

histidine at pD 8.6 and 12.4, where the imidazole ring is nonprotonated. The upfield shift is indicative of an increase in electron density about the ring protons at pD >11, as a result of ionization of the "pyrrole" hydrogen. P.m.r. data support amino coordination in the Zn(II)-histidine complex at pD 11.8, for the alanine resonances lie between those observed for free histidine at pD 8.6 and 12.4, where the amino group is protonated and nonprotonated, respectively. Furthermore, the alanine resonances lie upfield of those observed in the Zn(II) complex at pD 5.0. These facts are consistent with replacement of the proton on the amino group by the less electronegative Zn(II) ion in going from pD 5.0 to 11.8.

In summary, the particular value of the infrared spectral data, and the ancillary p.m.r. data, is that the structures of the coordinated species at a given pD are deducible from correlation of the spectra with the spectra of free ligand as a function of pD. The evaluation of formation constants is not very precise, but the agreement of the results in this area with the expectations based on more precise methods is sufficiently good to lend support to the interpretations placed on the infrared data. In addition, observation of certain frequency shifts in favorable cases, such as the anti-symmetric carboxylate stretching mode, yields information which relates closely to the relative enthalpies of coordinate bond formation.

CONTRIBUTION FROM THE INSTITUTE FOR ATOMIC RESEARCH AND DEPARTMENT OF CHEMISTRY
IOWA STATE UNIVERSITY, AMES, IOWA

A Redetermination of the Crystal Structure of CsCuCl₃^{1a}

By ALBERT W. SCHLUETER, ROBERT A. JACOBSON, AND ROBERT E. RUNDLE^{1b}

Received August 23, 1965

The crystal structure of CsCuCl₃ has been redetermined by single crystal X-ray diffraction techniques. Data were taken by scintillation counter methods and the atomic parameters were refined by anisotropic least-squares methods. The final parameters are essentially those originally reported by Wells. The structure consists of distorted hexagonally close-packed layers of cesium and chlorine atoms with the copper atoms occupying octahedral holes. The Cu-Cu distance is a very short 3.062 Å, but direct copper-copper bonding is ruled out by overlap integral calculations. The close approach of the coppers is considered to be a result of the copper atoms sharing faces of their chlorine octahedra instead of edges as has been found in all other halogen-bridged copper compounds. This interpretation differs from that of Wells. Several other metal-halogen systems having chain structures similar to that found in CsCuCl₃ are considered and the overlap integrals for metal-metal interaction calculated. From the calculations it is apparent that d-d metal-metal bonding for the elements at the right of the transition metal series is unlikely. The implications of this are related to the much discussed Cu₂(C₂H₃O₂)₂·2H₂O.

Introduction

The structure determination of CsCuCl₃ was first attempted by Klug and Sears in 1946,^{2a} but they were unsuccessful in solving the structure. Wells^{2b} next investigated the problem and obtained a model which gave an agreement factor, $R = \Sigma ||F_o| - |F_c|| / \Sigma |F_o|$, of 20% on the basis of 60 reflections.

Because of their interesting structural and magnetic features, several related copper compounds have recently been investigated in this laboratory, namely, LiCuCl₃·2H₂O,^{3,4} KCuCl₃, and NH₄CuCl₃.⁵ These structures all consisted of Cu₂Cl₆²⁻ dimers associated into chains by long unsymmetric Cu-Cl bridges and are quite different from the CuCl₃⁻ infinite chain found for CsCuCl₃ by Wells. Thus CsCuCl₃ should represent a

different and particularly simple system to study magnetically because of its essentially one-dimensional character. Preliminary to such a study we decided to carry out a more accurate determination of its structure, which is being reported here. A magnetic investigation of this compound is currently underway, by B. C. Gerstein of this laboratory.

Experimental Section

Dark red-brown crystals of CsCuCl₃ were obtained by evaporating to dryness a solution containing copper and cesium chlorides. A molar excess of copper chloride was used to prevent formation of Cs₂CuCl₄. The crystals appeared as hexagonal prisms often capped by hexagonal pyramids.

