

Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb) and Ba₂LnGaS₅ (Ln = Pr, Nd): Syntheses, Structures, and Magnetic and Optical Properties

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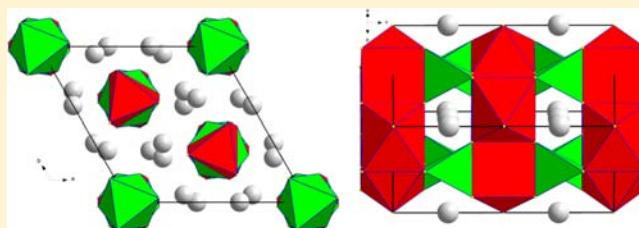
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Supporting Information

ABSTRACT: Six new quaternary rare-earth sulfides Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb) and Ba₂LnGaS₅ (Ln = Pr, Nd) have been synthesized for the first time. Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb) belong to the centrosymmetric space group *R*3̄c of the trigonal system. The structures contain infinite one-dimensional anionic chains $[\text{LnInS}_6]^{6-}$, which are built from face-sharing LnS₆ distorted triangular prisms and InS₆ octahedra. Ba₂LnGaS₅ (Ln = Pr, Nd) crystallize in the centrosymmetric space group *I*4/*mcm* of the tetragonal system. Their structures consist of (BaLn)S layers built from (BaLn)S₈ bicapped trigonal prisms. These layers are stacked perpendicular to the *c* axis and further connected by GaS₄ tetrahedra to form a three-dimensional framework with channels occupied by Ba²⁺ cations. As deduced from magnetic susceptibility measurements on Ba₂NdGaS₅, it is paramagnetic and obeys the Curie–Weiss law. Besides, the band gap of Ba₂NdGaS₅ is determined to be about 2.12(2) eV.



INTRODUCTION

Interest in rare-earth chalcogenides originates not only from their rich structures as a result of the diversity in the geometry of the Ln-centered coordination polyhedra and the connectivity between them but also from the important physical properties related to the 4f electrons.^{1,2} In recent years, extensive exploratory synthesis has led to the discovery of many new multinary rare-earth chalcogenides exhibiting interesting structures and physical properties.^{3–33} For example, CsCu₃Dy₂S₅ and CsCu₃Er₂S₅ have an isotypic three-dimensional network $[\text{Cu}_3\text{M}_2\text{S}_5]^-$, spreading along the *a* axis with large channels, that is well-suited to take up the highly coordinated Cs⁺ cations;³⁴ KCuCe₂Se₆ has an interesting structure consisting of two-dimensional $[\text{CuCe}_2\text{Se}_6]^-$ layers with orderly arranged Cu atoms;³⁵ in the compounds BaLn₂FeS₅ (Ln = Ce, Pr, Nd, Sm), the Fe²⁺ ion exhibits an antiferromagnetic ordering at around 40 K,³⁶ and AgPb_{*m*}LaTe_{*m+2*} shows high electrical conductivity and a relatively small Seebeck coefficient.³⁷

Most of these multinary rare-earth chalcogenides also contain a d-block transition metal. However, increasing attention has recently been paid to chalcogenides containing both f-block rare-earth metals and p-block main-group elements: for example, K₂Ln₂As₂Se₉ (Ln = Sm, Gd) contains interesting

chairlike As₂Se₄ units,³⁸ and the Eu₂Ga₂GeS₇,²⁸ La₄InSbS₉,³² and Sm₄GaSbS₉²⁷ compounds are found to show strong second-harmonic-generation responses in middle IR. In our research, we think that the quaternary A/Ln/M/Q (A = alkaline-earth metal; Ln = rare earth; M = group IIIA metals Ga and In; Q = S, Se, Te) system is attractive because of the combination of these microscopic nonlinear optical (NLO)-active MQ₄ units with magnetic rare-earth cations in one chalcogenide, and the interplay of the covalent M–Q bonding with the ionic Ln–Q or A–Q bonding may generate compounds with multifunctional properties and interesting structures. Our earlier exploratory efforts have led to the discovery of 12 new selenides in this family, namely, the Ba₂LnMSe₅ (M = Ga, In; Ln = Y, Nd, Sm, Gd, Dy, Er) compounds, which adopt two different structure types for the six Ga-containing selenides and six In-containing selenides, respectively, and exhibit interesting NLO and magnetic properties.³⁹

Here, we extend synthetic efforts to the sulfides in this quaternary A/Ln/M/Q (A = alkaline-earth metal; Ln = rare earth; M = group IIIA metals Ga and In; Q = S, Se, Te) system.

