## Inorganic Chemistry

## New Valence-Sandwich [Mn<sup>II</sup><sub>4</sub>Mn<sup>III</sup><sub>4</sub>Mn<sup>II</sup><sub>4</sub>] Aggregate Showing Single-Molecule Magnet Behavior

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The reaction of manganese(II) acetate, 1,1,1-tris(hydroxymethyl)methane (H<sub>3</sub>thme), and triethylamine in methanol leads to the formation of [Mn<sub>12</sub>O<sub>2</sub>(OMe)<sub>2</sub>(thme)<sub>4</sub>(OAc)<sub>10</sub>(H<sub>2</sub>O)<sub>4</sub>]·2MeOH. The [Mn<sup>III</sup><sub>4</sub>Mn<sup>II</sup><sub>8</sub>] core consists of a central [Mn<sup>III</sup><sub>4</sub>O<sub>6</sub>] rhombus sandwiched by two [Mn<sup>III</sup><sub>4</sub>O<sub>7</sub>] fragments. Frequency-dependent ac susceptibility and hysteresis loops in the magnetization indicate single-molecule magnet behavior with a pure quantum-tunneling regime of relaxation below 0.2 K.

A bottom-up approach to the synthesis and crystallization of high-nuclearity, magnetically coupled transition-metal aggregates is currently a major focus of research in the area of nanoscale materials.<sup>1</sup> In this field, nanoscale molecules exhibiting single-molecule magnet (SMM) behavior are of special interest owing to their potential applications in data storage and quantum computation as quantum bits.<sup>2,3</sup> Manganese carboxylate aggregates have been a fruitful source of new SMMs in recent years<sup>3a,4</sup> because these clusters include the necessary ingredients of SMMs, i.e., large-spin ground states (*S*) and a large uniaxial anisotropy (*D*) (here arising from the presence of Jahn–Teller distorted Mn<sup>III</sup> ions), giving rise to the superparamagnetic-like property governed by a barrier to magnetization relaxation. At lower

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temperatures, this thermally active regime is replaced by quantum tunneling of the magnetization (QTM) that is governed by a small transverse anisotropy term (E). Synthetic strategies have therefore concentrated on obtaining aggregates with a high-spin ground state and appropriate magnetic anisotropy (large and small values of D and Eparameters, respectively). The use of tripodal ligands such as 1,1,1-tris(hydroxymethyl)methane (H3thme) in manganese carboxylate chemistry has been successful in this respect, resulting in the syntheses of [Mn<sub>2</sub>],<sup>5a</sup> [Mn<sub>3</sub>],<sup>5b</sup> [Mn<sub>8</sub>],<sup>5c</sup>  $[Mn_9]$ <sup>5d</sup>  $[Mn_{12}]$ <sup>5c,e</sup> and  $[Mn_{22}]$ <sup>5f</sup> SMMs. It is noteworthy that these compounds have been prepared exclusively from the reactions of trinuclear starting materials  $[Mn_3O(RO_2)_6(L)_3]^n$ (n = 0 or 1+) with tripodal ligands in MeCN or methanol. We are investigating the reactions of simple Mn<sup>II</sup> salts with a range of carboxylate anions and tripodal ligands in basic methanol solutions in order to explore new routes to Mn aggregates that might show SMM behavior. We describe here a new mixed-valence [Mn<sup>III</sup><sub>4</sub>Mn<sup>II</sup><sub>8</sub>] cluster, [Mn<sub>12</sub>O<sub>2</sub>(OMe)<sub>2</sub>- $(\text{thme})_4(\text{OAc})_{10}(\text{H}_2\text{O})_4] \cdot 2\text{MeOH}$  (1).

The reaction of  $Mn(OAc)_2 \cdot 4H_2O$  (3 mmol) and the ligand  $H_3$ thme (1 mmol) with triethylamine (3 mmol) in MeOH

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**Figure 1.** (a) ORTEP representation of the aggregate in 1, with 50% thermal ellipsoids for Mn and O atoms. H atoms and lattice MeOH molecules have been omitted for clarity. Solid black lines indicate the Mn<sup>III</sup> Jahn–Teller elongation axes; Mn<sup>III</sup> ions are in purple, Mn<sup>II</sup> ions in yellow, O atoms in red, and C atoms in gray. (b) Metal–oxygen core of 1. The central [Mn<sup>III</sup><sub>4</sub>] rhombus (purple) is sandwiched by two [Mn<sup>II</sup><sub>4</sub>] fragments (yellow).

