# **Inorganic Chemistry**

# Structures of  $d^4$  MH<sub>3</sub>X: a Computational Study of the Influence of the Metal and the Ligands

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# **S** Supporting Information

[AB](#page-9-0)STRACT: [Density func](#page-9-0)tional theory (DFT, PBE0, and range separated DFT,  $RSH + MP2$ ) and coupled-cluster with single and double and perturbative triple excitations (CCSD- (T)) calculations have been used to probe the structural preference of  $d^4$  MH<sub>3</sub>X<sup>q</sup> (M = Ru, Os, Rh<sup>+</sup>, Ir<sup>+</sup>, and Re<sup>-</sup>; X = H, F, CH<sub>3</sub>, CF<sub>3</sub>, SiH<sub>3</sub>, and SiF<sub>3</sub>) and of  $MX_4$  (M = Ru; X = H,



F, CH<sub>3</sub>, CF<sub>3</sub>, SiH<sub>3</sub>, and SiF<sub>3</sub>). Landis et al. have shown that complexes in which the metal is sd<sup>3</sup> hybridized have tetrahedral and non-tetrahedral structures with shapes of an umbrella or a 4-legged piano stool. In this article, the influence of the metal and ligands on the energies of the three isomeric structures of  $d^4$  MH<sub>3</sub>X and MX<sub>4</sub> is established and rationalized. Fluoride and alkyl ligands stabilize the tetrahedral relative to non-tetrahedral structures while hydride and silyl ligands stabilize the non-tetrahedral structures. For given ligands and charge, 4d metal favors more the non-tetrahedral structures than 5d metals. A positive charge increases the preference for the non-tetrahedral structures while a negative charge increases the preference for the tetrahedral structure. The factors that determine these energy patterns are discussed by means of a molecular orbital analysis, based on Extended Hü ckel (EHT) calculations, and by means of Natural Bond Orbital (NBO) analyses of charges and resonance structures (NRT analysis). These analyses show the presence of through-space interactions in the non-tetrahedral structures that can be sufficiently stabilizing, for specific metals and ligands, to stabilize the non-tetrahedral structures relative to the tetrahedral isomer.

# **ENTRODUCTION**

Unsaturated transition metal complexes are a model of reactive intermediates, and their structures and electronic properties can provide useful information on how they could react. These species are often very reactive and hard to characterize by experimental means. Therefore, structure and bonding analyses and computational studies have been useful to establish the structures of these complexes. The structure of unsaturated species does not always follow the rules that apply for saturated systems. For instance, it has been found that the structure of unsaturated species cannot be always predicted by the valence shell electron pair repulsion (VSEPR) theory<sup>1</sup> or by a simple ionic model. WMe<sub>6</sub> and WH<sub>6</sub> are notable cases for the failure of the VSEPR theory since they have distorted [tri](#page-9-0)gonal prismatic structures with  $C_{3v}$  symmetry<sup>2</sup> while VSEPR would have predicted an octahedral coordination. The VSEPR also fails to predict the bent structure of m[o](#page-9-0)lecular dihalides of the heavy alkaline earth metals (Ca, Sr, and Ba) and the pyramidal structure of trihydrido complexes such as ScH<sub>3</sub>, LaH<sub>3</sub>, TiH<sub>3</sub><sup>+</sup> ,  $\text{ZrH}_{3}^{+}$ , and their alkyl analogues. These species and related  $\text{d}_{3}^{0}$ complexes have been described in a comprehensive review.<sup>3</sup> The main feature that emerges from the studies is that distortion away from the structure predicted by VSEPR [is](#page-9-0) prevalent for complexes with hydride or alkyl ligands, which have no  $\pi$ -donor capability, while the presence of  $\pi$ -donor ligand like halide lead more often to structures that follow the VSEPR rule. For instance, WCl<sub>6</sub> is octahedral<sup>4</sup> while WMe<sub>6</sub> and WH<sub>6</sub> are not. Analysis of the metal–ligand bonding within such

complexes shows that the electron delocalization between the  $\pi$ -donor ligand and the formally empty metal d orbital increases the electron density at the metal which, in turn, becomes less unsaturated. Therefore, the structures of polyhydride or polyalkyl complexes are considered as informative of the fundamental structural preferences associated with the metal coordination number and the number of d electrons. Unfortunately, few systems of this type have been characterized experimentally, a clear consequence of the poor stability of highly unsaturated species. For this reason, computational studies have been widely used.<sup>5</sup>

Landis et al. have shown that valence bond concepts can be used for describing the shape [of](#page-9-0) a variety of hydride and alkyl systems.<sup>6</sup> A primary feature of their studies is that the valence empty  $(n + 1)p$  orbitals of a *n*d metal do not contribute to the metal–l[ig](#page-9-0)and bonds, which are constructed from the  $(n + 1)s$ and available  $nd$  orbitals.<sup>7</sup> Consequently, the shape of the hydride complex, where the metal has  $\lambda$  nonbonding electron pairs, is determined by th[e](#page-9-0) optimal spatial distribution of the sd<sup>ω</sup> ( $\omega$  = 5 –  $\lambda$ ) hybrids set, that is, the spatial distribution that minimizes the interactions between the hybrid orbitals. In general, there is more than one optimal spatial distribution for a given  $sd^{\omega}$  set and quantum calculations are used to determine the relative energies of these isomeric structures.<sup>7a,8,9</sup> Similar results were obtained using the orbitally ranked symmetry

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analysis method (ORSAM) since for low coordinated complexes the metal orbital uses also the  $(n + 1)s$  and nd orbitals and the ranking of the structures is also done by quantum calculations.<sup>10</sup> In these studies, only hydride complexes were calculated but other ligands were not explicitly considered. Likewise t[he](#page-9-0) influence of the metal was not explored.

Well-defined experimental species are rather limited since the only known complexes of this type are the homoleptic neutral  $(M = Ru \text{ and } Os)$  and cationic  $(M = Ir^+)$  d<sup>4</sup>  $MR_4$   $(R = alkyl \text{ and }$ aryl) complexes synthesized by Wilkinson. These complexes,  $\text{Ru}(C_6H_{11})_4$ ,  $^{11}$  Ru(o-Tolyl)<sub>4</sub>,  $^{12'}$  Ru(Mesityl)<sub>4</sub>,  $^{13}$  Os( $C_6H_{11}$ )<sub>4</sub>,  $^{11}$  $Os(o-Tolyl)<sub>4</sub>$ <sup>14</sup> and Ir(Mesityl)<sub>4</sub><sup>+</sup>,<sup>13</sup> are stable and have a , tetrahedral [coo](#page-9-0)rdination at [the](#page-9-0) metal. How[eve](#page-9-0)r, the bondi[ng](#page-9-0) analysis of La[nd](#page-9-0)is and the ORSAM [an](#page-9-0)alysis of Hall predict that other isomers with non-tetrahedral structures should be possible, and calculations for  $RuH<sub>4</sub>$  and  $OsH<sub>4</sub>$  have shown that non-tetrahedral structures could even be preferred or at least competitive, but no rationale for these results is available. We were intrigued by these results, and we wanted to explore more widely the structural preferences of  $d^4$  MH<sub>3</sub>X and MX<sub>4</sub> for various metals, charges, and ligands. We present a study based on density functional theory (DFT) calculations, validated by ab initio calculations, for neutral  $MH_3X$  (M = Ru and Os;  $X = H$ , F, CF<sub>3</sub>, CH<sub>3</sub>, SiF<sub>3</sub>, and SiH<sub>3</sub>), cationic  $[MH<sub>3</sub>X]<sup>+</sup>$  (M = Rh and Ir), and anionic  $[Ref<sub>3</sub>X]<sup>-</sup>$  complexes. We also present results for homoleptic  $RuX_4$  (X = H, F, CH<sub>3</sub>,  $CF_3$ , SiH<sub>3</sub>, and SiF<sub>3</sub>) complexes.

