

## Probing the Vanadyl and Molybdenyl Bonds in Complex Vanadomolybdate Structures

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A solid solution was found to exist in the quaternary  $\text{Li}_2\text{O}-\text{MgO}-\text{V}_2\text{O}_5-\text{MoO}_3$  system between the two phases  $\text{Mg}_{2.5}\text{VMoO}_8$  and  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$ . Both  $\text{Mg}_{2.5}\text{VMoO}_8$  and  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  are isostructural with the mineral lyonsite, and substitution according to the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ , where  $\square$  denotes a cation vacancy) demonstrates that a complete solid solution exists coupling the addition of molybdenum and lithium with the subtraction of cation vacancies, magnesium, and vanadium and vice versa. Vibrational Raman spectroscopy indicates that molybdenum–oxo double bonds preferentially associate with the cation vacancies.

### Introduction

Lyonsite,  $\alpha\text{-Cu}_3\text{Fe}_4(\text{VO}_4)_6$ ,<sup>1</sup> is a naturally occurring mineral with a remarkable structure. The structure is ubiquitous in multicomponent vanadates and molybdates that form discreet  $\text{V}/\text{MoO}_4$  tetrahedra, and various homeotypes have been studied for their properties ranging from ionic conduction to catalysis.<sup>2,3</sup> Numerous mixed-metal vanadates, molybdates, and tungstates adopt the lyonsite structure including  $\text{Li}_3\text{M}(\text{MoO}_4)_3$  ( $\text{M} = \text{Fe, Al, Ga, Cr, Sc, In}$ ),<sup>3–6</sup>  $\text{Li}_2\text{M}_2(\text{MoO}_4)_3$  ( $\text{M} = \text{Fe, Co, Ni, Mg, Mn, Zn}$ ),<sup>7,8</sup>  $\text{Li}_2\text{Zr}(\text{MoO}_4)_3$ ,<sup>9</sup>  $\text{NaCo}_{2.31}(\text{MoO}_4)_3$ ,<sup>10</sup>  $\text{Cu}_{3.85}(\text{MoO}_4)_3$ ,<sup>11</sup>  $\text{Na}_{2-2x}\text{M}_{2+x}(\text{MoO}_4)_3$  ( $\text{M} = \text{Mg, Zn, Co, and Cu}$ ),<sup>7,12</sup>  $\text{M}_{2.5}\text{VMoO}_8$  ( $\text{M} = \text{Mg, Zn, Mn,$

$\text{Co}$ ),<sup>13–16</sup>  $\text{Mg}_{2.5}\text{VWO}_8$ ,<sup>17</sup>  $\text{Li}_2\text{Mg}_2(\text{WO}_4)_3$ ,<sup>18</sup>  $\text{Co}_4\text{Fe}_{3.33}(\text{VO}_4)_6$ ,<sup>19</sup>  $\text{Co}_{3.6}\text{Fe}_{3.6}(\text{VO}_4)_6$ ,<sup>20</sup>  $\text{Co}_5\text{Cr}_{2.667}(\text{VO}_4)_6$ ,<sup>21</sup>  $\text{Cu}_{4.05}\text{Cr}_{3.3}(\text{VO}_4)_6$ ,<sup>22</sup> and  $\text{Cu}_4\text{Fe}_{3.33}(\text{VO}_4)_6$ .<sup>22</sup>

The nominal lyonsite structure is orthorhombic (space group  $Pnma$ ) and consists of  $\text{V}/\text{Mo}/\text{W}$ -centered tetrahedra linking zigzag sheets and columns formed by  $\text{MO}_6$  octahedra and trigonal prisms (Figure 1). These isolated  $\text{V}/\text{Mo}/\text{W}$ -centered tetrahedra are joined to the inner wall of hexagonal tunnels formed by  $\text{MO}_6$  octahedra and  $\text{MO}_6$  trigonal prisms. Infinite columns of face-sharing  $\text{MO}_6$  octahedra pass through the center of the tunnels (Figure 2a).

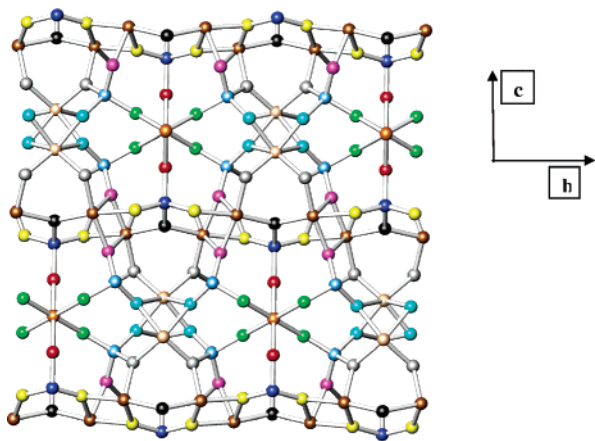
A striking feature of this structure is that it can incorporate a variable number of cationic vacancies for charge balance in order to accommodate numerous combinations of cations with different charges. For example, in  $\text{M}_{2.5}\text{VMoO}_8$ , one-sixteenth of the  $\text{M}$  ( $\text{M} = \text{Mg, Zn, Mn, Co}$ ) sites are vacant. In  $\text{Co}_{3.6}\text{Cr}_{3.6}(\text{VO}_4)_6$ , one-tenth of the  $\text{Co}/\text{Cr}$  sites are vacant,

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**Figure 1.** Structure of  $\text{Mg}_{2.5}\text{VMoO}_8/\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$ . Mo/V(1), dark blue; Mo/V(2), light blue; Mg/Li(1), dark orange; Mg/Li(2), orange (in the middle of the pseudo-hexagonal tunnel); Mg/Li(3), light orange; O(1), gray; O(2), red; O(3), yellow; O(4), black; O(5), green; O(6), pink; O(7), turquoise.

and in naturally occurring lyonsite,  $\alpha\text{-Cu}_3\text{Fe}_4(\text{VO}_4)_6$ , a full one-eighth of the Cu/Fe sites are not occupied. In contrast, although cation vacancies are common, many of the molybdate phases, such as  $\text{Li}_2\text{M}_2(\text{MoO}_4)_3$  ( $\text{M} = \text{Fe}, \text{Co}, \text{Ni}, \text{Mg}, \text{Mn}, \text{Zn}$ ) and  $\text{Li}_3\text{M}(\text{MoO}_4)_3$  ( $\text{M} = \text{Fe}, \text{Al}, \text{Ga}, \text{Cr}, \text{Sc}, \text{In}$ ), have no vacancies associated with the Li/M site. However, vacancies can be created in these molybdates, for example, by replacing lithium with magnesium  $\text{Li}_{2-2x}\text{Mg}_{2+x}(\text{MoO}_4)_3$  ( $0 \leq x \leq 0.3$ ).<sup>3</sup>

