3d Prussian blue analogues.⁸ Using

DMF as a ligand, a cyano-bridged 4f–3d dimer $[\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_4\text{Fe}(\text{CN})_6] \cdot \text{H}_2\text{O}$ was obtained, which can be described as a variation of 3D $\text{SmFe}(\text{CN})_6 \cdot 4\text{H}_2\text{O}$.⁴ Similar cyano-bridged binuclear $[\text{Ln}(\text{DMF})_4(\text{H}_2\text{O})_3\text{Fe}(\text{CN})_6] \cdot \text{H}_2\text{O}$ ($\text{Ln} = \text{lanthanide}$

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Table 1. Crystallographic Data for Complexes **2** and **3**

	2	3
formula	C ₁₈ H ₃₄ CrN ₁₀ O ₇ Sm	C ₁₂ H ₂₂ CrN ₈ O ₆ Sm
fw	704.90	576.73
$\lambda/\text{\AA}$	0.710 73	0.710 73
space group	<i>P</i> 2 ₁ / <i>c</i>	<i>P</i> 1
<i>a</i> / \AA	13.161(3)	7.7535(1)
<i>b</i> / \AA	12.928(3)	9.4307(3)
<i>c</i> / \AA	19.174(4)	16.2905(5)
α/deg	90	94.1590(14)
β/deg	109.82(3)	100.0597(18)
γ/deg	90	100.9154(18)
<i>V</i> / \AA^3	3069.1(12)	1144.75(5)
<i>Z</i>	4	2
$\rho_{\text{calc}}/(\text{g cm}^{-3})$	1.526	1.673
$\mu(\text{Mo K}\alpha)/\text{mm}^{-1}$	2.298	3.055
reflins collected	10423	17175
independent reflins	5405	4015
data/restraints/params	5405/0/358	4015/0/286
<i>R</i> [<i>I</i> > 2 σ (<i>I</i>)]	0.0372	0.0330
wR2 (all data)	0.0895	0.0828

ions except Sm and Nd) complexes and [Pr(DMF)₄(H₂O)₃Cr(CN)₆] \cdot H₂O have been reported subsequently.⁵ However, cyano-bridged two-dimensional (2D) 4f–3d assemblies have been rarely characterized so far, especially an ordering one.⁷ We just communicated a 2D brick wall-like bimetallic complex [Gd(DMF)₂(H₂O)₃Cr(CN)₆] \cdot H₂O⁹ (**1**) with long-range anti-ferromagnetic ordering below 3.5 K and a small coercive field of less than 10 Oe at 1.8 K. As an extension of this work, two bimetallic Sm–Cr complexes were synthesized and characterized, because the Sm³⁺ ion is anisotropic and a magnet with a larger coercive field was expected. We report here the synthesis, structures, and magnetic investigations of [Sm(DMF)₄(H₂O)₂Cr(CN)₆] \cdot H₂O (**2**) and [Sm(DMF)₂(H₂O)₃Cr(CN)₆] \cdot H₂O (**3**).

Experimental Section

Elemental analyses of carbon, hydrogen, and nitrogen were carried out with an Elementar Vario EL. The infrared spectroscopy on KBr pellets was performed on a Magna-IR 750 spectrophotometer in the 4000–400 cm⁻¹ region. Variable-temperature magnetic susceptibility, zero-field ac magnetic susceptibility measurements, and field dependence magnetization were performed on a Maglab System 2000 magnetometer. Effective magnetic moments were calculated by the equation $\mu_{\text{eff}} = 2.828(\chi_{\text{m}}T)^{1/2}$, where χ_{m} is the molar magnetic susceptibility. The experimental susceptibilities were corrected for the diamagnetism of the constituent atoms (Pascal's tables).

Syntheses. All of the chemicals and solvents used in the synthesis were reagent grade. The precursors K₃[Cr(CN)₆] and (Bu₄N)₃[Cr(CN)₆] were prepared by the literature methods.^{10,11}

[Sm(DMF)₄(H₂O)₂Cr(CN)₆] \cdot H₂O (2**).** To SmCl₃ \cdot 6H₂O (1 mmol) in DMF (5 mL) was added dropwise, with stirring, an equal molar amount of (Bu₄N)₃[Cr(CN)₆] in absolute MeOH (5 mL). Slow evaporation of the resultant yellow mixture in the dark at room temperature gave well-shaped single crystals. Anal. Calcd (%) for C₁₈H₃₄N₁₀O₇CrSm (704.90): C, 30.67; N, 19.87; H, 4.86. Found: C, 30.14; N, 19.94; H, 4.68. IR (cm⁻¹, KBr): $\nu_{\text{C}\equiv\text{N}}$ 2141, 2126, $\nu_{\text{C}=\text{O}}$ 1653.