# **COMPUTATIONAL METHODS**

Two methods (1 and 2) were used to obtain the geometries and energies at the DFT level. Additional methods were used to increase the accuracy of the relative energies of the isomers, methods 3 and 4. The validity of the DFT method was established with method 5.

Method 1. The geometry optimizations were performed at the DFT level of theory with Gaussian 09<sup>15</sup> using the B3PW91 functional.<sup>16</sup> The metal centers were described with the Stuttgart/ Dresden (SDD) quasi-relativistic effective co[re](#page-9-0) potentials (RECP) and the associ[ate](#page-9-0)d basis sets.<sup>17</sup> The 6-31G(d,p) basis set were used for the other atoms.<sup>18</sup>

Method 2. The geo[me](#page-9-0)tries obtained with the method 1 were used as initial inp[ut](#page-9-0) for optimization with the PBE0 functional<sup>19</sup> and triple- $\zeta$ quality basis sets on all atoms, def2-TZVPP,<sup>20</sup> with the corresponding RECP for the metal. $21$  With rare exceptions, the optimi[ze](#page-9-0)d structures obtained with method 2, and method 1 are [sim](#page-9-0)ilar. Bond lengths only differ by about 0.01 [Å](#page-9-0), and bond angles differ by less than 5°. The structures and energies obtained with method 2 are used for describing the results in this Article. All stationary points were fully characterized by analytical frequency calculations as either a minimum or a transition state. Intrinsic reaction coordinate (IRC) calculations followed by geometry optimizations were used to associate a transition state to the corresponding two minima. The energies reported in this article do not include the ZPE correction, but it has been verified that the ZPE correction does not modify significantly the relative energies of the isomers. The optimized geometries for all stationary points and the ZPE corrected energies are given in the Supporting Information.

Method 3. Single-point short-range (sr-) DFT/long-range (lr-) ab initio calculations on the DFT (PBE0/def2-TZVPP) geometries were carried out for all structures using the M[olpro](#page-9-0) [2010.1](#page-9-0) [program.](#page-9-0)<sup>22</sup> The short-range exchange-correlation terms were treated by the sr-PBE variant functional. $23^{\circ}$  The long-range portion of the exchange e[ner](#page-9-0)gy is explicitly treated (HF). This step defines a "range-separated hybrid" (RSH) scheme, w[hi](#page-9-0)ch is corrected in a second step for the long-range correlation effects by a second order perturbation theory, leading to MP2-like correction. This method is referred to as  $RSH+MP2<sup>24</sup>$  The range-separation parameter  $\mu$  has been set to 0.4  ${a_0}^{-1}$ . For these singlepoint calculations, all atoms were described with quadruple-ζ quality basis sets, def2-QZVPP<sup>25</sup> with the corresponding RECP for the  $metal.<sup>2</sup>$ 

Method 4. Single-poi[nt](#page-9-0) coupled-cluster with single and double and pertu[rba](#page-9-0)tive triple excitations (CCSD(T)) calculations on the DFT (PBE0/def2-TZVPP) optimized geometries were carried out for  $MH_3X$  (M = Ru and Os; X = H, F, CF<sub>3</sub>, CH<sub>3</sub>, SiH<sub>3</sub>, and SiF<sub>3</sub>) using the Orca 2.9 program.<sup>26</sup> In this case, all atoms were described with def2-QZVPP<sup>25</sup> basis sets with the corresponding RECP for the metal.<sup>21</sup>

Method [5.](#page-9-0) Single-[po](#page-9-0)int MR-CI and MR-ACPF-2 $a^{27}$  calculations on th[e D](#page-9-0)FT (PBE0/def2-TZVPP) optimized geometries were carried out for  $\text{RuH}_4$  using the Orca 2.9 program.<sup>26</sup> In this case, [al](#page-9-0)l atoms were described with a def2-QZVPP<sup>25</sup> basis sets with the corresponding RECP for the metal.<sup>21</sup> Resolution of the [id](#page-9-0)entity (RI) approximation was used for the integral transf[orm](#page-9-0)ation with def2-QZVPP/C auxiliary basis sets.<sup>28</sup> For th[es](#page-9-0)e MRCI and MR-ACPF-2a calculations, the reference space was built from full valence CASSCF orbitals (12 electrons, [10](#page-9-0) orbitals) using a selection threshold (Tpre) set to  $10^{-6}$ . . The selection of the excited configurations from the reference space was done with a selection threshold (Tsel) of 10<sup>−</sup><sup>9</sup> . The effect of the rejected configurations was estimated using second order perturbation theory.

Analysis of the Results. The Extended Hückel (EHT) calculations were carried out with the Yaehmop program.<sup>29</sup> The Ru 4d, 5s, and 5p orbitals are  $-14.9$ ,  $-10.4$ , and  $-6.87$  eV, respectively.<sup>30</sup> The energy of the "normal" hydrogen 1s orbital is −13.6 e[V.](#page-9-0) A better electron-donating ligand than "normal" hydrogen is characterized [by](#page-9-0) Hii higher than −13.6 eV and a poorer electron-donating ligand is characterized by H<sub>ii</sub> lower than  $-13.6$  eV. Ligands of variable electrondonating power were represented by varying the  $H<sub>ii</sub>$  value of the four hydrogen atoms between −15.6 and −12.6 eV. The Natural Bond Orbital (NBO) charges and the Natural Resonance Theory (NRT) analysis were obtained using the NBO methodology with the NBO 5.0 program.<sup>31</sup>

#### ■ RE[SU](#page-10-0)LTS

In  $d^4$  ML<sub>4</sub> species, where two d orbitals are used to host the four metal electrons, the  $(n + 1)$ s and three *n*d metal orbitals are available to form the four metal−ligand bonds. This leads to  $sd<sup>3</sup>$  hybridization at the metal, and the shape of the complex is determined by the optimal spatial distribution of these four  $sd<sup>3</sup>$ hybrids. The neutral polyhydride complex,  $d^4$  MH<sub>4</sub>, is an ideal system in that the H−M−H angles are very close to those that minimize the overlap between the metal hybrid orbitals at the metal which are  $71^{\circ}$  and  $109^{\circ}$ .<sup>7a,8,32</sup> It results that the structures can be constructed by placing the metal at a center of a cube and the four ligands at four [ver](#page-9-0)[tic](#page-10-0)es such that no ligands are trans (Figure 1). This gives three possible structures, which



Figure 1. Dispositions of ligands for  $sd<sup>3</sup>$  hybridization in idealized  $d<sup>4</sup>$  $ML_4$  (M in blue at the center of the cube and L in black).<sup>8</sup>

have the shape of a tetrahedron, td, an umbrella, umb, and a 4 legged-piano stool, pst, as shown in the top part of Figure 2. In td, all coordination sites are equivalent and all bond angles are 109 $^{\circ}$ . In umb, there are three basal hydrides,  $H_b$ , and one [api](#page-2-0)cal hydride, H<sub>a</sub>, with H<sub>a</sub>−M−H<sub>b</sub> and H<sub>b</sub>−M−H<sub>b</sub> angles of 71<sup>°</sup> and 109°, respectively. In pst, the four hydrides are equivalent, and the cis H−M−H angle is equal to 71°. Changing the metal and

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Figure 2. Possible structures for  $d^4$  MH<sub>4</sub> (top) and  $d^4$  MH<sub>3</sub>X (bottom) with labeling.

the ligands lead to different angles, which influence the opening of the umbrella and the height of the piano stool.

