

Elucidation of a Sc(I) Complex by DFT Calculations and Reactivity Studies[†]

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Sc(BrMgL)₂Br (L = (R₂NCH₂CH₂NCMe)₂CH, R = H) was studied by DFT methods leading to the conclusion that this diamagnetic formal scandium(I) system enjoys stabilization of its Sc-based filled d_{yz} orbital by a δ -acceptor linear combination of BrMgL ring orbitals. Investigation of the reactivity of Sc(BrMgL)₂Br (L = (R₂NCH₂CH₂NCMe)₂CH, R = Et) with H₂O·B(C₆F₅)₃ and (HOCH₂)₂CMe₂, respectively, led to decomposition, with LMgBr being isolated in the latter case.

Introduction

In the organometallic chemistry of the lanthanides, molecular complexes incorporating low-oxidation-state metals hold a particular interest. The availability of compounds with the formal oxidation state +2 for samarium, europium, and ytterbium due to the stability of the f⁶, f⁷, and f¹⁴ electronic configurations is no longer a problem.¹ In recent years, reports on lanthanum and scandium increased the number of the rare earth +2 compounds.²⁻⁴ In most of the metal (0) or (+1) containing compounds, stabilization is ensured by using a considerably bulky ligand with π -acceptor properties⁵ that supports back-donation of electron density from the metal.

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Incontestably, the principally used technique for accessing low oxidation states in different ligand environments for these metals apart from samarium, europium, and ytterbium is metal vapor synthesis.⁵ Our recent interest in the chemistry of β -diketiminato derivatives of scandium⁶ led us to an unexpected sandwich-like scandium (+1) complex (Scheme 1 and Figure 1).⁷

Scheme 1

$$LScBr_2 + 2$$
 (\longrightarrow)MgBr \longrightarrow 1 + \bigcirc

In comparison with the first reported Sc(I) compound,³ this complex is obtained in a facile route in solution, and it is diamagnetic. Earlier this year another subvalent scandium compound with a 1,2,4-triphosphacyclopentadienyl ligand, [Sc{($P_3C_2tBu_2$)_2}], was reported by Cloke and co-workers.⁴ This compound presented as a mixed valence dimer Sc(I)—Sc(III) displaying a magnetic behavior similar to that of **1**. A reasonable extension of our work is the investigation through calculations on the basis of the unusual magnetic behavior of compound **1** and the investigation of the reactivity of such a scandium +1 center. In this paper, we describe the outcome of the theoretical study and the results of redox reactions undergone by **1** with H₂O·B(C₆F₅)₃ and (HOCH₂)₂CMe₂. This work is designed to shed light on the

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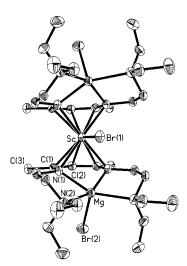


Figure 1. Molecular structure of 1. Solvent and protons are omitted for clarity.

question of the best formal oxidation state assignment for the scandium center in complex **1**.

Experimental Section

General Comments. All manipulations were performed on a high-vacuum line or in a glovebox under a purified N_2 atmosphere, using Schlenk techniques with rigorous exclusion of moisture and air. All the necessary glassware was oven-dried at 150 °C for a minimum period of 12 h, assembled hot, and cooled under high vacuum with intermittent flushing of nitrogen or argon. The samples for spectral measurements were prepared inside a MBraun MB 150-GI glovebox where the O_2 and H_2O levels were normally maintained below 1 ppm. Toluene was distilled from Na/benzophenone, THF from Na prior to use. The melting points of 1, 2, and 3 were measured in sealed capillaries on a Bühler SPA-1 instrument.

¹H, ¹³C, and ⁴⁵Sc NMR spectra (benzene- d_6 , THF- d_8 , toluene- d_8) were recorded on Bruker MSL-400, AM-250, and AM-200 instruments. Solid-state ⁴⁵Sc and ¹³C NMR spectra were recorded on Bruker A-600. The solvents for NMR measurements were dried over K or CaH₂ and trap-to-trap distilled prior to use. Mass spectra were obtained on a Finnigan MAT 8230 by EI technique. EPR spectra were recorded on a Varian Century-Line 9 GHz. UV–vis spectrum was recorded on a Perkin-Elmer 320. Magnetic measurements were performed on a Squid-Magnetometer (Quantum Design, California) at different magnetic fields in the range of temperatures between 2 and 300 K.