(10 mL) results in the crystallization of **1** after 3 days in  $\sim 60\%$  yield.<sup>6</sup> Complex **1** (Figure 1a) crystallizes in the monoclinic space group  $P2_1/n$ ,<sup>7</sup> and the structure of the aggregate is shown in Figure 1. In the core, all of the Mn ions are six-coordinate except Mn6, which is seven-coordinate. Mn1 and Mn2 can be assigned as Mn<sup>III</sup>, with the remaining Mn centers as Mn<sup>II</sup>, based on charge-balance consideration, bond-valence sum calculations,<sup>8</sup> and the Jahn–Teller elongation axes on Mn1 and Mn2 (O4–Mn1–O7' and O1–Mn2–O11). Notably, all of the elongation axes of the

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 $Mn^{III}$  centers are close to co-parallel. On the basis of the above oxidation states, the mixed-valent [ $Mn_{12}O_{20}$ ] core can be described as a central rhomboid { $Mn^{III}_4O_6$ } layer (Mn1, Mn2, O1, O4, O7, and their symmetry equivalents) sandwiched by two { $Mn^{II}_4O_7$ } layers (Mn3, Mn4, Mn5, Mn6, O2, O3, O5, O6, O8, O13, O15, and their symmetry equivalents) (Figure 1b). Such a valence-sandwich-type topology is only the second reported example in Mn aggregates.<sup>10</sup> With the exception of Mn6, which has been pulled out of position by the chelating acetates, the structure of the metal—oxygen core approximates closely to a fragment of the rock salt lattice type.

The two oxo ligands (O1 and O1') adopt a  $\mu_5$ -bridging mode, which is still rather unusual in manganese carboxylate aggregates. The methoxo ligands adopt a typical  $\mu_3$ -bridging mode. The four thme<sup>3-</sup> ligands are all fully deprotonated, and each of them uses two arms to form  $\mu_3$ -OR bridges, while the third arm acts as a  $\mu_2$  bridge. Of the 10 acetate ligands, 4 form syn,syn bridges, another 4 adopt a less common chelate-bridging mode,<sup>4d</sup> and the remaining 2 bond terminally to Mn<sup>II</sup> centers (Mn5). The four aqua ligands are coordinated to the peripheral Mn<sup>II</sup> centers (Mn4). These form H bonds to acetate ligands, either in the same molecule, stabilizing the monodentate acetates, or to acetates in neighboring molecules, forming a planar two-dimensional packing of [Mn<sub>12</sub>] molecules in the crystal (Figure S1 in the Supporting Information).

Solid-state dc magnetic susceptibility ( $\gamma$ ) data were collected in the range 1.8-300 K at 0.1 T (Figure S2 in the Supporting Information). Above 100 K, the magnetic susceptibility follows a Curie–Weiss behavior with C = 41.5 $cm^3 mol^{-1} K$  and  $\theta = -47 K$ . The Curie constant is slightly smaller than the expected value for 4 Mn<sup>III</sup> and 8 Mn<sup>II</sup> noninteracting spins (~47 cm<sup>3</sup> mol<sup>-1</sup> K for g = 2). The negative Weiss constant indicates the presence of dominant antiferromagnetic interactions within the aggregates. The effects of these interactions are also observable on a plot of  $\chi T$  vs T, where  $\chi T$  decreases continuously from 35.4 cm<sup>3</sup> mol<sup>-1</sup> K at 300 K to 8.2 cm<sup>3</sup> mol<sup>-1</sup> K at 1.8 K. Even at this temperature, the  $\gamma T$  product is still decreasing, indicating that low-lying excited states are still significantly populated (Figures S2 in the Supporting Information). Therefore, fitting of the M vs H/T data cannot be used to determine unambiguously the spin ground state S and the D term of this  $Mn_{12}$ complex.

To investigate whether **1** might be a SMM, ac susceptibility measurements were performed in a 3.5-G ac field oscillating at 10–1000 Hz and with a zero applied dc field (Figure 2). The  $\chi'T$  product (Figure 2, top) steadily decreases

<sup>(6)</sup> Calcd for C<sub>42</sub>H<sub>80</sub>O<sub>40</sub>Mn<sub>12</sub> (i.e., loss of both lattice MeOH): C, 26.77; H, 4.28. Found: C, 26.65; H, 4.45. Selected IR data (KBr, cm<sup>-1</sup>): 1573vs, 1423vs, 1344m, 1212w, 1121s, 1050s, 1029s, 1000w, 953w, 914w, 674m, 657w, 617m, 579s, 506m, 469s.