When one of the three hydrogen atoms is replaced by an X ligand, tetrahedral, umbrella, and 4-legged piano stool isomeric structures are still found (Figure 2 bottom). Only the umbrella complex has nonequivalent sites. The presence of X gives rise to two isomers, namely,  $umb_a$  where  $X$  is at the apical site and  $umb<sub>b</sub>$  where X is at the basal site.

Validation of the Use of Single Reference Methods for  $d^4$  MH<sub>3</sub>X. The geometries of MH<sub>3</sub>X were optimized at a DFT level and the energetics were obtained with methods 2, 3, and 4, DFT (PBE0 and RSH+MP2), and CCSD(T), respectively. However, the use of single reference (SR) methods for the description of unsaturated species could be questionable. Even with truncation of excitations at the double and triple level, single reference methods, such as  $CCSD(T)$ , recover a large portion of the dynamic correlation energy, but may fail when a significant amount of nondynamical correlation energy is present. In contrast, multireference methods, such as MRCI or MR-ACPF, recover both dynamical and nondynamical correlation energies.

To check the validity of the SR calculations for  $d^4$  MH<sub>3</sub>X, single points MRCI and MR-ACPF-2a calculations were carried out for the three isomers of  $RuH_4$  (td, umb, and pst). For these MRCI and MR-ACPF-2a calculations, the reference space was built from full valence CASSCF orbitals (12 electrons, 10 orbitals) and the effect of the rejected configurations was estimated using second order perturbation theory (Table 1).

All calculations show that the umb and pst isomers are lower in energy than the td isomer. Furthermore, all methods give a similar energetic pattern. Umb and pst are lower than td by 9.1

# Table 1. Energetics<sup> $a$ </sup> of RuH<sub>4</sub> Using the td Isomer as Reference



to 12.9 and 9.4 to 14.7 kcal mol<sup>-1</sup>, using the PBE0, RSH+MP2, and CCSD(T) methods, respectively. There is a marginal tendency for pst to have a lower energy than umb. It is also worth noting that the range separated hybrid RSH+MP2 calculations and the single reference coupled cluster calculations give values that are close. This overall good agreement between the single reference and multireference methods suggests that these unsaturated metal species do not contain a significant amount of nondynamical correlation energy. This indicates that a single reference method can be used for the computational study of  $d^4$  MH<sub>3</sub>X.

In addition, several diagnostics have been suggested to evaluate the SR/MR character for molecular systems. The  $T_1$ and  $D_1$  diagnostics<sup>33</sup> (the Frobenius norm and matrix 2-norm of coupled cluster amplitudes for single excitations respectively) are arguably the [mos](#page-10-0)t widely used diagnostics for SR coupled cluster calculations. Alternatively, the weight  $(C_0^2)$  of the leading configuration state function in a MRCI can be used to determine the multireference character. The  $T_1$  and  $D_1$ diagnostics were tested primarily on small organic molecules, and it was suggested that any  $T_1$  diagnostic larger than 0.02 and any  $D_1$  diagnostic larger than 0.05 invalidate the use of SR methods. However, a recent study of Jiang et al. pointed out that the criteria established for organic molecules ( $T_1$  < 0.02 and  $D_1 < 0.05$ ) are no longer valid for molecular systems with transition metals element.<sup>34</sup> On the basis of a statistical analysis of a set of 225 species, these authors proposed  $T_1$  < 0.05 and  $D_1$  < 0.15 as amplitude [crit](#page-10-0)eria for the validation of SR-based methods in the case of d-block energetics. The  $T_1$  and  $D_1$ diagnostics of the single reference CCSD(T) calculations are reported in Table 2 for the three isomers of  $RuH_4$  (td, umb,

Table 2.  $T_1$  and  $D_1$  Diagnostics of the Single Reference CCSD(T) Calculations and Weight of the Leading Configuration  $(C_0^2)$  of MRCI and MRACPF2 Calculations for the Three Isomers of RuH4

td	umb	pst
0.048	0.045	0.043
0.091	0.100	0.089
0.82	0.84	0.84
0.81	0.83	0.83

and  ${\sf pst)}$ ; the weight of the leading configuration  $({C_0}^2)$  from the multireference MRCI or MR-ACPF-2a calculation is also indicated. For the three isomers of RuH<sub>4</sub>, both  $T_1$  and  $D_1$ diagnostics from  $SR CCSD(T)$  calculations are less than the values suggested by Jiang et al. as acceptable limits. In addition, the weight of the leading configuration is higher than 0.80 for the MRCI and MR-ACPF-2a calculations. This confirms that these unsaturated molecular systems do not contain any substantial MR character. Single references methods are thus validated for the calculations of RuH<sub>4</sub>. This validation applies to all  $d^4$  MH<sub>3</sub>X and MX<sub>4</sub> studied in this article (T<sub>1</sub> and D<sub>1</sub> diagnostics given in the Supporting Information for  $RuH_3X$  and  $OsH_3X$ ).

**RuH<sub>3</sub>X.** (a). RuH<sub>4</sub>. [The DFT calculations sh](#page-9-0)ow that RuH<sub>4</sub> has td, umb, and pst minima. The umb and pst isomers have similar energies and both species are more stable than the tetrahedral structure by about 9 kcal mol<sup>−</sup><sup>1</sup> . These results differ from those obtained by Landis. Molecular mechanic (MM) calculations including hybridization and resonance in a valence bond (VB) approach (Hypervalent-VALBOND (HV-VB) MM <span id="page-3-0"></span>method) $9$  give the same three minima but show that **td** is the most stable isomer. DFT calculations using a different ECP and a differe[nt](#page-9-0) functional for the metal<sup>7a,9</sup> from the one used in this work give td and pst as the only minima, with pst being 8.6 kcal mol<sup>−</sup><sup>1</sup> lower in energy relative to [td](#page-9-0). The Hartree−Fock and MP2 calculations by Hall et al. show several isomers with relative energives that depend on the computational method used. $10$ 

The essential structural features of the three isomers of  $RuH<sub>4</sub>$ are s[ho](#page-9-0)wn in Figure 3. The Ru−H bond is 1.55 Å in td. In umb,



Figure 3. DFT optimized structures of minima and transition states of interconversion between minima for  $RuH<sub>4</sub>$ . The distances are in Å and the angles in degrees. The DFT energies are given in kcal mol<sup>−</sup><sup>1</sup> relative to the td structure, using method 2.

the apical Ru−Ha bond of 1.51 Å is slightly shorter than the basal Ru−H<sub>b</sub> bond of 1.54 Å. The H<sub>b</sub>−Ru−H<sub>a</sub> and the H<sub>b</sub>− Ru−Hb angles of 64° and 102°, respectively, are close to the corresponding ideal values of 71° and 109° predicted for a sd<sup>3</sup> hybridized metal. The acute  $H_b-Ru-H_a$  angle leads to a nonbonded  $H_a \cdots H_b$  distance of 1.62 Å, which is less than the sum of the van der Waals (vdW) radii of the two hydrogens (1.20 Å).<sup>35</sup> In pst, the Ru−H bond of 1.53 Å and the angles of 67° and 103° for cis H−Ru−H and trans H−Ru−H angles, respectiv[ely](#page-10-0), are similar to the ideal values of 71° and 109°. In pst, the shorter nonbonded H $\cdots$ H distance is 1.69 Å, which is less than the sum of the vdW radii of the two hydrogens although slightly longer than in umb. Therefore, both umb and pst have short nonbonded H···H distances that are not present in the td structure wherein the H···H nonbonded distance is 2.54 Å.