[Sm(DMF)₂(H₂O)₃Cr(CN)₆] \cdot H₂O (3**).** To an aqueous solution of Sm(NO₃)₃ \cdot 6H₂O (0.5 mmol) and DMF (1 mmol) was added K₃[Cr(CN)₆] (0.5 mmol) in water (5 mL). The resultant yellow mixture was kept undisturbed in the dark at room temperature. After 1 week, large platelike yellow single crystals suitable for X-ray diffraction analysis were obtained. IR (cm⁻¹, KBr): $\nu_{\text{C}\equiv\text{N}}$ 2158, 2136, $\nu_{\text{C}=\text{O}}$ 1676_{sh}, 1654.

The complexes are insoluble in inorganic and organic solvents except water.

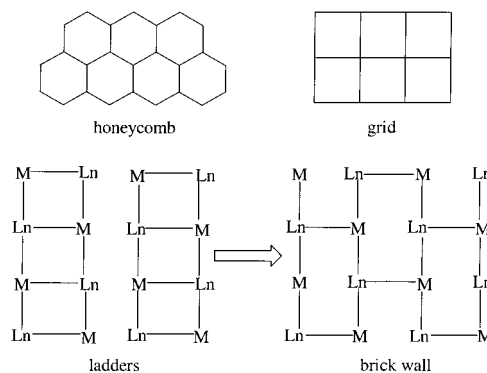


Figure 1. Schematic view of the possible topologies for LnM complexes and transformation of parallel ladders into brick wall.

Crystallographic Data Collection and Structure Determination. The data collections of **2** and **3** were made at 293 K on a Rigaku R-Axis RIPID IP and a Nonius κ -CCD diffractometer, respectively. The structures were solved by the direct method (SHELXS-97) and refined by full-matrix least-squares (SHELXL-97) on F^2 . Anisotropic thermal parameters were used for the non-hydrogen atoms and isotropic parameters for the hydrogen atoms. Hydrogen atoms were added geometrically and refined using a riding model. Weighted *R* factors (*wR*) and all of the goodness-of-fit (*S*) values are based on F^2 ; conventional *R* factors (*R*) are based on F , with F set to zero for negative F^2 . The weighting scheme is $w = 1/[s^2F_o^2 + (0.0433P)^2 + 0.0000P]$, where $P = (F_o^2 + 2F_c^2)/3$. The final Fourier map of **3** showed some residual peak and hole in the vicinity of the Sm atom (1.765/–2.051 e \AA^{-3}), which is quite normal for complexes containing heavy metals. The crystal data are summarized in Table 1.

Results and Discussion

Synthesis and General Characterization. We have previously reported a cyano-bridged 4f–3d dimer [Sm(DMF)₄(H₂O)₄Fe(CN)₆] \cdot H₂O prepared by the reaction of Sm(NO₃)₃ \cdot 6H₂O and K₃[Fe(CN)₆] in a DMF–H₂O mixture. Attempts are being made to synthesize polymeric 4f–3d species with the hope of long-range magnetic ordering. Considering that the polar solvent water has an unfavorable effect on the formation of the network structure,⁷ we tried a nonaqueous synthetic procedure or, rather, a reaction with the existence of the least amount of water. When equimolar amounts of SmCl₃ \cdot 6H₂O and (Bu₄N)₃[Cr(CN)₆] are mixed in a DMF–MeOH solution at room temperature, a light-yellow solution is obtained. Slow evaporation of the solution results in the isolation of the one-dimensional (1D) chain of the formula [Sm(DMF)₄(H₂O)₂Cr(CN)₆] \cdot H₂O (**2**), as determined by single-crystal X-ray diffraction analysis.