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Table 1. Crystal Data and Structure Refinements for Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb)^a and Ba₂LnGaS₅ (Ln = Pr, Nd)^b

	Ba ₃ PrInS ₆	Ba ₃ SmInS ₆	Ba ₃ GdInS ₆	Ba ₃ YbInS ₆	Ba ₂ PrGaS ₅	Ba ₂ NdGaS ₅
fw	860.11	869.55	876.45	892.24	645.61	648.94
<i>a</i> (Å)	12.1196(2)	12.063(2)	12.072(2)	12.057(2)	8.16340(10)	8.1061(11)
<i>c</i> (Å)	13.9803(5)	13.853(3)	13.816(3)	13.713(3)	13.6093(3)	13.541(3)
<i>V</i> (Å ³)	1778.38(8)	1745.8(5)	1743.6(5)	1726.2(5)	906.94(3)	889.8(2)
ρ_c (g/cm ³)	4.819	4.962	5.008	5.150	4.728	4.844
μ (cm ⁻¹)	27.95	17.930	18.607	21.158	17.845	18.548
<i>R</i> (<i>F</i>) ^c	0.0122	0.0199	0.0225	0.0177	0.0204	0.0286
<i>R</i> _w (<i>F</i> _o ²) ^d	0.0293	0.0383	0.0581	0.0399	0.0467	0.0736

^aFor all four structures, *Z* = 6, space group = *R* $\bar{3}c$, *T* = 153 (2) K, and λ = 0.71073 Å. ^bFor both structures, *Z* = 4, space group = *I4/mcm*, *T* = 153 (2) K, and λ = 0.71073 Å. ^c $R(F) = \sum ||F_o| - |F_c|| / \sum |F_o|$ for $F_o^2 > 2\sigma(F_o^2)$. ^d $R_w(F_o^2) = \{ \sum [w(F_o^2 - F_c^2)^2] / \sum wF_o^4 \}^{1/2}$ for all data. $w^{-1} = \sigma^2(F_o^2) + (zP)^2$, where $P = (\max(F_o^2, 0) + 2F_c^2) / 3$.

To our surprise, the first sulfides that we obtained, with the formula Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb), possess stoichiometries and structures different from those of the 12 selenides reported earlier, which clearly demonstrated the richness of the phases in the quaternary A/Ln/M/Q system. The Ba₃LnInS₆ structures feature one-dimensional [LnInS₆]⁶⁻ chains of ordered arrays of face-sharing LnS₆ distorted trigonal prisms and InS₆ octahedra. Actually, they represent the first sulfide analogues of the extensively studied hexagonal perovskite-related (A₃A'BO₆) oxides, which have displayed very flexible compositions and interesting magnetic and optical properties.^{40–42} The investigation into related chalcogenides of the (A₃A'BO₆) oxides would be worthwhile because the increased covalence of the metal–chalcogen bond compared to the metal–oxygen bond may have a different influence on the electronic structures and physical properties. Even more interestingly, when we try to replace the In atom of these Ba₃LnInS₆ compounds with a Ga atom, the obtained products are not the hypothetical Ba₃LnGaS₆ compounds but the Ba₂LnGaS₅ (Ln = Pr, Nd) compounds. Although Ba₂LnGaS₅ (Ln = Pr, Nd) possesses the same stoichiometry as the 12 selenides reported earlier, it adopts a structure type totally different from those of the selenides. In this paper, we report the syntheses, structural characterizations, and magnetic and optical properties of Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb) and Ba₂LnGaS₅ (Ln = Pr, Nd).

EXPERIMENTAL SECTION

Syntheses. BaS (99%, Sinopharm Chemical Reagent Co., Ltd.), Ga (99.99%, Sinopharm Chemical Reagent Co., Ltd.), In (99.99%, Sinopharm Chemical Reagent Co., Ltd.), S (99.99%, Sinopharm Chemical Reagent Co., Ltd.), and Ln (Ln = Pr, Nd, Sm, Gd, Yb; 99.9%, Alfa Aesar China (Tianjin) Co., Ltd.) were used as received. The binary starting materials Ga₂S₃ and In₂S₃ were prepared from direct reactions of the elements at high temperatures in sealed silica tubes evacuated to 10⁻³ Pa.

Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb). The mixtures of BaS (0.339 g, 2 mmol), In₂S₃ (0.163 g, 0.5 mmol), Ln (Ln = Pr, Sm, Gd, Yb; 1 mmol), and S (0.048 g, 1.5 mmol) were ground, loaded into fused-silica tubes under an argon atmosphere in a glovebox, then flame-sealed under a high vacuum of 10⁻³ Pa. The tubes were then placed in computer-controlled furnaces, heated to 1323 K in 24 h, left for 48 h, cooled to 693 K at a rate of 3 K/h, and finally cooled to room temperature by switching off the furnace. Dark-red crystals were found in the ampules. The crystals were stable in air.

The chip-shaped crystals were manually selected for structure characterization and determined as Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb). Analyses of the crystals with an energy-dispersive X-ray (EDX)-equipped Hitachi S-4800 scanning electron microscope showed the presence of Ba, Ln, In, and S in the approximate ratio of 3:1:1:6.

Ba₂LnGaS₅ (Ln = Pr, Nd). The mixtures of BaS (0.339 g, 2 mmol), Ga₂S₃ (0.118 g, 0.5 mmol), Ln (Ln = Pr, Nd; 1 mmol), and S (0.048 g, 1.5 mmol) were ground, loaded into fused-silica tubes under an argon atmosphere in a glovebox, and then flame-sealed under a high vacuum of 10⁻³ Pa. The tubes were then placed in computer-controlled furnaces, heated to 1323 K in 24 h, left for 48 h, cooled to 693 K at a rate of 3 K/h, and finally cooled to room temperature by switching off the furnace. Dark-red crystals were found in the ampules. The crystals were stable in air.

The chip-shaped crystals were manually selected for structure characterization and determined as Ba₂LnGaS₅ (Ln = Pr, Nd). Analyses of the crystals with an EDX-equipped Hitachi S-4800 scanning electron microscope showed the presence of Ba, Ga, Ln, and S in the approximate ratio of 2:1:1:5. Inductively coupled plasma (ICP) measurement on the crystals indicated that the molar ratio of Ba/Ln/Ga is close to 2:1:1.

Structure Determination. The single-crystal X-ray diffraction measurements were performed on a Rigaku AFC10 diffractometer equipped with graphite-monochromated K α (λ = 0.71073 Å) radiation at 153 K. The *CrystalClear* software⁴³ was used for data extraction and integration, and the program *XPREP*⁴⁴ was used for face-indexed absorption corrections.

The structure was solved with direct methods implemented in the program *SHELXS* and refined with the least-squares program *SHELXL* of the *SHELXTL*.PC suite of programs.⁴⁴ For Ba₃LnInS₆ (Ln = Pr, Sm, Gd, Yb), one Ba position (Wyckoff site 18e), one Ln position (Wyckoff site 6a), one In position (Wyckoff site 6b), and one S position (Wyckoff site 36f) were found. It was straightforward to refine the structure of the compounds containing the larger rare-earth cations, namely, Pr³⁺, Sm³⁺, and Gd³⁺: the anisotropic displacement parameters of all atoms are normal, and refinement of the occupancy of all atoms led to values very close to 100% occupancy, which indicated no detectable disorder between the Ln³⁺ and In³⁺ cations in these compounds. However, in the initial structure refinement of Ba₃YbInS₆, which contains the smaller rare-earth element Yb, the isotropic displacement parameter of Yb was unusually large, while that of In was close to zero, which was indicative of possible disorder between the Yb and In cations. Subsequently, disorder of Yb and In was introduced at Yb (Wyckoff site 6a) and In (Wyckoff site 6b) positions, and the refinements showed that Yb and In were disordered at the Yb (Wyckoff site 6a) and In (Wyckoff site 6b) positions with Yb/In ratios of 0.840:0.160 and 0.158:0.842, respectively. The resultant *R* index *R*₁/*wR*₂ decreased from 0.0261/0.0885 to 0.0177/0.0399 and the largest peak/hole in the difference electron density map was reduced from 4.57/−2.82 to 0.88/−1.16 e/Å³ with this disorder model. Considering that virtually the same total amounts of Yb and In cations in the structure were present, the final stoichiometry was also determined as Ba₃YbInS₆.