The transition states between these minima were searched using method 2 to evaluate the depth of the associated wells on the potential energy surface and the possibility of exchange between the three minima (Figure 3). A transition state of 8.7 kcal mol<sup>−</sup><sup>1</sup> above td was located between td and umb and another transition state between umb and pst with an energy of 2.9 kcal mol<sup>−</sup><sup>1</sup> above umb. No transition state was identified between td and pst. The tetrahedral structure is thus separated from the other minima by a significant barrier, but there is a low energy barrier between umb and pst. The structural features of the two transition states are shown in Figure 3. The transition state between td and umb isomers inverts the configuration at the metal as it inverts the umbrella. At the transition state, the ruthenium atom is in the equatorial plane formed by the three  $H<sub>b</sub>$  hydrogens and the Ru−H<sub>a</sub> bond is perpendicular to the equatorial plane; this inversion at the metal occurs without significant change in the Ru−H bond distances. The transformation of umb to pst is a concerted swinging motion of two  $H<sub>b</sub>$  relative to the plane defined by Ru−H<sub>a</sub> and the third Ru−  $H<sub>b</sub>$  bonds. At the transition state, Ru–H<sub>a</sub> and two Ru–H bonds are coplanar. This transformation also occurs without significant change of the Ru−H bond distances.

(b). RuH<sub>3</sub>X  $(X = H, F, CH_3, CF_3, SiH_3, and SiF_3)$ . The structures of RuH<sub>3</sub>X (X = H, F, CH<sub>3</sub>, CF<sub>3</sub>, SiH<sub>3</sub>, and SiF<sub>3</sub>) were determined at the DFT level with method 2. The results are shown in Figure 4 as a histogram in which the energies of the



Figure 4. DFT energies, in kcal mol<sup>-1</sup>, for RuH<sub>3</sub>X (X = F, CF<sub>3</sub>, CH<sub>3</sub>, H,  $SiF_3$ , and  $SiH_3$ ) complexes, relative to the tetrahedral isomer, td. The energies are in blue for the umbrella with apical  $X$ , umb<sub>a</sub> in red for the umbrella with basal  $X$ , umb<sub>b</sub>, and in green for the 4-legged pianostool structure, pst. Light green is used to indicate structures with hydride bridging Ru−Si.

non-tetrahedral structures are plotted relative to the td structure, which is a minimum for all X groups. Figure 4 shows that X influences the relative energies of the various minima and also the number of minima. For instance, RuH3F has only two minima,  $\mathbf{td}$  and  $\mathbf{umb}_{\mathbf{b}}$ , while the four minima shown in Figure 2 are obtained for all other X.

Single point calculations with method 3, RSH+MP2, and method 4,  $CCSD(T)$  $CCSD(T)$  $CCSD(T)$ , are shown in Figure 5. These calculations show energetic patterns similar to that obtained with method 2. Methods 3 and 4 stabilize more the non-[tet](#page-4-0)rahedral structures relative to the tetrahedral references than method 2. In the case where method 2 gives a non-tetrahedral structure less stable than td, the difference in energy between td and nontetrahedral structures is decreased. In the case where method 2 gives a non-tetrahedral structure more stable than td, the difference in energy is increased.

In RuH<sub>3</sub>F, the isomers **td** and  $umb<sub>b</sub>$  are isoenergetic. The Ru−H bond lengths and the angles between the Ru−H bonds are similar to that of  $RuH<sub>4</sub>$  in the corresponding isomers. The Ru–F is longer at the basal site of **umb**<sub>b</sub> than in td (1.85 and 1.80 Å, respectively). The F–Ru–H<sub>a</sub> angle is nearly 92<sup>°</sup> in

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Figure 5. RSH+MP2 (top) and CCSD(T) (bottom) energies, in kcal mol<sup>-1</sup>, for RuH<sub>3</sub>X (X = F, CF<sub>3</sub>, CH<sub>3</sub>, H, SiF<sub>3</sub>, and SiH<sub>3</sub>) complexes, relative to the tetrahedral isomer, td. The energies are in blue for the umbrella with apical X,  $umb_a$ , in red for the umbrella with basal X,  $umb_{b}$ , and in green for the 4-legged piano-stool structure, pst.

 $umb<sub>b</sub>$ , which is notably different from the H<sub>b</sub>−Ru−H<sub>a</sub> angle of 63°. For RuH<sub>3</sub>CH<sub>3</sub>, the td, umb<sub>b</sub>, and pst have similar energies and  $umb_a$  is significantly higher in energy. Although the four minima have shapes that are overall similar to the corresponding minima of RuH<sub>4</sub>, some structural aspects need to be mentioned. The Ru−C bond length of 1.95 Å in td is slightly shorter than in the other isomers (1.96 to 2.1 Å). In  $umb<sub>b</sub>$ , the methyl group moves away from the axis of the umbrella as indicated by the CH3−Ru−Ha angle of 80°, which is larger than the  $H_b$ -Ru- $H_a$  angle of 66 $^{\circ}$ . Likewise, in pst, the CH3−Ru−H angle of 81° is larger than that of 63° between two cis Ru−H bonds. RuH<sub>3</sub>CF<sub>3</sub> also has four minima, the most stable ones being **umb**<sub>b</sub> and **pst** by about 8 kcal mol<sup>-1</sup> at the DFT level. The structural features of the  $RuH_3CF_3$  and  $RuH_3CH_3$  complexes are similar.

 $\rm RuH_3SiH_3$  and  $\rm RuH_3SiF_3$  also have the four minima like were found for  $RuH_3CH_3$  and  $RuH_3CF_3$ . However, the energy pattern for the alkyl and the silyl ligands is significantly different. For the silyl ligand, all non-tetrahedral structures are at significantly lower energy than the tetrahedral structure by about 18 kcal mol<sup>−</sup><sup>1</sup> at the DFT level. In addition, the influence on the geometries is also different for the silyl and the methyl ligands. For instance, in the  $umb_a$  isomer of  $RuH_3SiH_3$ , the Ru−Si distance of 2.26 Å is slightly shorter than in td, 2.30 Å, and the Si−Ru−H angles are similar to the corresponding H− Ru−H angles. In umb<sub>b</sub> of RuH<sub>3</sub>SiH<sub>3</sub> the Si–Ru–H<sub>a</sub> angle is equal to the  $H_b-Ru-H_a$  angle. Thus, the Si–Ru–H angles are

systematically smaller than the corresponding C−Ru−H angles in the non-tetrahedral geometry. For the pst structure, one of the hydrogen of the SiH<sub>3</sub> group bridges the Ru–Si bond. The ability for the silyl group and notably the  $SiH<sub>3</sub>$  group to enter in the bridging situation is well documented and is not the focus of this study.<sup>36</sup>

As a partial summary, these calculations show that all  $RuH_3X$ have several [iso](#page-10-0)meric structures, among which the tetrahedral structure is not always the most stable isomer. For  $X = F$  and  $CH<sub>3</sub>$ , the td structures are isoenergetic with the other isomers. In the case of  $X = H$ ,  $CF_3$ ,  $SiH_3$ , and  $SiF_3$ , there is a clear preference for non-tetrahedral structures especially with the silyl ligands. The structures of the td,  $umb_a$ ,  $umb_b$ , pst isomers are rather similar for all X and, thus, the structure of  $RuH<sub>4</sub>$  is a good representative of the entire set. There is a noticeable tendency for the X groups that favor significantly the nontetrahedral structures ( $SiH_3$  and  $SiF_3$ ) to have a structure close to that of  $RuH_4$ . In contrast, for the X groups that do not favor the non-tetrahedral structures (F and  $CH<sub>3</sub>$ ), the X-Ru–H angle is systematically larger than the corresponding H−Ru−H angle. This structural pattern will reveal through-space interaction in the non-tetrahedral structures, as it will be discussed later.