While most of the known lyonsite-type molybdates have no associated cation vacancies, all of the known lyonsite-type vanadates have associated cation vacancies. Thus, while a large number of quaternary molybdates crystallize in the lyonsite structure, it is much less common for vanadates. Instead, many mixed-metal vanadates avoid cation vacancies and adopt either the triclinic  $\beta\text{-Cu}_3\text{Fe}_4(\text{VO}_4)_6$ <sup>23</sup> or the triclinic  $\text{GaZn}_2\text{V}_3\text{O}_{11}$ <sup>24</sup> structure types. When cation vacancies are present in the lyonsite structure, the vacancies are localized within the infinite chains of face-sharing  $\text{MO}_6$  octahedra (Figure 2a) in order to minimize the Coulombic repulsions associated with the close positioning of cations at the center of these face-sharing octahedra.<sup>25</sup>

In addition to, and in part a result of, the large concentration of cation vacancies that can be incorporated, the lyonsite structure can accommodate a wide variety of metals. A large number of mono-, di-, tri-, and tetravalent metals have each been shown to occupy the  $\text{MO}_6$  octahedral and trigonal prismatic positions, while only three metals, vanadium ( $\text{V}^{5+}$ ), molybdenum ( $\text{Mo}^{6+}$ ), and tungsten ( $\text{W}^{6+}$ ), are known to occupy the tetrahedral positions. With this in mind, there exists the general possibility to form various solid solutions by coupling the substitution of  $\text{M}^+$ ,  $\text{M}^{2+}$ ,  $\text{M}^{3+}$ , or  $\text{M}^{4+}$  among themselves with charge balance achieved by a combination of cation vacancies and the substitution of  $\text{M}^{5+}$  for  $\text{M}^{6+}$ , or vice versa, in the tetrahedral positions.

To demonstrate this phenomenon,  $\text{Mg}_{2.5}\text{VMoO}_8$  and  $\text{Li}_2\text{-Mg}_2(\text{MoO}_4)_3$  were chosen to form a solid solution according to the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ , where  $\square$  denotes a cation vacancy). The formula is remarkably complex but can be described simply by one variable. More specifically, lithium ( $\text{Li}^+$ ) was substituted for magnesium ( $\text{Mg}^{2+}$ ) in  $\text{Mg}_{2.5}\text{VMoO}_8$ , and charge neutrality was maintained by the simultaneous substitution of molybdenum ( $\text{Mo}^{6+}$ ) for vanadium ( $\text{V}^{5+}$ ), concurrent with a reduction in the number of cation vacancies.

Both  $\text{Mg}_{2.5}\text{VMoO}_8$  and  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  have interesting properties. The surface of  $\text{Mg}_{2.5}\text{VMoO}_8$  has been shown to effectively catalyze the oxidative dehydrogenation of propane and butane.<sup>2,26</sup> Recently,  $\text{Mg}_{2.5}\text{VMoO}_8$  has been used as a photocatalyst for  $\text{O}_2$  evolution under visible light,<sup>27</sup> while  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  is part of the NASICON family, which are well known for their conductive properties.<sup>3</sup> A solid solution between the compounds is interesting because the addition of alkali metals to oxidation catalysts often reduces the acidity and improves selectivity. Additionally, the creation of cationic vacancies in ionic conductors is advantageous because it facilitates lithium mobility. This work illustrates the remarkable function of cation vacancies to preserve solid-state structures, which can lead to new and significant solid-state chemistry.

The structures of various bulk and supported metal oxides (e.g., vanadates, molybdates, and tungstates) have been characterized extensively by Raman spectroscopy.<sup>28,29</sup> Raman spectroscopy is sensitive to oxidation states and the short-range coordination environment around metal and oxygen ions and is particularly useful in operando<sup>30</sup> or in-situ studies for surface species where crystallinity and homogeneity are in question. Raman spectra of surface species are frequently ambiguous and difficult to assign clearly. For example, higher-order metal–oxo bonds on the surface of metal oxides have been identified by comparing observed Raman frequencies with the nominal  $\text{M}=\text{O}$  frequency range. Primarily, the ambiguity results from the overlap of the nominal  $\text{M}=\text{O}$  stretching region,  $\sim 900\text{--}1040\text{ cm}^{-1}$ ,<sup>31</sup> with the  $\nu_1$  stretching frequency region for  $\text{MO}_x$  polyhedra (e.g.,  $\nu_1$  band for  $\text{WO}_4$  tetrahedra with  $T_d$  symmetry appears at  $931\text{ cm}^{-1}$ ). Also, the vibrational frequency of the  $\text{M}=\text{O}$  stretching depends on the local environments of the metal and oxygen, as well as the metal type and the oxidation state of the metal, i.e., the coordination number of the metal, the degree of angular distortion of the  $\text{MO}_x$  polyhedra,  $\text{M}-\text{O}$  bond length, and the coordination of other chemical species (cations) to the oxygen.<sup>31,32</sup> The detailed structural information such as

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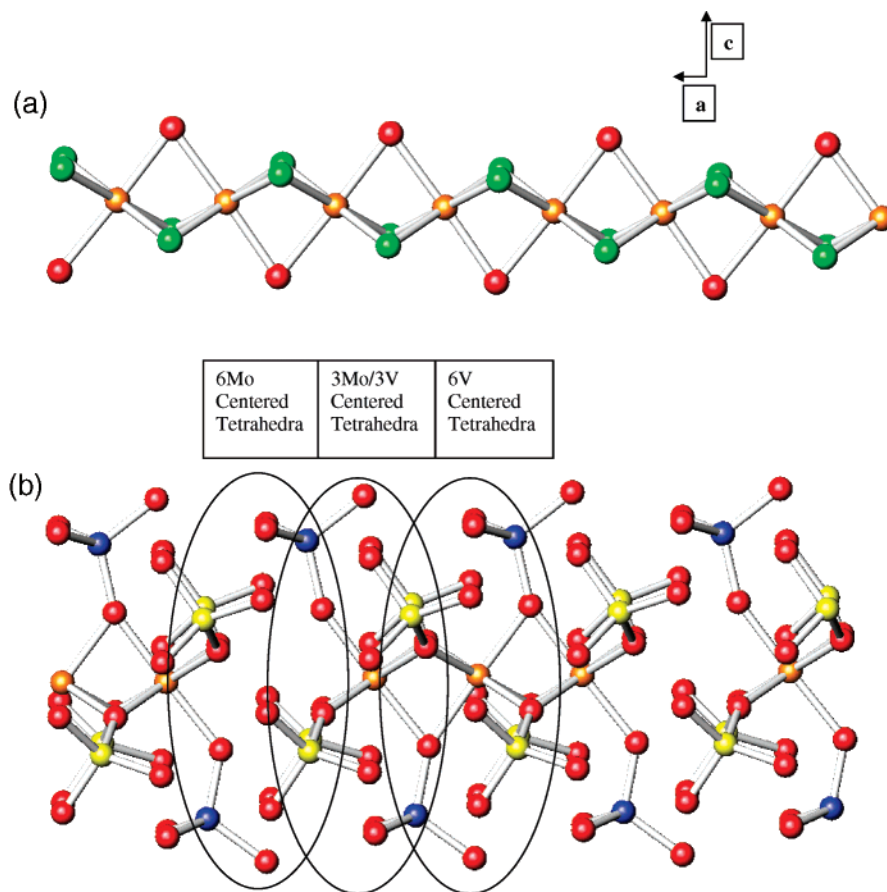
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**Figure 2.** (a) Face-sharing octahedra. Mg/Li(2), orange; O(2), red; O(5), green. (b) Face-sharing octahedra surrounded by six MoO<sub>4</sub> around the vacancy (left oval) and six VO<sub>4</sub> tetrahedra around the middle Mg of a Mg trimer (right oval). The V/Mo(1) position is blue, and the V/Mo(2) position is yellow.

