

Syntheses, Characterization, and Computational Study of WSF₄ and WSF₄·CH₃CN

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A new and improved synthetic route to WSF₄ was developed from the reaction of WF₆ and Sb₂S₃ in anhydrous HF. Tungsten sulfide tetrafluoride was characterized by Raman and ¹⁹F NMR spectroscopy for the first time in HF solvent. Both studies provided evidence for its monomeric form in HF solution. In the solid state, WSF₄ was also characterized by Raman and infrared spectroscopy. The WSF₄ · CH₃CN adduct was prepared from WSF₄ and CH₃CN in anhydrous HF solvent and by the direct combination of WSF₄ with excess CH₃CN, and was characterized by Raman and infrared spectroscopy in the solid state and ¹⁹F NMR spectroscopy in CH₃CN solution. The crystal structure of WSF₄·CH₃CN was obtained and showed that CH₃CN coordinates to W in an end-on fashion and trans to the W-S bond. Quantumchemical calculations using B3LYP and PBE1PBE methods were used to calculate the gas-phase geometries and vibrational frequencies of WSF₄ and WSF₄ · CH₃CN.

Introduction

Although oxide fluorides of transition metals have attracted considerable attention,¹ reports related to the analogous high-oxidation-state sulfide fluorides have been scarce. Among all transition-metal sulfide fluorides, WSF₄ has been the most studied system. Tungsten sulfide tetrafluoride has been characterized by electron-diffraction in the gas-phase;² by EXAFS³ and ¹⁹F NMR spectroscopy⁴ in solution; and in the solid state by Raman and infrared spectroscopy,⁴ and X-ray crystallography.⁵ While WSF₄ was found to be monomeric in the gas-phase² and in CH₃CN and CH₂Cl₂ solutions,³ it was found to adopt a fluorine-bridged chain structure in the solid state.⁵

The primary synthetic route to WSF₄ has been the solid state substitution of two fluorine ligands in WF_6 by a sulfur ligand using high-temperature reactions and reagents such as $Sb_2S_3 (300 \degree C), ^4B_2S_3 (260 \degree C)$, and sulfur $(300 \degree C), ^5$ followed by purification of WSF₄ by vacuum sublimation.² A solution synthesis utilizing S(Si(CH₃)₃)₂ in CH₃CN and CH₂Cl₂ solvents (eq 1) has also been reported.³ However, variable amounts of a dark brown solid were obtained as a side product. A solvent-assisted dismutation of WSF₄ in CH₃CN

was suggested resulting in the formation of WF₆ and WS₃ according to eq 2.

$$WF_6 + S(Si(CH_3)_3)_2 \xrightarrow{CH_3CN \text{ or } CH_2Cl_2} WSF_4 + FSi(CH_3)_3$$
(1)

$$3WSF_4 \xrightarrow{CH_3CN} 2WF_6 + WS_3$$
 (2)

In spite of these studies, a facile, high-yield synthetic route to pure WSF₄ is still missing. This is primarily a consequence of side-reactions resulting in a dark brown solid that is presumed to be WS_3 . In addition, WSF_4 was found to be photosensitive, changing its color to pale brown upon exposure to direct sunlight.4

The solubility of WSF₄ in CH₃CN suggested the possible formation of a WSF₄·CH₃CN adduct.⁵ Preliminary crystallographic data confirmed adduct formation; however, the data quality was insufficient to determine the correct space group and to provide detailed structural information.

The goal of this study is to provide a reliable and facile high-yield synthesis of pure WSF₄ and WSF₄·CH₃CN, and their full structural characterizations.

Results and Discussion

Synthesis of WSF₄. Pure WSF₄ was synthesized from WF₆ and Sb₂S₃ in anhydrous HF (aHF) solvent at room temperature with an approach analogous to that used for ReSF_4 (eq 3).⁶

$$3WF_{6(HF)}+Sb_2S_{3(s)} \xrightarrow{HF, 14 \ h} 3WSF_{4(HF)}+2SbF_{3(s)} \ (3)$$

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The reaction proceeded slowly at room temperature, yielding a stable yellow HF solution above insoluble SbF3 and incompletely reacted antimony sulfide fluorides. Decantation of the supernatant followed by the removal of all volatiles under vacuum yielded yellow WSF₄. The solubility of isolated solid WSF4 in aHF solvent was somewhat less than that of WSF₄ that was generated in solution, suggesting the presence of monomeric WSF4 in HF solvent and the formation of fluorine-bridged chains in the solid state. It was found that solid WSF₄ changes color from yellow to brown after a day at room temperature in the absence of light. Maintaining the sample at low temperature (below -10 °C) prevented the color change. The color change is possibly the result of (a) slow decomposition or (b) particle coagulation over a period of time that results in the formation of larger aggregates. Although fluorescence of freshly prepared WSF₄ was minimal in the Raman spectrum, a marked increase in fluorescence was observed for older samples. Because the reactivity of brown WSF₄ samples that were stored at room temperature for extended periods of time did not differ from freshly prepared samples, decomposition of the bulk sample seems an unlikely cause of the color change.

