# **Inorganic Chemistry**

# A Binuclear Mn<sup>III</sup> Complex of a Scorpiand-Like Ligand Displaying a Single Unsupported Mn<sup>III</sup>−O−Mn<sup>III</sup> Bridge

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**S** Supporting Information

[AB](#page-7-0)STRACT: [The crystal s](#page-7-0)tructure of a binuclear Mn<sup>III</sup> complex of a scorpiand-like ligand (L) displays an unsupported single oxo bridging ligand with a Mn<sup>III</sup>-O-Mn<sup>III</sup> angle of 174.7°. Magnetic susceptibility measurements indicate strong antiferromagnetic coupling between the two metal centers. DFT calculations have been carried out to understand the magnetic behavior and to analyze the nature of the observed Jahn−Teller distortion. Paramagnetic <sup>1</sup>H NMR has been applied to rationalize the formation and magnetic features of the complexes formed in solution.



### **■ INTRODUCTION**

Manganese is an element of biological, chemical, and pharmaceutical relevance. Living systems use this essential element for important functions dealing with redox reactions such as water splitting in photosynthesis<sup>1</sup> or reactive oxygen species  $removal<sup>2</sup>$  among others. Although many manganese enzymes such as superoxide dismutase or [m](#page-7-0)anganese oxygenase contain a singl[e](#page-7-0) manganese atom $3$  in their active center, a growing number of manganese enzymes with binuclear metal sites<sup>4</sup> such as catalases,<sup>5,6</sup> ribo[n](#page-7-0)ucleotide reductase,<sup>7</sup> or arginase<sup>8</sup> are being discovered. In all these metalloenzymes, man[g](#page-7-0)anese is cycling betw[een](#page-7-0) different oxidation states.

Mn<sup>III</sup> unsaturated polynuclear complexes have a strong tendency to form bridges between two neighboring metal atoms. The most common bridging ligand types are acetate, carbonate, water, hydroxo, and oxo anions.<sup>6,9–13</sup> However, the participation of a single unsupported oxo bridge ligand is rather unusual. Indeed, the fo[r](#page-7-0)mation of linear  $\mu$ [-o](#page-7-0)xo complexes requires of a particular steric hindrance for preventing the bending of the angle or the formation of multibridged complexes.6,9−11,13−<sup>18</sup> Since MnIII is a paramagnetic atom, such binuclear complexes may exhibit interesting magnetic properties [like sup](#page-7-0)e[r-a](#page-7-0)ntiferromagnetic behavior.<sup>1</sup>

Scorpiand-like ligands have a structure consisting of a macrocyclic core and a flexible pending ar[m](#page-7-0) containing additional donor atoms. Compound L, depicted in Scheme 1,



which belongs to this class of ligands, is the result of the condensation of 2,6-bis(bromomethyl)pyridine with tris(aminoethyl)amine.<sup>20</sup> In this paper, we report that the interaction of  $L$  with a  $Mn^{II}$  salt followed by atmospheric oxidation leads to th[e f](#page-7-0)ormation of a binuclear complex in which two  $[Mn^{III}L]^{3+}$  subunits are interconnected by a single, almost linear oxo bridge, which has been studied by a variety of experimental techniques both in solution and in the solid state.

## **EXPERIMENTAL SECTION**

Synthesis and Chemicals Used. The synthetic procedure for the obtention of L has been already described.<sup>20</sup> The  $Mn^{\text{II}}$  source was a

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Figure 1. Unit cell showing the arrangement around the 61 screw axis. Br<sup>−</sup> atoms, shown in yellow; water molecules, shown in red; hydrogen atoms, not displayed.

MnSO4·H2O standardized stock solution. All other chemicals were purchased from several commercial sources and used as received.

The formation of the  $[Mn_2OL_2]^{4+}$  complex is achieved by dissolving in water equivalent amounts of L-4HBr-2H<sub>2</sub>O and  $Mn<sup>H</sup>$  and increasing the pH up to ca. 7 with diluted NaOH. The formation and further oxidation of the complex can be followed by a change in the color of the solution from colorless to purple.

X-Ray Structures. A solution containing one equivalent of both L-4HBr-2H<sub>2</sub>O and MnSO<sub>4</sub>·H<sub>2</sub>O at neutral pH was left to slowly evaporate in an open vessel. After several days, hexagonal dark magenta crystals suitable for X-ray diffraction were obtained.

The single crystal X-ray diffraction studies were carried out using Agilent Gemini equipment with Mo K $\alpha$  monochromatic radiation ( $\lambda$  = 0.71073 Å) and a CCD detector at room temperature. The data reduction and scaling were made with the CrysAlisPro suite software. $21$ The structure was solved and refined with SHELXS97 and SHELXL97, respectively.<sup>22</sup> The crystallographic data are collected [in](#page-7-0) Table S1 (Supporting Information).