It can be seen that the aforementioned 4f–3d chain and dimers were prepared in an excess of DMF and that there are four coordinated DMF molecules around each lanthanide ion. We infer that a 2D or 3D array might be obtained when fewer bulky DMF molecules, for example, Ln/DMF = 1:2, were employed in the synthesis. Such an LnM array might assume a grid, honeycomb, or brick wall-like topology dependent on the mode of the M–C≡N–Ln linkages (Figure 1). Naturally, a 1D ladder motif might form, as observed in an Eu(II)–Ni(II) complex [Eu(DMF)₄Ni(CN)₄]₈.^{6c} We found that the reaction between Sm(NO₃)₃ \cdot 6H₂O, DMF, and K₃[Cr(CN)₆] in the molar ratio of 1:2:1 in H₂O afforded a 2D brick wall-like complex [Sm(DMF)₂(H₂O)₃Cr(CN)₆] \cdot H₂O (**3**).

The IR spectra of the complexes show two sharp bands (2141 and 2126 cm⁻¹ for **2** and 2158 and 2136 cm⁻¹ for **3**) in the range 2000–2200 cm⁻¹ that are attributed to C≡N stretching modes. The splitting of $\nu(\text{C}\equiv\text{N})$ suggests the presence of both bridged and nonbridged CN⁻ ligands. The strong bands at 1653

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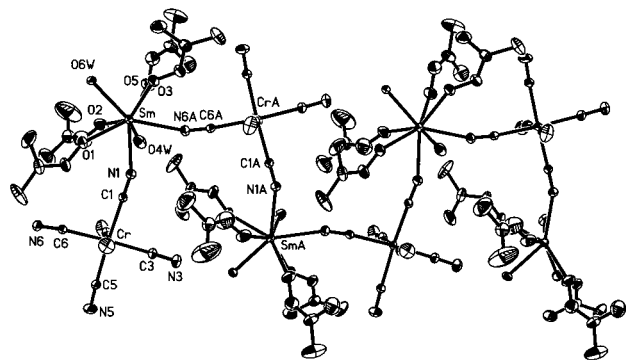


Figure 2. Chainlike structure of **2**.

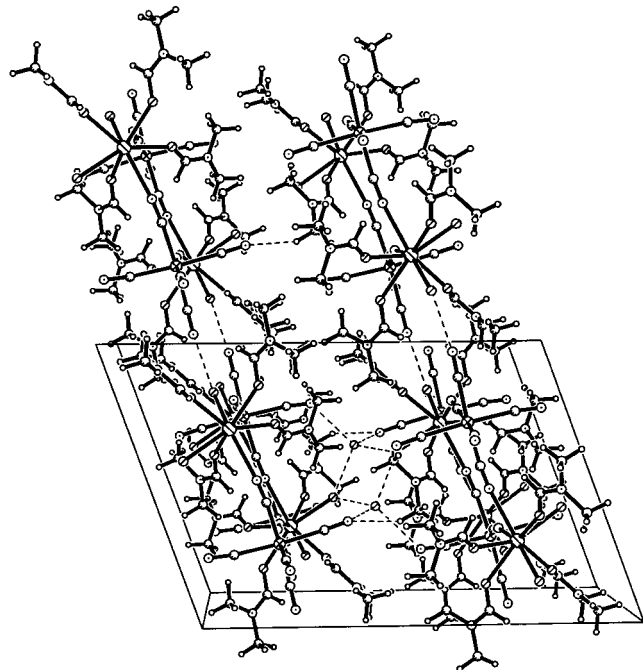


Figure 3. Unit cell packing diagram of **2**, showing the hydrogen-bonded chains.

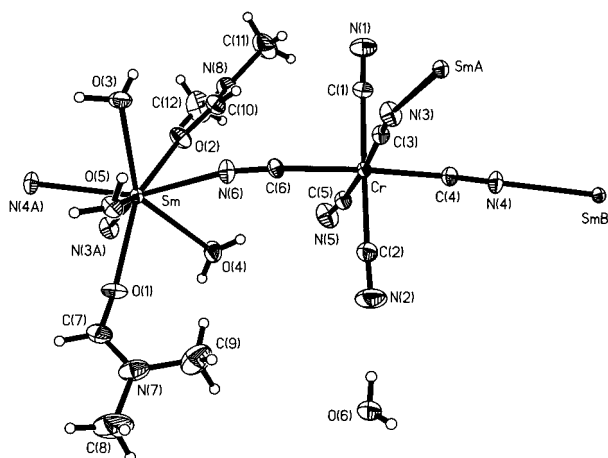


Figure 4. ORTEP plot of **3** with the numbering of the unique atoms.

cm^{-1} for **2** and 1676_{sh} and 1654 cm^{-1} for **3** are assigned to the $\text{C}=\text{O}$ stretching vibration (DMF).