coordination geometry, bond lengths, and bond angles that can be obtained by single-crystal X-ray crystallography<sup>17</sup> or X-ray absorption spectroscopy<sup>33</sup> combined with vibrational spectroscopic data can provide additional insight and more reliable vibrational assignments.

In a companion paper, high-frequency Mo=O stretching vibrations were associated with cation vacancies.<sup>17</sup> In this work, the assignments are further supported by the Li<sub>2</sub>Mg<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub>–Mg<sub>2.5</sub>VMoO<sub>8</sub> solid solution, wherein lower concentrations of cation vacancies correlate with a systematic decrease in the oxygen pπ–molybdenum dπ bonding. In addition, because there is no definitive evidence for a V=O double bond, local ordering that places primarily MoO<sub>4</sub> tetrahedra around the cation vacancies is discussed.

## Experimental Section

Compounds were synthesized by the stoichiometric addition of Li<sub>2</sub>CO<sub>3</sub> (Aldrich, 99+%), MgO (Alfa Aesar, 99.95%), V<sub>2</sub>O<sub>5</sub> (Alfa Aesar, 99.6%), and MoO<sub>3</sub> (Alfa Aesar, 99.95%). Powders were heated in platinum crucibles at 850–1000 °C.

Powder X-ray diffraction (PXRD) data for the refinements were collected every 0.01° for 10° < 2θ < 70° on a Scintag XDS 2000 diffractometer with Ni-filtered Cu Kα radiation (λ = 1.5418 Å) with 1 s per step. Lattice parameters were refined using EXPGUI

graphical interface for GSAS<sup>34</sup> with the structural parameters of Mg<sub>2.5</sub>VMoO<sub>8</sub><sup>35</sup> and Li<sub>2</sub>Mg<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> as a starting reference.

Unpolarized Raman spectra (100–1200 cm<sup>-1</sup>) of the polycrystalline samples were collected on a Bio-Rad FT-Raman spectrophotometer with 0.5 cm<sup>-1</sup> resolution and 200 scans.

## Results and Discussion

**Formation of the Solid Solution of  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ).** Limited solid solutions of Mg<sub>2.5</sub>VMoO<sub>8</sub> with differing amounts of Mg, V, and Mo have been reported previously.<sup>13,25</sup> The concentration of cation vacancies varies with the vanadium and molybdenum concentrations according to the formula Mg<sub>2.5+x</sub>V<sub>1+2x</sub>Mo<sub>1-2x</sub>O<sub>8</sub> (–0.05 ≤ x ≤ 0.05).<sup>25</sup> At the same time, the charge is balanced by a corresponding gain or loss of Mg<sup>2+</sup> ions. Recently, a solid solution has been reported for Li<sub>2</sub>Mg<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> as Li<sub>2-2x</sub>Mg<sub>2+x</sub>(MoO<sub>4</sub>)<sub>3</sub> (0 ≤ x ≤ 0.3), which also results in cation vacancies.<sup>3</sup> Owing to the previously reported versatility of the structure with regards to cation vacancies and solid solutions, the solid solution should exist between Mg<sub>2.5</sub>VMoO<sub>8</sub> and Li<sub>2</sub>Mg<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub>.

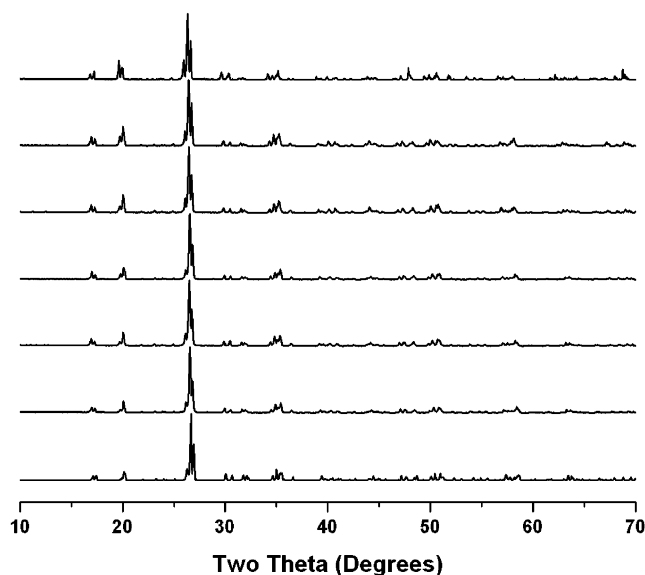
Both Mg<sub>2.5</sub>VMoO<sub>8</sub> and Li<sub>2</sub>Mg<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> adopt the same lyonsite-type structure consisting of vanadium or molybdenum tetrahedra linking zigzag sheets and columns formed

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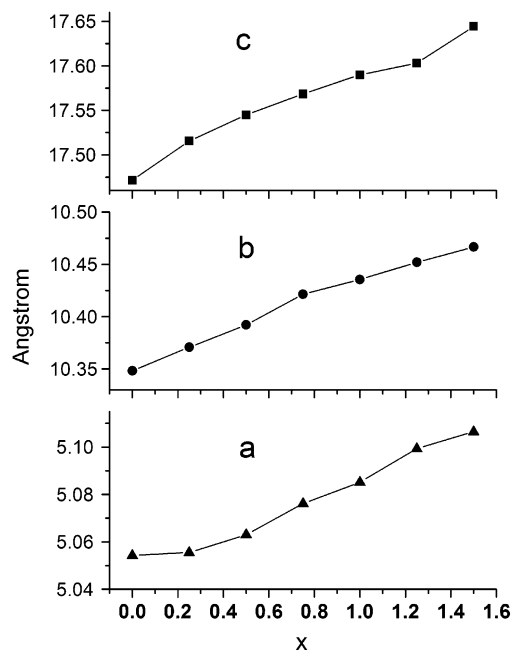
**Figure 3.** Powder X-ray diffraction patterns.  $\text{Mg}_{2.5}\text{VMoO}_8$  ( $x = 0$ ), bottom;  $x = 0.25$ , 2nd from bottom;  $x = 0.50$ , 3rd from bottom;  $x = 0.75$ , 4th from bottom;  $x = 1.00$ , 5th from bottom;  $x = 1.25$ , 6th from bottom;  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  ( $x = 1.5$ ), top.