Magnetic Measure[me](#page-7-0)nts. Magnetic susceptibility measurements on polycr[ystalline samples of com](#page-7-0)pound  $[Mn_2OL_2]\cdot 4Br\cdot 4H_2O$  were carried out in the 1.9−300.0 K temperature range with a Quantum Design SQUID magnetometer under an applied magnetic field of 5000 G. Corrections for the diamagnetic contribution of the constituent atoms and for the magnetization of the sample holder were applied to the raw data.

Theoretical Computations. Calculations were performed with the Gaussian 09 package using the B3LYP functional and the quadratic convergence approach.<sup>23</sup> Triple-ξ and double-ξ all electron basis sets, as proposed by Ahlrichs et al., were used for the metal ions and the rest of atoms, respec[tiv](#page-7-0)ely.<sup>24</sup> The broken symmetry approach was employed to describe the unrestricted solutions of the antiferromagnetic spin states, $25$  which [we](#page-8-0)re obtained from the guess functions generated with the fragment tool implemented in the Gaussian code. The full geome[try](#page-8-0) was used for  $[Mn_2OL_2]^{4+}$ . We also used a polarizable continuum model with the parameters corresponding to the acetonitrile solvent.<sup>26</sup> However, since similar results for the magnetic coupling constant J were found, these data have not been included in the discussio[n.](#page-8-0) A study on molecular models was done to verify our results, to know which magnetic orbitals are responsible for the strong observed magnetic coupling and which is the most stable arrangement of the Jahn−Teller axis. These molecular models are simple, being constituted only by two Mn<sup>III</sup> ions, one oxo group preserving the linearity of the Mn−O−Mn axis, and ammonia groups to complete the ideal octahedral metal coordination sphere (N−Mn−  $N = 90^{\circ}$ ). The broken-symmetry methodology allows the determination of all the exchange coupling constants present in a polynuclear transition metal complex.<sup>25</sup> Since our system contains only two  $Mn$ <sup>III</sup>  $(S = 2)$  paramagnetic centers, the magnetic coupling constant in  $[Mn_2OL_2]^{4+}$  can be inferr[ed](#page-8-0) from the energies of a nonet and a singlet spin function, displaying a parallel and antiparallel alignment of the local spin moments. The electronic spectra in acetonitrile and water solutions have been partially calculated using the time dependent formalism applied to the density functional theory (DFT).<sup>27</sup> In both solvents, the obtained results are similar. The inclusion of the solvent in the electronic calculations was done by using a PCM [mo](#page-8-0)del with parameters corresponding to both acetonitrile and water.<sup>28</sup>

Paramagnetic NMR Measurements. The paramagnetic NMR measurements were acquired on a Bruker Avance 400 spectrometer operating at 399.91 MHz. One-dimensional spectra were recorded in  $D_2O$  solvent with presaturation of the  $H_2O$  signal during part of the relaxation delay to eliminate the  $H_2O$  signal. A relaxation delay time of 200 ms, a 75 kHz spectral width, and an acquisition time of 200 ms were used. 1D spectra were processed using exponential linebroadening weighting functions as apodization with a value of 60 Hz. Chemical shifts were referenced to residual solvent protons of  $D_2O$  resonating at 4.76 ppm (298 K) relative to TMS. The concentration of the samples for paramagnetic <sup>1</sup>H NMR was 5 mmol·dm<sup>−</sup><sup>3</sup> . The longitudinal relaxation times of the hyperfine shifted resonances were determined using the inversion recovery pulse sequence  $(d_1-180^{\circ}-\tau-90^{\circ}-acq$ , where  $d_1$  is the relaxation delay and acq is the acquisition time); 14 values of  $\tau$  were selected between 0.8 and 500 ms.<sup>29</sup>  $(d_1 + acq)$  values were at least 5 times the longest expected  $T_1$  ranging from 200 to 300 ms, and the number of scans was 9000. The  $T_1$  $T_1$  values were calculated from the inversion−recovery equation. Transversal relaxation times were obtained measuring the line broadening of the isotropically shifted signals at half-height through the equation  $T_2^{-1} = \pi \Delta \nu_{1/2}$ .

Electrochemical Measurements. The electrochemistry of the monomer systems was studied in water using  $0.15$  M NaClO<sub>4</sub> as a supporting electrolyte. The electrochemical behavior was recorded in the 4.0−9.0 pH range under an argon atmosphere. A glassy carbon electrode was used as a working electrode, a platinum wire as an auxiliary, and Ag/AgCl as a reference electrode in a three-electrode cell coupled to a CH 420I potentiostat.