Weissenberg photographs exhibited 6/mmm Laue symmetry. The only systematic extinction observed was $l \neq 6n$ for 00 l reflections, indicating the space group P6₃22 and confirming Wells' observation. The lattice constants were determined by a least-squares extrapolation treatment of twelve reflections in the back-reflection region observed using a single crystal orienter. The orienter was previously aligned with an aluminum single crystal. The lattice constants found are: $a = b = 7.2157 \pm 0.0005$ and $c = 18.1777 \pm 0.0010$ Å.

Complete three-dimensional data to $\sin \theta/\lambda = 0.904$ were taken with zirconium-filtered molybdenum K α radiation using a General

(1) (a) Work was performed at the Ames Laboratory of the U. S. Atomic Energy Commission. Contribution No. 1780. (b) Deceased, Oct. 9, 1963.

(2) (a) H. P. Klug and G. W. Sears, Jr., *J. Am. Chem. Soc.*, **68**, 1133 (1946); (b) A. F. Wells, *J. Chem. Soc.*, 1662 (1947).

(3) P. H. Vossos, L. D. Jennings, and R. E. Rundle, *J. Chem. Phys.*, **32**, 1590 (1960).

(4) P. H. Vossos, D. R. Fitzwater, and R. E. Rundle, *Acta Cryst.*, **16**, 1037 (1963).

(5) R. D. Willet, et al., *J. Chem. Phys.*, **38**, 2429 (1963).

TABLE I
OBSERVED AND CALCULATED STRUCTURE FACTORS (L, 10Fo, 10Fc)

Table with multiple columns of structure factor values and their corresponding reflections, organized by L, 10Fo, and 10Fc values.

* 0+0+, 20L(35;21) 0+0+2, 32L(19;31) 0+1-0, 227L(31;32) 2+0+3, 2+1L(35;39) 2+0+4, 2+0L(20;23) 2+0+9, 237L(35;52) 2+0+0, 2+3L(29;54)

These reflections affected by extinction, Fo (Fc), are not on an absolute scale.

Electric XRD-5 X-ray unit equipped with a single crystal orienter. For each of the 1092 reflections a total count was obtained from a 100-sec. scan using a scan rate of 1°/min. A correction for background and streak was made from the recorder trace for each reflection. The recorder was calibrated and zeroed frequently to obtain good accuracy. An absorption correction was made using Busing's polyhedral absorption correction program rewritten for the IBM 7074 computer.

Structural Refinement

The parameters of Wells were tested and an agreement factor of R = 45% was obtained. Using a full-matrix least-squares program two cycles of refinement reduced the R factor to 35%. A model constructed at this point indicated that the cesium atom, on the basis of a hexagonal close-packed arrangement, was displaced by 0.2 Å. and too close to two chlorine atoms. It was shifted to a more nearly close-packed position. Further isotropic refinement of the structure with an R factor of 11%. At this point seven of the strongest, low-order reflections, suspected of being strongly affected by ex-

tinction, were removed. Anisotropic refinement lowered the R factor for the 1062 nonzero reflections to a final agreement factor, R = 5.05%. The final parameters and standard deviations are given in Table II. In Table I the values of 10Fo and 10Fc are listed on an absolute scale.

Discussion

The CsCuCl3s structure consists of distorted hexagonally close-packed layers of cesium ions and chlorine atoms, with the copper atoms, located on twofold axes, occupying all the octahedral holes which have six chlorine nearest neighbors. Two chlorines are at a distance of 2.35 Å., two at 2.28 Å., and two at 2.78 Å. from the coppers (Figure 1). The distortion from perfect close packing is slight with Cl(2') 0.149 Å. above the plane and Cl(2) the same distance below (Figure 2). The copper atoms are 0.42 Å. from the 61c axis and are oriented to form spiraling chains by sharing faces of the