Reaction of 1 with H₂O·B(C₆F₅)₃. A 0.1 g (1.8 mmol) portion of H₂O·B(C₆F₅)₃ was added to a solution of 0.075 g (0.07 mmol) of 1 (obtained from reaction of 0.3 g (0.6 mmol) of LScBr₂ with 1.2 mL of (C₃H₅)MgBr (1 M in ether, 1.2 mmol) in a yield of 25%) in toluene (10 mL). The reaction mixture immediately turned to brown. After stirring for a further hour, the solution was concentrated to approx 5 mL until it became turbid and 5 mL THF was added to dissolve the precipitate. Colorless crystals of 2 suitable for X-ray analyses appeared after one week at -26 °C (approx 0.050 g). Mp 167-178 °C. Anal. Calcd for C₄₁H₃₅N₄B₂F₂₀O₂Sc: C, 46.36; H, 3.32; N, 5.27. Found: C, 46.86; H, 3.98; N, 5.06. ¹H NMR (500 MHz, C₆D₆, 300 K): $\delta = 3.51$ (s, 1 H, C(Me)CHC-(Me), 2.96 (t, 4 H; Et₂NCH₂CH₂), 2.75 (m, 8 H, NCH₂CH₃), 2.58 (t, 4 H, Et₂NCH₂CH₂), 1.40 (s, 6H, CCH₃), 0.72 (t, NCH₂CH₃). ¹⁹F NMR (188 MHz, ext C₆F₆, C₆D₆, 300 K): -136.7 (q, 8F, BC₆F₅ ortho), -155.8 (t, 4F, BC₆F₅ para), -162.7 (m, 8F, BC₆F₅ meta). ⁴⁵Sc NMR (121 MHz, referenced to $[Sc(H_2O)_6]^{3+}$ in D₂O, C₆D₆,

300 K): δ 244.11. EI-MS: *m*/z (rel int %) 1062 [C₄₁H₃₅N₄B₂F₂₀O₂-Sc⁺, 5], 976 [C₄₁H₃₅N₄B₂F₂₀O₂Sc⁺ - C₅H₁₂N, 50], 86 [C₅H₁₂N, 100].

Reaction of 1 with (HOCH₂)₂CMe₂. A solution of 0.125 g of (HOCH₂)₂CMe₂ (2 mmol) in toluene (5 mL) was added to a solution of **1** (obtained from reaction of 0.5 g (1 mmol) of LScBr₂ with 2 mL of (C₃H₅)MgBr (1 M in ether, 2 mmol) of a yield of 25%) in toluene (15 mL). The reaction takes place instantaneously with formation of a brown solution that was stirred for 6 h which was then concentrated to approx 10 mL, and 1 mL THF was added. Crystals of **3** were obtained after several days at -26 °C.

The reaction was repeated in a NMR tube with 0.020 g of **1** (0.02 mmol) in C₆D₆ (0.5 mL) and 0.002 g of (HOCH₂)₂CMe₂ in C₆D₆ (0.5 mL). The ¹H NMR spectrum was recorded after 5 min and 24 h. The second ¹H NMR spectrum showed no changes in comparison with the first one. Spectral data for **3** follow. ¹H NMR (500 MHz, C₆D₆, 300 K): δ 4.74 (s, 2 H, C(Me)*CH*C(Me)), 3.62 and 2.63 br signals assigned to 16H, CH₂, 1.80 (s, 6 H, CCH₃), 0.72 (t, 24 H, NCH₂CH₃).

X-ray Crystallography. Data for the crystal structure of 1-toluene were collected on a Stoe Image Plate IPDS II-System, for 2 on a Stoe-Siemens-Huber four circle diffractometer, and for 3 on Bruker Smart Apex CCD diffractometer. For refinement of 2 and 3 as nonmerohedric twins, the two matrices of the two domains were determined, and every domain was integrated on its own. Then a new file with the reflections of the dominant domain was written, without reflections that were strongly overlapped with reflections of the other domain. With these data, the structures were solved and refined.

The structures were solved by direct methods (SHELX-97) and refined against F^2 using SHELXL-97.⁸ *R* values were defined as $R1 = \sum ||F_o| - |F_c||/\sum |F_o|$, $wR2 = [\sum w(F_o^2 - F_c^2)^2/\sum w(F_o^2)^2]^{0.5}$, $w = [\sigma^2(F_o^2) + (g_1P)^2 + g_2P]^{-1}$, $P = 1/3[\max(F_o^2, 0) + 2F_c^2]$. Heavy atoms were refined anisotropically. Hydrogen atoms were included using the riding model with U_{iso} tied to U_{iso} of the parent atoms. Crystal data collection details, and the solution and refinement procedures, are summarized in Table 1.