by lithium or magnesium octahedra and trigonal prisms, as previously described. It is possible to understand how the formulas are related by considering cation vacancies.  $\text{Mg}_{2.5}\text{VMoO}_8$ , with  $Z = 6$ , can be written as  $\text{Mg}_{15}\text{V}_6\text{Mo}_6\text{O}_{48}$ . Magnesium cations occupy three crystallographically distinct positions: an 8d, and two separate 4c positions, for a total of 16 magnesium positions. Therefore, one-sixteenth of the Mg positions are vacant, and  $\text{Mg}_{2.5}\text{VMoO}_8$  can be rewritten as  $\square_{1/6}\text{Mg}_{2.5}\text{VMoO}_8$  or  $\square_{1/4}\text{Mg}_{15/4}\text{V}_{3/2}\text{Mo}_{3/2}\text{O}_{12}$ .

Further inspection of the end member formulations revealed that it is easier to determine the substitution pattern of lithium into  $\text{Mg}_{2.5}\text{VMoO}_8$  when the formulations are normalized with the same oxygen stoichiometry as  $\square_{1/4}\text{Mg}_{15/4}\text{V}_{3/2}\text{Mo}_{3/2}\text{O}_{12}$  and  $\text{Li}_2\text{Mg}_2\text{Mo}_3\text{O}_{12}$ . A more general formula can be written as  $\text{A}_4\text{B}_3\text{O}_{12}$ , where A = some combination of  $\square$ , Li, or Mg, and B = some combination of V or Mo. On the basis of these observations, the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) was derived.

Substitution according to the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) was successful. X-ray diffraction shows no evidence of other phases (Figure 3). Additionally, a lattice shift is observed as  $x$  increases. The peaks shift to a lower angle, indicating a larger  $d$ -spacing, which is expected owing to the slightly larger size of lithium and molybdenum.

The unit cell for  $\text{Mg}_{2.5}\text{VMoO}_8$  has been determined to be  $a = 5.0543(2)$  Å,  $b = 10.3483(4)$  Å, and  $c = 17.4713(7)$  Å



**Figure 4.** Variation of unit cell parameters  $a$ ,  $b$ , and  $c$  as a function of  $x$  in the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ).

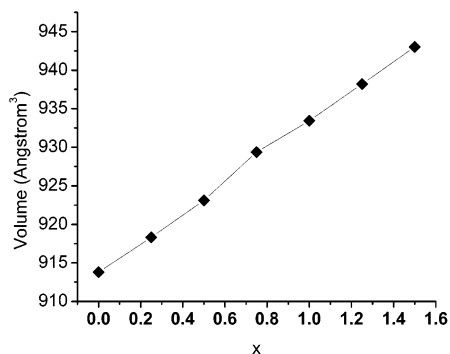
and is in good agreement with the published values  $a = 5.0515(1)$  Å,  $b = 10.3455(2)$  Å,  $c = 17.4683(4)$  Å from PXRD.<sup>35</sup> Likewise, the unit cell for  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  has been determined to be  $a = 5.1064(7)$  Å,  $b = 10.4666(14)$  Å, and  $c = 17.6443(23)$  Å. The published values for  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  are  $a = 5.1167(2)$  Å,  $b = 10.4646(4)$  Å, and  $c = 17.6228(8)$  Å from single-crystal X-ray diffraction.<sup>3</sup> Owing to the slightly larger sizes of both lithium and molybdenum compared to magnesium and vanadium respectively, the lattice constants of  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  are noticeably larger when compared to those of  $\text{Mg}_{2.5}\text{VMoO}_8$ .

The unit cell parameters and volume obtained for compounds corresponding to  $x = 0, 0.25, 0.50, 0.75, 1.00, 1.25$ , and  $1.5$  for the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) are listed in Table 1. The lattice parameters increase as more lithium and molybdenum are incorporated and magnesium and vanadium are removed, corresponding to the aforementioned shift in peaks to lower  $2\theta$ . Graphs displaying the increasing lattice parameters as a function of  $x$  are shown in Figure 4, and Figure 5 displays the increase in volume as a function of  $x$ .

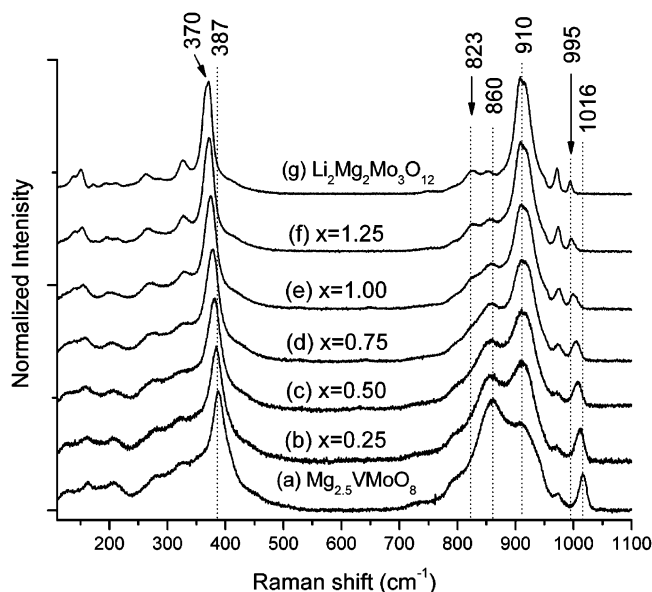
Initial attempts at forming the solid solution based on the formula  $\text{Li}_x\text{Mg}_{2.5-x}\text{V}_{1-x}\text{Mo}_{1+x}\text{O}_8$  failed and resulted in the formation of a mixture of the phases  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  and  $\text{MgMoO}_4$  when  $x = 1$ . This is expected as the lithium-rich end member would be 'LiMg<sub>1.5</sub>Mo<sub>2</sub>O<sub>8</sub>' or 'Li<sub>2</sub>Mg<sub>3</sub>Mo<sub>4</sub>O<sub>16</sub>',