Crystal Structures. ORTEP drawings of complexes **2** and **3** are shown in Figures 2 and 4, respectively. Figures 3 and 5 show the unit cell packing diagrams of **2** and **3**, respectively. Selected bond distances and angles for **2** and **3** are listed in Tables 2 and 3, respectively.

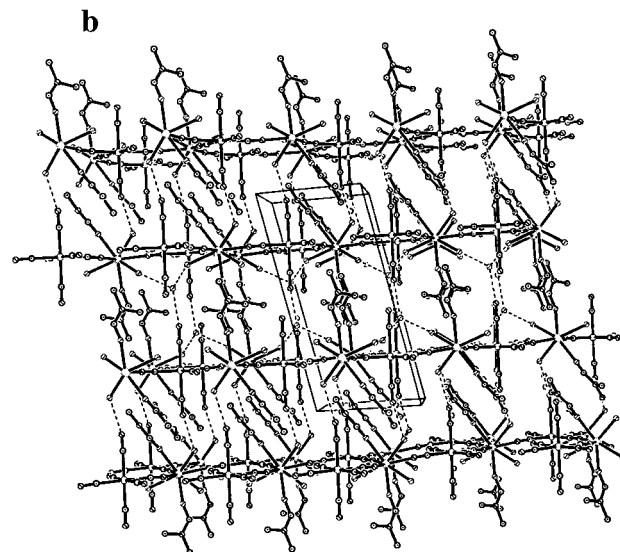
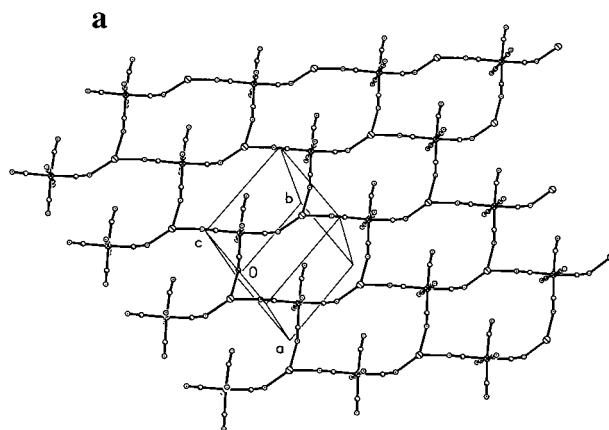


Figure 5. (a) Projection of **3** along the *c* axis showing the brick wall-like backbone (DMF and H_2O molecules are omitted for clarity). (b) Projection of **3** along the *b* axis showing the H-bonded layers.

Table 2. Selected Bond Lengths (Å) and Angles (deg) for **2**^a

Bond Length (Å)			
Sm–O(1)	2.378(4)	Sm–O(2)	2.386(4)
Sm–O(3)	2.391(3)	Sm–O(4W)	2.438(4)
Sm–O(5)	2.408(4)	Sm–O(6W)	2.374(4)
Sm–N(1)	2.538(4)	Sm–N(6)#1	2.537(4)
Cr–C(1)	2.082(5)	Cr–C(2)	2.067(6)
Cr–C(3)	2.068(5)	Cr–C(4)	2.063(5)
Cr–C(5)	2.077(5)	Cr–C(6)	2.098(5)
C(2)–N(2)	1.149(7)	C(1)–N(1)	1.143(6)
C(3)–N(3)	1.134(7)	C(4)–N(4)	1.145(6)
C(5)–N(5)	1.129(6)	C(1)–N(1)	1.144(6)
Angle (deg)			
N(2)–C(2)–Cr	172.0(5)	N(1)–C(1)–Cr	174.6(4)
N(3)–C(3)–Cr	178.1(5)	N(4)–C(4)–Cr	173.3(5)
N(5)–C(5)–Cr	177.4(5)	N(6)–C(6)–Cr	171.2(4)
C(6)–N(6)–Sm#2	163.0(4)	C(1)–N(1)–Sm	159.4(4)

^a Symmetry transformations used to generate equivalent atoms: #1, $-x + 1, y + 1/2, -z + 1/2$; #2, $-x + 1, y - 1/2, -z + 1/2$.