Theoretical Calculations. NMR shielding calculations for complexes **1** and Cp₂Sc(BH₄)⁹ were carried out using density functional theory (DFT) and gauge-including atomic orbitals (GIAOs) using the Amsterdam Density Functional package (version *ADF2000.02*).¹⁰ The method involved has been described in detail by Schreckenbach in his description of the ⁵⁷Fe NMR shieldings in ferrocene.¹¹ For the purpose of minimizing the calculation time, the model system Sc(BrMgL)₂Br (L = (R₂NCH₂CH₂NCMe)₂CH, R = H) was used where ethyl groups remote from scandium in the real system are replaced by hydrogen atoms. The remainder of the molecule was left intact. Since the real system is known to crystallize with $C_{2\nu}$ molecular point group symmetry, geometry

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Table 1. Summary of Crystallograp

	1-toluene	2·toluene	3.0.5toluene	
formula	C41H78Br3Mg2N8Sc	$C_{48}H_{43}N_4B_2F_{20}O_2Sc$	C _{20,50} H ₃₉ BrMgN ₄	
fw	1016.42	1154.44	445.78	
space group	Cmcm	$P\overline{1}$	$P2_1/c$	
cryst syst	orthorhombic	triclinic	monoclinic	
a, Å	21.5760(12)	10.865(2)	17.508(4) Å	
b, Å	12.9704(7)	13.518(3)	7.2970(14)	
c, Å	17.7340(8)	19.172(4)	19.674(4)	
α, deg		71.83(3)		
β , deg		85.12(3)	110.493(3)	
γ, deg		69.11(3)		
V, Å ³	2148.51(17)	2498.6(9)	2354.4(8)	
Z	4	2	4	
$D_{\text{calcd}}, \text{g/cm}^3$	1.360	1.534	1.258	
temp, K	133(2)	133(2)	173(2)	
2θ range, deg	1.83-24.71	1.12-23.21	2.21-27.10	
index ranges	$-25 \le h \le 25$	$-11 \le h \le 12$	$-22 \le h \le 21$	
	$-15 \le k \le 14$	$-13 \le k \le 14$	$0 \le k \le 9$	
	$-20 \le l \le 18$	$0 \le l \le 21$	$0 \le l \le 25$	
reflns collected/unique	26177/2256 [<i>R</i> (int) = 0.0720]	$\frac{12159}{12159}$ [<i>R</i> (int) = 0.0000]	23318/5067 [<i>R</i> (int) = 0.0768]	
completeness to θ (%)	99.9	98.3	97.5	
data/restraints/params	2256/0/199	12159/856/767	5067/0/266	
GOF	1.107	1.080	1.052	
final <i>R</i> indices $[I > 2\sigma(I)]$	R1 = 0.0315,	R1 = 0.0454,	R1 = 0.0567,	
	wR2 = 0.0849	wR2 = 0.1364	wR2 = 0.1372	
R indices (all data)	R1 = 0.0346,	R1 = 0.0486,	R1 = 0.0874,	
	wR2 = 0.0867	wR2 = 0.1403	wR2 = 0.1570	
largest diff peak and hole (e/A^3)	0.946 and -0.608	0.329 and -0.435	1.966 and -0.596	

optimization on the model system was carried out with imposition of this symmetry constraint. Since no structural data are available for Cp₂Sc(BH₄), we carried out geometry optimization taking the η^3 -BH₄ bonding mode as a starting point, employing no symmetry constraints. Since the Cp₂Sc(BH₄) system exhibited a tendency to $C_{\rm s}$ symmetry, final optimization was carried out with incorporation of this constraint. Single-point calculations, preceding the NMR calculations, were run using the Zora(V) basis set for all atoms, as implemented in the ADF suite. Full electronic configuration was used for all atoms. Relativistic effects were included by virtue of the zero order regular approximation (ZORA).12 However, no spinorbit coupling effects were taken into account in the derivation of the isotropic shielding for the scandium atoms. The local density approximation (LDA) by Vosko, Wilk, and Nusair (VWN)13 was used together with the exchange and correlation corrections published by Perdew in 1991 (PW91).14

Results and Discussion

Our aforementioned work⁷ formulated, accordingly to the X-ray and NMR studies, compound **1** as $Sc(BrMgL)_2Br$ (L = (R₂NCH₂CH₂NCMe)₂CH, R = Et). But magnetic and EPR measurements exposed an interesting and unanticipated magnetic behavior (diamagnetic) for the compound, a detailed understanding of which warranted a computational study as follows.

A DFT study performed on the model system Sc-(BrMgL)₂Br (L = $(R_2NCH_2CH_2NCMe)_2CH$, R = H) re-

Figure 2. (a) A contour plot of the molecule's highest occupied molecular orbital (HOMO) in the yz plane, with Sc located at the origin. The Sc-Br vector defines the z axis (vertical) and the HOMO in the C_{2v} point group has b_2 symmetry. (b) This shows the same 10×10 Å² area in the yz plane, but it displays contours corresponding to the total SCF electron density.

а

h

vealed that the HOMO (Figure 2a) is comprised largely of the scandium d_{yz} orbital, which is antibonding with respect to the filled bromine p_y orbital. Nodal surfaces cutting through the yz plane are indicated in Figure 2a by dash dot—dash (-•-) lines, while solid (—) lines and dashes (- -) denote positive and negative contours, respectively.