Synthesis of WSF₄·CH₃CN. The WSF₄·CH₃CN adduct was synthesized by the reaction of freshly prepared WSF₄ with CH₃CN in aHF solvent below -20 °C (eq 4) and by the reaction of solid WSF₄ with neat CH₃CN solvent at -40 °C (eq 4). Crystallization of WSF₄·CH₃CN from aHF gave WSF₄·CH₃CN in high purity as bright yellow needles. The reaction in neat CH₃CN initially produced a red-brown solution, yielding a gray solid upon removal of all volatiles. The WSF₄·CH₃CN adduct is stable at room temperature under anhydrous conditions and can be used as a stable synthetic substitute for WSF₄. The WSF₄·CH₃CN adduct is moderately soluble in CH₂Cl₂.

$$WSF_4 + CH_3CN \xrightarrow{aHF \text{ or } CH_3CN} WSF_4 \cdot CH_3CN \quad (4)$$

Vibrational Spectroscopy. (a). WSF₄. The Raman spectrum of a solution of WSF_4 in aHF solvent is shown in Figure 1, and the Raman and infrared spectra of solid WSF_4 are shown in Figure 2. The observed Raman frequencies of WSF_4 in HF solution and their detailed assignments are summarized in Table 1 together with the calculated vibrational frequencies for monomeric WSF_4 . The observed vibrational frequencies of solid WSF_4 are listed in Table 2 along with their general assignments. Overall, there is better agreement between experimental and calculated data using the B3LYP method.

The Raman spectrum of an aHF solution of WSF₄ is in agreement with monomeric WSF₄ having $C_{4\nu}$ symmetry as predicted by computational studies and indicated in a previous EXAFS study in CH₃CN and CH₂Cl₂ solvents and a gas-phase electron-diffraction study.^{2,3} The vibrations of monomeric WSF₄ span the irreducible representations $\Gamma = 3A_1 + 2B_1 + B_2 + 3E$ under $C_{4\nu}$ symmetry. All vibrational modes are Raman active, while the A₁ and E modes are infrared active.

The band observed at 690 cm⁻¹ is attributable to the $v_s(WF_4)$ mode. The $v_{as}(WF_4)$ modes, $v_7(E)$ and $v_4(B_1)$, are



Figure 1. Raman spectrum of an aHF solution of WSF₄ in a FEP tube recorded at -15 °C using 1064 nm excitation. Asterisks (*) and dagger (†) denote FEP bands and a laser line, respectively.



Figure 2. Vibrational spectra of WSF₄: Raman spectrum (lower trace) recorded at -100 °C using 1064 nm excitation and the infrared spectrum (upper trace) recorded at room temperature in a KBr pellet. The dagger (†) denotes a laser line.

calculated to be weaker and were not observed in the Raman spectrum. The value reported for matrix-isolated WSF₄, 671 cm⁻¹,¹⁰ is, however, in agreement with the calculated value for $\nu_7(E)$. As predicted by the calculations, the W–S stretching mode appears as the most intense band in the Raman spectrum at 563 cm⁻¹. Another, weaker band is observed at 544 cm⁻¹ and has been assigned to the $\nu(W-^{34}S)$ mode by analogy with the WSF₄·CH₃CN adduct (vide infra). A similar isotopic effect on the W–S stretching frequency was observed for matrix-isolated WSF₄.¹⁰ The bands at 303, 313 cm⁻¹ are assigned to $\delta_{scissoring}(WF_4)$ and that at 241 cm⁻¹ is assigned to $\delta(SWF)$, based on its relative intensity. The absence of bands in the W–F_{bridging} stretching region between 500 and 540 cm⁻¹ confirms the monomeric structure of WSF₄ in HF solution.

The Raman spectrum of pure solid WSF₄ is more complex, reflecting the polymeric nature of WSF₄ in the solid state. The present infrared values are in agreement with those reported earlier.⁴ Previously, only one Raman band had been reported at 577 cm⁻¹, which was attributed to the W–S stretching mode. The nine bands between 635 and 718 cm⁻¹, and the five bands between 554 and 585 cm⁻¹, observed in the present study, can be attributed to W–F_t and W–S stretching modes, respectively, by analogy with those observed for monomeric

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Table 1. Experimental Raman Frequencies and Assignments for WSF₄ Dissolved in aHF and Calculated Vibrational Frequencies for Monomeric WSF₄

	calcd ^{ac}				exptl ^{a,b}		
	$W^{34}SF_4$		² SF ₄	W ³²	4	WSF4	
assgnt $(C_{4v})^{f}$	PBE1PBE	B3LYP	PBE1PBE	B3LYP	matrix ^e	Raman ^d	
$v_1(A_1), v_s(WF_4)$	729(11)[95]	710(13)[92]	729(12)[95]	709(13)[91]	707	690(11)	
$v_7(E), v_{as}(WF_4)$	697(1)[202]	682(1)[196]	697(1)[201]	680(1)[196]	671		
$v_4(B_1), v_{as}(WF_4)$	640(3)[<0.1]	625(3)[0]	640(3)[0]	624(3)[0]			
$v_2(A_1), v(W^{32}S)$			598(27)[53]	579(28)[47]	577.3	563(77)	
$v_2(A_1), v(W^{34}S)$	583(26)[50]	565(26)[44]			562.9	544(4)	
$v_6(B_2), \delta_{scissoring}(WF_4)$	338(2)[0]	331(2)[0]	338(2)[0]	331(2)[0]		313(6) ^g	
o(2), beisbeinig(i)						$303(8)^{g}$	
$v_3(A_1), \delta_{umbrella}(WF_4)$	247(1)[4]	243(1)[5]	247(1)[4]	244(1)[5]			
$v_8(E), \delta_{in-plane}(WF_4)$	238(<1)[19]	235(<1)[18]	238(<1)[19]	235(<1)[19]			
$v_9(E), \delta(SWF)$	205(3)[<0.1]	203(3)[<0.1]	208(3)[0]	206(4)[0]		241(40)	
$v_5(B_1), \delta_{out-of-plane}(WF_4)$	123(<1)[0]	120(<1)[0]	123(<1)[0]	122(<1)[0]			
		. /	. /	. /		94(100)	
modes associated with						72(24)	
solvation						60(54)	

^{*a*} Frequencies are given in cm⁻¹. ^{*b*} Values in parentheses denote relative Raman intensities. ^{*c*} Stuttgart f (W) and aug-cc-pVTZ (H,C,N,O,F,S) basis sets. Values in parentheses denote Raman intensities (Å⁴ u⁻¹). Values in square brackets denote infrared intensities (km mol⁻¹). Values in parentheses denote relative Raman intensities. ^{*d*} The Raman spectrum of WSF₄ dissolved in aHF was recorded in an FEP sample tube at -15 °C using 1064 nm excitation. FEP bands were observed at 293(6), 382(4), 386(4), 574sh, 733(16), 752(3) cm⁻¹; a laser line was observed at 85 cm⁻¹. ^{*e*} Values from ref 10 in a nitrogen-matrix. ^{*f*} The abbreviations denote symmetric (s), asymmetric (as), stretch (ν), bend (δ), and rock (ρ_{rock}). The in-plane and out-of-plane mode descriptions are relative to the WF₄ plane. ^{*g*} These bands overlap with the FEP band at 293 cm⁻¹.