For all systems, the various non-tetrahedral structures are separated by low energy barriers and therefore non-tetrahedral structures are fluxional. However, the non-tetrahedral structures are separated from the td structure by a barrier that is significantly higher than their difference in energy with the exception of  $RuH_3SiF_3$  where the energy barrier is very close to the difference in energy. See Supporting Information for further details.

OsH<sub>3</sub>X (X = H, F, CH<sub>3</sub>, CF<sub>3</sub>, SiH<sub>3</sub>, and SiF<sub>3</sub>). Three minima were located for  $OsH_4$  (Figure 6). In contrast to the preference for non-tetrahedral structures found for  $RuH<sub>4</sub>$ , the three isomers of  $OsH<sub>4</sub>$  are isoener[ge](#page-5-0)tic. Similar results have been obtained by Landis et al. with umb and pst being 3.5 and 5.2 kcal mol<sup>-1</sup> above the td.<sup>37</sup> The structures of td, umb, and pst are similar to those found for RuH4. An energy barrier of 10.9 kcal mol<sup>−1</sup> relative to **td** [sep](#page-10-0)arates  $\mathbf{t}\overline{\mathbf{d}}$  and  $\mathbf{umb},^{37}$  and an energy barrier of 5.8 kcal mol<sup>-1</sup> relative to **umb** separates **umb** and **pst**.

The DFT energies for OsH<sub>3</sub>X (X = F, CF<sub>3</sub>, CH<sub>3</sub>, H, SiF<sub>3</sub>, and  $SiH<sub>3</sub>$ ), relative to the td isomers, are shown as a histogram in Figure 7 using the convention already used in Figure 4 for  $RuH<sub>3</sub>X$ . There is a significant difference between the isoelectro[nic](#page-5-0) ruthenium and osmium complexes. For  $X = F$  $X = F$  $X = F$ ,  $CH<sub>3</sub>$  and  $CF<sub>3</sub>$ , the tetrahedral structure is more stable than any non-tetrahedral isomer. For  $X = H$ , the tetrahedral and nontetrahedral isomers have similar energies. For  $X = SiH<sub>3</sub>$  and  $SiF<sub>3</sub>$ , the preference for the non-tetrahedral structure remains, but the difference in energy between td and non-tetrahedral structures decreases by about 10 kcal mol<sup>-1</sup>, compared to the ruthenium complexes. Thus, everything being equal, replacing ruthenium by osmium increases the stability of the tetrahedral isomer relative to the non-tetrahedral.

[ReH<sub>3</sub>X]<sup>-</sup>, [RhH<sub>3</sub>X]<sup>+</sup>, and [IrH<sub>3</sub>X]<sup>+</sup>. To further understanding of the effect of the metal on the structural preference of  $d^4$  MH<sub>4</sub>, the isoelectronic anionic [ReH<sub>3</sub>X]<sup>−</sup> and cationic  $[RhH_3X]^+$ ,  $[IrH_3X]^+$  species were optimized with method 2, and single points RSH+MP2 energies were obtained with method 3.

 $\left[\text{Re}\text{H}_{4}\right]^{-}$  and  $\left[\text{Ir}\text{H}_{4}\right]^{+}$  have three minima while  $\left[\text{Rh}\text{H}_{4}\right]^{+}$  has two minima, Figure 8. The anionic  $[\rm{Re}H_4]^-$  has a preference for the tetrahedral structure while the cationic  $[\mathrm{RhH}_4]^+$  and

<span id="page-5-0"></span>

Figure 6. DFT optimized structures of minima and transition states of interconversion between the minima for  $OsH_4$ . The distances are in Å and the angles in degrees. The energies are given in kcal mol<sup>−</sup><sup>1</sup> relative to the td structure, using method 2.



Figure 7. DFT energies, in kcal mol<sup>-1</sup>, for OsH<sub>3</sub>X (X = F, CF<sub>3</sub>, CH<sub>3</sub>, H,  $SiF_3$ , and  $SiH_3$ ) complexes, relative to the tetrahedral structure, td, with similar conventions as in Figure 4.

 $[IrH<sub>4</sub>]$ <sup>+</sup> have a preference for non[-t](#page-3-0)etrahedral geometries. The preference for the non-tetrahedral relative to the tetrahedral structures is stronger for the cationic than for the neutral species; for  $[RhH_4]^+$  the preference for non-tetrahedral structure increases by over 25 kcal mol<sup>-1</sup> relative to RuH<sub>4</sub>. For  $[\mathrm{IrH}_4]^+$ , the preference for the non-tetrahedral structure is more marked than for  $\rm{OsH_4}$  but less so than for  $\rm{RhH_4^{+}}$ . The influence of the total charge and the nature of metal on the preferential structures of  $d^4 [MH_4]^q$   $(q = -1, 0,$  and 1) is thus important. Consequently, going from the left to the right of the periodic table increases the preference for non-tetrahedral geometries but going down a column decreases this preference.

The geometries of the minima for  $\mathrm{[Re H_4]}^-$ , and  $\mathrm{[Ir H_4]}^+$  are similar to those previously found for RuH<sub>4</sub>. The H-M-H



Figure 8. DFT optimized geometries (distances are given in Å and angles in degrees) of  $\left[\text{Re}\text{H}_{4}\right]^{-}$ ,  $\left[\text{Rh}\text{H}_{4}\right]^{+}$ , and  $\left[\text{Ir}\text{H}_{4}\right]^{+}$ . The energies are given in kcal mol<sup>-1</sup> relative to td. N/A signifies that no minimum of this type could be located.

angle shows a pattern that parallels the preference for the nontetrahedral structures. In the **umb** isomer of  $[{\rm Re}H_4]^-$ , the  $H_a$ –  $M-H_b$  and the  $H_b-M-H_b$  angles are larger than in the neutral systems. A similar pattern is seen in the pst isomer. The opposite trend is found in the cationic systems. For  $[\mathrm{IrH}_4]^+$ , the minima are those of the type shown in Figure 2. In the case of  $[RhH<sub>4</sub>]<sup>+</sup>$ , the pst isomer is a highly fluxional structure where the decrease of a cis H−Rh−H angle from 6[3](#page-2-0) to 38° occurs with essentially no energy cost. No umb type structure was located as minimum. Overall, the structures of isoelectronic  $d<sup>4</sup>$  $MH<sub>4</sub>$  are thus sensitive to the total charge; a negative charge opens the umbrella and flattens the 4-legged piano stool and a positive charge closes the umbrella and raises the height of the 4-legged piano stool. Going down a column of the periodic table also opens the umbrella and flattens the 4-legged piano stool.