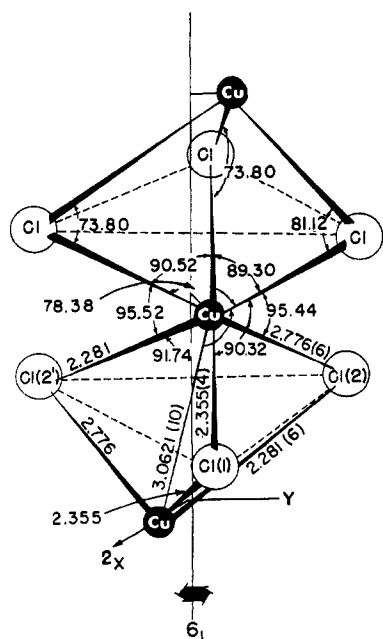
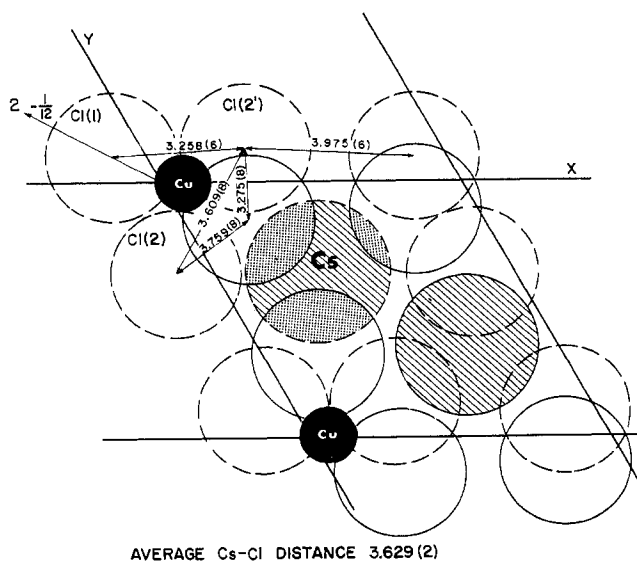
TABLE II

Final Atom Positional Parameters with Standard Deviations

Atom	x/a	y/b	z/c	σ_x	σ_y	σ_z
Cs^a	0.35458 (0.345) ^b	0.70916	0.25000	0.0001	0.0002	...
Cu	0.0616 (0.07)	0.0000	0.0000	0.0003
$\text{Cl}(1)^a$	0.8877 (0.90)	0.7754	0.2500	0.0003	0.0006	...
$\text{Cl}(2)$	0.3540 (0.35)	0.2095 (0.22)	0.2418 (0.25)	0.0005	0.0005	0.0002

Final Thermal Parameters with Deviations

Atom	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Cs	0.01311 (14)	0.01113 (7)	0.00222 (2)	0.00570 (3)	-0.00094 (3)	0
Cu	0.01026 (22)	0.01123 (11)	0.00101 (2)	0.00562 (6)	0.00020 (2)	0.00040 (4)
$\text{Cl}(1)$	0.0125 (4)	0.0094 (2)	0.0012 (1)	0.0047 (1)	0.0001 (1)	0
$\text{Cl}(2)$	0.0097 (2)	0.0128 (3)	0.0014 (1)	0.0051 (2)	-0.0006 (1)	-0.0002 (1)

^a Special positions: $y = 2x, z = 1/4$. ^b Wells' parameters in parentheses.Figure 1.—The molecular structure of CuCl_3^- in CsCuCl_3 .Figure 2.—The hexagonal close packing in CsCuCl_3 .

chlorine octahedra about them. One of the bridging chlorines, $\text{Cl}(1)$, is bound symmetrically to two coppers—a feature that has been found in all the garnet-colored