Table 2. Experimental Vibrational Frequencies a and Assignments for SolidWSF4

D. b.c	1b.d	: c1b.e	
Kaman	infrared	intrared	assgnt
718(8)	716 1		
715(8)	/16 sh		
707(10)	<0.5	600	
697(18)	695 s	699	
681(4)			$\gamma V(W-F_{term})$
672(7)	675 sh	673	
654(4)			
642(5)		643	
635(4)	630 vs)
585sh			
578(100)	576 s	577	
568(23)			ν (W–S)
560sh			
554sh	556 s)
535(13)	538 sh	534)
513(17)	514 s	514	$(W_{E_{1}})$
473(2)	463 m		(V I bridge)
457(2)			J
319sh			
309(4)			
293(2)			
263(10)			
248(14)			
236(14)			
225(44)			
218sh			
208(10)			
195(8)			
190sh			
130(6)			
118(5)			
101(26)			
59(74)			

^{*a*} Frequencies are given in cm⁻¹. ^{*b*} Abbreviations denotes shoulder (sh), very strong (vs), strong (s), medium (m), weak (w), very weak (vw). ^{*c*} The Raman spectrum of solid WSF₄ was recorded in a sealed mp capillary at -100 °C using 1064-nm excitation. A laser line was observed at 84 cm⁻¹. ^{*d*} The infrared spectrum of WSF₄ was recorded on a KBr pellet at room temperature. ^{*e*} Values from ref 4.

WSF₄. The most intense Raman band at 578 cm⁻¹ occurs 15 cm⁻¹ higher than that of WSF₄ in solution which is counterintuitive because monomeric WSF₄ should

exhibit more covalent W–S bonding when compared with a fluorine-bridged chain. However, vibrational coupling in the fluorine-bridge chain can easily result in a high-frequency shift of a W–S stretching band. The modes around 473 and 513 cm⁻¹ in the Raman spectrum (463 and 514 cm⁻¹ in the infrared spectrum) do not appear in the solution spectrum and are confidently assigned to the W–F_b stretching modes. Such modes are indeed expected in view of the chain structure found by X-ray crystallography. The modes observed below 319 cm⁻¹ are all associated to deformation modes which are likely strongly coupled because of the chain structure of WSF₄. The intense Raman band at 225 cm⁻¹ can tentatively be assigned to δ (SWF) in comparison with the Raman spectrum of WSF₄ in HF solvent.

The W-F_t stretching modes for solid WSF₄ appear at lower frequency than those in WOF₄ (660 and 725 cm⁻¹)⁸ which adopts a fluorine-bridged tetrameric structure in the solid state,⁷ in agreement with the greater electronegativity of oxygen, which is expected to result in an electron-poorer tungsten center and a more covalent W-F bond. The W-F_b stretches in WSF₄ appear at even lower frequency than those in WOF₄ (520 and 560 cm⁻¹), indicating much weaker fluorine bridging in WSF₄.

(b). $WSF_4 \cdot CH_3CN$. The Raman and infrared spectra of $WSF_4 \cdot CH_3CN$ are shown in Figure 3. The observed vibrational frequencies of $WSF_4 \cdot CH_3CN$, their assignments, and the calculated frequencies are summarized in Table 3. The assignments of the modes are based on the calculated frequencies and Raman intensities of the energy-minimized geometry. Because some of the bands exhibit splittings, a factor-group analysis was carried out (Supporting Information, Table S1).

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Figure 3. Vibrational spectra of WSF₄·CH₃CN: Raman spectrum (lower trace) recorded at -100 °C using 1064 nm excitation and the infrared spectrum (upper trace) recorded at room temperature in a KBr pellet. The dagger (†) denotes a laser line.

The WSF₄·CH₃CN adduct in the crystal structure adopts C_s point symmetry (vide infra) with a slightly bent W-N-C moiety. When neglecting the hydrogen atoms of the methyl group, the idealized symmetry of the F₄SWNCC framework is $C_{4\nu}$ point symmetry, and its vibrational modes span the irreducible representations $\Gamma = 6A_1 + C_1$ $2B_1 + B_2 + 6E$, all modes being Raman active, while only the A_1 and E modes are infrared active. The vibrational analysis using C_{4v} point symmetry is warranted because no splitting of the experimental and calculated E modes is observed and the lowering of the overall symmetry by the slight deviation from linearity of the SWNCC moiety and the presence of the CH_3 group to C_s or lower symmetry does not have a significant influence on the vibrational modes. The symmetry labels for the F₄SWNCC framework under $C_{4\nu}$ symmetry are given in Table 3.