The structure of **2** consists of a cyano-bridged alternating array of $\text{Cr}(\text{CN})_6$ and $\text{Sm}(\text{DMF})_4(\text{H}_2\text{O})_2$ fragments. The Sm atom is eight-coordinated with a distorted square antiprism geometry. Six oxygen atoms of two water molecules and four DMF molecules and two nitrogen atoms of the bridging CN ligands are bound to Sm, with an Sm–O distance ranging from 2.374(4) to 2.438(4) Å. Two bridging cyanides coordinate to the samarium(III) ion [$\text{N}(1)\text{–Sm} = 2.538(4)$ Å and $\text{N}(6)\text{–Sm}\#2 = 2.537(4)$ Å] in a bent fashion, with bond angles of $159.4(4)^\circ$

Table 3. Selected Bond Lengths (Å) and Angles (deg) for **3^a**

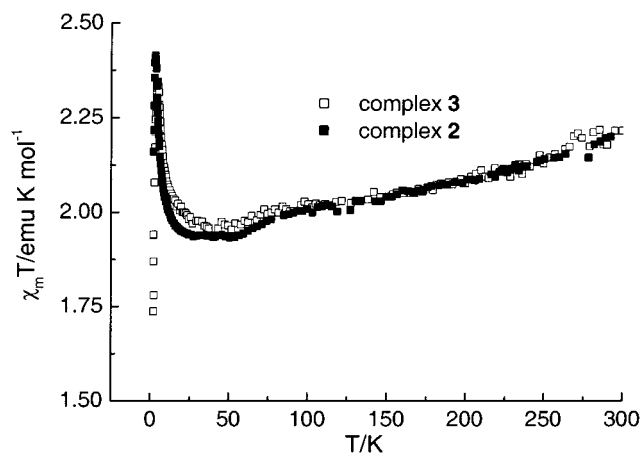
Bond Length (Å)			
Sm–O(1)	2.361(3)	Sm–O(2)	2.345(3)
Sm–O(3)	2.445(4)	Sm–O(4)	2.442(3)
Sm–O(5)	2.447(3)	Sm–N(6)	2.520(4)
Sm–N(4)#1	2.514(4)	Sm–N(3)#2	2.524(4)
Cr–C(1)	2.059(4)	Cr–C(2)	2.074(5)
Cr–C(3)	2.071(4)	Cr–C(4)	2.081(4)
Cr–C(5)	2.072(4)	Cr–C(6)	2.075(4)
C(2)–N(2)	1.138(6)	C(1)–N(1)	1.144(6)
C(3)–N(3)	1.149(5)	C(4)–N(4)	1.129(5)
C(5)–N(5)	1.152(6)	C(1)–N(1)	1.131(5)
Angle (deg)			
N(2)–C(2)–Cr	177.9(4)	N(1)–C(1)–Cr	177.3(5)
N(3)–C(3)–Cr	179.5(4)	N(4)–C(4)–Cr	177.0(4)
N(5)–C(5)–Cr	172.6(4)	N(6)–C(6)–Cr	171.2(4)
C(6)–N(6)–Sm	152.7(3)	C(4)–N(4)–Sm#3	175.0(3)
C(3)–N(3)–Sm#4	167.1(3)		

^a Symmetry transformations used to generate equivalent atoms: #1, $x - 1, y - 1, z$; #2, $x, y - 1, z$; #3, $x + 1, y + 1, z$; #4, $x, y + 1, z$.

for C(1)–N(1)–Sm and 163.0(4)° for C(6)–N(6)–Sm#2 (#2 denotes the symmetry transformation $-x + 1, y - 0.5, -z + 0.5$). Each Cr(CN)₆ coordinates to two Sm(III) ions using two *cis*-cyanide ligands, while each Sm(DMF)₄(H₂O)₂ group connects two Cr(CN)₆ moieties in a *cis* fashion, giving rise to a zigzag chainlike structure. The average value of N–Sm–O_{H₂O} angles is 112.14(15)°, larger than that of the average N–Sm–O_{DMF} angles [94.92(15)°]. A similar result is observed in complex **3** and in other cyano-bridged LnM species with both H₂O and DMF in the coordination sphere of Ln.^{4–6} This reflects that small water molecules are more polar than bulky DMF and that the electrostatic repulsion force is dominant and decisive. The adjacent intrachain metal–metal distances are 5.660(1), 5.667(1), and 7.326(1) Å for Cr···Sm, Cr···Sm#2, and Cr···Sm#1 (#1: $-x + 1, y + 0.5, -z + 0.5$), respectively. The nearest interchain metal–metal separation is 7.623(1) Å for Cr···Sm#3 (#3: $x + 1, y, z$).