The sample was prepared by direct mixing of the ligand as hydrobromide salt and one equivalent of  $MnSO<sub>4</sub>·H<sub>2</sub>O$  in 0.15 M of NaClO<sub>4</sub>. The pH was adjusted with either  $HClO<sub>4</sub>$  or NaOH aqueous solutions.

Experiments on solutions of  $[Mn_2OL_2]$  +Br·4H<sub>2</sub>O in 0.10 M  $Bu_4NPF_6/DMSO$  and 0.10 M  $Bu_4NPF_6/MeCN$  were performed using a Pt disk pseudoreference electrode, the potentials being measured relative to the ferrocenium/ferrocene couple using a 0.50 mM ferrocene (Fluka) internal standard.

Spectroscopic Measurements. UV-vis spectra were recorded in an Agilent 8453 spectrometer.

#### ■ RESULTS AND DISCUSSION

Description of Crystal Structure of [Mn<sub>2</sub>OL<sub>2</sub>]·4Br·4H<sub>2</sub>O. The crystal system is hexagonal, and the unit cell contains six [LMn-O−MnL]<sup>4+</sup> units helicoidally arranged according to a  $6<sub>1</sub>$  screw axis (Figure 1). The c axis is unusually long, 55 Å, showing long-range organization of anions and water molecules in the crystal structure. Four bromide anions and four lattice water molecules complete the structure forming the hydrogen bond network.

Each  $\text{Mn}^{\text{III}}$  in a  $[\text{Mn}_2\text{OL}_2]^{4+}$  moiety is coordinated in a distorted octahedral geometry (Figure 2). The pyridine nitrogen, the primary amino group of the tail, the tertiary nitrogen, and the bridging oxo group define the equatorial plane (Table 1), while the distorted axial po[sit](#page-2-0)ions are occupied

<span id="page-2-0"></span>

Figure 2. ORTEP plot for the structure  $[Mn_2OL_2]^{4+}$ . Bromide and water moieties and hydrogen atoms are not displayed. Thermal ellipsoids are plotted at the 50% probability level.

Table 1. Selected Bond Lengths and Angles

bond length/Å			angles/deg			
$Mn1 - O1$	1.753(2)	$Mn1 - O1 - Mn1$	174.7(3)			
$Mn1-N2$	2.308(4)	$O1 - Mn1 - N5$	93.1(2)			
$Mn1-N1$	2.088(4)	$O1-Mn1-N1$	92.8(2)			
$Mn1-N3$	2.101(4)	$N5-Mn1-N3$	83.3(2)			
Mn1-N4	2.268(4)	N1-Mn1-N3	91.0(2)			
$Mn1-N5$	2.034(4)	$O1-Mn1-N4$	97.3(1)			
		N5–Mn1–N4	107.1(2)			
		$N1-Mn1-N4$	77.0(2)			
		$N3 - Mn1 - N4$	81.0(2)			
		$O1 - Mn1 - N2$	102.6(1)			
		$N5-Mn1-N2$	97.8(2)			
		$N1-Mn1-N2$	76.1(2)			
		$N3 - Mn1 - N2$	80.90(2)			

by the nitrogen atoms of the two secondary amino groups close to the pyridine aromatic spacer (Mn1−N2 = 2.269(3)Å, Mn1−  $N4 = 2.304(6)$ Å; Table 1). As such an orientation of the elongation axis was already observed for related Cu<sup>II</sup> complexes,<sup>20</sup> the direction of the distortion should be somehow imposed by the ligand topology. The oxo bridge is almost linear  $(Mn1-O1-Mn1 = 174.7(3)°)$ . This value agrees with the angles reported for the other few cases in which single Mn<sup>III</sup>− O-Mn<sup>III</sup> bridges have been described.<sup>10,14,16,17</sub><sup>e</sup>,18 Steric</sup> congestion was put forward to explain a more acute angle of 146.15° reported in one of the examples.1[7c](#page-7-0) [Evidence](#page-7-0) for the oxo bridge is provided by the infrared spectrum (Figure S1, Supporting Information) in which two c[hara](#page-7-0)cteristic peaks of metal–oxo–metal bonds can be found at around 870 cm<sup>-1,30</sup> .

The  $Mn$ <sup>III</sup>L [subunits a](#page-7-0)re not placed one above the other but rotated 38.18° around the Mn1−O1−Mn1 axis. Moreov[er,](#page-8-0) both pyridine rings are not strictly parallel to each other, with an angle of 8.08° between their mean planes.