The vibrational frequencies that are attributable to the CH₃CN ligand all exhibit complexation shifts, compared to free CH₃CN, that are in good agreement with other adducts of CH₃CN and transition metal oxide fluorides, such as $\text{ReO}_2\text{F}_3 \cdot \text{CH}_3\text{CN}$.⁹ The strong bands observed at 674 and 662 cm⁻¹ are assigned to the $\nu_s(\text{WF}_4)$ mode, and are significantly lower than that in monomeric WSF₄ in HF solution (690 cm^{-1}) or in the gas-phase (707 cm^{-1}), reflecting the increase in bond polarity upon adduct formation. The infrared band at 619 cm^{-1} is assigned to the doubly degenerate $\nu_{as}(WF_4)$ mode and is very strong, as predicted from the calculations. This band also appears at lower frequency than in WSF₄ in the gas-phase (671 cm⁻¹). The bands at 593 and 599 cm⁻¹ are assigned to the $\nu_{as}(WF_4)$ mode of B₁ symmetry. The bands at 544, 555, and 561 cm⁻¹ are attributed to W–S stretching modes. Several spectra were recorded for the adduct, and they all revealed the presence of three bands and an invariant intensity ratio, confirming that the band at 561 cm⁻¹ is not due to an excess of WSF₄, but originates from vibrational coupling in the crystallographic unit cell, consistent with the factor-group analysis, which predicts the splitting of all Raman modes into two Raman-active components (Supporting Information, Table S1). The

strongest band at 555 cm⁻¹ is readily assigned to a W-S stretching band, and the band at 544 cm⁻ is assigned to an isotopic shift resulting from the $W^{34}SF_4 \cdot CH_3CN$ isotopomer. The assignment is confirmed by the experimental shift, 11 cm^{-1} , which is well reproduced by a calculated shift of 14 cm^{-1} at both levels of theory. The W-S stretching frequency is also lower when compared with the W-S stretching frequency (563 cm^{-1}) of monomeric WSF₄. The W–N stretching mode is predicted to occur at very low frequency with a very low intensity, and was consequently not observed. The low frequency reflects the weakness of the coordinate bond. Microcrystalline bulk samples of WSF₄·CH₃CN that were produced by reactions in HF and CH₃CN solvents frequently contained unidentified impurities that gave rise to broad Raman and infrared bands at 800 cm^{-1} .

NMR Spectroscopy. The solubility of WSF₄ in HF solvent allowed for its study by ¹⁹F NMR spectroscopy in that solvent. The $^{19}\mathrm{F}$ NMR spectrum of $\bar{W}SF_4$ in HF showed a singlet at 77.4 ppm with satellites due to a ${}^{1}J({}^{183}W-{}^{19}F)$ coupling of 42 Hz. The chemical shift is significantly lower than that observed for WSF₄ in the donor solvent CH₃CN (85.4 ppm), which is in good agreement with previously reported values,^{3,4} corroborating the formation of the WSF₄·CH₃CN adduct characterized by Raman spectroscopy and X-ray diffraction. The chemical shift of WSF4·CH3CN dissolved in CH₂Cl₂ exhibited a similar chemical shift at 86.1 ppm; ${}^{1}J({}^{1\bar{8}3}W-{}^{19}F)$ coupling in CH₂Cl₂ could not be resolved as a consequence of the line width, $\Delta v_{1/2} = 21$ Hz. A smaller chemical shift difference between the adduct and the uncomplexed species is observed for WOF_4 (aHF solvent: 60.8 ppm, ${}^{1}J({}^{183}W-{}^{19}F) = 68$ Hz; CH₃CN solvent: 66.7 ppm), which was observed in some of the NMR samples as a result of small degree of hydrolysis.

Crystal Structures of WSF₄ and $WSF_4 \cdot CH_3CN$. Details of the data collection parameters and other crystallographic information for WSF₄ and WSF₄ \cdot CH₃CN are given in Table 4 while important bond lengths and bond angles are listed in Table 5.

WSF₄. Although the published structure is crystallographically correct, it was of average quality, with some issues as indicated by the severe distortion of several of the thermal ellipsoids.⁵ The present structure of WSF₄ consists of fluorine-bridged chains and contains two crystallographically independent bridged WSF₄ moieties (Figure 4a). The fluorine bridge is asymmetric, with one short bridging W-F_b distance, 1.932(5) and 1.947(5) A, and one long bridging W--- F_b distance, 2.323(4) and 2.308(5) Å. The bridging fluorine forms the long contact trans to the adjacent W=S bond. The cis-bridging results in meandering chains that run along the *c*-axis; however, there are no intermolecular contacts between adjacent chains (Figure 4b). The W=S bond lengths in solid WSF₄ (2.084(2) and 2.091(2) Å) are not significantly different from those in the gas-phase, indicating the weakness of the W---F_b bridging contact. The F_t-W---F_b angles are significantly smaller compared to the $F_t - W - F_t$ angles due to the long W---F_b distances. The $W-F_{b}--W$ angles range from 160.5(2) to 147.0(3)°, indicating the weakness of the fluorine bridge and the flexibility of the angle.

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Table 3.	Experimental and	Calculated R	aman Frequencies	and Intensities for	r WSF₄∙CH ₃ CN	and CH ₃ CN an	d Calculated	Infrared Intensi	ties for WSF ₄	·CH ₃ CN and
CH ₃ CN	-		_							
, V	VSE. CH. CN and	1 CH-CN								