**Table 1.** Unit Cell Parameters and Volume for  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ )

$x$	formulation	$a$ (Å)	$b$ (Å)	$c$ (Å)	$V$ (Å <sup>3</sup> )
0	$\text{Mg}_{2.5}\text{VMoO}_8$	5.0543(2)	10.3483(4)	17.4713(7)	913.81(6)
0.25	$\square_{0.21}\text{Li}_{0.33}\text{Mg}_{3.46}\text{V}_{1.25}\text{Mo}_{1.75}\text{O}_{12}$	5.0555(4)	10.3707(6)	17.5156(12)	918.33(29)
0.50	$\square_{0.17}\text{Li}_{0.67}\text{Mg}_{3.17}\text{V}_{1.00}\text{Mo}_{2.00}\text{O}_{12}$	5.0630(4)	10.3921(7)	17.5448(12)	923.12(34)
0.75	$\square_{0.13}\text{Li}_{1.00}\text{Mg}_{2.86}\text{V}_{0.75}\text{Mo}_{2.25}\text{O}_{12}$	5.0761(5)	10.4213(8)	17.5684(2)	929.36(8)
1.00	$\square_{0.08}\text{Li}_{1.33}\text{Mg}_{2.58}\text{V}_{0.50}\text{Mo}_{2.50}\text{O}_{12}$	5.0851(9)	10.4355(16)	17.5899(27)	933.42(389)
1.25	$\square_{0.04}\text{Li}_{1.67}\text{Mg}_{2.29}\text{V}_{0.25}\text{Mo}_{2.75}\text{O}_{12}$	5.0993(3)	10.4519(5)	17.6030(8)	938.19(12)
1.50	$\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$	5.1064(7)	10.4666(14)	17.6443(23)	943.03(225)



**Figure 5.** Variation of unit cell volume as a function of  $x$  in the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ).



**Figure 6.** Raman spectra of the solid solution  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ). Intensities are normalized to a strong band attributed to the bending vibration centered at  $370\text{--}387\text{ cm}^{-1}$ .

which is not a known compound. The end member of the solid solution is  $\text{Li}_2\text{Mg}_2\text{Mo}_3\text{O}_{12}$ .

In contrast to solid solutions where one or two metals and the ratio therein is varied, changing  $x$  in the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) affects the stoichiometry of four metals and corresponding vacancies in order to maintain compositions where one phase is observed. The solid-state crystal chemistry with other combinations of metal oxides that can be envisioned with only one variable in solid solutions that parallel  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) appears to be extensive.

**Raman Spectra of the Solid Solution.** Figure 6 displays the Raman spectra for the solid solution of  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) at formulations corresponding to  $x = 0.25, 0.50, 0.75, 1.00, 1.25$ , as well as the end members  $\text{Mg}_{2.5}\text{VMoO}_8$  (when  $x = 0$ ) and  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  (when  $x = 1.5$ ). On the basis of known trends of positions and relative intensities of the  $\nu_1$ -( $A_1$ ) Raman bands for  $\text{VO}_4$ ,  $\text{MoO}_4$ , and  $\text{WO}_4$  tetrahedra and by comparing Raman spectral features for  $\text{Mg}_{2.5}\text{VMoO}_8$  and  $\text{Mg}_{2.5}\text{VWO}_8$  in a companion paper,<sup>17</sup> the bands centered at

$910$  and  $860\text{ cm}^{-1}$  in the Raman spectrum of  $\text{Mg}_{2.5}\text{VMoO}_8$  are assigned to the symmetric stretching  $\nu_1$  mode of  $\text{MoO}_4$  and  $\text{VO}_4$  tetrahedra, respectively. The bands at  $370$  and  $387\text{ cm}^{-1}$  are assigned to the bending ( $\nu_2(E)$  or  $\nu_4(T_2)$ ) vibration of  $\text{MoO}_4$  and  $\text{VO}_4$  tetrahedra, respectively. Also, bands at  $823$  and  $791\text{ cm}^{-1}$  are assigned to  $\nu_3(T_2)$  asymmetric stretching vibration of  $\text{MoO}_4$  and  $\text{VO}_4$  tetrahedra, respectively. The intensity of the bands varies with the content of vanadium (See Tables 1 and 2, and Figure 6), which strongly supports these assignments. The  $\nu_2$  band positions of the aqueous  $\text{MoO}_4^{2-}$  and  $\text{VO}_4^{3-}$  tetrahedra have been observed at  $317$  and  $336\text{ cm}^{-1}$ , respectively, and are estimated to overlap with the  $\nu_4$  band positions.<sup>28</sup> A relatively weak, but distinct band (see Figure 6d–g) at  $326\text{ cm}^{-1}$  is likely due to the  $\nu_2(E)$  or  $\nu_4(T_2)$  bending vibration of  $\text{MoO}_4$  tetrahedra because it is stronger when the Mo concentration is relatively higher.

**M=O Stretching Vibrations.**  $\text{M}=\text{O}$  ( $\text{M} = \text{V}, \text{Mo}, \text{W}$ , etc.) stretching Raman bands generally appear at  $\sim 900\text{--}1040\text{ cm}^{-1}$ .<sup>36</sup> While the high-frequency metal–oxo bands observed in Figure 6 are of similar frequency to surface  $\text{M}=\text{O}$  species, it is important to note that none of the oxygen atoms are likely to be singly coordinated. Each oxygen atom is three coordinate in  $\text{Li}_2\text{Mg}_2\text{Mo}_3\text{O}_{12}$ , when no cation vacancies are present. However, as vacancies are incorporated at the Li/Mg(2) position, the coordination environments of O(2) and O(5) are changed. Instead of bonding to two Li/Mg(2) and one V/Mo(1) or V/Mo(2) (three total bonds), O(2) and O(5) are bonded to only one Li/Mg(2) and one V/Mo(1) or V/Mo(2) (two total bonds), respectively. The double-bond character that is reflected in the high-frequency Raman bands is a result of the two-coordinate oxygen atoms being more strongly bonded to the higher-valent transition metal ion.

In a companion paper, the Raman spectrum for isostructural  $\text{Mg}_{2.5}\text{VWO}_8$  showed a band at  $1035\text{ cm}^{-1}$ , which was assigned to the  $\text{W}=\text{O}$  stretching vibration, and showed no other bands in the  $\sim 970\text{--}1030\text{ cm}^{-1}$  range.<sup>17</sup> Because vanadium is common to both  $\text{Mg}_{2.5}\text{VMoO}_8$  and  $\text{Mg}_{2.5}\text{VWO}_8$ ,  $\text{V}=\text{O}$ -related vibrations are expected to appear at the same frequency in the vibrational spectra for the respective materials. Therefore, the Raman bands observed at  $\sim 970\text{--}1016\text{ cm}^{-1}$  in Figure 6 should be attributed to MoO-related, and not VO-related, vibrations. The band centered at  $1016\text{ cm}^{-1}$  in the Raman spectrum of  $\text{Mg}_{2.5}\text{VMoO}_8$  was assigned to the  $\text{Mo}=\text{O}$  stretching vibration of the molybdenum–oxo bond located next to Mg(2) cation vacancies.<sup>17</sup> The Raman spectrum of  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  does not show the band at  $1016\text{ cm}^{-1}$  (Figure 6g), which strongly supports the assignment because  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  does not contain vacancies. Additionally, the band position at  $1016\text{ cm}^{-1}$  matches well with surface terminal  $\text{Mo}=\text{O}$  stretching vibrations seen at  $1012\text{ cm}^{-1}$ .<sup>37</sup> The relatively lower intensity of the band at  $1016$  than  $910\text{ cm}^{-1}$  (the symmetric stretching  $\nu_1$  mode of  $\text{MoO}_4$ )