Electronic Spectrum. The UV−vis spectrum displayed in Figure 3 shows that, as it was reported for a related system,  $14$ practically there is not absorbance above 500 nm. A summary of the energy and nature of the main low-energy ban[ds](#page-7-0) obtained in a preliminary theoretical study is displayed in Table S2 (Supporting Information). The study was done on the electronic spectrum of the broken-symmetry singlet by means of the time dependent formalism applied to DFT. The molecular orbit[als](#page-7-0) [involved](#page-7-0) [in](#page-7-0) [the](#page-7-0) [elect](#page-7-0)ronic transitions associated with the low-energy bands are depicted in Figure 4. The theoretical calculation indicates that around 500 nm appear several very weak bands corresponding to d-d transitions [i](#page-3-0)n the Mn<sup>III</sup> ions. Such bands are observed in the experimental spectrum at around 520 nm. A more intense band is postulated at 470 nm that involves a d–d transition and a charge transfer from the  $\pi$ electronic density on the pyridine group to a d metal orbital (LMCT). The last transition is the cause of the higher intensity. This transition probably corresponds to the band observed at 420 nm in the experimental spectrum. At higher energies, very intense bands are observed that can be associated with  $\pi-\pi^*$ 



Figure 3. Electronic spectrum of  $[Mn_2OL_2]$ <sup>-4Br·4H<sub>2</sub>O in water.</sup>

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Figure 4. Main selected molecular orbitals involved in the electronic transitions shown in Table S2 (Supporting Information).

transitions in the ligand. These intense bands hide other less intense electronic transitions such as those of moderate intensity that, as predicted by the theoretical study, should appear at 365 and 325 nm. These bands can be attributed to a mixture between metal−metal (MM'CT) and metal−ligand (MLCT) charge transfers.

Magnetic Properties. The molar magnetic susceptibility,  $\chi_{\text{M}}$ , for  $[\text{Mn}_2\text{OL}_2]\cdot 4\text{Br}\cdot 4\text{H}_2\text{O}$  was measured on grounded crystals as a function of the temperature, T. The results are shown in Figure 5 in the form of a  $\chi_{\rm M}T$  versus T plot. At room temperature, the value of  $\chi_M T$  of 0.65 cm<sup>3</sup> mol<sup>-1</sup> K is very much lower than that expected for two magnetically noninteracting high-spin Mn<sup>III</sup> ions  $[\chi_{\text{M}}T = (2N\beta^2 g^2/3k)S(S + 1) =$ 6.0 cm<sup>3</sup> mol<sup>-1</sup> K for S = 2 and g = 2]. Upon cooling,  $\chi_M T$ continuously decreases, and it vanishes below 50 K. This behavior is characteristic of a strong antiferromagnetic coupling between two  $S = 2$  spins. The data were fitted by the theoretical expression derived from an isotropic spin exchange Hamiltonian  $H = -JS_1S_2$  [eq 1].

$$
\chi_{\rm M} = \frac{2N\beta^2 g^2}{kT} \frac{e^x + 5e^{3x} + 14e^{6x} + 30e^{10x}}{1 + 3e^x + 5e^{3x} + 7e^{6x} + 9e^{10x}} \tag{1}
$$



Figure 5. Temperature dependence of  $\chi_{\rm M}T$  for  $[\text{Mn}_2\text{OL}_2]\cdot 4\text{Br}\cdot 4\text{H}_2\text{O}$ . The solid line is the best fit (see text).

The best fit was obtained with  $g = 2$  (kept constant) and  $J =$ −249(2) cm<sup>−</sup><sup>1</sup> . Both the magnitude and sign of the J value are comparable to those observed for other quasi-linear Mn<sup>III</sup>-O−  $Mn^{II}$  cores.<sup>14,15</sup>

However, since, as above stated, there are very few examples of Mn<sup>III</sup> di[nuclea](#page-7-0)r comples with only one oxo bridging ligand, a magneto-structural correlation cannot be established.<sup>10,14,17c</sup>

Moreover, unexpected weak antiferromagnetic coupling has been reported in two cases. In the other cases, the values of the Mn<sup>III</sup>-O-Mn<sup>III</sup> angle fall in a small range close to linearity (−168.5 to 180°), and there is not a clear-cut correlation with the antiferromagnetic coupling constants (J from −165.9 to −251.2 cm<sup>−</sup><sup>1</sup> ). Nevertheless, a discussion about the superexchange pathway in this compound is given below.