	exptl ^a	una erișeri			calcd ^{<i>a</i>}	b				
WSF ₄ ·CH ₃ CN		CH ₃ CN	W ³² SF ₄ ·C	$H_3CN(C_1)$	W ³⁴ SF ₄ ·C	$H_3CN(C_1)$	CH ₃ C	$N(C_{3v})$	-	assgnt $\{C_{dr}\}^f$
Raman ^c	ir ^{d,e}	Raman	B3LYP	PBE1PBE	B3LYP	PBE1PBE	B3LYP	PBE1PBE	-	
3019(3)		3001(7)	3124(69)[<1] 3124(69)[<1]	3154(67)[1] 3153(67)[1]	3124(69)[<1] 3124(69)[<1]	3154(67)[1] 3153(67)[1]	3115(59)[<1]	3147(56)[<1]	-}	$\nu_{as}(\rm CH_3)$
2953(34)	2983 vw 2916 vw 2364 w 2343 sh	2943(100) 2888(3) 2847(2) 2732(2)	3052(256)[<0.1]	3070(256)[<1]	3052(256)[<0.1]	3070(256)[<1]	3048(195)[3]	3067(193)[2]	J	$v_{s}(CH_{3})$ combination bands
2313(24)	2322 w 2308w	2293(5)								
2286(25)	2280 w 2253 w 2217 w 1696 w 1580 w	2253(60) 2204(1)	2409(228)[88]	2437(229)[98]	2409(228)[88]	2437(229)[98]	2363(81)[11]	2389(80)[13]		$v(C=N) \{v(A_1)\}$ combination bands
1364(7)	1405 m 1367w 1259 m	1443(3) 1416(2) 1375(6)	1468(6)[12] 1467(6)[12] 1411(12)[1]	1457(6)[13] 1457(6)[13] 1397(10)[2]	1468(6)[12] 1468(6)[12] 1411(12)[1]	1457(6)[13] 1457(6)[13] 1398(10)[2]	1475(5)[10] 1413(6)[2]	1465(5)[12] 1400(6)[3]	}	$\delta_{as}(CH_3)$ $\delta_s(CH_3)$ combination bands
	1191 w 1027 m		1062(<1)[3] 1062(<1)[3]	1051(<1)[4] 1051(<1)[4]	1062(<1)[3] 1062(<1)[3]	1051(<1)[4] 1051(<1)[4]	1063(<1)[2]	1053(<1)[2]	}	ρ _r (CH ₃)
939(5)	996 m 935 sh	919(18)	944(7)[13]	968(8)[12]	945(8)[13]	968(8)[12]	928(5)[1]	948(5)[1]	J	combination band $\nu(C-CH_3) \{\nu(A_1)\}$
674(23) 662 (2)	692 s	}	685(14)[71]	704(13)[73]	685(14)[73]	704(13)[76]				$\nu_s(WF_4)~\{\nu(A_1)\}$
	619 vs		654(<1)[207] 654(<1)[207]	668(<1)[212] 668(<1)[212]	654(<1)[207] 654(<1)[207]	669(<1)[213] 669(<1)[213]				$\nu_{as}(WF_4)~\{\nu(E)\}$
599(2) 593(4) 561(20)		}	607(3)[<0.1]	622(3)[<0.1]	607(3)[<0.1]	622(3)[<0.1]				$\nu_{as}(\mathrm{WF_4})\;\{\nu(B_1)\}$
555(100) 544(9)	558 s 539 sh 450 s	}	563(51)[103]	580(50)[114]	549(49)[99]	566(48)[110]				$\begin{array}{l} \nu(W^{32}S) \; \{\nu(A_1)\} \\ \nu(W^{34}S) \; \{\nu(A_1)\} \\ combination \; band \end{array}$
406(4)		380(10)	409(1)[1] 409(1)[1]	415(1)[1] 415(1)[1]	409(1)[1] 409(2)[1]	415(1)[1] 415(1)[1]	382(1)[<1]	387(1)[<1]		$\delta(CCN) \{v(E)\}$
301(7)			308(2)[0] 263(<1)[23]	313(2)[0] 269(<1)[23]	307(2)[0] 263(<1)[23]	313(2)[0] 269(<1)[23]				$\begin{array}{l} \delta_{scissoring}(WF_4) \; \{\!\nu(B_2)\} \\ \delta_{umbrella}(WF_4) \; \{\!\nu(A_1)\} \end{array}$
249(21)			236(2)[3] 236(2)[3] 222(<1)[14]	242(2)[4] 242(2)[4] 222(<1)[14]	235(2)[4] 235(2)[4] 221(51)[12]	241(2)[4] 241(2)[4] 222(<1)[12]				$\delta_{in \text{-plane}}(WF_4) \ \{v(E)\}$
220(11)			222(<1)[14] 222(<1)[14]	223(<1)[14] 223(<1)[14]	221(<1)[13] 221(<1)[13]	222(<1)[13] 222(<1)[13]				$\delta(SWF_4) \{v(E)\}$
1(1(20)			180(<1)[<0.1] 147(3)[<0.1]	185(<1)[<0.1] 152(3)[<0.1]	181(<1)[<0.1] 146(3)[<0.1]	185(<1)[<0.1] 151(3)[<0.1]				$\partial_{out \circ of plane} (WF_4) \{V(B_1)\}$
161(32)			147(3)[<0.1]	152(3)[<0.1]	146(3)[<0.1]	151(3)[<0.1]				o(SWN) {V(E)}
			140(<1)[15] 37(<0.1)[4]	153(<1)[16] 38(<0.1)[4]	140(<1)[5] 37(<0.1)[5]	153(<1)[16] 38(<0.1)[4]				$v(w N) \{v(A_1)\}$
			37(<0.1)[4]	38(<0.1)[4]	37(<0.1)[5]	38(<0.1)[4]				ρ _i (CH ₃)
104 (3)			10(0)[<0.1]	24(0)[≤0.1]	J(≦0.1)[≤0.1]	∠4(0)[<0.1]				ρ _t (CH ₃) lattice mode

^{*a*} Frequencies are given in cm⁻¹. ^{*b*} Stuttgart f (W) and aug-cc-pVTZ (H,C,N,O,F,S) basis sets. Values in parentheses denote Raman intensities (Å⁴ u⁻¹). Values in square brackets denote infrared intensities (km mol⁻¹). ^{*c*} The Raman spectrum of WSF₄·CH₃CN was recorded in a sealed mp capillary at -100 °C using 1064 nm excitation. Values in parentheses denote relative Raman intensities. ^{*d*} The infrared spectrum of WSF₄·CH₃CN was recorded on a KBr pellet at ambient temperature; broad. ^{*e*} Broad infrared band at 885 from unidentified impurity. ^{*f*} The abbreviations denote symmetric (s), asymmetric (as), stretch (ν), bend (δ), and rock (ρ_{rock}). The in-plane and out-of-plane mode descriptions are relative to the WF₄ plane.