TABLE III

DISTANCES AND ANGLES IN COPPER CHLORIDE COMPOUNDS

Formula	Cu-Cu, Å.	Av. angle Cl-Cu-Cl, deg.	Av. angle Cu-Cl-Cu, deg.
KCuCl_3	3.443	84.1	95.9
$\text{LiCuCl}_3 \cdot 2\text{H}_2\text{O}$	3.47	85.	95.
NH_4CuCl_3	3.42	84.9	
$\text{Cu}_2\text{Cl}_4(\text{CH}_3\text{CN})_2$	3.39	86.	94.
$\text{Cu}_3\text{Cl}_6(\text{CH}_3\text{CN})_2$	3.346	88.	94.2
$\text{Cu}_5\text{Cl}_{10}(\text{C}_3\text{H}_7\text{OH})_2$	3.310	87.	92.8
$\text{CuCl}_2 \cdot 2\text{H}_2\text{O}$	3.73		92.9
CuCl_2	3.30	87.	
CsCuCl_3	3.0621	90.2	76.2

copper compounds.^{4,5} The three-dimensional character of the compound can be considered to be infinite parallel chains held together by cesium ions. The more important bond distances and angles along with their standard deviations can be seen in Figures 1 and 2.

A very interesting feature of the structure is the very short copper-copper distance. As can be seen in Table III this distance is much shorter than that found in any other halogen-bridged copper compound. Two possible explanations for this short distance can be proposed: (1) The copper atoms form copper-copper bonds along the chain axis. Thus each copper is bonded strongly to four chlorines and weakly to two others and to two neighboring copper atoms. (2) In this compound the copper octahedra are connected together by sharing faces while in the other compounds listed in Table III the octahedra are connected together by sharing edges. The sharing of faces would force the copper-copper distance to be shorter than in the other compounds previously investigated. In this description no direct copper-copper bonding is suggested.

Evidence seems to support the latter explanation. As can be seen from Table III, the angles in CsCuCl_3 are distorted from the angles observed in the other copper compounds in such a way as to suggest the sharing of faces. Further evidence for the latter explanation is found in overlap integral calculations. In order to consider the likelihood of the formation of a metal-metal bond, D_{3d} symmetry was assumed at a copper atom. Under this symmetry the representation for σ bonding by copper breaks down into the irreducible representations: $2A_{1g} + E_g + 2A_{2u} + E_u$. Considering only the chlorine atoms the σ bonds transform as

$A_{1g} + E_g + A_{2u} + E_u$. This leaves A_{1g} and A_{2u} for metal-metal bonding in this simple treatment. The A_{1g} transforms as d_{z^2} which is directed right at the coppers along the chain while A_{2u} is a nonbonding orbital. Thus metal-metal bonds could possibly be formed by σ overlap of $3d_{z^2}$ orbitals along the chain axis. However, the value of the overlap integral for σ $3d_{z^2}$ overlap on copper atoms 3.062 Å. apart is less than 10^{-5} . (This value and the values to follow were determined using the improved orbital exponents calculated by Clementi.⁶) It must be realized of course that these calculations only give an approximate value of the overlap because of the uncertainty in the orbital exponent for the isolated atom and more important the unknown change in the exponent in considering a complexed metal. Nevertheless the value of the overlap in the above case is so extremely small that the metal-metal bond order is essentially zero. This result is confirmed by the magnitude of the magnetic moment for $CsCuCl_3$, which is normal for one unpaired electron down to 80°K.⁷

Similar calculations for other metal-halogen systems lend credence to this type of approach. Thus for $CsNiCl_3$, which crystallizes in a close-packed structure isomorphous with $CsCuCl_3$ with a nickel-nickel distance of 2.96 Å,⁸ the overlap integral for $3d_{z^2}$ orbitals between nickel atoms is less than 10^{-4} , and $CsNiCl_3$ also has a normal moment for two unpaired electrons.⁹ Overlap calculations have also been performed for $TiCl_3$ and $ZrCl_3$, the structures of which have recently been deduced from powder diffraction studies.¹⁰ The compounds have D_{3d} symmetry at the metal and form chains similar to $CsCuCl_3$ by sharing faces of chlorine octahedra. The metal-metal distance in both of these isomorphous compounds is 3.07 Å.¹⁰ The calculation of the overlap integrals for σ d_{z^2} overlap gave a value of approximately 0.01 for the titanium-titanium bonds and 0.10, almost the maximum possible overlap for $4d_{z^2}$ orbitals, for the zirconium-zirconium bonds. Subnormal magnetic moments have been observed for these substances. $TiCl_3$ has a moment of 1.3 B.M.¹¹ (compared to the normal moment for one unpaired electron of 1.75 B.M.) indicating some interaction between the metals, and $ZrCl_3$ is almost diamagnetic with a moment of only 0.4 B.M.¹¹ In the latter case