From the standpoint of energy considerations, the HOMO is located nearly equidistant between the HOMO - 1 and LUMO orbitals, a characteristic typically identified with a nonbonding (lone pair) orbital (Figure 5). This is consistent with assigning a formal +1 oxidation state to the scandium center. Examination of the HOMO in three dimensions suggests a substantial stabilization by virtue of δ -backbonding to the two heterocyclic rings sandwiching the scandium atom (Figure 3).

The bromine atom (top) and the scandium atom (bottom) are the only atoms that have significant electron density in the yz plane within this area.

Schreckenbach has elucidated the electronic mechanism of the ⁵⁷Fe NMR shielding in ferrocene and in iron pentacarbonyl.¹¹ He makes note that the diamagnetic contribution to the shielding, σ^d , does not contribute to relative chemical shifts because it is due to contributions from core

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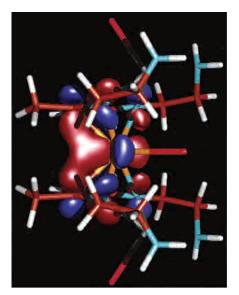


Figure 3. HOMO orbital of **1** in three dimensions.

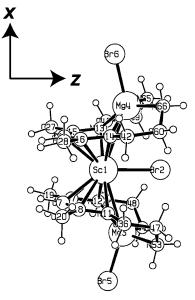


Figure 4. Calculated structure of the model compound $Sc(MgBrL)_2Br$, where L is $(R_2NCH_2CH_2NCMe)_2CH$ with R = H. Shown also are the Cartesian axis directions used for the NMR calculation. The molecule is oriented with Sc at the origin, the Sc–Br2 vector defining the *z* axis, and the *xz* plane being a mirror plane containing Sc, the three bromine atoms, and the two Mg atoms. Selected bond distances (Å): Sc1–Br2, 2.799; Sc–N13, 2.319; Sc–C15, 2.494; Sc–C8, 2.510; Mg4–N13, 2.163; Mg4–Br6, 2.525; Sc1–Mg4, 3.276.

electrons and is thus essentially invariant from one compound to another. Further, Schreckenbach points out that this truism has been noted previously for a variety of nuclei. By far the greatest contributor to the total shielding, in the absence of large spin—orbit effects,¹⁵ typically is the occupied-virtual component of the paramagnetic shielding term, $\sigma^{p.oc-vir}$. Additionally, it may turn out, as for the ferrocene ⁵⁷Fe case, that only a few occupied-virtual molecular orbital (MO) couplings determine the paramagnetic shielding almost completely. In such a case, the NMR shift and its associated anisotropy represent a sensitive probe of frontier MO electronic structure. Herein we use Schreckenbach's method to analyze via density functional theory (DFT) calculations the electronic structure of 1. The major MO mixings contributing to $\sigma^{p,oc-vir}$ are elucidated and are compared with those found for a reference scandium(III) sandwich system, Cp₂Sc(BH₄).¹⁶ Both systems have been studied by solution ⁴⁵Sc NMR methods,^{15,16} and the present work suggests that it would be of interest to probe their predicted anisotropies by solid-state ⁴⁵Sc NMR techniques. A priori expectations for any scandium(I) system include a small HOMO-LUMO gap and therefore a strong downfield shift relative to scandium(III) systems. This is because the MO mixings that comprise $\sigma^{p,oc-vir}$ become stronger as the gap between the involved orbitals diminishes. A small HOMO-LUMO gap is expected for scandium(I) because this oxidation state assignment implies a lone pair of electrons localized on scandium, i.e., a high-lying HOMO. To the extent that the HOMO becomes delocalized onto the ligands, the effect on the ⁴⁵Sc NMR downfield shift may be expected to diminish; Schreckenbach found that the most substantive MO mixings comprising $\sigma^{p,oc-vir}$ for ferrocene are between MOs having substantial iron 3d character.11 Calculations were carried out on the same model system $Sc(BrMgL)_2Br$ (L = (R₂NCH₂- CH_2NCMe_2CH , R = H) where ethyl groups remote from scandium in the real system are replaced by hydrogen atoms (Figure 4).

Depicted in Figure 5a is a set of three energy level diagrams (corresponding to Rx, Ry, and Rz) for Sc-(BrMgL)₂Br indicating via vertical dashed lines the five strongest occupied-virtual MO couplings as a function, respectively, of the three Cartesian axis directions. Separation of the occupied-virtual MO couplings in this manner comes about because the magnetic operator has the same symmetry properties as the Rx, Ry, and Rz rotations about the Cartesian axes. Note that the HOMO is, as anticipated for scandium-(I), energetically high-lying.