Table 4. Crystallographic Data for WSF_4 and $WSF_4 \cdot CH_3CN$

chem formula	WSF_4	$WSF_4 \cdot CH_3CN$
space group	<i>Pca</i> 2 ₁ (No. 29)	<i>Pnma</i> (No. 62)
a (pm)	16.886(3)	1099.5(3)
b (pm)	5.3754(8)	787.4(3)
c (pm)	9.4435(14)	804.6(2)
molecules/unit cell	8	4
mol wt (g mol^{-1})	291.91	332.96
calcd density (g cm ^{-3})	4.524	3.175
T (°C)	-120	-120
$\mu (mm^{-1})$	27.381	16.87
R_1^a	0.0201	0.0123
wR_2^{b}	0.0412	0.0301

^{*a*} R_1 is defined as $\sum_{i=1}^{n} ||F_0| - |F_c|| / \sum_{i=1}^{n} |F_0|$ for $I > 2\sigma(I)$. ^{*b*} wR_2 is defined as $[\sum_{i=1}^{n} w(F_0^{-2} - F_c^{-2})^2 / \sum_{i=1}^{n} w(F_0^{-2})^2]^{1/2}$ for $I > 2\sigma(I)$.

WSF₄•**CH**₃**CN.** The WSF₄•**CH**₃**CN** adduct crystallizes in the orthorhombic space group *Pnma*. The data quality in a previous study was insufficient to distinguish between the *Pna*2₁ or *Pnam* (conventional space group, *Pnma*) space groups.⁵ In the crystal structure (Figure 5a), the WSF₄•**CH**₃**CN** adduct is located on a crystallographic mirror plane and the CH₃**CN** ligand is coordinated to tungsten trans to the sulfur with an essentially linear S–W---N arrangement and a slightly bent W---N–C geometry (174.6(3)°). The W---N distance is rather long, implying rather weak coordination, in agreement with the small complexation shift observed for the ν (CN) mode in the adduct compared to that of free CH₃CN.

The W–S bond length (2.0959(11) Å) is slightly elongated or equal within $\pm 3\sigma$ when compared with the W–S bond lengths of solid WSF₄ (2.084(2) and 2.091(2) Å); the elongation reflects the greater Lewis basicity of N in CH₃CN compared to the "fluoride ion" of the bridge in solid WSF₄. A similar trend is observed between the W–F_t bond lengths in WSF₄·CH₃CN (1.856(2) and 1.857(2) Å) and the W–F_t bonds in WSF₄ (1.835(5) and 1.852(4) Å). Both the W–S and W–F_t bond lengths are comparable to those in the gas-phase.² The S–W–F angles with 100.74(6) and 100.12(6)° are significantly greater than 90° as a consequence of the greater steric demand of the sulfide group relative to that of the weakly coordinated CH₃CN ligand.

Computational Results. The electronic structures of WSF₄ and WSF₄·CH₃CN were optimized under $C_{4\nu}$

Table 5. Experimental and Calculated Metric Parameters for WSF₄ and WSF₄·CH₃CN

		WSF ₄	·CH ₃ CN					
	exptl (polymeric)	exptl (mo	nomeric)	calcd (n	nonomeric)	exptl	с	alcd
	X-ray	gas-phase ^a	$EXAFS^b$	B3LYP ^c	PBE1PBE ^c	X-ray	B3LYP ^c	PBE1PBE ^c
			Bond	Lengths (Å))			
$W-F_t$ $W-F_b$ WF_b	1.835(5) to 1.852(4) 1.932(5), 1.947(5) 2.308(5), 2.323(4)	1.847(3)	1.863(3)	1.859	1.845	1.8564(15),1.8572(17)	1.871	1.857
W-S WN N-C1 C1-C2	2.084(2), 2.091(2)	2.104(7)	2.026(8)	2.108	2.094	2.0959(11) 2.369(3) 1.130(5) 1.453(5)	2.120 2.494 1.144 1.450	2.107 2.450 1.143 1.444
			Bond	Angles (deg)			
S-W-F _t S-WF _b S-W-F _t	101.37(18) to 102.59(19) 177.86(15), 177.93(15) 99.48(16), 99.86(16)	104.5(11)		105.5	105.7	100.74(6), 100.12(6)	101.7	101.8
$S W T_{b}$ $F_{t}-W-F_{t}$ $cis-F_{t}-W-F_{b}$ $F_{t}-W-F_{b}$ $F_{b}-W-F_{b}$ $W-F_{b}-W-W$	84.0(2) to 91.9(2) 155.0(2) to 158.8(2) 83.9(2) to 85.1(2) 77.2(2) to 80.1(2) 78.07(19), 78.89(18) 160 5(2) 147 0(3)	86.6(5)	85.9 148.9	85.8 148.6		87.99(11), 88.30(14) 155.0(2), 158.8(2)	87.6 156.5	87.6 156.3
S-WN F _t -WN WN-C N-C-C	100.0(2), 1110(3)					179.32(9) 79.39(8), 79.74(8) 174.6(3) 179.0(4)	180 78.3 179.6 180	180 78.2 179.9 180

^a Values from ref 2. ^b Values from ref 3. ^c Stuttgart f (W) and aug-cc-pVTZ (H,C,N,O,F,S) basis sets.