Theoretical Computations. DFT calculations of the experimental geometry of complex  $[Mn_2OL_2]^{4+}$  were performed in order to estimate the magnetic coupling constant between the two Mn<sup>III</sup> centers. The obtained *J* value of −244.1 cm<sup>-1</sup> supports a strong antiferromagnetic nature for the coupling. A detailed analysis of the molecular geometry of  $[Mn_2OL_2]^{4+}$  shows that the Mn<sup>III</sup> ions have, because of a second order Jahn−Teller effect, an elongation in an axis noted as the Jahn−Teller axis. Thus, the electronic configuration of the high-spin Mn $^{\text{III}}$ , 3d $^4$ , involves the three  $\bm{{\mathsf{t}}}_{\text{2g}}$  orbitals and the  $\mathrm{d}_{z^2}$  orbitals. Only these orbitals, not all of them, act as magnetic orbitals and contribute to the magnetic coupling. Since the  $d_z$ <sup>2</sup> orbital delocalizes in the  $\sigma$  pathway, it is usually thought that this orbital is mainly responsible for the strong antiferromagnetic coupling. However, the Jahn−Teller axis is placed perpendicular to the exchange pathway direction. Thus, only the electronic density of the  $d_{z^2}$  ring, and not the larger electronic density on the axial lobe, will be delocalized on the exchange pathway. This situation should not favor a large magnetic coupling. However, the exchange pathway is short enough to magnify the magnetic coupling, and moreover, the contribution from the  $t_{2g}$  orbitals to magnetic coupling can be significant. A search in the Cambridge Database shows that, even if only a few examples are reported of  $Mn$ <sup>III</sup> binuclear systems with a single oxo group acting as bridging ligand between the metal ions, in all of them an elongated or compressed Jahn-Teller axis is aligned perpendicular<sup>6,9</sup> or parallel<sup>9,10,14,15,17</sup> to the linear or quasi-linear Mn<sup>III</sup>-O-Mn<sup>III</sup> exchange pathway, respectively (models A and D in Fig[ure](#page-7-0) 4), the latt[er case bei](#page-7-0)ng the most usual. There is also a case with an octahedral-like geometry and another one with a squ[ar](#page-3-0)e pyramidal coordination sphere where exceptionally the  $Mn$ <sup>III</sup> −O−Mn<sup>III</sup> angle is far from linearity.<sup>12,16</sup> In those systems in which magnetic measurements were performed, a strong antiferromagnetic coupling was observe[d.](#page-7-0)[10,1](#page-7-0)4,15 In order to understand the preference for this arrangement and the qualitative contribution of each magneti[c orbi](#page-7-0)tal, we have performed electronic calculations on simplified binuclear models built up as explained in the caption to Figure 6.

The Mn−N(ammonia) and Mn−O bond lengths have been optimized forcing axial elongation and compression for different Jahn−Teller orientations (see Figure 6). In agreement with the experimental data, the most stable conformation is that showing a compression in a Jahn−Teller axis that involves the Mn−O bond (conformation D). However, conformation A with an axial elongation perpendicular to the Mn−O axis, observed in  $[Mn_2OL_2]^{4+}$  and also in some complexes of the literature, is just at 0.75 kcal/mol. Hence, both conformations can be found, and only the adequate design of the peripheral ligand benefits one of them. Since in our models there are not restrictions in peripheral ligands, conformation B (contraction in a perpendicular direction to the Mn−O axis) converges to the most stable conformation A. Finally, model C, where axial elongation is placed on the Mn−O axis, is clearly unstable (26.75 kcal/mol). These results allow for an explanation of the observation of two preferential conformations in the literature.



Figure 6. Molecular models used in the DFT study. Thin and thick dotted lines are employed for elongated and compressed Jahn−Teller axes, repectively. The oxidation state of the metal ion is indicated as a superscript character.

In the two most stable conformations (A and D), strong antiferromagnetic coupling is expected  $(-143.3 \text{ and } -163.4 \text{)}$ cm<sup>−</sup><sup>1</sup> , respectively) but moderately stronger in the last case. It makes sense if we bear in mind that the electronic delocalization from the magnetic orbitals to the pathway is larger for shorter Mn−O bond lengths and that this electronic promotion is larger for  $d_{z^2}$  than for  $d_{x^2-y^2}$ . While the electronic density on the equatorial ring lobe of the  $d_z^2$  is pointing to the Mn−O axis, the  $d_{x^2-y^2}$  orbital does not overlap with the oxo orbitals. On the other hand, a similar model has been built for  $Mn^W$  ions, E, where only the  $t_{2g}$  orbitals acts as magnetic orbitals. Also, a strong antiferromagnetic coupling is found  $(J =$ −112.3 cm<sup>−</sup><sup>1</sup> ). Even though, to our knowledge, there are not any experimental data for such systems, the obtained result is supported by those obtained for mixed  $Mn^{III}-Mn^{IV}$   $\mu$ -oxo complexes that also present strong antiferromagnetic coupling.<sup>14</sup> In the last cases, the magnetic coupling is stronger probably because it can be associated with partial charge trans[fer](#page-7-0)s assisted by the mixed valence character and the nature of the peripheral ligands.<sup>11</sup> In any case, to compare qualitatively the contribution to the magnetic coupling of the  $t_{2g}$  and  $e_g$ orbitals, we must use the  $n^2J$  $n^2J$  product, where *n* is the number of unpaired electrons per center. Notwithstanding that this procedure is purely qualitative because some metal−ligand distances have been modified, it can be intuitive and useful. This product takes the absolute values 2293, 2614, and 1011 for A, D, and E, respectively. The large value for E proves that the contribution of  $t_{2g}$  orbitals is important, as it was suggested and indirectly well proved by Girerd et al. $11,14$  However, the large difference between this value and that for modes A and D supports that [the](#page-7-0) contribution of the  $d_{z^2}$  orbital is also important.