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**Table 2.** Assignments and Raman Wavenumber Shifts in  $\text{cm}^{-1}$  of  $\text{VO}_4^{3-}$  and  $\text{MoO}_4^{2-}$  Tetrahedra in the Solid Solution<sup>d</sup>

Point group, symmetry species, vibrational mode (all modes are Raman active), and the correlation			$\text{Mg}_{2.5}\text{VMoO}_8$ (when $x = 0$ )	$\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$ (when $x = 1.50$ ), & $x = 0.25$ -1.50	
$T_d$ : 4 frequencies		$C_s$ : 9 frequencies	$\text{MoO}_4^{2-}$	$\text{VO}_4^{3-}$	$\text{MoO}_4^{2-}$
$\nu_s(\text{M-O})$	$\nu_1, A_1$	$\nu_s(\text{M-O})$	910	860↓	910↑
		$\nu_s(\text{M-O})^a$	1016 (Mo=O), 974 (Mo <sup>•••</sup> O---Mg)	944 (V <sup>•••</sup> O)	995-1012↑ (Mo <sup>•••</sup> O---Li), 974↑ (Mo <sup>•••</sup> O---Mg)
$\nu_{as}(\text{M-O})$	$\nu_3, T_2$	$\nu_s(\text{M-O})^a$	807 <sup>b</sup>	799 <sup>b</sup>	853 <sup>c</sup>
		MO <sub>2</sub> scissor	370,	387↓	370↑
Bending (O-M-O)	$\nu_2, E$	MO <sub>2</sub> wag	326		326↑
		Deformation			
		MO <sub>2</sub> rock			
	$\nu_4, T_2$	MO <sub>2</sub> twist			
		$\nu_{as}(\text{M-O})$		823	791↓

<sup>a</sup> O<sub>s</sub> and O<sub>l</sub> is the oxygen with a shorter and a longer M–O bond length, respectively, than any others. M = V, Mo. <sup>b</sup> Estimated for the V/Mo–O bond length of 1.77 Å. <sup>c</sup> Estimated for the Mo–O bond length of 1.79 Å. <sup>d</sup> † and ‡ means that the band is stronger and weaker, respectively, as  $x$  increases from 0 to 1.5. The  $\text{MoO}_4^{2-}$  column on the left refers only to  $\text{Mg}_{2.5}\text{VMoO}_8$ , while the  $\text{MoO}_4^{2-}$  column on the right refers to the phases corresponding from  $x = 0.25$  to  $x = 1.50$ . The  $\text{VO}_4^{3-}$  column in the middle is common to the phases from  $x = 0$  to  $x = 1.25$ .

in the Raman spectrum of  $\text{Mg}_{2.5}\text{VMoO}_8$  (Figure 6a) can be explained by the low concentration of the vacancies (6.25%).

The band centered at  $1016 \text{ cm}^{-1}$  in the Raman spectrum of  $\text{Mg}_{2.5}\text{VMoO}_8$  ( $x = 0$ ) seen in Figure 6a gradually shifts to lower wavenumbers, and the intensity becomes weaker as  $x$  in the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) increases. The gradual red-shifts and the weakening of the shifted bands as  $x$  increases from 0 to 1.50 are likely due to a systematic decrease in the oxygen  $p\pi$ -molybdenum  $d\pi$  bonding related to the decrease in the number of the Mg(2) cation vacancies. The weaker bonds are reflected in the increase of the average bond lengths of Mo(1)–O(2) and Mo(2)–O(5) and the expansion of the unit cell dimensions (see Tables 1 and 3).

Structural distortion of the tetrahedral  $\text{MO}_4$  ion that lowers the  $T_d$  symmetry to  $C_s$  symmetry was discussed in a companion paper<sup>17</sup> and the band at  $974 \text{ cm}^{-1}$  seen in Figure 6a was assigned to the symmetric stretching ( $A'$  symmetry,  $C_s$  group) vibration of Mo=O-type (i.e., Mo<sup>•••</sup>O) bonds where  $\text{Mg}^{2+}$  binds to the oxygen. This assignment is supported by the appearance of the band at  $\sim 974 \text{ cm}^{-1}$  in the Raman spectra for all compositions of the solid solution because  $\text{Mg}^{2+}$  is common to each phase and by the absence of the band in the Raman spectrum of  $\text{Mg}_{2.5}\text{VVO}_8$ , where the Mo is replaced by W.<sup>17</sup> The band position is in good agreement with a band observed at  $967 \text{ cm}^{-1}$  from the Mo–O stretching in the distorted  $\text{MoO}_4$  tetrahedra of  $\text{MgMoO}_4$ .<sup>38</sup> The shortest Mo–O bond length in  $\text{MgMoO}_4$  is  $1.73 \text{ \AA}$ ,<sup>39</sup> while the shortest Mo–O bond length in

**Table 3.** Comparison of Bond Lengths Reproduced from Refs 3 and 13 for All the Atoms that Are on the Same Positions in Isostructural  $\text{Li}_2\text{Mg}_2\text{Mo}_3\text{O}_{12}$  and  $\text{Mg}_{2.5}\text{VMoO}_8$ <sup>a</sup>

$\text{Mg}_{2.5}\text{VMoO}_8$ ( $x = 0$ )		$\text{Li}_2\text{Mg}_2\text{Mo}_3\text{O}_{12}$ ( $x = 1.5$ )	
V/Mo(1)–O(2)	1.713(4)	Mo(1)–O(2)	1.745(2)
V/Mo(1)–O(3) × 2	1.741(3) × 2	Mo(1)–O(3) × 2	1.765(2) × 2
V/Mo(1)–O(4)	1.770(5)	Mo(1)–O(4)	1.793(3)
V/Mo(2)–O(1)	1.758(3)	Mo(2)–O(1)	1.780(2)
V/Mo(2)–O(5)	1.716(3)	Mo(2)–O(5)	1.768(2)
V/Mo(2)–O(6)	1.745(3)	Mo(2)–O(6)	1.786(2)
V/Mo(2)–O(7)	1.744(3)	Mo(2)–O(7)	1.749(2)
Mg(1)–O(1)	2.038(3)	M(1)–O(1)	2.055(2)
Mg(1)–O(3)	2.071(4)	M(1)–O(3)	2.099(3)
Mg(1)–O(3)	2.142(4)	M(1)–O(3)	2.186(2)
Mg(1)–O(4)	2.066(3)	M(1)–O(4)	2.110(2)
Mg(1)–O(6)	2.072(4)	M(1)–O(6)	2.076(2)
Mg(1)–O(6)	2.095(4)	M(1)–O(6)	2.117(2)
Mg(2)–O(2) × 2	2.106(6), 2.119(6)	M(2)–O(2) × 2	2.129(4) × 2
Mg(2)–O(5) × 2	2.049(4) × 2	M(2)–O(5) × 2	2.050(3) × 2
Mg(2)–O(5) × 2	2.066(4) × 2	M(2)–O(5) × 2	2.058(3) × 2
Mg(3)–O(1) × 2	2.059(4) × 2	M(3)–O(1) × 2	2.133(3) × 2
Mg(3)–O(7) × 2	2.117(4) × 2	M(3)–O(7) × 2	2.172(4) × 2
Mg(3)–O(7) × 2	2.178(4) × 2	M(3)–O(7) × 2	2.206(3) × 2

<sup>a</sup> All of the atomic labels from ref 13 were converted to the numbering system used in ref 3 for clarity.