Figure 4. (a) X-ray crystal structure of WSF₄; thermal ellipsoids are shown at the 50% probability level. (b) The WSF₄ chain viewed along the *c*-axis. (c) Calculated geometry of monomeric WSF₄.

symmetry at the B3LYP and PBE1PBE levels and resulted in stationary points with all frequencies real (Figures 4c and 5b). The energy-minimized geometries and vibrational frequencies of WF₆ were also calculated (Supporting Information, Table S2) to serve as a benchmark. The calculated geometrical parameters for WSF₄ and WSF₄·CH₃CN are listed in Table 5.

The calculated parameters for WSF_4 agree very well with the experimental gas-phase parameters obtained



Figure 5. (a) X-ray crystal structure of WSF₄·CH₃CN; thermal ellipsoids are shown at the 50% probability level. (b) The calculated geometry of the WSF₄·CH₃CN adduct.

from an electron-diffraction study, and those for $WSF_4 \cdot CH_3CN$ agree well with the present experimental structure. As expected, the largest discrepancy occurs for the $W \cdot \cdot \cdot N$ contact, which is overestimated at both levels.

Conclusions

Pure WSF₄ can be prepared by means of a facile synthetic route in aHF solvent. Solutions of WSF₄ in aHF were studied for the first time and were shown to contain monomeric $C_{4\nu}$ symmetric WSF₄. The WSF₄·CH₃CN adduct was formed in the donor-solvent CH₃CN and in aHF solvent in the presence of CH₃CN. For the first time, the WSF₄·CH₃CN adduct could be unambiguously characterized by vibrational spectroscopy in

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conjunction with quantum-chemical calculations and X-ray diffraction. The $WSF_4 \cdot CH_3CN$ adduct can serve as the synthetic equivalent to monomeric WSF_4 , with its reactivity currently under investigation in our laboratory.

Experimental Section

Materials and Apparatus. All volatile materials were handled (a) on a Pyrex vacuum line equipped with glass/Teflon J. Young valves and (b) a vacuum line constructed of nickel, stainless steel, and FEP. Nonvolatile materials were handled in the dry nitrogen atmosphere of a drybox (Omni Lab, Vacuum Atmospheres). Reaction vessels and NMR sample tubes were fabricated from $\frac{1}{4}$ in. o.d. thin-wall ($\frac{3}{16}$ in. i.d.) and 4 mm o.d., 2.8 mm i.d. FEP tubing, respectively, and outfitted with Kel-F valves. All reaction vessels and sample tubes were rigorously dried under dynamic vacuum prior to passivation with 1 atm of F₂ gas.

Anhydrous HF (Air Products, 99.9%) was dried over K_2NiF_6 . The solvents, CH_3CN (Sigma-Aldrich, HPLC grade) was dried according to standard literature methods¹¹ and CH_2Cl_2 was dried using molecular sieves (type 4 Å). Tungsten hexafluoride (Elf Atochem) was used without further purification, whereas Sb₂S₃ (Alfa Aesar, 99.5%) was purified by drying under dynamic vacuum at 160 °C for about 6 h.

Preparation of WSF₄. Inside the drybox, dried Sb₂S₃ (0.136 g, 0.400 mmol) was added to a 1/4 in. FEP T-reactor equipped with a Kel/F valve. After distillation of 1.76 g of aHF, WF₆ (0.326 g, 1.09 mmol) was condensed onto the frozen reaction mixture at -196 °C. The reactor was backfilled with dry N_{2(g)} and allowed to warm to room temperature. After warming to room temperature, and with intermittent agitation over the course of 16 h, the reaction progressed slowly, going from a colorless solution above black Sb_2S_3 to a deep yellow solution. The deep yellow solution was decanted to the side arm of the FEP T-reactor where the solution was cooled to about -78 °C causing yellow WSF₄ to precipitate. More WSF₄ was extracted from the solid reaction mixture by condensing HF solvent back onto the solid and washing the solid again at room temperature, followed by decanting the yellow solution into the side arm. The washing procedure was repeated 2 to 3 times. Anhydrous HF was removed under vacuum at -78 °C over several hours, leaving some yellow crystalline material that was pumped on for an additional 10 min at room temperature to ensure total HF removal. The amount of recovered WSF₄ was 0.258 g (0.867 mmol) with a yield of 79.3%.

Preparation of WSF₄·**CH₃CN.** (a) Inside a glovebox, 0.055 g (0.16 mmol) of Sb₂S₃ was transferred into the straight arm of a T-reactor fabricated out of 1/4 in. o.d. FEP tubing. After vacuum distillation of about 0.27 mL of aHF onto Sb₂S₃, 0.172 g (0.577 mmol) of WF₆ was condensed onto the Sb₂S₃ at -196 °C by vacuum-distillation. The reaction was allowed to proceed at room temperature overnight, producing a bright yellow solution above a dark gray precipitate. After decanting the solution into the side arm, 0.044 g (1.07 mmol) of CH₃CN was then vacuumdistilled into the side arm at -196 °C. The reaction mixture was vigorously agitated at -35 °C, resulting in a bright yellow solution with some undissolved bright yellow solid. The mixture was then slowly cooled to $-61 \,^{\circ}$ C to allow complete precipitation of WSF₄·CH₃CN. The remaining HF was decanted back into the main arm, frozen at -196 °C, and the main arm was heat sealed under dynamic vacuum. Volatiles in the side arm were removed under dynamic vacuum while the cooling bath was allowed to slowly warm to 13 °C. Greenish-yellow WSF₄·CH₃CN (0.1283 g, 0.3853 mmol) was collected inside the drybox in 79.3% yield. (b) Inside a dry nitrogen atmosphere drybox, 0.016 g of WSF₄ (0.055 mmol) was transferred into a $\frac{1}{4}$ in. thin wall FEP reactor equipped with a Kel-F valve. Approximately 0.35 mL of dry CH₃CN was vacuum distilled onto the solid. Slow melting of the solvent at -43.3 °C resulted in a dark brown solution. After 5 min of agitation, volatiles were removed under dynamic vacuum while the reaction mixture was allowed to slowly warm to -30 °C. Gray-colored WSF₄·CH₃CN (0.018 g, 0.054 mmol) was recovered in 98% yield.