Paramagnetic NMR Measurements. Paramagnetic  ${}^{1}$ H NMR spectroscopy is a powerful technique for evaluating the coordination geometry and electronic structure of metal complexes in solution.31−<sup>33</sup> Usually, mono- and binuclear  $Mn^{III}$  have short electronic relaxation times and, thus, are expected to present wel[l resol](#page-8-0)ved isotropically shifted <sup>1</sup>H NMR spectra.<sup>31</sup> The isotropic shifts of proton resonances in

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Figure 7. 400 MHz proton NMR spectra in D<sub>2</sub>O at pD = 7 and 298 K of  $[Mn_2OL_2]$ <sup>4Br·H<sub>2</sub>O 5 mM. The asterisks mark the residual solvents (\*,</sup>  $H<sub>2</sub>O; **, HOD).$ 

Table 2.  $\rm ^1H$  NMR Hyperfine-Shifted Resonances of  $\rm [Mn_2OL_2]$  Complex in  $\rm D_2O$  at 298 K and pD  $\rm 7$ 

signal	$\delta$ (ppm)	no. of protons	assignments	temperature dependence	$T_1$ (ms)	$\Delta\nu_{1/2}$ (Hz)	$T_2^a$ (ms)
a	57.4			Curie	$\leq$ 1	2210	0.14
b	38.0			Curie	$\leq$ 1	1360	0.23
d	$-21.3$			Curie	$\leq$ 1	2260	0.14
	$-42.9$	32	$\alpha$ CH <sub>2</sub>	Curie	$\leq$ 1	2260	0.14
g	$-57.4$			Curie	$\leq$ 1	b	b
h	$-63.0$			Curie	$\leq$ 1	b	b
$\mathbf{C}$	31.3	4	$H_m$ -py	Curie	$\sim$ 1.2	290	1.1
e	$-32.2$		$H_p$ -py	Curie	$\sim1.5$	240	1.3
$\sim$ - -	. .	$\sim$ $\sim$ $\sim$ $\sim$ $\sim$ $\sim$ $\sim$ .		.			

<sup>a</sup>Measured from the linewidth at half-height. <sup>b</sup>Overlap prevents measurement of this value.



Figure 8. Temperature dependence of the <sup>1</sup>H NMR isotropic shifts of the  $[Mn_2OL_2]^{4+}$  complex at pD = 7. Temperature range 288–323 K.

paramagnetic systems may be of contact or dipolar origin or a combination of both,  $(\Delta H/H_0)^{iso} = (\Delta H/H_0)^{con} + (\Delta H/H_0)^{dip}$ . As already described in the literature,  $34,36$  high-spin Mn<sup>III</sup> and binuclear Mn<sup>III</sup> complexes have large [cont](#page-8-0)ributions by contact shift.35,36

In this work, we have studied the electronic properties of the  $[Mn_2OL_2]$  $[Mn_2OL_2]$  $[Mn_2OL_2]$  $[Mn_2OL_2]$ <sup>-4</sup>Br·4H<sub>2</sub>O complex dissolved in water. We have recorded the paramagnetic  $^1\mathrm{H}$  NMR spectra, measured the  $^1\mathrm{H}$ longitudinal and transversal relaxation times,  $T_1$  and  $T_2$ , and analyzed the temperature dependence of the chemical shifts.