The coordinated water molecule, O(4W), is hydrogen-bonded to the lattice water molecule, O(1W), with an O(1W)···O(4W) contact of 2.811(5) Å. In addition, the terminal cyanide ligands are also hydrogen-bonded to the water molecules O(6W) and O(1W), with the O–N separations ranging from 2.772 to 2.897 Å. Therefore, the molecular structure can be described as a hydrogen-bonded 3D network derived from SmCr chains (Figure 3).

Complexes **3** and **1**⁹ are isostructural and consist of 2D brick wall-like layers with slightly distorted Cr₃Sm₃ rectangles. Each [Cr(CN)₆]³⁻ unit uses three cyanide groups in the meridional arrangement to connect with three [Sm(DMF)₂(H₂O)₃]³⁺ units, and each [Sm(DMF)₂(H₂O)₃]³⁺ unit, in turn, links with three [Cr(CN)₆]³⁻ (Figure 4). The metal ions within a layer lie in a plane with the largest deviation of 0.17 Å for Sm. This local molecular disposition extends, generating a flat brick wall-like topology. A similar linking mode was observed in the ladder-like complex [Eu(DMF)₄Ni(CN)₄]₈.^{6c} Actually, the brick wall network can be generated by linking the ladders via the displacement of every other rung of the parallel ladders (Figure 1). The bridging cyanide ligands coordinate to the Sm(III) ions in two different fashions: two nearly linear, [Sm–N(4)#1–C(4)#1 = 175.0(3)° and Sm–N(3)#2–C(3)#2 = 167.1(3)°] and one bent [Sm–N(6)–C(6) = 152.7(3)°] (#1 and #2 denote the symmetry transformations $x - 1, y - 1, z$ and $x, y - 1, z$, respectively), leading to a slight distortion of the brick wall. This observation supports the proposition by Kautz et al. that the Ln–N≡C–M linkage angles are most probably affected

**Figure 6.** Temperature dependences of $\chi_m T$ in an applied field of 10 kOe for **2** and **3**.

by steric factors and crystal packing forces.⁵ In complex **3**, the bent Sm–N(6)–C(6) linkage is due to the steric hindrance of the coordinated DMF molecules (Figure 4). The Sm(III) ion exists in a slightly distorted square antiprism coordination geometry with a dihedral angle between the two faces of 2.7(2)°. The Sm–O bond distances range from 2.345(3) to 2.447(3) Å, and the Sm–N bond lengths range from 2.514(4) to 2.524(4) Å. The adjacent Cr···Sm distances are 5.487(1) Å for Cr···Sm, 5.708(1) Å for Cr···Sm#3, and 5.718(1) Å for Cr···Sm#4, respectively (#3, $x, y + 1, z$; #4, $x + 1, y + 1, z$). The shortest interlayer metal–metal distance is 7.743(1) Å for Cr···Sm#5 (#5, $1 - x, 1 - y, -z$). The unit cell packing diagram (Figure 5) shows four parallel layers. The flat layers align parallel with separations of ca. 7.4(1) and 8.5(1) Å, respectively. The uncoordinated water molecules are positioned between the distantly separated layers and linked to one terminal CN⁻ ligand of [Cr(CN)₆]³⁻ and the coordinated water molecules via hydrogen bonding. Between the closer layers, the nonbridging cyanide ligands interact with the coordinated water molecules through hydrogen bonds to connect the two layers.

Magnetic Properties. The variable-temperature (2–300 K) magnetic susceptibilities of **2** and **3** have been measured on a bunch of small crystals in a field of 10 kOe. A plot of $\chi_m T$ versus T for **2** and **3** is shown in Figure 6, where χ_m is the magnetic susceptibility per SmCr unit. The $\chi_m T$ value at 300 K is ca. 2.24 emu K mol⁻¹ (the calculated value is in the range of 2.16–2.24 emu K mol⁻¹ assuming $g_{Cr} = 2.0$),¹² which decreases smoothly with a decrease of the temperature, reaching a minimum value of 2.0 and 1.9 emu K mol⁻¹ at 52 and 45 K for **3** and **2**, respectively. With a further decrease of the temperature, $\chi_m T$ increases sharply, reaching a maximum value of 2.4 emu K mol⁻¹ at 4.5 K before decreasing again. The magnetic susceptibilities of both of the complexes deviate from the Curie–Weiss law, which is caused by first- and second-order Zeeman effect contributions of the ground (⁶H_{5/2}) and first excited (⁶H_{7/2}) states of Sm(III), which are close in energy and thermally populated at room temperature because of the relatively low spin–orbit coupling effect. Therefore, the nature of the magnetic coupling between the adjacent Sm(III) and Cr(III) ions cannot simply be determined from the decrease of $\chi_m T$, which should be due to the superposition of thermal depopulation of the low-lying excited states of Sm(III) and the Sm(III)–Cr(III) magnetic interaction. Hulliger et al. have demonstrated that the Sm(III)–Cr(III) magnetic interaction in