the fourteen σ -bonding electrons (1 from Zr^{3+} , 12 from six chlorines and $1/2 \times 2$ from two neighboring zirconium atoms) enter the A_{1g} , E_g , E_u , A_{1u} , and A_{1g}' bonding orbitals, the last being the metal-metal d_{z^2} - d_{z^2} bonding orbital.

This treatment thus consistently explains the magnetic properties of the above materials and also gives evidence that no direct metal-metal bonding occurs in $CsCuCl_3$. In fact from these overlap calculations it was found that for the metals at the right of the first transition series the single bond length for d_{z^2} - d_{z^2} σ bonding is so extremely short that such metal-metal bonding cannot possibly occur unless there are gross changes in the orbital exponent of the combined metal.

The orbital exponents of Clementi were used in these calculations rather than those determined using Slater's rules. The difference between these values is quite significant (for copper $\mu = 2.7$ by Slater's rules and $\mu = 4.4$ as calculated by Clementi), with Clementi's values indicating much more compact d orbitals for the atoms on the right of the transition metal series.

The choice of orbital exponents also has a strong bearing on any discussion of the bonding in copper acetate monohydrate where many approaches have been used to attempt to explain the short copper-copper distance.¹²⁻¹⁸ All until recently have invoked some sort of direct 3d-3d bonding, the δ bond being very popular; however, none has been completely successful in explaining all the observed data. Our calculations of overlap rule out any direct 3d-3d bonding. Since our calculations were made, Kokoszka^{19,20} and Hansen and Ballhausen²¹ have given consistent interpretations of the available data assuming no direct copper-copper interaction. These results are consistent with our work and further indicate that overlap integrals obtained using the improved orbital exponents present a reliable method of evaluating the existence of direct metal-metal interactions.

(11) J. Lewis, *et al.*, *J. Chem. Soc.*, 2036 (1962).

(12) B. N. Figgis and R. L. Martin, *ibid.*, 3837 (1956).

(13) B. N. Figgis and R. L. Martin, *ibid.*, 2545 (1957).

(14) I. G. Ross, *Trans. Faraday Soc.*, **55**, 1057 (1959).

(15) I. G. Ross and J. Yates, *ibid.*, **55**, 1064 (1959).

(16) M. L. Tonnet, S. Yamada, and I. G. Ross, *ibid.*, **60**, 840 (1964).

(17) L. S. Forester and C. J. Ballhausen, *Acta Chem. Scand.*, **16**, 1385 (1962).

(18) E. A. Boudreaux, *Inorg. Chem.*, **3**, 506 (1964).

(19) G. F. Kokoszka, H. C. Allen, and G. Gordon, *J. Chem. Phys.*, **42**, 3693 (1965).

(20) C. W. Reimann, G. F. Kokoszka, and G. Gordon, *Inorg. Chem.*, **4**, 1082 (1965).

(21) A. E. Hansen and C. J. Ballhausen, *Trans. Faraday Soc.*, **61**, 631 (1965).

(6) E. Clementi and D. L. Raimondi, *J. Chem. Phys.*, **38**, 2686 (1963).

(7) B. N. Figgis and C. M. Harris, *J. Chem. Soc.*, 855 (1959).

(8) G. N. Tishehenko, *Tr. Inst. Krist. Akad. Nauk SSSR*, **11**, 93 (1955).

(9) R. W. Asmussen and H. Soling, *Z. anorg. allgem. Chem.*, **283**, 3 (1956).

(10) L. F. Dahl, *et al.*, *Inorg. Chem.*, **3**, 1236 (1964).