In the point group $C_{2\nu}$, for example, Rx has B2 symmetry, and thus, the magnetic field can induce mixing about the xaxis between occupied and virtual orbitals whose direct product is also of B2 symmetry. The strongest such MO mixing associated with Rx is accordingly seen to be between 37B2 (the HOMO) and 59A1 (the LUMO + 1). Associated with this pair of orbitals is a small energy gap and a B2symmetric direct product. In contrast to the delocalization of the HOMO, the LUMO + 1 with which the HOMO may mix strongly via Rx is a nearly pure $d_{x^2-v^2}$ orbital. From a physical point of view, in the presence of an applied magnetic field a current is induced about the x axis involving the lobes of scandium's d_{yz} and $d_{x^2-y^2}$ orbitals that reside in the *yz* plane. Even though the LUMO + 12 (64A1) is energetically wellseparated from the HOMO, it is this orbital (scandium localized and of d_{z^2} symmetry) with which the HOMO couples second-strongest. It is worthwhile at this point to make mention of the calculated chemical shielding anisotropy associated with Sc(BrMgL)2Br. The Cartesian axes indicated

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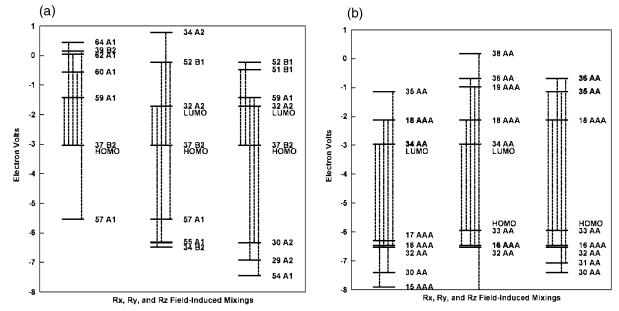


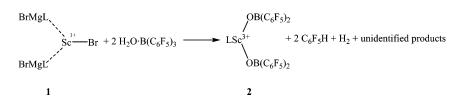
Figure 5. (a) Plot of the 5 strongest field-induced mixings of occupied and virtual energy levels of $Sc(MgBrL)_2Br$ for each of the three rotations Rx, Ry, and Rz, from left to right (left side). (b) Plot of the 5 strongest field-induced mixings of occupied and virtual energy levels of Cp_2ScBH_4 for each of the three rotations Rx, Ry, and Rz, from left to right (right side).

in Figure 4 map onto the magnetic shielding in the following manner: x, δ_{11} ; y, δ_{33} ; z, δ_{22} . What this means is that the strongest $\sigma^{p,oc-vir}$ and largest downfield chemical shifts are associated with the Cartesian x axis direction. The system is found to be rhombic; i.e., $\delta_{11} \neq \delta_{22} \neq \delta_{33}$. The orbitals that are most strongly mixed by the Ry magnetic operator are 37B2 (HOMO) with 32A2 (LUMO), and the second being 34B2 (HOMO - 13) with 32A2 (LUMO), and by the Rz magnetic operator are 37B2 (HOMO) with 51B1 (LUMO + 3) and 37B2 (HOMO) with 52B1 (LUMO + 5), respectively. Note that although HOMO-LUMO mixing is facilitated by Ry, it is associated with δ_{33} and the direction of weakest paramagnetic shielding. This is because the LUMO is mostly ligand-based with only a very small contribution from the scandium d_{xy} orbital. LUMO + 3 and LUMO + 5 both enjoy strong contributions from scandium's d_{xz} orbital, leading to the intermediate coupling and downfield shift associated with δ_{22} and R_z (see Supporting Information).

Now it is appropriate to ask if the calculated shieldings for $Sc(BrMgL)_2Br$ really do signify a downfield shift relative to a typical scandium(III) reference compound. For this purpose we choose $Cp_2Sc(BH_4)$, a sandwich complex geometrically similar to $Sc(BrMgL)_2Br$ but with what most would agree to be an unambiguous example of scandium in the +3 oxidation state.