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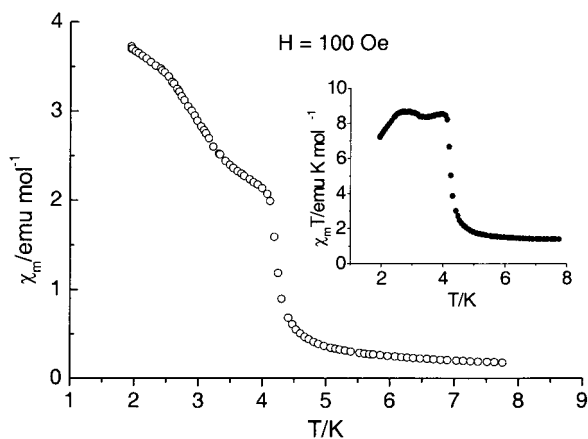


Figure 7. FCM for **3**. Inset: Corresponding $\chi_m T$ product as a function of temperature in a 100 Oe field for **3**.

3D $\text{SmCr}(\text{CN})_6 \cdot 4\text{H}_2\text{O}$ is antiferromagnetic.² To ascertain the nature of magnetic coupling between magnetic-lanthanide and transition-metal ions, an experimental methodology has been developed in which the intrinsic rare-earth ion magnetic properties are dismissed by comparing the $\chi_m T$ versus T curves for two isostructural analogues [i.e., LnM and LnM' where M and M' represent a paramagnetic and a diamagnetic transition-metal ion, respectively (for example, Cu^{2+} and Zn^{2+})]. The difference of the $\chi_m T$ products of these two species, for a given lanthanide ion, affords information about the nature of the magnetic interaction between $\text{Ln}(\text{III})$ and paramagnetic $\text{M}(\text{II})$ ions.^{7,13,14} In our experiment, we were not able to synthesize SmCo analogues of complexes **3** and **2**; therefore, such an approach could not be used.

The FCM curve of **3** measured in a low field of 100 Oe shows an abrupt increase in χ_m at ca. 4.2 K (Figure 7). The absence of a peak down to 2.0 K suggests long-range magnetic ordering due to the ferromagnetic interaction between the adjacent layers. The derivative curve ($d\text{FCM}/dT$) presents an extremum at 4.2 K, corresponding to the transition temperature (T_c). The corresponding $\chi_m T$ product as a function of temperature for **3**, shown in the inset of Figure 7, presents two maximums at 2.7 and 3.9 K, suggesting the existence of two transitions. This is evidenced by the ac magnetic susceptibility measurements.

The temperature dependence of ac magnetic susceptibilities for complex **3** verifies the magnetic-phase transition temperature at 4.2 K defined by the inflection point in χ' as well as the point where distinct χ'' appears (Figure 8). This ordering temperature is in good agreement with that of the FCM result. Below T_c , the frequency dependence of ac magnetic susceptibilities suggests the presence of a degree of spin-glass behavior, and the increase of χ' below 3.5 K indicates an unidentified transition, which may be tentatively attributed as a reentrant spin-glass state.¹⁵

The field dependence of the magnetization (0–70 kOe) for complex **3**, measured at 1.8 and 3.0 K, shows a rapid saturation of magnetization, as expected for a magnet. The saturation magnetization reaches a value of $3.5 N\beta$ at 1.8 K, higher than the literature value of $2.2 N\beta$ for the 3D SmCr ferrimagnet.² A hysteresis loop at 1.85 K (Figure 9) for complex **3** is observed with a coercive field of 100 Oe, larger than that of the GdCr analogue **1** (less than 10 Oe).⁹ This indicates that the coercive