The <sup>1</sup>H NMR spectrum of the binuclear  $[Mn_2OL_2]^{4+}$ complex cation recorded in  $D_2O$  at  $pD = 7$  shows, in the downfield region, three well resolved isotropically shifted signals (a−c; see Figure 7). In addition, it displays five upfield shifted signals (d−h). Hyperfine-shifted resonances, linewidths at half-height, and longitudinal relaxation time values  $(T_1)$  are reported in Table 2. The isotropically shifted signals show linewidths, measured at half-height, of around 1300−2200 Hz, except signals c and e with linewidths of 290 and 240 Hz, respectively. Longitudinal relaxation times vary from less than 1 ms to ∼1.5 ms for signal e.  $T_1$  measurements along with the



**Figure 9.** Cyclic voltammograms at a glassy carbon electrode for a 1.0 mM Mn/L solution in 0.15 M NaClO<sub>4</sub> at pH 3.0 (a) and 7.0 (b). Potential scan rate 50 mV/s.

observed chemical shift values and  $\Delta \nu_{1/2}$  are characteristic of binuclear  $Mn^{III}$  complexes.<sup>34,36</sup> The group of signals a, b, d, and f−h, integrates 32 protons and can be assigned on the basis of the extremely short  $T_1$  va[lues](#page-8-0) and large chemical shifts to the CH<sub>2</sub> protons closest to the manganese centers in  $Mn^{\text{III}}_{2}OL_{2}$  ( $\alpha$ -CH<sub>2</sub>, see inset in Figure 7). As described in the literature,<sup>36</sup>  $\alpha$ - $CH<sub>2</sub>$  protons in spin-coupled Mn<sup>III</sup> complexes have similar large shifts, broad linewidths, a[nd](#page-5-0) short  $T_1$  values. The other gro[up](#page-8-0) of signals (c, e), which have a longer relaxation time ( $T_1 \sim 1.2$  and  $\sim$ 1.5 ms respectively) and relatively narrow signals ( $\Delta \nu_{1/2}$  of 290 and 240 Hz respectively), agree with the assignment to meta and para pyridine protons, respectively (see Table 2).

The above-mentioned isotropic shift resonances in paramagnetic systems can have a contact or dipolar origi[n o](#page-5-0)r a combination of both. The contact term presents a linear dependency with the inverse of the temperature  $(T^{-1})$ , whereas the dipolar term is a function of  $T^{-2,3}$ <sup>7</sup> Thus, it is possible to . use the dependency on the isotropic shift resonances with the temperature to estimate the relati[ve](#page-8-0) importance of these contributions. Nevertheless, due to the fact that the studied interval of temperatures is narrow because of the limitations imposed by the solvent, typically an inversely proportional change of the isotropic shift resonances is observed with the temperature. Deviations due to a small  $T^{-2}$  component cannot be detected. The extrapolation to infinite temperature in the plot of  $\delta$  as a function of the  $T^{-1}$  is the method used for the study of the dependency type of the isotropic shifts with the temperature. Important deviations from the zero extrapolation out of the diamagnetic region will indicate the existence of a dipolar contribution to the isotropic shifts. Variable temperature <sup>1</sup>H NMR spectra of  $[Mn_2OL_2]$ ·4Br·4H<sub>2</sub>O at pD = 7 were registered from 288 to 323 K. In Figure 8, the observed isotropic shift resonances are plotted as a function of T<sup>−</sup><sup>1</sup> . All isotropic shift resonances follow a Curie [be](#page-5-0)havior as they invariably decrease with increasing temperature. The isotropic shift signals h and e are more temperature dependent and show an appreciable deviation from the zero extrapolation at 29.6 and 15.9 ppm, respectively. As described above, these results are indicative of some dipolar contribution to the hyperfine shifted resonances originated by the presence of a zero-field splitting and consequently of certain magnetic anisotropy. The other signals show  $\delta$  intercept values at infinite temperature within or close to the diamagnetic region and are in general indicative of a predominant contact contribution. These results agree with the existence of a magnetically coupled binuclear Mn<sup>III</sup> system with antiferromagnetic coupling. The paramagnetic NMR results are consistent with the magnetic studies described and show that the  $[Mn_2OL_2]$  4Br·4H<sub>2</sub>O complex present in

solution has similar electronic properties as that in the solid state.

Electrochemistry. Figure 9 collects cyclic voltamperograms of aqueous solutions of  $Mn^{\text{II}}$  and L recorded at pH < 6. The electrochemical response is essentially identical to that of uncomplexed  $Mn^{II}$  solutions, consisting of an irreversible oxidation signal at +0.70 V vs Ag/AgCl, corresponding to the oxidation of the free  $Mn^{\text{II}}$  to  $MnO_2^{38}$  preceded by a weak reversible wave at a midpeak potential of +0.40 V attributable to one-electron oxidat[io](#page-8-0)n of Mn<sup>II</sup> ions adsorbed on the electrode surface to adsorbed  $Mn$ <sup>III</sup> ones. In the region of negative potentials, there is a cathodic wave at −0.58 V corresponding to a two-electron reduction of  $Mn^{\text{II}}$  (see Figure 9a).