For Ba₂LnGaS₅ (Ln = Pr, Nd), three metal positions are found: M1 (Wyckoff site 4a) surrounded by about 10 S atoms with a M1–S distance of about 3.4 Å, M2 (Wyckoff site 8h) coordinated to 8 S atoms at a distance of around 3.0 Å, and M3 (Wyckoff site 4b) coordinated to 4 S atoms at a distance of around 2.2 Å. In addition,

two S atoms (Wyckoff sites 4c and 16l) were decided. On the basis of the bond lengths and the electron density, Ba can be assigned to the M1 position and Ga to the M3 position, respectively. From the bond length, the Ln atom should take the M2 position. However, if the M2 position is totally occupied by the Ln elements, the resultant composition $\text{BaLn}_2\text{GaS}_5$ neither agrees with the EDX or ICP results regarding the elemental ratios in these compounds nor satisfies the charge-balance requirements. A disordered model was then tried, in which the M2 position was occupied by both Ba and Ln atoms. The resultant occupancy of Ba at the M2 position is close to 50%. Hence, the M2 position is set to be occupied by Ba and Ln in a ratio of 1:1 in the final refinement, leading to the formula of $\text{Ba}_2\text{LnGaS}_5$ (Ln = Pr, Nd), which agrees with the elemental analysis results by EDX and ICP and achieves charge balance. Similar structural features were found in the $\text{Sr}_2\text{EuAlO}_5$ and $\text{Sr}_2\text{EuFeO}_5$ oxides.⁴⁵

The final refinements of all structures include anisotropic displacement parameters and secondary extinction correction. The program *STRUCTURE TIDY*⁴⁶ was then employed to standardize the atomic coordinates. Additional experimental details are given in Table 1, and selected metrical data are given in Tables 2 and 3. Further information may be found in the Supporting Information.

Table 2. Selected Interatomic Distances (Å) for $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb)

	$\text{Ba}_3\text{PrInS}_6$	$\text{Ba}_3\text{SmInS}_6$	$\text{Ba}_3\text{GdInS}_6$	$\text{Ba}_3\text{YbInS}_6$
Ba–S×2	3.171(1)	3.151(1)	3.154(2)	3.147(1)
Ba–S×2	3.233(1)	3.225(1)	3.232(1)	3.251(1)
Ba–S×2	3.300(1)	3.286(1)	3.288(2)	3.287(1)
Ba–S×2	3.335(1)	3.316(1)	3.318(1)	3.311(1)
Ln–S×6	2.836(1)	2.799(1)	2.784(1)	2.735(1)
In–S×6	2.633(1)	2.624(1)	2.622(2)	2.620(1)

Table 3. Selected Interatomic Distances (Å) for $\text{Ba}_2\text{LnGaS}_5$ (Ln = Pr, Nd)

	$\text{Ba}_2\text{PrGaS}_5$	$\text{Ba}_2\text{NdGaS}_5$
Ba1–S1×2	3.402(1)	3.385(1)
Ba1–S2×8	3.442(1)	3.421(2)
Ln–S2×2	2.922(1)	2.901(2)
Ln–S1×2	3.056(1)	3.036(1)
Ln–S2×4	3.200(1)	3.174(2)
Ga–S2×4	2.246(1)	2.239(2)

Magnetic Susceptibility Measurements. Because of the low yields in our synthesis and the tiny size of the obtained crystals, only enough $\text{Ba}_2\text{NdGaS}_5$ crystals can be picked for measuring the magnetic susceptibility. A SQUID magnetometer (Quantum Design) operating at 10 kOe was used to carry out magnetic susceptibility measurements on $\text{Ba}_2\text{NdGaS}_5$. The zero-field-cooled and field-cooled (ZFC/FC) magnetic susceptibility was obtained in the temperature range of 2–300 K on ground crystals in a gelatin capsule.

Diffuse-Reflectance Spectroscopy. A Cary 5000 UV–visible–near-IR (NIR) spectrophotometer with a diffuse-reflectance accessory was used to measure the spectrum of $\text{Ba}_2\text{NdGaS}_5$ over the range 400 nm (3.10 eV) to 1300 nm (0.95 eV).