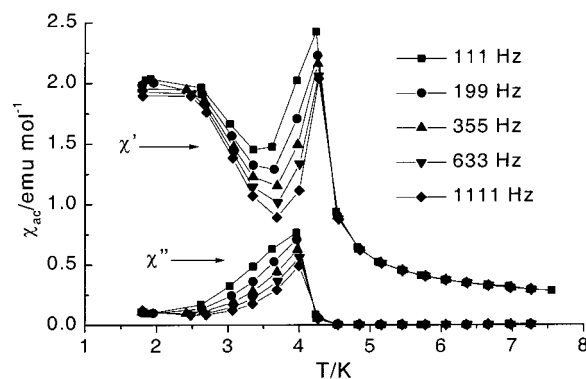


Figure 8. Real (χ') and imaginary (χ'') ac magnetic susceptibilities in zero applied dc field and an ac field of 2 Oe at different frequencies for **3**. The sample was zero-field cooled and data taken upon warming.

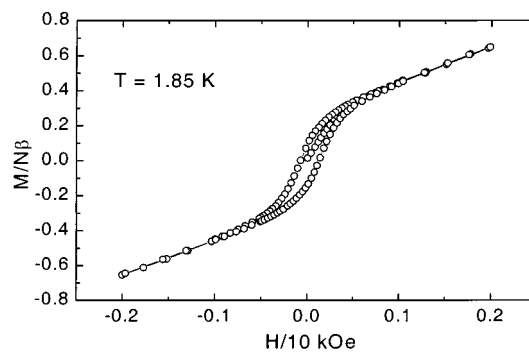


Figure 9. Hysteresis loop at 1.85 K for **3**.

field might be tuned by the choice of lanthanide ions with large magnetic anisotropy. It is worth mentioning that the structure parameters of **3** are very close to that of isostructural complex **1**. However, different bulky magnetic properties are observed: ferromagnetic-like ordering for **3** and antiferromagnetic ordering for **1**. This suggests that the nature of the intra- and interlayer magnetic interactions relate to the interacting metal ions. The ferromagnetic-like behavior of **3** might be from the stronger magnetic anisotropy of Sm compared with that of Gd in **1**.

The temperature dependence of zero-static field ac magnetic susceptibilities for complex **2** shows that the in-phase component (χ') has no maximum down to 2.0 K for frequencies of 199, 355, 633, and 1111 Hz and that the out-of-phase component (χ'') remains at zero, indicating that no magnetic ordering occurs down to 2.0 K. This reveals that the interchain magnetic interaction through H bonds in 1D species usually results in lower or no magnetic-phase transition temperature. The field dependence of the magnetization (0–70 kOe) for complex **2** measured at 1.8 and 4.2 K shows the saturation of magnetization reaching a value of 3.1 (2.9) $N\beta$ at 2.0 (4.2) K, close to that for complex **3** and higher than the literature value of $2.2 N\beta$ for $\text{SmCr}(\text{CN})_6 \cdot 4\text{H}_2\text{O}$.² On the basis of the present results, it follows that the magnetic exchange between adjacent $\text{Sm}(\text{III}) \cdots \text{Cr}(\text{III})$ is ferromagnetic-like rather than antiferromagnetic. Therefore, the decrease in $\chi_m T$ for complexes **2** and **3** should be due to the dominant thermal depopulation of the low-lying excited states of $\text{Sm}(\text{III})$ to which $\text{Sm}(\text{III})$ – $\text{Cr}(\text{III})$ ferromagnetic coupling is superimposed. Obviously, this supposition should be supported by further experimental results, for example, by comparison of magnetic properties for isostructural $\text{Sm}(\text{III})$ – $\text{Cr}(\text{III})$ and $\text{Sm}(\text{III})$ – $\text{Co}(\text{III})$ complexes. Further work along this line is in progress in our group.

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Conclusion

We have synthesized a new cyano-bridged 2D brick wall-like lanthanide transition-metal complex (**3**) exhibiting long-range magnetic ordering at 4.2 K and a new cyano-bridged 1D zigzag chainlike complex (**2**). The absence of magnetic ordering down to 2.0 K for complex **2** is in part due to the weak interchain magnetic interaction propagated through hydrogen bonds. In contrast to the 3D parent complex, the low critical temperature of the 2D complex (**3**) is connected with the weak interlayer magnetic interaction in a 2D system. Therefore, the synthesis of new cyano-bridged 3D 4f–3d arrays is our next target.

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Supporting Information Available: An X-ray crystallographic file (CIF). Field dependence (0–70 kOe) of magnetization for **3** and **2**. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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