The histograms for the relative energies of the isomers for  $[{\rm Re}H_3X]^-$ ,  $[{\rm Rh}H_3X]^+$ , and  $[{\rm Ir}H_3X]^+$  are shown in the Supporting Information. The pattern of energies can be understood from the pattern obtained for the ruthenium and osmium complexes. For  $[ReH_3X]^-(X = F, CF_3, CH_3, H, SiH_3,$ and  $SiF_3$ ) the preference for the tetrahedral structure decreases from  $CH_3$  to  $SiF_3$ . However, the combined effect of the negative charge and the 5d metal results in a general preference for the tetrahedral coordination for any X; in the case of  $[ReH<sub>3</sub>F]$ <sup>-</sup> the tetrahedral structure is the only minimum. For X =  $CH_3$ ,  $CF_3$ , and  $SiH_3$ , the tetrahedral structures are more stable than the non-tetrahedral structures, and, in the case of [ReH<sub>3</sub>SiF<sub>3</sub>]<sup>−</sup>, **pst** and **td** are isoenergetic. For [IrH<sub>3</sub>X]<sup>+</sup> the preference for non-tetrahedral structures is found for  $X = SiF<sub>3</sub>$ and SiH<sub>3</sub>. For  $X = CH_3$  and  $CF_3$ , non-tetrahedral structures are preferred except for the umb structure with apical X group that is significantly less stable than td. In the case of  $[IrH_3F]^+$ , the  $umb<sub>b</sub>$  structure with equatorial F is marginally more stable than  $td$ , but the  $umb<sub>a</sub>$  structure with apical F is significantly less stable than **td**. In the case of  $\overline{[RhH_3X]}^+$ , non-tetrahedral structures are more stable than td for  $X = H$ ,  $SiF_3$  and  $SiH_3$ . For  $X = CF_3$  and  $CH_3$ , structures of different natures than found for

all other metal complexes are found to be more stable than tetrahedral structure. No comparison can be done with the other cases, and this was not explored further.

 $\text{RuX}_4$ . The diamagnetic homoleptic  $\text{Ru}(C_6H_{11})_4$ ,  $\text{Ru}(O_6H_{11})_5$  $\text{Tolyl})_{4}^{12} \text{ Ru(Mesityl)}_{4}^{13} \text{ Os}(C_6H_{11}^{1})_{4}^{11} \text{ Os}(O-Tolyl)_{4}^{14}$  and Ir(Mesityl)<sub>4</sub><sup>+</sup>,<sup>13</sup> which are the only 12-electron  $d^4$  MR<sub>4</sub> comple[xes](#page-9-0) to have bee[n](#page-9-0) synthesized [and](#page-9-0) fully charac[ter](#page-9-0)ized, are remarkabl[y s](#page-9-0)table. They are colored as expected from the presence of low-lying empty metal d orbitals. In all cases, X-ray diffraction studies show that these complexes have a tetrahedral structure in the solid state. The NMR study of the tetramesityl complex of ruthenium shows a fluxional behavior attributed to the rotation of the ligand about the Ru−C bond.<sup>13</sup> To better compare with the experimental structures,  $RuX_4$  complexes were calculated for  $X = F$ ,  $CH_3$ ,  $CF_3$ ,  $SiH_3$ , a[nd](#page-9-0)  $SiF_3$ . The optimization of  $RuF_4$  and  $Ru(CF_3)_4$  yields only tetrahedral isomers. Two minima are found for  $Ru(CH_3)_4$  with the td structure being 32.2 kcal mol<sup>−</sup><sup>1</sup> below the pst structure. A preference for non-tetrahedral isomers is obtained with the silyl substituted complexes. For  $Ru(SiH_3)_4$  and  $Ru(SiF_3)_4$ , the umb isomer is 12.6 kcal mol<sup>-1</sup> and 17.1 kcal mol<sup>-1</sup> more stable than the td isomer, respectively. In addition, structures with a pst shape in which the  $SiH_3$  and  $SiF_3$  ligands have hydrides or fluoride bridging the Ru−Si bond are also located as minima lower than the td isomer  $(-29.8 \text{ and } -23.5 \text{ kcal mol}^{-1}$  for  $Ru(SiH<sub>3</sub>)<sub>4</sub>$  and  $Ru(SiF<sub>3</sub>)<sub>4</sub>$ , respectively), but they will not be considered. These calculations on homoleptic complexes show that non-tetrahedral geometries are possible even with ligands more bulky than H. Furthermore, the energy trend obtained for  $RuH<sub>3</sub>X$  (Figure 4) is also present for  $RuX<sub>4</sub>$ . Fluorine is the ligand that most disfavors the non-tetrahedral geometry in the  $RuH<sub>3</sub>X$  series; [co](#page-3-0)nsequently  $RuF<sub>4</sub>$  has only a tetrahedral structure. The  $CH<sub>3</sub>$  ligand disfavors slightly less the nontetrahedral structures; consequently a non-tetrahedral structure is found as a high-lying secondary minimum. Finally,  $SiH<sub>3</sub>$  and  $SiF<sub>3</sub>$  ligands, which most favor non-tetrahedral structures in the  $RuH<sub>3</sub>X$  series, lead to homoleptic complexes with preference for non-tetrahedral structures. The only exception to this trend is the case of  $Ru(CF_3)_4$  whose tetrahedral structure cannot be predicted from the influence of a single  $CF<sub>3</sub>$  ligand on the structural preference of  $RuH_3(CF_3)$ . Clearly, the structures of the alkyl and aryl complexes that have been synthesized do not display the diversity of possible structures for  $d^4$  RuX<sub>4</sub>.

#### ■ DISCUSSION

The bonding analysis of Landis and Weinhold $^8$  as well as the ORSAM analysis of Hall<sup>10</sup> show that  $d^4$  MH<sub>3</sub>X and MX<sub>4</sub> have [se](#page-9-0)veral structures. In this analysis, the metal uses the  $(n + 1)s$ and the nd orbitals, whic[h](#page-9-0) are not occupied by the electrons of the metal lone pairs to establish the covalent metal−ligand bond. For an  $sd^3$  hybridization that applies to  $d^4$  tetracoordinated complexes, the tetrahedral structure is one of the possible structures. The non-tetrahedral structures have umbrella and 4 legged piano stool shapes. However, the valence-bond analysis of Landis and Weinhold<sup>8</sup> and the ORSAM model of Hall<sup>10</sup> do not provide any information on the relative energies of the various isomeric forms. [T](#page-9-0)he DFT and ab initio calcul[atio](#page-9-0)ns show that the metal and the ligands influence significantly the relative energies of the several isomers. In  $MH<sub>3</sub>X$ , the nontetrahedral structures have nonbonded distances between proximate atoms that are relatively short. This is the case in particular for  $MH_4$  wherein the nonbonding  $H \cdots H$  distances in the non-tetrahedral geometries are shorter than the sum of the

vdW distances. The increased steric hindrance resulting from the replacement of the hydrogens in  $RuH_4$  by relatively bulky alkyls and aryl ligands could have been the reason for all experimentally known systems to be tetrahedral. However, the calculations show that  $CH<sub>3</sub>$  and SiH<sub>3</sub> ligands lead to opposite structural preference, the former ligand increasing the preference for the tetrahedral structure and the latter for the non-tetrahedral structures. These trends apply to  $RuH<sub>3</sub>X$  and  $RuX<sub>4</sub>$  showing that the steric effects of groups larger than hydrides do not determine the structural preferences. Calculations show that the metal also plays an important role on the structural preference. Going from the left to the right of the periodic table, that is, from anionic to cationic isoelectronic complexes increases the preference for the non-tetrahedral structures while going down a column of the periodic table increases the preference for the tetrahedral structures. Factors that also play a role are, among others, the relative electronegativity of the metal and the electron-donating/ electron-withdrawing ability of the ligands. To understand better these factors, a molecular orbital analysis, based on EHT calculations, has been carried out for  $RuH<sub>4</sub>$  and a NBO charge analysis complemented by a NRT analysis based on the calculated DFT densities have been carried out for  $RuH<sub>4</sub>$  and  $OsH_4$ .