RESULTS AND DISCUSSION

Synthesis. Six sulfides $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb) and $\text{Ba}_2\text{LnGaS}_5$ (Ln = Pr, Nd) have been synthesized by traditional high-temperature solid-state reactions for the first time. The yields range from 10% to 20% based on Ln. Great efforts have been made to synthesize analogues containing other rare-earth elements available to us, namely, Y, La, Ce, Nd, Dy, Er, and Lu for $\text{Ba}_3\text{LnInS}_6$ and Y, La, Ce, Sm, Gd, Dy, Er,

Yb, and Lu for $\text{Ba}_2\text{LnGaS}_5$, but it did not work. Thus, we only report here the six members that we obtained.

Structure of $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb). The four compounds $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb) are isotypic and crystallize in the K_4CdCl_6 structure type,⁴⁷ which is in the centrosymmetric trigonal space group $R\bar{3}c$. The asymmetric unit of $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb) contains one crystallographically independent Ba atom, one Ln atom, one In atom, and one S atom, which are at Wyckoff positions 18e, 6a, 6b, and 36f, respectively. Without metal–metal or S–S bonds in the structures, the oxidation states of 2+, 3+, 3+, and 2– can be assigned to Ba, In, Ln, and S, respectively.

The structure of $\text{Ba}_3\text{GdInS}_6$ is illustrated in Figure 1. The basic structural unit of $\text{Ba}_3\text{GdInS}_6$ is an infinite one-dimensional

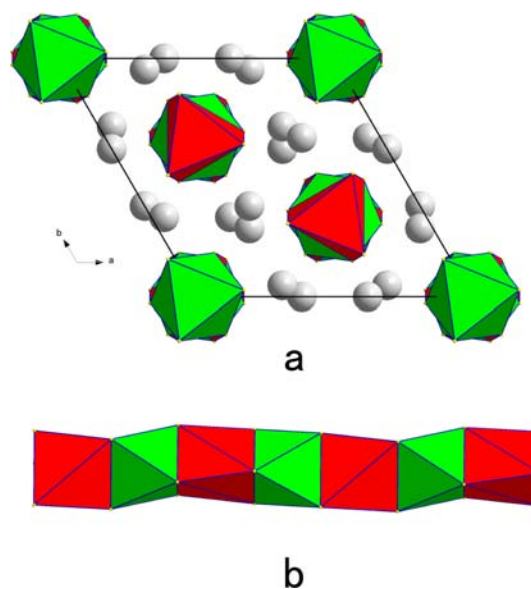


Figure 1. Crystal structure (a) and the $[\text{GdInS}_6]^{6-}$ anionic chain (b) of $\text{Ba}_3\text{GdInS}_6$ (GdS_6 and InS_6 polyhedra are shaded in red and green, respectively; the white balls are Ba atoms).

$[\text{GdInS}_6]^{6-}$ anionic chain built from InS_6 octahedra (oct) and GdS_6 distorted triangular prisms (tp), which are connected alternately via face-sharing to form chains along the *c*-axis direction, which are separated by Ba atoms in a bicapped trigonal-prismatic geometry. The compound Ba_4CrUS_9 ⁴⁸ possesses a similar structure with one-dimensional $[\text{Cr}_2\text{US}_9]^{8-}$ chains separated by Ba^{2+} cations. $[\text{Cr}_2\text{US}_9]^{8-}$ is a chain built from CrS_6 octahedra and US_6 trigonal prisms in the sequence *oct oct tp oct oct tp*, while for $[\text{GdInS}_6]^{6-}$, InS_6 octahedra and GdS_6 distorted triangular prisms arrange in the sequence of *oct tp oct tp*. In addition, there are two crystallographically unique one-dimensional chains in Ba_4CrUS_9 , but only one exists in $\text{Ba}_3\text{GdInS}_6$.