$\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$ ,  $1.745 \text{ \AA}$  for Mo(1)–O(2) (see Table 3), is similar to that observed for  $\text{MgMoO}_4$ . The Mo–O bond order in Mo<sup>•••</sup>O–Mg<sup>2+</sup> for a frequency of  $971$ – $974 \text{ cm}^{-1}$  is estimated to be  $\sim 1.8$ .<sup>38</sup>

Similarly, the bands at  $995$ – $1012 \text{ cm}^{-1}$  seen in Figure 6b–g are assigned to the symmetric stretching ( $A'$  symmetry,  $C_s$  group) vibration of Mo=O-type (i.e., Mo<sup>•••</sup>O) bonds where  $\text{Li}^+$  binds to the oxygen. The  $\text{Li}^+$  coordination to the oxygen increases the Mo=O bond length, which

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lowers the Mo=O stretching frequency. As  $x$  in the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) increases, the number of  $\text{Li}^+$  bound to O(7) increases because it has been shown that the M(3) site is primarily occupied by  $\text{Li}^+$  in  $\text{Li}_2\text{Mg}_2\text{Mo}_3\text{O}_{12}$ .<sup>3</sup> Additionally, the number of  $\text{Li}^+$  bound to O(2) will increase as the M(2) cation vacancies are gradually filled. The gradual increase of  $\text{Li}^+$  binding to O(2) and O(7) and subsequent decrease in  $\pi$  bonding between the oxygen p and molybdenum d orbitals are likely responsible for the red-shift of the Mo–O stretching frequency from 1016 to 995  $\text{cm}^{-1}$ . This assignment is supported by the lack of the bands at 995–1012  $\text{cm}^{-1}$  in the Raman spectrum of  $\text{Mg}_{2.5}\text{VMoO}_8$  because it does not contain lithium.

Although tetrahedrally coordinated  $\text{Li}_2\text{MoO}_4$  might be expected to be a model compound for estimating the stretching frequencies of  $\text{Li}^+$ -coordinated Mo–O (or Mo=O-type) bonds, the bond lengths and angles of  $\text{Li}_2\text{MoO}_4$ <sup>40</sup> are not a good match for modeling  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$ . The Mo–O bond lengths in  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  are shorter compared to the Mo–O bond lengths in  $\text{Li}_2\text{MoO}_4$ , and the  $\text{MoO}_4$  tetrahedral angles in  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  show a greater degree of angular distortion compared to the  $\text{MoO}_4$  tetrahedral angles in  $\text{Li}_2\text{MoO}_4$ .<sup>40</sup> Therefore, the highest Mo–O Raman stretching frequency for  $\text{Li}_2\text{Mg}_2(\text{MoO}_4)_3$  is estimated to appear at a higher frequency than the highest observed Mo–O stretching band at 903  $\text{cm}^{-1}$  for the  $\text{Li}_2\text{MoO}_4$ .<sup>41</sup>

A more accurate model for  $\text{Li}^+$  coordinated to Mo=O-type bonds may be the coordination of a hydrogen atom from a water molecule to the oxygen in terminal M=O bonds on surfaces. A band at 975  $\text{cm}^{-1}$  has been attributed to surface terminal stretching vibrations of Mo=O where hydrogen in water is coordinated to the oxygen (i.e.,  $\text{Mo}=\ddot{\text{O}}-\text{H}$  in  $\text{H}_2\text{O}$ ),<sup>37</sup> and is close in position to the observed band centered at 995  $\text{cm}^{-1}$  in Figure 6. The Mo–O bond orders for the bands at 1016 and 995  $\text{cm}^{-1}$  are estimated to be near 2.0 (i.e., ideal Mo=O double bond) and 1.9, respectively, from a correlation relating Mo–O bond order to Mo–O stretching frequency.<sup>38</sup>

Because  $\text{MoO}_3$  is a strong Raman scatterer and was used as a starting material in the synthetic procedure, the two bands at 995 and 823  $\text{cm}^{-1}$  appearing together in the Raman spectra of the  $x = 1.0$ – $1.5$  samples (Figure 6e–g) may be due to a trace amount of  $\text{MoO}_3$ . However, this possibility can be excluded owing to the lack of other characteristic strong bands of  $\text{MoO}_3$ <sup>42</sup> at 663 and 289  $\text{cm}^{-1}$ , as well as the lack of any XRD diffraction peaks from  $\text{MoO}_3$  and to the reproducible Raman spectra (i.e., appearance of the two bands at 995 and 823  $\text{cm}^{-1}$ ) when the samples were calcined at 900 °C, a temperature that is well above the melting point of  $\text{MoO}_3$  (795 °C). The band at 823  $\text{cm}^{-1}$  is assigned to the Mo–O asymmetric stretching vibration (see Table 2).

Since vanadium and molybdenum atoms equally occupy each V/Mo-centered tetrahedral position in the crystal-

lographic model,<sup>13</sup> a Raman band attributed to the VO stretching vibration associated with the vacancies is expected to appear. The VO stretching vibration in the  $\text{Mg}_{2.5}\text{VMoO}_8$  structure should appear above 1000  $\text{cm}^{-1}$  since the terminal V=O stretching on surfaces typically appears in this region. However, the assignment of the band at 1016  $\text{cm}^{-1}$  to V=O was excluded in the companion paper because a high-frequency band appeared only at 1035  $\text{cm}^{-1}$  for  $\text{Mg}_{2.5}\text{VVO}_8$ , and the V=O vibration is expected to be located at the same frequency in these materials.<sup>17</sup>

The observed V/Mo(1)–O(2) bond length, which is dependent upon the crystallographic model, of 1.713 Å (see Table 3) should be an average of four bonds: The V(1)–O(2) that does not border a vacancy, the Mo(1)–O(2) that does not border a vacancy, the V(1)–O(2) that borders a vacancy, and the Mo(1)–O(2) that borders a vacancy. The bond lengths of the V(1)–O(2) and Mo(1)–O(2) bonds that border a vacancy are estimated to be  $\sim 1.66$  and  $\sim 1.70$  Å, respectively. The estimation is obtained from the average V/Mo(1)–O(2) crystallographic bond length of 1.713 Å, an estimated V/Mo(1)–O(2) bond length of  $\sim 1.741$  Å based on the V/Mo(1)–O(3) bond length that does not border a vacancy, as well as the degree of difference in V–O and Mo–O bond lengths in compounds<sup>13</sup> that are isostructural to  $\text{Mg}_{2.5}\text{VMoO}_8$  but consist of ordered (i.e., V-only or Mo-only) tetrahedra. The estimated VO and MoO bond lengths of 1.66 and  $\sim 1.70$  Å correlate with a bond order of 1.5–1.6 and 1.8–2.0, respectively. This correlates with a vibrational frequency of 929 and 1016  $\text{cm}^{-1}$ , respectively, on the basis of Wachs' correlation<sup>43</sup> and a correction for the degree of tetrahedral angular distortion assuming that the Mo(1) $\text{O}_4$  and V(1) $\text{O}_4$  tetrahedra have the same degree of tetrahedral angular distortion.