(a). Molecular Orbital Analysis. The EHT total energies of the td, umb, and pst structures for  $RuH<sub>4</sub>$  give a preference for a non-tetrahedral structure. Increasing the electronwithdrawing ability of the ligand by lowering the hydrogen  $H<sub>ii</sub>$  Coulombic integral gives a preference for a tetrahedral structure. This modeling of ligands reproduces the results of the DFT calculations. For instance (i) F and  $CH<sub>3</sub>$  are more electron-withdrawing than H and  $SiH<sub>3</sub>$  and thus favor the tetrahedral structure, (ii) Os 5d orbitals are higher in energy than Ru 4d orbitals and thus ligands appear to be more electron-withdrawing relative to Os than Ru. Likewise, in a cationic complex where the metal is more electron-attracting, all ligands appear more electron-donating. The reverse is true for anionic complexes.

A Walsh diagram is used to gain further insight into why EHT calculations reproduce qualitatively the structural preference obtained with higher-level calculations. Only the occupied molecular orbitals of  $RuH<sub>4</sub>$  are necessary in this analysis (Figure 9). The  $C_{3v}$  or  $C_{4v}$  group notations are used for labeling the molecular orbitals for the three structural forms. The td structur[e h](#page-7-0)as two nonbonding d orbitals of e symmetry to host four electrons. The four Ru−H bonds are represented by two molecular orbitals of  $a_1$  symmetry and two orbitals of  $e$ symmetry. Only the molecular orbitals of  $a_1$  symmetry have a contribution on the hydrogen located on the z axis. In umb, all molecular orbitals of e symmetry have the same energy as in td because they have no contribution on the hydrogen on the  $z$ axis. The only molecular orbitals, which have different energies for  $td$  and  $umb$ , are the molecular orbitals of  $a_1$  symmetry. They are constructed in a similar manner in td and umb. The lower one,  $1a<sub>1</sub>$ , is the in-phase combination of the metal 5s orbital and the 1s orbital of the four hydrogens. The higher orbital of  $a_1$ symmetry,  $2a_1$ , is mostly made of the in-phase combination of a dz <sup>2</sup> with the four hydrogens. The contribution of the 5p orbitals has been found to be negligible in all molecular orbitals even if permitted by symmetry; it will not be mentioned further. The key point of the analysis is that, in  $2a_1$ , the coefficient of the hydrogen on the  $z$  axis has the sign opposite to that of the three other hydrogens, which are in-phase with the torus part of  $d_{z^2}$ .

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Figure 9. Walsh diagram for the td to umb and pst transformation. The green orbitals are stabilized from td to either umb or pst and the red orbitals are destabilized. The green and red wriggling lines indicate positive and negative overlaps, respectively.

The energy of  $1a<sub>1</sub>$  is thus lower in umb than in td because the in-phase relationship between all hydrogens stabilizes the orbital when the hydrogen atoms are closer (umb vs td). In contrast, the energy of  $2a_1$  is higher in umb than in td because the out-of-phase relationship between the apical and basal hydrogens raises the energy of this molecular orbital when the distances between the nonbonded hydrogens are shorter and the overlap increases. Two antagonist effects are thus at work and umb is more stable than td if the energy lowering of  $1a<sub>1</sub>$  is more important than the energy rise of  $2a_1$ . The  $1a_1$  orbital is more located on the hydrogens than on the metal because the energy level H<sub>ii</sub> of the hydrogen 1s orbital is always lower than the energy value  $H_{ii}$  of the 5s of the metal. In contrast, the  $2a_1$ orbital is more located on the metal when the energy of the hydrogen 1s orbital is higher than that of the metal d orbital, and it is more located on the hydrogen atoms when it is lower. Therefore, the energy rise of the  $2a_1$  orbital upon going from td to umb is small in the first case and large in the second case. Consequently, electron-donating ligands favor umb and electron-withdrawing ligands favor td.

A similar reasoning applies to the comparison between td and pst. In pst, the four nonbonding metal electrons are hosted in  $2\mathsf{a}_1\ (\mathrm{d}_{z^2})$  and  $\mathsf{b}_2\ (\mathrm{d}_{\mathrm{xy}})$  orbitals (Figure 9). The four Ru–H bonding orbitals are described by an  $1a_1$  orbital made mostly from the metal s orbital, two orbitals of e symmetry made mostly from the metal  $d_{xz}$  and  $d_{yz}$  orbitals, and a  $\mathbf{b}_1$  orbital made from  $d_{x^2-y^2}$ . The contribution of the 5p orbitals has been found to be negligible even if permitted by symmetry, as already noted above. The  $1a_1$  orbital favors the pst structure because of the in-phase relationship between all hydrogens. The  $b_1$  orbital disfavors the pst structure because of the out-of-phase relationship between the hydrogens. The two degenerate orbitals of e symmetry have essentially no influence on the structural preference because of the long distance between the two trans hydrogens. The influence of the energy level  $H_{ii}$  of the hydrogens on the variation of energy of the  $1a_1$  and  $b_1$ orbitals between td and pst is similar to that obtained for the

case of td vs umb. Electron-donating ligands favor pst and electron-withdrawing ligands favor td.

(b). NBO Charge Analysis. The NBO charges for  $RuH<sub>4</sub>$ and  $OsH<sub>4</sub>$  are shown in Figure 10. For these two species in





which the metal is formally at a high oxidation state,  $M^V$ , the NBO charge at the metal is negative indicating, as currently recognized, that the calculated charge has no relation with formal oxidation state. The negative charge is small for the td complexes but increases to a maximum value of −0.3 for the non-tetrahedral species. Accordingly, the charge on the hydrogens is very small and positive in RuH4. In the nontetrahedral structures, the hydrogens carry a positive charge; the largest positive charge is carried by the apical hydrogen of umb. The results are qualitatively similar for  $OsH<sub>4</sub>$ , but the electronic density on the metal is smaller for Os than for Ru and the charge is even marginally positively charge in td. The charge distribution in the Ru and Os complexes does not give a rationale for the stronger energetic preference for nontetrahedral structures in the case of Ru. However, the charge distribution indicates that electron-donor ligands are beneficial for the non-tetrahedral structures.

(c). NRT Analysis. A natural resonance (NRT) analysis<sup>38</sup> of the DFT density was carried out for the three isomers, td, umb, an[d](#page-10-0) pst of  $RuH_4$  and  $OsH_4$ . This analysis is aimed at determining the valence structures that describe the total electron density of these complexes in the various isomeric forms.<sup>8</sup> This analysis has met some criticism<sup>38a</sup> but has been found to be qualitatively useful for understanding electronic struct[ur](#page-9-0)es.<sup>8</sup> The results are shown in Figu[re 1](#page-10-0)1. In the td isomer, the dominant structure has four covalent M−H bonds and two l[on](#page-9-0)e pairs at the metal. This valence stru[ctu](#page-8-0)re accounts for 97% of the density for  $RuH_4$  and 98% for  $OsH_4$ . For the umb isomer, the same structure accounts for only 80% of the total density for  $\text{RuH}_4$  and 90% for  $\text{OsH}_4$ .<sup>39</sup> The remaining part of the electron density is represented by structures with two covalent M−H bonds, three lone pairs [a](#page-10-0)t the metal, and a proton and a hydride characterizing an ionic interaction between the metal fragment and the hydrogen atoms. The structure with a proton at the apical site has higher weight (12% for RuH<sub>4</sub> and 6% for OsH<sub>4</sub>) than that with a hydride at the apical site (7% for  $RuH_4$  and 3% for  $OsH_4$ ). The structure with the proton and the hydride at the basal sites has the smallest weight 1%. For the pst isomer, the main structure is that with four M−H bonds and two lone pairs at the metal (79% of the total electron density for  $RuH_4$  and 89% for  $OsH_4$ ). The remaining part of the electron density is represented by a valence structure with two covalent M−H bonds, three lone pairs at the metal, and a proton and a hydride on cis hydrogen

<span id="page-8-0"></span>

Figure 11. Main valence bond structures for  $RuH_4$  and  $OsH_4$ . The value indicated below each structure represents the percentage of this structure that contributes to the total density. The values indicated are summed over all equivalent sites.

sites. It accounts for 21% of the density for  $RuH<sub>4</sub>$  and 11% of the density for  $OsH_4$ .