The selected bond distances of $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb) are listed in Table 2. The Ba–S distances of 3.171(1)–3.335(1) Å are similar to those of $\text{Ba}_2\text{AgInS}_4$ [3.128(2)–3.314(2) Å] for eight-coordinated Ba;⁴⁹ The In–S distances lie in the range of 2.620(1)–2.633(1) Å, close to those of $\text{Bi}_3\text{In}_4\text{S}_{10}$ [2.51(1)–2.82(2) Å] for six-coordinated In atoms.⁵⁰ The Ln–S distances are also normal for the six-coordinated Ln cations: the Pr–S distance of 2.836(1) Å is comparable to those in $\text{BaPr}_2\text{MnS}_5$ [2.773(7)–3.081(7) Å],⁵¹ the Sm–S distance of 2.799(1) Å is similar to those in $\text{BaSm}_2\text{FeS}_5$ [2.752(1)–

3.017(1) Å],⁵² the Gd–S distance of 2.784(1) Å resembles those in Ba₄Gd₂Cd₃S₁₀ [2.765(3)–2.799(3) Å],⁵³ and the Yb–S distance of 2.735(1) Å agrees with those in CaYbInS₄ [2.678(1)–2.694(2) Å].³ In addition, it is interesting to note that twist angles between the pair of triangular faces (15.604°, 15.745°, 15.817°, and 15.987° for Ln = Pr, Sm, Gd, and Yb, respectively) increase with a decrease of the Ln³⁺ ion radius.

Structure of Ba₂LnGaS₅ (Ln = Pr, Nd). The two compounds Ba₂LnGaS₅ (Ln = Pr, Nd) are isostructural and crystallize in the Cs₃CoCl₅ structure type of centrosymmetric tetragonal space group *I4/mcm*.⁵⁴ In the asymmetric unit of Ba₂LnGaS₅ (Ln = Pr, Nd), Ba1, Ga, S1 and S2 atoms are at Wyckoff positions 4a, 4b, 4c, and 16l with 100% occupancy, respectively, while the Ba2 and Ln atoms are disordered at Wyckoff position 8h with 50% occupancy for each atom. Because there are no metal–metal or S–S bonds in the structures, the oxidation states of 2+, 3+, 3+, and 2– can be assigned to Ba, Ga, Ln, and S, respectively.

The structure of Ba₂NdGaS₅ is shown in Figure 2. Each Ga atom is coordinated to four S atoms to form perfect GaS₄ tetrahedra with four identical Ga–S bonds. The Nd(Ba2) atom is joined to eight S atoms to form a bicapped trigonal prism (btp). Ba1 atoms are coordinated by 10 S atoms. As shown in Figure 2, one Nd(Ba2)S₈ bicapped trigonal prism (btp) is

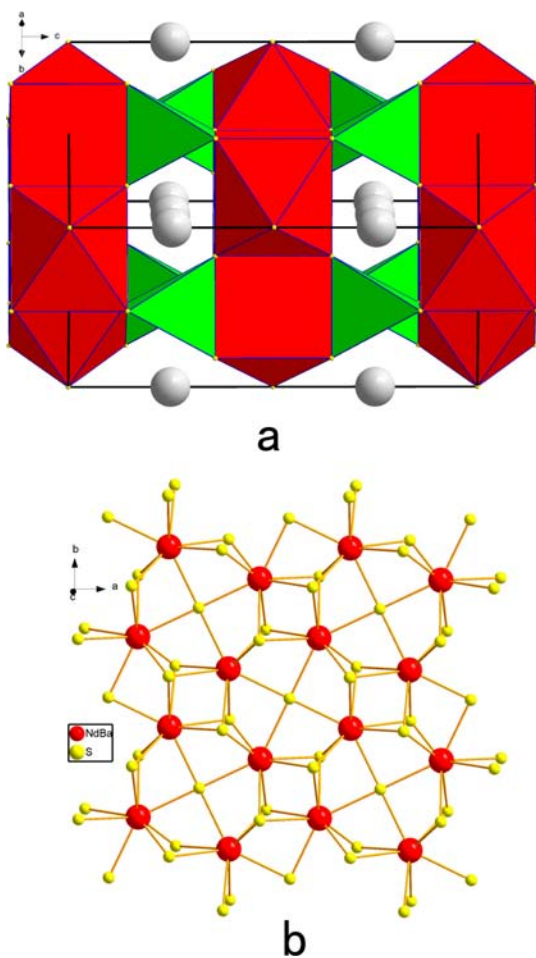


Figure 2. Crystal structure seen along the [110] direction (a) and the Nd(Ba2)S layer (b) of Ba₂NdGaS₅ [Nd(Ba2)S₈ and GaS₄ polyhedra are shaded in red and green, respectively; the white balls are Ba atoms].