McGlinchey and co-workers assigned to the borohydride ligand of Cp₂Sc(BH₄) an η^3 bonding mode, despite the fact that their orbital analysis suggested BH₄⁻ to be only a 4e donor to the Cp₂Sc⁺ fragment; i.e., they call the complex a 16e system.¹⁶ What we found was that the system converged nicely to an η^2 -BH₄ (see Supporting Information). Interestingly, as expected for scandium(III) relative to scandium(I), the energy level diagrams (Figure 5b) are reflective of a large HOMO–LUMO gap, essentially double the value obtained for Sc(BrMgL)₂Br. Also, while the symmetry of Cp₂Sc(BH₄) is low consistent with an expectation of rhombic magnetic symmetry about Sc, we find instead that $\delta_{22} \approx \delta_{33}$, such that the system is (accidentally) nearly axial about z. The magnetic operator of Rz symmetry permits strong mixing between HOMO -3 and LUMO +2, on one hand, and between HOMO -2 and LUMO +1 on the other. Respectively, these orbitals are strong in scandium $d_{x^2-v^2}$, d_{xy} and d_{xz} , d_{yz} character (see Supporting Information). The HOMO for $Cp_2Sc(BH_4)$ is quite delocalized, involving metal-ring π bonding as well as terminal B-H bonding. The LUMO is nearly a pure scandium d_{r^2} orbital consistent with assignment of this molecule as a Lewis acidic 16e complex. This orbital plays a role in the strongest occupiedvirtual mixings associated with δ_{22} and δ_{33} (see Supporting Information). Having discussed the magnetic anisotropy associated with the two scandium systems considered here, as well as the underlying molecular orbital basis for the anisotropy, we turn attention now to the total isotropic shielding for these systems. The compound $Cp_2Sc(BH_4)$ has been reported to have an isotropic solution ⁴⁵Sc NMR shift in C_6D_6 of +67.5 ppm relative to saturated ScCl₃ in D_2O . Taking Cp₂Sc(BH₄) as our reference compound, we can convert a given calculated shielding value $\sigma(\text{calc})$ into a chemical shift δ via the equation: $\delta(\text{calc}) = \delta(\text{ref, calc}) - \delta(\text{ref, calc})$ δ (calc) + 67 ppm. The expectation is that Sc(BrMgL)₂Br should exhibit its ⁴⁵Sc NMR signal ca. 40 ppm downfield relative to Cp₂Sc(BH₄). On the other hand, the compound $Sc(BrMgL)_2Br$ (L = (R₂NCH₂CH₂NCMe)₂CH, R = Et) has been reported to resonate at $\delta = 167.5$ ppm relative to $[Sc(H_2O)_6]_3^{+,7}$ presumably the same reference used in the case of Cp₂Sc(BH₄). The discrepancy between the calculated and observed δ_{iso} is rather substantial, given that the ⁴⁵Sc chemical shift range has been estimated to be only ca. 340 ppm.¹⁷ One possible explanation is that the experimental

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reference samples used might possibly differ in their chemical shifts; the pH of the aqueous scandium solutions was not stated in either case. A recent ⁴⁵Sc NMR study reported the use of $Sc(NO_3)_3$ in water at pH 1 as the reference standard.¹⁸ On the other hand, Schreckenbach has pointed out that calculated metal NMR shifts are very sensitive to small changes in structure (i.e., bond lengths) and to the absolute calculated MO energies, both of which can vary according to the specific choice of computational method.¹⁴ Regardless of our ability at present to accurately calculate absolute isotropic NMR shieldings, we assert that the type of study carried out here has great value in being able to relate the anisotropy of the magnetic shielding to the molecular electronic structure. We expect that the gross features of the anisotropy should be relatively independent of the computational method employed, and accordingly propose that solid-state ⁴⁵Sc NMR may be the best method firmly to connect the molecular orbitals discussed herein to the NMR observable.19

The foregoing discussion substantiates the formal oxidation state of +1 for Sc in complex 1, while recognizing that the phenomenon of δ back-bonding to the LMgBr ring LUMOs draws some of electron density away from Sc. It is difficult to conceive a reaction complex 1 might undergo in which scandium maintains its unusual oxidation state(I).¹⁹ This is because the disposition and the electronic configuration of LMgBr have a decisive contribution to the stability of complex 1. Attempts to exchange LMgBr for other ligands were not successful. Accordingly, investigation of the reactivity of 1 has proceeded in the direction of redox reactions with substrates possessing acidic hydrogens, including alcohols and water.

Reaction of 1 with the Water Adduct H₂O·B(C_6F_5)₃. For a controlled hydrolysis, H₂O·B(C_6F_5)₃ was used, and product **2** was isolated (Scheme 2). The volatiles of this reaction were analyzed by GC-MS, and C_6F_5H was identified as a byproduct (Scheme 2).

With these findings, we can propose a path involving the reduction of water by scandium, which in turn is oxidized to the (3+) state. In compensation to the fragmentation of molecule **1**, the ligand (L) migrates for stabilizing the oxidized scandium, and the $B(C_6F_5)_3$ molecule is cleaved.²⁰ In compound **2**, scandium adopts a pseudo-octahedral geometry (O(1)-Sc-O(2) 166.1°) forming a perfect planar

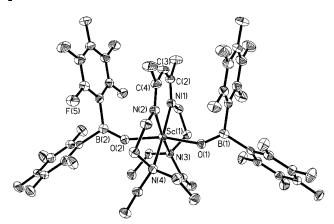


Figure 6. Molecular structure of 2. Solvent and protons are omitted for clarity.