Recently, combined EXAFS and Raman data<sup>44</sup> for the tetrahedrally coordinated V=O terminal bond/stretching on alumina supported vanadia catalysts demonstrate that a terminal V=O Raman band at 1022  $\text{cm}^{-1}$  corresponds to a very short V=O bond length of 1.58 Å. Wachs' correlation<sup>38,43,45</sup> indicates that the V–O bond length should be much shorter than the corresponding Mo/W–O bond lengths for the M–O (M = V, Mo, W) Raman bands with the same vibrational frequency. For example, the V–O, Mo–O, and W–O bond lengths are estimated to be 1.60, 1.69, 1.71 Å, respectively, for a Raman frequency at 1000  $\text{cm}^{-1}$ . Likewise, the correlation shows that the V–O stretching should appear at a much lower frequency than the Mo/W–O stretching at the identical V/Mo/W–O bond length. For example, Raman bands for the V–O, Mo–O, and W–O stretching are estimated to appear at 902, 1076, and 1120  $\text{cm}^{-1}$ , respectively, at a single M–O bond length of 1.65 Å.

The difference in bond length/strength between V–O and Mo/W–O bonds is primarily due to the difference in valence. To satisfy the bond valence, a tetrahedrally coordinated

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$\text{VO}_4^{3-}$  unit should have three 'single' bonds and one 'double' bond, while tetrahedrally coordinated  $\text{MoO}_4^{2-}$  or  $\text{WO}_4^{2-}$  units should have two 'single' bonds and two 'double' bonds. Therefore, the Mo/W–O bond orders tend to be higher than the corresponding V–O bond order for the same tetrahedral geometry. For example, a Mo=O double bond was observed by Raman spectroscopy in the mixed V/MoO<sub>4</sub> tetrahedra of FeVMoO<sub>7</sub> and CrVMoO<sub>7</sub>. However, no vanadium–oxo double bond was observed even though the shortest V–O bond length of 1.65–1.66 Å is shorter than the shortest Mo–O bond length of 1.69 Å in the tetrahedral structure.<sup>46</sup>

Using 1.66 Å as an estimation for the V(1)–O(2) bond length associated with cation vacancies, a shoulder band appearing at 944 cm<sup>-1</sup>, closest to the estimated frequency of 929 cm<sup>-1</sup>, in the Raman spectrum of Mg<sub>2.5</sub>VMoO<sub>8</sub> is assigned to the V(1)–O(2) stretching vibration. Supporting evidence is that the band at 944 cm<sup>-1</sup> observed as a shoulder in the Raman spectrum of Mg<sub>2.5</sub>VMoO<sub>8</sub> is not seen in the Raman spectrum of vacancy-filled and vanadium-missing Li<sub>2</sub>Mg<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub>. However, the shoulder is very weak, suggesting that the concentration of V(1)–O(2) bonds surrounding the vacancies are much less than the corresponding Mo(1)–O(2) bonds, and thus MoO<sub>4</sub> preferentially associates with the cation vacancies. If this V(1)–O(2) bond length is somewhat longer than ~1.66 Å, then the vibrational assignment should be revised to the Raman band at 860 cm<sup>-1</sup>, which shows similar spectral variation (red-shifted and weaker as *x* increases) to the band at 1016 cm<sup>-1</sup>. The observed Raman frequencies and their assignments, as we understand them at the present time, are summarized in Table 2.

Because there is no definitive evidence of vanadium–oxo double bonds in the structure of Mg<sub>2.5</sub>VMoO<sub>8</sub>, short-range ordering may place predominately the molybdenum-centered tetrahedra nearest the cation vacancies. No long-range ordering of either the cation vacancies or the V/Mo tetrahedra has been observed by X-ray<sup>13</sup> or neutron diffraction,<sup>25</sup> but the relatively large thermal parameters associated with Mg(2) along the *a* axis have been attributed to the displacement of Mg<sup>2+</sup> cations toward cation vacancies arising from columbic repulsions between the face-sharing Mg<sup>2+</sup> cations.<sup>13</sup> The Mg(2) site is in the hexagonal tunnel on the 4*c* position,

and one-fourth of these sites are empty on average. Clusters of trimers of face-sharing Mg<sup>2+</sup> cations bordered by vacancies effectively relax the two Mg<sup>2+</sup> cations at the end of the trimer which displace toward the vacancies. Such clusters may also couple to additional order of the tetrahedral sites that comprise the wall of the hexagonal tunnel. Specifically, the vacancies are surrounded by six Mo<sup>6+</sup> centered tetrahedra. The two Mg<sup>2+</sup> ions on the end of the trimer are surrounded by three V<sup>5+</sup>- and three Mo<sup>6+</sup>-centered tetrahedra, while the Mg<sup>2+</sup> in the center of the trimer is surrounded by six V<sup>5+</sup>-centered tetrahedra. The model is further illustrated in Figure 2b.

## Conclusions

A complete solid solution between the phases Mg<sub>2.5</sub>VMoO<sub>8</sub> and Li<sub>2</sub>Mg<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> governed by the formula  $\square_{1/4-x/6}\text{Li}_{4x/3}\text{Mg}_{15/4-7x/6}\text{V}_{3/2-x}\text{Mo}_{3/2+x}\text{O}_{12}$  ( $0 \leq x \leq 1.5$ ) was discovered in which remarkably complex solid-state chemistry can be described by one variable. Cation vacancies are required to maintain electrical neutrality. High-frequency Mo=O Raman bands shift to lower frequencies, correlating with a decrease in the molybdenum–oxygen bonding, as the cation vacancies are filled. The apparent absence of V=O bands suggest local ordering of the VO<sub>4</sub> and MoO<sub>4</sub> tetrahedra surrounding the vacancies that places primarily MoO<sub>4</sub> tetrahedra around the cation vacancies.

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**Note Added after ASAP Publication.** This article was released ASAP on December 27, 2005, with arrows missing in Table 2. The correct version was published on January 6, 2006.

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