linked to one nearest-adjacent Nd(Ba2)S₈ btp by sharing the four-membered face and to four next-nearest-neighboring Nd(Ba2)S₈ btp's by sharing the three-membered faces, and in this way, a two-dimensional anionic (BaNd)S layer was generated. The layers are stacked along the *c* axis and further connected by isolated GaS₄ tetrahedra to form a three-dimensional framework with channels occupied by Ba atoms along the [110] direction. The structure of Ba₂GaNdS₅ is similar to that of BaLn₂FeS₅,⁵³ except that the Wyckoff 8h metal position was occupied by Ln metals only in the BaLn₂FeS₅ series of compounds. Another isostructural series of compounds are the Sr₂EuAlO₅ and Sr₂EuFeO₅ oxides,⁴⁵ in which the Wyckoff 8h metal position was half-occupied by Eu atoms and half-occupied by Sr atoms, just as in the case of Ba₂LnGaS₅ (Ln = Pr, Nd).

In comparison, the corresponding selenides Ba₂LnGaSe₅ (Ln = Y, Nd, Sm, Gd, Dy, Er) crystallize in the space group *P* $\bar{1}$ of the triclinic system, and the structure features a one-dimensional ∞ [LnGaSe₅]⁴⁻ chain built from LnSe₆ octahedra and GaSe₄ tetrahedra and separated by Ba²⁺ cations. Because of the smaller size of the S atoms, Ln atoms can be coordinated by two more S atoms in the sulfides than in the selenides. The more surrounding S atoms, in turn, will enhance the opportunity of connecting among these LnS₈ polyhedra, leading to a denser framework, such as that in Ba₂LnGaS₅ (Ln = Pr, Nd).

The selected bond distances of Ba₂LnGaS₅ (Ln = Pr, Nd) are listed in Table 3. The distances of the Ga–S bonds are 2.239(2) Å for Ba₂NdGaS₅ and 2.246(1) Å for Ba₂PrGaS₅, which agree with those of 2.212(1)–2.239(1) Å in BaGa₂SiS₆ for tetrahedrally coordinated Ga atoms.⁵⁴ The Ba–S distances of 3.385(1)–3.442(1) Å are similar to those of BaGa₂SiS₆ [3.474(1)–3.6320(8) Å].⁵⁴ The Pr–S [2.922(1)–3.200(1) Å] and Nd–S [2.901(2)–3.174(2) Å] distances are a little larger than the 2.773(7)–3.081(7) Å distances of Pr–S in BaPr₂MnS₅⁵² and the 2.754(8)–3.079(8) Å distances of Nd–S in BaNd₂MnS₅.⁵⁵ Such a deviation may result from the occupation of Ba in these Ln positions because the Ba²⁺ cation has a larger radius than the Ln³⁺ cation.

Magnetic Susceptibility Measurement. Figure 3 shows the temperature variation of the molar magnetic susceptibility (χ_m) and the inverse magnetic susceptibility ($1/\chi_m$) for the Ba₂NdGaS₅ compound. The ZFC and FC magnetic susceptibility data are essentially superimposable at all temperatures.

The magnetic susceptibility data of the Ba₂NdGaS₅ sample were fitted by a least-squares method to a modified Curie–

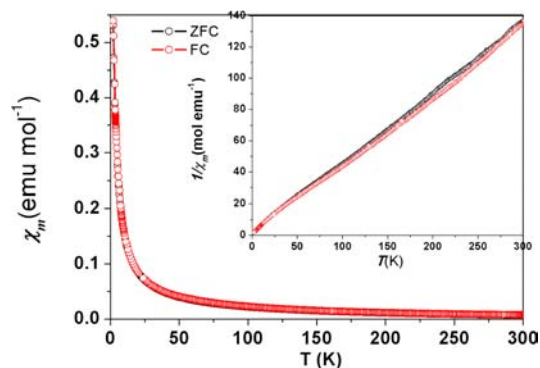


Figure 3. χ_m versus temperature of Ba₂NdGaS₅ for FC and ZFC data. Inset: plot of $1/\chi_m$ versus temperature.