The total electron density is well described by these valence structures for all isomers of  $RuH_4$  and  $OsH_4$  since the sum of the contributions is 100%. In the valence structures displaying the ionic contributions, the proton and the hydride interact with the metal and also together since the two species with opposite charges are on atoms that are at relatively short distances. This gives rise to an electrostatic interaction that contributes to the stability of the umb and pst isomers. The structures with these ionic components have a higher weight for  $RuH_4$  than for  $OsH_4$ , which could rationalize the preference for the umb and pst isomers in the case of RuH4. The same analysis was carried out for  $RuH_3X$  (X = F,  $CH_3$ , and  $SiH_3$ ) (see Supporting Information). In all cases, the tetrahedral structure is described by essentially a single structure with three M−H bonds, one M−X bond, and two lone pairs on the metal. The [electron](#page-9-0) [densities](#page-9-0) [of](#page-9-0) umb and pst isomers are still dominated by the same valence structures but structures with two covalent bonds to the ruthenium, three lone pairs on the metal, and a positively charged group (H or X) and a negatively charged group (H or X) have relatively high percentages. It has not been possible to establish a quantitative correlation between the percentage of the valence structures with zwitterionic contributions and the relative energies of the non-tetrahedral and tetrahedral structures. However, it is likely that they play a key role in stabilizing the non-tetrahedral structures.

(d). Through-Space Interactions in the non-tetrahedral Structures from MO and NRT Analyses. The Walsh diagram and the NRT analyses provide a consistent interpretation of the stability of the non-tetrahedral structures of  $d^4$  RuH<sub>4</sub> although they use different languages and properties of the electronic structures. In the molecular orbital analysis, it was shown that the out-of-phase contribution between nearby "hydrogen" orbitals, which disfavors the umb and pst isomers relative to the td isomers, depends on the electron-donating power of the "hydrogen" ligands to the metal center. Ligands

that are good electron donor contribute less to this orbital than ligands that are less good electron donor. When this destabilizing interaction is low, the through-space stabilizing interaction arising from the all in-phase low-lying occupied orbital dominates. Structures with the shortest distances between nonbonded atoms, that is, umb and pst, are favored relative to td. These stabilizing interactions do not create any covalent bonds between these hydrogens but favor interatomic distances shorter than the sum of the vdW radii. The NRT analysis gives a complementary interpretation of the results. It confirms the absence of covalent bond between hydrogens that are at distance shorter than the sum of the vdW radii (axial  $H_a$ ) and basal  $H<sub>b</sub>$  in **umb** and two cis basal H in **pst**). However, the NRT establishes the presence of an electrostatic interaction between positively and negatively charged hydrogens.

The analysis established for tetrahydride complexes can readily be generalized to complexes with any type of ligand. It provides an interpretation of the effect of the ligand using only the electron-withdrawing/electron-donating property of the ligands, but other factors can also contribute. For instance, fluorine is an electron-withdrawing ligand and thus the nontetrahedral structures are disfavored. This result is suggested independently from the presence of the lone pairs on the fluorine lone pairs that would disfavor the presence of any group at short distance from fluorine as in the non-tetrahedral structures. Likewise, an alkyl group is reasonably electronwithdrawing and thus disfavors a non-tetrahedral structure. Steric effects between the alkyl group and the other ligands also disfavor the non-tetrahedral structures. At the other extreme, a good electron-donating group like a silyl group favors nontetrahedral structures. In the case of a complex like  $RuH_3(SiH_3)$ the stabilizing interactions that have been described above are supplemented by the well-known through space interaction between a hydride and a nearby silyl group, originating from the ability of silicon to become hypervalent. This interaction, which has been known as SISHA<sup>40</sup> or IHI,<sup>41</sup> is magnified when the silyl group is substituted by halide. This accounts for the preference for non-tetrahedral s[tru](#page-10-0)ctures [for](#page-10-0)  $RuH_3(SiH_3)$  and  $RuH_3(SiF_3)$ . The fact that  $Ru(SiF_3)_4$  and  $Ru(SiF_3)_4$  also prefer non-tetrahedral structures show that interactions different from SISHA and IHI are at work.

#### ■ **CONCLUSIONS**

DFT (PBE0 and range separated DFT, RSH+MP2) and CCSD(T) calculations have been used to determine the preferred structures of  $MH_3X^q$  (M = Ru, Os, Rh<sup>+</sup>, Ir<sup>+</sup>, and Re<sup>-</sup>;  $X = H$ , F, CH<sub>3</sub>, CF<sub>3</sub>, SiH<sub>3</sub>, and SiF<sub>3</sub>) and MX<sub>4</sub> (M = Ru; X = H, F,  $CH_3$ ,  $CF_3$ ,  $SiH_3$ , and  $SiF_3$ ). Tetrahedral and non-tetrahedral structures are found as possible minima. The non-tetrahedral structures (umbrella and 4-legged piano stool shaped complexes) are those predicted for an  $sd<sup>3</sup>$  hybridization following the analysis of valence bond of Landis et al. and the ORSAM analysis of Hall et al. The calculations show that non-tetrahedral structures can be energetically preferred to the expected tetrahedral geometry for certain metals and ligands. The non-tetrahedral isomers are more energetically preferred for 4d metal than for for 5d metals. Cationic complexes also tend to favor the non-tetrahedral isomers while anionic complexes do not. The relative energies of the non-tetrahedral and the tetrahedral structures are significantly influenced by the nature of the ligands. Electron-withdrawing ligands like halide and alkyl favor the conventional tetrahedral structure while ligands that are more electron donating like hydride and silyl <span id="page-9-0"></span>favor the non-tetrahedral structures. The factors that determine these energy patterns have been discussed by means of a molecular orbital analysis based on EHT calculations and NBO charge analysis supplemented by an NRT analysis. All analyses show that electron-donating ligands favor the non-tetrahedral structures. Furthermore, the EHT and the NRT analyses suggest the presence of possible weak attractive interactions between atoms that are closer in umb and pst than in td. These weak attractive interactions are magnified when the ligands are good electron donor to the metal center.

# ■ ASSOCIATED CONTENT

#### **S** Supporting Information

List of coordinates of all calculated structures with energies E and ZPE corrected energies  $(E + ZPE)$ . Histograms of energies (methods 2 and 3) for  $[ReH_3X]^-$ ,  $[RhH_3X]^+$ , and  $[IrH_3X]^+$ . Potential energy surfaces for the transformations between the various isomers of RuH<sub>3</sub>X. NBO charges and NRT analysis for  $RuH<sub>3</sub>X$  (X = F, CH<sub>3</sub>, and SiH<sub>3</sub>). This material is available free of charge via the Internet at http://pubs.acs.org.

# ■ AUTHOR INFORMATI[ON](http://pubs.acs.org)

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### Notes

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