Vibrational Spectroscopy. The Raman spectra of WSF₄ and WSF₄·CH₃CN were recorded on a Bruker RFS 100 FT Raman spectrometer with a quartz beam splitter, a liquid-nitrogencooled Ge detector, and a low-temperature accessory. The backscattered (180°) radiation was sampled. The actual usable Stokes range was 50 to 3500 cm^{-1} with a spectral resolution of 2 cm⁻¹. The 1064 nm line of an Nd:YAG laser was used for excitation of the sample. All Raman spectra were corrected for effects arising form the optics and the frequency dependence of the Raman scattering by using the white light spectrum of a tungsten lamp. The low-temperature spectra of WSF4 and WSF₄·CH₃CN were recorded on a powdered sample in powdered samples in melting point capillaries using laser powers of 150 and 100 mW, respectively. The FT-infrared spectra were recorded on a Nicolet Avatar 360 FTIR spectrometer at ambient temperature. A KBr sandwich was formed in a Wilks minipress inside the drybox by sandwiching the sample between two layers of KBr. The spectra were acquired in 64 scans at a resolution of 2 cm^{-1} .

Nuclear Magnetic Resonance Spectroscopy. All NMR spectra were recorded unlocked on a 300 MHz Bruker Avance II NMR spectrometer equipped with a 5 mm broadband probe. Fluorine-19 (282.404 MHz) NMR spectra were referenced externally to neat CFCl₃ at 25 °C. The ¹⁹F NMR spectra were typically acquired in 128 K memory with spectral settings of 56 kHz, yielding an acquisition time of 1.15 s and a data point resolution of 0.433 Hz/data point. The number of transients accumulated was 100 using a pulse width of 10.3 μ s.

X-ray Crystal Structure Determination of WSF₄ and WSF₄·CH₃CN. (a). Crystal Growth and Crystal Mounting. Crystals of WSF₄ were grown directly from a HF solution upon slow removal of solvent HF at -78 °C under dynamic vacuum. Crystals of WSF₄·CH₃CN were grown in aHF solvent at a temperature slowly varied from -53 to -59 °C. Crystals of WSF₄ and WSF₄·CH₃CN having the dimensions $0.10 \times 0.06 \times$ 0.03 mm³ and $0.55 \times 0.33 \times 0.24$ mm³, respectively, were selected at -80 °C for low-temperature X-ray structure determination under a flow of cold nitrogen and mounted as previously described.¹²

(b). Collection and Reduction of X-ray Data. The crystal was centered on a Bruker SMART APEX II diffractometer, equipped with an APEX II 4K CCD area detector and a triple-axis goniometer, controlled by the APEX2 Graphical User Interface (GUI) software,¹³ and a sealed source emitting graphite-monochromated Mo K α radiation ($\lambda = 0.71073$ A). Diffraction data collection at -120 °C consisted of four ω scans at various φ settings of 366 frames each at a fixed $\chi = 54.74^{\circ}$ with a width of 0.5° . The data collection was carried out in a 512×512 pixel mode using 2×2 pixel binning. Processing of the raw data was completed by using the APEX2 software,¹³ which applied Lorentz and polarization corrections to three-dimensionally integrated diffraction spots. The program SADABS¹⁴ was used for the scaling of diffraction data, the application of a decay correction, and an empirical absorption correction on the basis of the intensity ratios of redundant reflections.

(c). Solution and Refinement of the Structure. The XPREP program was used to confirm the unit cell dimensions and the

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crystal lattice. The solutions were obtained by direct methods which located the positions of the atoms defining two bridged WSF₄ moieties for WSF₄, and defining WSF₄ and the carbon and nitrogen atoms of the CH₃CN molecules for WSF₄·CH₃CN. The final refinement was obtained by introducing anisotropic thermal parameters and the recommended weightings for all of the atoms. The positions of the hydrogen atoms were calculated. The maximum electron densities in the final difference Fourier map were located near the heavy atoms. All calculations were performed using the SHELXTL-plus package for the structure determination and solution refinement and for the molecular graphics.¹⁵

Computational Methods. The optimized geometries and frequencies of WF₆, WSF₄, and WSF₄·CH₃CN were calculated by the density functional theory (DFT) method at the PBE1PBE and B3LYP¹⁶ levels. The Stuttgart basis set augmented by one f-type polarization function (α_f W 0.823)¹⁷ for tungsten and aug-cc-pVTZ basis sets for oxygen, fluorine, nitrogen,

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carbon, and hydrogen was used. Pseudopotentials were used for tungsten.

Quantum-chemical calculations were carried out using Gaussian 98 and Gaussian 03.^{16,18} The levels and basis sets were benchmarked by calculating WF_6 and comparing with the experimental geometries¹⁹ and vibrational frequencies.²⁰ The geometries were fully optimized using analytical gradient methods. After optimization at one level of theory, the geometries were calculated at the other level of theory to ensure an equivalent energy-minimized geometry had been achieved. The vibrational frequencies were calculated at the PBE1PBE and B3LYP levels using the appropriate minimized geometry, and the vibrational mode descriptions were assigned with the aid of Gaussview.²¹

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Supporting Information Available: Factor-group analysis for $WSF_4 \cdot CH_3CN$ (Table S1); benchmark computational study for WF_6 (Table S2); full version of refs 16 and 18; X-ray crystallographic file in CIF format for the structure determination of $WSF_4 \cdot CH_3CN$. This material is available free of charge via the Internet at http://pubs.acs.org.

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