Table 2. Selected Bond Distances (Å) and Angles (deg) for 2

bond lengths		angles			
$\begin{array}{c} C(1)-C(2)\\ C(4)-C(5)\\ C(2)-C(3)\\ N(1)-C(2)\\ B(1)-O(1)\\ B(2)-O(2)\\ Sc(1)-O(1)\\ Sc(1)-O(1)\\ \end{array}$	$\begin{array}{c} 1.508(4) \\ 1.521(4) \\ 1.402(4) \\ 1.331(3) \\ 1.296(3) \\ 1.290(3) \\ 2.0449(17) \\ 2.0449(17) \end{array}$	$\begin{array}{c} O(1)-Sc-O(2)\\ C(2)-C(3)-C(4)\\ Sc(1)-O(1)-B(1)\\ Sc(1)-O(2)-B(2)\\ N(1)-Sc(1)-N(2)\\ N(1)-Sc(1)-N(2)\\ N(1)-C(2)-C(3)\\ N(2)-C(4)-C(3)\\ N(2)\\ N(2)-C(4)-C(3)\\ N(2)-C(4)-C(4)\\ N(2)-C(4)-C(4)\\$	166.1(7) 129.5(2) 145.84(17) 149.90(17) 85.33(8) 123.9(2) 124.1(2) 70.15(8)		
$\begin{array}{l} Sc(1)-O(2) \\ Sc(1)-N(1) \\ Sc(1)-N(2) \\ Sc(1)-N(3) \\ Sc(1)-N(4) \end{array}$	2.0481(17) 2.163(2) 2.171(2) 2.402(2) 2.392(2)	$\begin{array}{l} N(1) - Sc(1) - N(3) \\ N(2) - Sc(1) - N(4) \\ N(3) - Sc(1) - N(4) \\ C(2) - N(1) - Sc(1) \\ C(4) - N(2) - Sc(1) \end{array}$	79.15(8) 79.15(8) 116.59(7) 128.53(17) 128.51(18)		

six-membered ring with the β -diketiminato backbone (Figure 6). The differences in bond lengths within the β -diketiminato moiety among 1, 2, and $LScBr_2^{21}$ are negligible (C(1)-C(2)) 1.411 Å in 1, 1.401 Å in 2, and 1.394 Å in LScBr₂; N(1)-C(1) 1.385 Å in 1, 1.330 Å in 2, and 1.340 Å in LScBr₂), and the B–O bond length in 2 (av 1.292 Å) is also very close to a covalent B–O bond (1.311 Å),²⁰ although much shorter than the coordinative B-O bond found in H₂O. $B(C_6F_5)_3$ (1.597 Å)²² (Table 2). An interesting feature of this structure is the arrangement of two C₆F₅ groups almost parallel contiguous to the ScNCCCN ring in 1. Obviously this has a direct influence on the electron density on the proton from the C(Me)CHC(Me) position whose resonance in the ¹H NMR spectrum is significantly upfield shifted from 4.82 ppm in LScBr₂ to 3.51 ppm in 2. As expected, the crystal packing reveals the incidence of $\pi - \pi$ stacking interactions between the C₆F₅ groups of different molecules.

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⁽¹⁹⁾ See ring displacement reactions with participation of Ln(0) compounds in ref 5a.

⁽²⁰⁾ Splitting of the B(C₆F₅)₃ molecule is not new. See: Neculai, D.; Roesky, H. W.; Neculai, A. M.; Magull, J.; Walfort, B.; Stalke, D. Angew. Chem. **2002**, 114, 4470–4471; Angew. Chem., Int. Ed. **2002**, 41, 4294–4295.

⁽²¹⁾ Neculai, A. M.; Neculai, D.; Nikiforov, G.; Roesky, H. W.; Magull, J.; Schlicker, C.; Herbst-Irmer, R.; Noltemeyer, M. Eur. J. Inorg. Chem. 2003, 3120–3126.

⁽²²⁾ Bergquist, C.; Bridgewater, B. M.; Harlan, C. J.; Norton, K. R.; Friesner, R. A.; Parkin, G. J. Am. Chem. Soc. 2000, 122, 10581– 10590.

Scheme 3

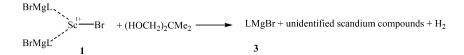


Table 3. Selected Bond Distances (Å) and Angles (deg) for **3** and 1^a

compound 3			compound 1				
Mg = N(1)	2.087(3)	C(2)-C(3)-C(4)	128.0(3)	Mg-N(1)	2.142(2)	C(1)-C(2)-C(4)	129.1(4)
Mg-N(2)	2.094(3)	N(2)-Mg-N(1)	88.22(12)			N(1)-Mg-N(1)#2	82.14(12)
Mg-N(3)	2.309(3)	N(4)-Mg-N(3)	99.29(11)	Mg - N(2)	2.345(2)	N(2)-Mg-N(2)#2	110.62(12)
Mg - N(4)	2.335(3)	N(3)-Mg-N(1)	79.32(11)			N(2)-Mg-N(1)	79.76(8)
C(1) - C(2)	1.532(5)	N(2)-Mg-N(4)	78.87(11)	C(1) - C(3)	1.511(4)		
C(2) - C(3)	1.411(5)			C(1) - C(2)	1.411(3)		
C(2) - N(1)	1.341(5)			C(1) - N(1)	1.385(3)		
C(4) - N(2)	1.334(5)						

^a Symmetry transformations used to generate equivalent atoms for 1: #1 - x, y, z; #2 x, y, $-z + \frac{3}{2}$; #3 - x, y, $-z + \frac{3}{2}$.