Weiss law: $\chi_m = C/(T - \theta)$, where C is the Curie constant and θ is the Weiss constant. The effective magnetic moment $[\mu_{\text{eff}}(\text{total})]$ was calculated from the equation $\mu_{\text{eff}}(\text{total}) = (7.997C)^{1/2} \mu_B$.⁵⁶

As shown, it is paramagnetic and obeys the Curie–Weiss law over the entire experimental temperature range. As deduced from the fitting results, the values of C and θ for $\text{Ba}_2\text{NdGaS}_5$ are 2.25 and 3.69, respectively. The calculated effective magnetic moment is $4.24 \mu_B/\text{Nd}$ atom, which is a little bigger than the calculated theoretical value for the Nd^{3+} ion ($3.62 \mu_B/\text{Nd}$ atom).⁵⁷ The divergence between ZFC and FC data within the 200–250 K range may be from a tiny amount of an impure phase, which is attached on the surfaces of the picked crystals and is impossible to get rid of.

Experimental Band Gap. On the basis of the UV–visible–NIR diffuse-reflectance spectra of $\text{Ba}_2\text{NdGaS}_5$ (Figure 4), the

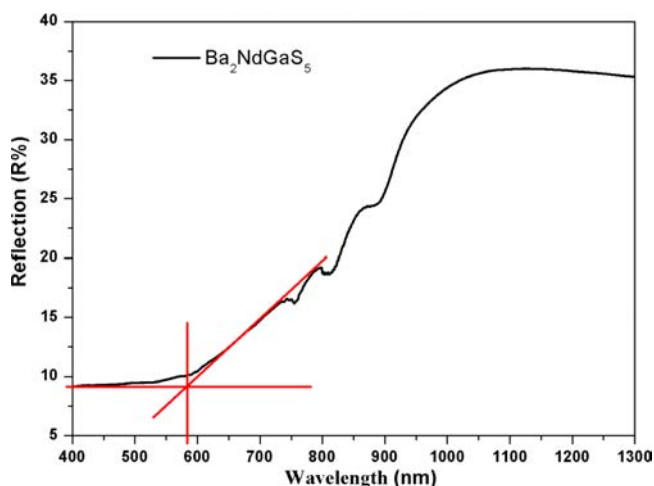


Figure 4. Diffuse-reflectance spectra of $\text{Ba}_2\text{NdGaS}_5$.

band gap can be deduced by a straightforward extrapolation method.⁵⁸ The absorption edge of $\text{Ba}_2\text{NdGaS}_5$ is about 585 nm (2.12 eV) and agrees with the dark-red color. As can be seen, the existence of typical f – f transitions of Nd^{3+} ions brings several broad absorption bands below the optical band gaps between 500 and 950 nm.²²

CONCLUSIONS

In summary, six new rare-earth sulfides in the quaternary A/Ln/M/Q (A = alkaline earth; Ln = rare earth; M = group IIIA metals Ga and In; Q = S, Se, Te) system, $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb) and $\text{Ba}_2\text{LnGaS}_5$ (Ln = Pr, Nd), have been discovered and characterized. $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb) crystallizes in the centrosymmetric space group $R\bar{3}c$ of the trigonal system with infinite one-dimensional anionic chains $\infty[\text{LnInS}_6]^{6-}$ consisting of alternating LnS_6 distorted triangular prisms and InS_6 octahedra in the sequence of *oct tp oct tp*. $\text{Ba}_2\text{LnGaS}_5$ (Ln = Pr, Nd) crystallizes in the centrosymmetric space group $I4/mcm$. They possess a three-dimensional framework structure built from $(\text{Ba}/\text{Nd})\text{S}_8$ bicapped trigonal prisms and GaS_4 tetrahedra. The Ba atoms reside in the channels in the framework. $\text{Ba}_2\text{NdGaS}_5$ exhibits paramagnetic behavior, obeying the Curie–Weiss law down to 2 K. In addition, the band gap of $\text{Ba}_2\text{NdGaS}_5$ is 2.12(2) eV according to the diffuse-reflectance measurement. These compounds may have potential use as magnetic semiconductors or optical filters.

ASSOCIATED CONTENT

Supporting Information

Crystallographic file in CIF format for $\text{Ba}_3\text{LnInS}_6$ (Ln = Pr, Sm, Gd, Yb) and $\text{Ba}_2\text{LnGaS}_5$ (Ln = Pr, Nd). This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Notes

The authors declare no competing financial interest.

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