Above  $pH = 6$ , the complex starts forming, and the electrochemical response changes dramatically. The couple at +0.40 V and the oxidation peak at +0.70 V disappear, while a new oxidation peak appears at +0.55 V (see Figure 9b) followed, in the subsequent cathodic scan, by weak reduction peaks at +0.44 and −0.08 V. On the basis of the variation of peak potentials with the potential scan rate and the pH, it was concluded that the oxidation process at +0.55 V corresponds to a reversible one electron transfer process followed by a rapid chemical reaction (EC mechanism),<sup>39</sup> possibly due to dissociation or ligand reorganization. Accordingly, cathodic peaks at +0.44 and −0.08 V correspon[d](#page-8-0) to the reduction of different  $Mn^{III}$  species electrochemically generated.

Figure 10 shows the cyclic voltammetric response at a glassy carbon electrode of a deaerated 1.0 mM solution of  $[Mn_2OL_2]$ ·4Br·4H<sub>2</sub>O in 0.10 M Bu<sub>4</sub>NPF<sub>6</sub>/DMSO. In the initial an[odi](#page-7-0)c scan, two overlapping peaks at  $+0.25$  (A<sub>1</sub>) and +0.53 V  $(A_2)$  vs coupled  $Fe^+/Fe$  appear in the subsequent cathodic scan, with weaker reduction peaks at  $+0.65$  (C<sub>2</sub>) and +0.35 V  $(C_1)$ . This voltammetric profile corresponds to two successive one-electron transfer processes which can be judged as nearly reversible on the basis of the anodic-to-cathodic peak potential separation, which tends to 60 mV at low potential scan rates for both couples. No initial reduction processes were detected in the investigated potential window. The observed voltammetric response can be described in terms of two successive one-electron transfer processes as

$$
\text{couple } A_1 / C_1: \left[ \text{Mn}^{\text{III}} \mathbf{L} - \mathbf{O} - \mathbf{L} \text{Mn}^{\text{III}} \right]^{4+}
$$
\n
$$
\rightarrow \left[ \text{Mn}^{\text{IV}} \mathbf{L} - \mathbf{O} - \mathbf{L} \text{Mn}^{\text{III}} \right]^{5+} + e^- \tag{2}
$$

$$
\text{couple } A_2/C_2: \left[ \text{Mn}^{\text{III}} \mathbf{L} - \mathbf{O} - \mathbf{L} \text{Mn}^{\text{IV}} \right]^{5+}
$$

$$
\rightarrow \left[ \text{Mn}^{\text{IV}} \mathbf{L} - \mathbf{O} - \mathbf{L} \text{Mn}^{\text{IV}} \right]^{6+} + e^{-}
$$
(3)

<span id="page-7-0"></span>

Figure 10. Cyclic voltammogram (a) and its deconvoluted curve (b) recorded at a glassy carbon electrode of a deaerated 1.0 mM solution of  $Mn_2OL_2$  in 0.10 M Bu<sub>4</sub>NPF<sub>6</sub>/DMSO. Potential scan rate 20 mV/s.

Oxidized forms are stable on the time scale of voltammetric experiments, as denoted by the maintenance of the voltammetric pattern in repetitive voltammetry, thus suggesting that a similar  $Mn^V$  O-bridged complex is finally formed.

#### ■ CONCLUSIONS

The crystal structure of the binuclear  $Mn^{III}$  complex  $[Mn_2OL_2]$ ·4Br·4H<sub>2</sub>O exhibits a single oxo bridging ligand with a large Mn<sup>III</sup>-O−Mn<sup>III</sup> angle. Magnetic susceptibility measurements indicate a strong antiferromagnetic coupling between the metal centers  $(J = -249(2) \text{ cm}^{-1})$ . The paramagnetic NMR studies support the same structure as well as an antiferromagnetic coupling for the binuclear complex in solution. Electrochemical data suggest that the complex is reversibly oxidized in DMSO via two successive one-electron transfer processes forming, presumably, a similar binuclear  $Mn^{\text{IV}}$  complex.

#### ■ ASSOCIATED CONTENT

#### **S** Supporting Information

Crystallographic data, selected experimental and calculated electronic transition energy data, and a crystallographic information file. This material is available free of charge via the Internet at http://pubs.acs.org.

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#### Notes

The auth[ors declare no competin](mailto:enrique.garcia-es@uv.es)g fi[nancial interest.](mailto:francisco.lloret@uv.es)

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