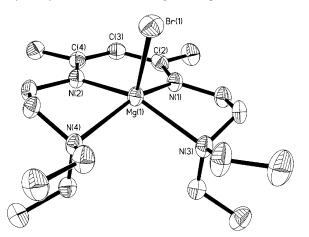


Figure 7. Molecular structure of 3. Solvent and protons are omitted for clarity.

Reaction of 1 with (HOCH₂)₂CMe₂; Formation of LMgBr (3). The reduction using alcohols led as expected to the cleavage of the sandwich molecule **1** with isolation only of the protective groups LMgBr (Scheme 3).

The reaction was monitored by ¹H NMR spectroscopy. Because of the complexity of the reaction mixture, only representative resonances were followed. The observations are consistent with a redox reaction; namely, the HOCH₂ resonances fully disappeared from the initial position (2.04 ppm) in the spectrum of the alcohol in the deuterated solvent. The resonance corresponding to the proton from C(Me)CHC-(Me) position was shifted from 2.82 ppm in 1 to 4.74 ppm in the reaction mixture, analogous with the resonance found for the same proton from the final isolated product LMgBr. The characteristic resonances of LMgBr can be found in the final spectrum. A spectral subtraction of the resonances of **3** gave no useful information about the scandium product(s). The only useful information extracted from the ⁴⁵Sc NMR spectrum (due to the multiple peaks) was related to the disappearance of the characteristic resonance for complex 1 (167 ppm).

Examination of the bond lengths and angles revealed by the X-ray structure of 3 (Figure 7) is important because it facilitates the understanding of the influence of the coordination of ScBr to the LMgBr molecule.

Moreover, the structure of 3 indicates that during the

transformation L has not changed (preservation of its monoanionic character) (Scheme 1). The pentacoordinated magnesium atom is in a pyramidal environment having the Mg-N bond distances slightly shorter compared to those in **1** (Mg $-N(\beta$ -diketiminato) 2.087 Å; 2.094 Å in **3** and 2.142 Å in 1; Mg–N(coordinated arms) 2.309 Å, 2.335 Å in 3 and 2.345 Å in 1) attributable most likely to the higher coordination number of Mg in 1 versus 3. These values accommodate with the literature data for β -diketiminato magnesium complexes.²³ Furthermore, we were interested in a comparison of distances and angles within the β -diketiminato backbone (C(1)-C(2) in 1 1.411 Å, C(2)-C(3) 1.411 Å, C(4)-C(3) 1.415 Å in **3**; C(1)-N(1) in **1** 1.385 Å, C(2)-N(1) in **3** 1.341 Å, the internal angle being 129.1° in 1 and 128.03° in 3). A closer look shows that in fact the dimensions do not change so much to conclude that the LMgBr in **1** is negatively charged (anion or radical-anion). The slightly shorter C–N bonds in 3 compared to 1 can be attributed to the different coordination numbers around the β -diketiminato nitrogen atoms, together with removal of the δ -back-bonding understood to be operative in the case of **1** (Table 3).

Note that the deviation of the Mg from the N_2C_3 plane is significantly lower in **3** (0.22 Å) compared to **1** (1.09 Å). In general, the magnesium ion coordination of L may be understood to be flexible.²⁴ In compound **1**, the conformation is presumably forced by sandwiching Sc together with repulsive interactions involving the Sc and Mg-bound bromine atoms.

Conclusion

Herein we have provided computational support for the statement that in compound **1** the formal oxidation state of

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scandium is +1. The DFT calculations are in good agreement with the experimental data and disclosed that HOMO has substantial Sc d_{yz} character, but it is additionally delocalized onto the ligands via a δ back-bonding mechanism. This delocalization represents stabilization of the unusual +1 oxidation state for scandium. Moreover, although the skeletal arrangement of the molecule was not maintained, redox processes in which **1** was involved helped to understand the chemistry of such a system and are in concert with the initial formulation of **1**. Also the fact that the L ligand was not activated during these reactions lends a bias toward the formulation of the ligand as having consistent monoanionic character.²⁵ Importantly, the synthetic method for obtaining compound 1 relies on solution chemistry, thus illustrating an alternative to metal atom vapor methods for preparing low-valent scandium compounds.

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Supporting Information Available: The CIF files for compounds **2** and **3** and details of the computational study. This material is available free of charge via the Internet at http://pubs.acs.org.

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