<sup>(7)</sup> Crystal data for **1** ( $C_{44}H_{88}Mn_{12}O_{42}$ ): red-brown block ( $0.23 \times 0.21 \times 0.08 \text{ mm}$ ), M = 1948.41, monoclinic,  $P2_1/n$ , a = 11.8465(12) Å, b = 14.2914(14) Å, c = 20.564(2) Å,  $\beta = 96.596(2)^\circ$ , V = 3458.5(6) Å<sup>3</sup>, T = 100 K, Z = 2,  $D_{calc} = 1.871 \text{ Mg m}^{-3}$ , F(000) = 1976,  $\mu(Mo K\alpha) = 2.207 \text{ mm}^{-1}$ . 16 134 data were measured on a Bruker SMART Apex ( $4.5 < 2\theta < 56.3^\circ$ ), of which 7823 were unique ( $R_{int} = 0.0297$ ); semiempirical absorption correction, <sup>11</sup>  $T_{max}/T_{min} = 0.901/0.592$ . Refinement<sup>12</sup> (466 parameters) to wR2 = 0.1038, S = 1.007 (all data), R1 = 0.0398 [6395 data with  $I > 2\sigma(I)$ ]; final difference peak/hole +1.00/-0.74 e Å<sup>-3</sup>. CCDC 286909.

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**Figure 2.** Temperature dependence of the in-phase ( $\chi'$ ) (as  $\chi'T$ , top) and out-of-phase ( $\chi''$ ) (bottom) ac susceptibility in a 3.5-G oscillating field at the indicated frequencies.



**Figure 3.** Field hysteresis loops of the normalized magnetization for 1 at a field scan rate of 0.14 T s<sup>-1</sup> in the temperature range of 0.8-0.04 K.

with decreasing temperature, confirming the population of low-lying excited states. From these data, a rough estimation of the spin ground state can be obtained by extrapolating the plot to 0 K, where only the ground state would be populated. A  $\chi'T$  value of ~5 cm<sup>3</sup> mol<sup>-1</sup> K is found and could indicate a ground state around S = 2 or 3.<sup>4b</sup> Although the  $\chi'T$  vs *T* plot shows no obvious frequency dependence above 1.8 K, the  $\chi''$  vs *T* plot (Figure 2, bottom) clearly shows nonzero and frequency-dependent  $\chi''$  signals below 3 K, indicating a slow relaxation of the magnetization.

To confirm SMM behavior, measurements of magnetization vs dc field were made on single crystals of **1** using a micro-SQUID apparatus.<sup>9</sup> Hysteresis loops, the characteristic property of a magnet, are indeed observed (Figures 3 and S3 in the Supporting Information). The coercivity is strongly temperature- and time-dependent, increasing with decreasing temperature or increasing field sweep rate, demonstrating the superparamagnetic-like behavior of **1**. The observation of very "rounded" hysteresis loops that do not fully saturate further indicates the presence of very low-lying spin states,



**Figure 4.**  $\tau$  vs 1/T plot for **1**.

as might be expected given the high Mn<sup>II</sup> content of **1**. These closely spaced spin levels, together with probable intermolecular interactions in 1 (Figure S1 in the Supporting Information) may well have smeared out the step features (resonant QTM) usually observed on M vs H curves for a SMM. Nevertheless, hysteresis loops become temperatureindependent below about 0.2 K (while staying sweep-fieldrate-dependent), indicating the presence of a pure quantum regime of magnetization relaxation as expected for SMMs. To determine the relaxation time of this system, magnetization decay data were collected in the 0.04-0.8 K range (Figure S4 in the Supporting Information). For complex 1, however, the  $M/M_s$  vs t curves cannot be fitted by a simple exponential law, and a scaling procedure<sup>10</sup> was applied to extract the temperature dependence of the mean relaxation time  $\tau(T)$ . The result is shown in Figure 4. Above ca. 0.2 K, the relaxation rate is temperature-dependent. The straight line in Figure 4 is a fit to an Arrhenius law, yielding the values  $\tau_0 = 3 \times 10^{-11}$  s and  $U_{\rm eff} = 12.7$  K. Below ca. 0.2 K, the relaxation rate becomes temperature-independent, suggesting the occurrence of ground-state QTM, expected for SMM behavior.

In conclusion, complex **1** is a rare example in the manganese carboxylate—tripodal ligand system of an aggregate assembled from simple  $Mn^{II}$  ions showing an unusual valence-sandwich structural motif. dc and ac susceptibilities confirm its SMM behavior with hysteresis loops and QTM. Efforts to substitute the acetate anion with other carboxylate ligands (such as PhCO<sub>2</sub><sup>-</sup>, Cl<sub>3</sub>CCO<sub>2</sub><sup>-</sup>, or Me<sub>3</sub>CCO<sub>2</sub><sup>-</sup>) have led to a series of new Mn clusters, which will be reported later.

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**Supporting Information Available:** Crystallographic details (CIF) for **1** and plots of the packing,  $\chi T$  vs *T*, and  $M/M_s$  vs *t* for **1**. This material is available free of charge via the Internet at http://pubs.acs.org.

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