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The Crystal Structure of $Cs(B_9C_2H_{11})_2Co^1$

BY ALLAN ZALKIN, TED E. HOPKINS, AND DAVID H. TEMPLETON

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The crystal structure of $(B_{9}C_{2}H_{11})_{2}C_{0}C_{5}$ has been determined from an X-ray diffraction study on a single-crystal specimen. The monoclinic cell, space group C2/c, with a = 11.698, b = 7.381, c = 21.207 A, and $\beta = 101.73^{\circ}$, contains four formula units. The structure consists of a cesium cation and a $(B_9C_2H_{11})_2C_0$ anion. The latter ion consists of two icosahedra sharing one common apex, which is the cobalt atom on a center of symmetry. The carbon atoms are disordered among the pentagon of atoms adjacent to the cobalt. All 11 hydrogens, one for each boron and carbon atom, were located as well.

Introduction

The salt $Cs(B_9C_2H_{11})_2Co$ was synthesized in the laboratory of Professor M. F. Hawthorne² at the University of California at Riverside. A crystal structure study was undertaken to confirm the proposed structure of two icosahedra joined at the cobalt apex; this we proved to be true. The synthesizers of this material had hoped they had produced a compound in which the carbon atoms were not adjacent to each other; unfortunately, we found the carbon atoms to be disordered and therefore could neither affirm nor deny this hypothesis.

Experimental Section

A light brown parallelepiped-shaped crystal of dimensions $\sim 0.15 \times 0.15 \times 0.30$ mm was selected and glued to the end of a Pyrex fiber in air.

A General Electric XRD-5 X-ray diffraction apparatus equipped with a molybdenum X-ray tube (Mo $K\alpha_1$, $\lambda 0.70926$ A), a scintillation counter, a pulse height discriminator, and a quarter-circle Eulerian-cradle type of goniostat was used. The X-ray tube was operated at 50 kv and 20 ma; a Zr filter was used on the receiving slit. The crystal was oriented such that the b axis was parallel to the ϕ axis of the instrument. The cell dimensions were obtained from carefully measured 2θ (θ is the Bragg angle) values of the 0k0, h00, and 00l reflections whose α_1 and α_2 components were resolved. The β angle was measured directly from the angle between the h00 and 00l set of reflections.

A total of 2051 independent intensities were measured, of which 144 were recorded as zero. A stationary-crystal, stationary-counter technique using a 10-sec count for every reflection was used. The maximum 2θ angle was $55^{\circ} [(\sin \theta)/\lambda = 0.651]$. The maximum observed raw intensity was 46,900 cps for the 200 reflection. Background was plotted as a function of 2θ and these values were used for most of the intensities; in the cases where background was seriously affected by streaking, individual backgrounds were measured. The absorption parameter is 30 cm^{-1} , which we estimate causes the absorption factor to vary by a factor 1.5 between the most extreme cases. Variations of the intensity of 020 when the ϕ setting was changed 90° ($\chi = 90^{\circ}$) were about 20%. No corrections were made for either absorption or extinction.

Our unpublished least-squares program minimizes the function $\Sigma w(|F_o| - |F_c|)^2 / \Sigma w F_o^2$, where F_o and F_c are the observed and calculated structure factors and w is the weighting factor. All of the nonzero reflections were given equal weights, and zero observed intensities were given zero weight.

Scattering factors for Cs⁺, Co, O, C, B, and H were taken from the "International Tables."3 The Cs⁺ and Co scattering factors were given $\Delta f'$ corrections, the real part of the anomalous dispersion,^{4,5} by the amounts -0.5 and +0.4 electron, respectively. The imaginary components of the anomalous dispersion, $\Delta f''$, +2.49 and +1.05 electrons for Cs⁺ and Co, were introduced into the least-squares calculations as well. The anisotropic temperature factors used have the form $\exp(-\beta_{11}h^2 - \beta_{22}k^2 \beta_{33}l^2 - 2\beta_{12}hk - 2\beta_{13}hl - 2\beta_{23}kl$). B_{ij} values, in units of A^2 , where $B_{ij} = 4\beta_{ij}/a_i^*a_j^*$, are reported in this paper.

Results

Unit Cell and Space Group.-Four formula units of $(B_9C_2H_{11})_2CoCs$ are in the monoclinic C-centered unit cell whose dimensions are: $a = 11.698 \pm 0.005$ A, $b = 7.381 \pm 0.004$ A, $c = 21.207 \pm 0.006$ A, and $\beta =$ $101.73 \pm 0.05^{\circ}$; the errors are estimates. The calculated X-ray density is 1.692 g/cc. The density measured by flotation is 1.69 g/cc. The observed extinctions correspond to C2/c or Cc, and good agreement of observed and calculated structure factors was achieved in the centric space group, C2/c.

Determination of the Structure.-The Patterson function readily revealed the locations of the cesium and cobalt atoms. The cesium occupies the special fourfold position on the twofold axis, and the cobalt is on a center of symmetry at the origin. A leastsquares refinement on these two atoms alone using anisotropic temperature factors gave an R factor of 0.25, where $R = \Sigma ||F_0| - |F_c|| / \Sigma |F_o|$. A difference function phased by these two heavy atoms revealed the positions of the 11 boron and carbon atoms as the 11 largest peaks. Three cycles of full-matrix leastsquares refinement of the parameters of the havy atoms with anisotropic temperature factors and the other 11 atoms as boron with isotropic temperature factors resulted in R = 0.069.

In attempting to locate the two carbon atoms, we used the two criteria of temperature factors and interatomic distances. Figure 1 shows the framework of the $(B_9C_2H_{11})_2C_0$ group. When all of these atoms with the exception of Co were refined as borons, the isotropic temperature factors of the pentagon of atoms adjacent to the cobalt, *i.e.*, atoms 2-6, ranged from 2.46 to 2.71 A^2 ; those of the remaining atoms ranged from 2.96 to 3.24 A². A carbon atom refined with boron (3) J. A. Ibers in "International Tables for X-Ray Crystallography," Vol. III, The Kynoch Press, Birmingham, England, 1962, p 202.

⁽¹⁾ Work done under the auspices of the U.S. Atomic Energy Commission. (2) M. F. Hawthorne and T. D. Andrews, Chem. Commun., 19, 443 (1965).

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		Atomic	2 ^a and Theri	MAL PARAMI	eters ^b in (]	$B_9C_2H_{11})_2C_0$	Cs		
Atom	x	У	z	B_{11}	B_{22}	B_{33}	B_{12}	B_{13}	B_{23}
Cs	O^c	0.4060	$1/4^{\circ}$	7.46	3.05	4.30	0°	-1,11	0°
Co(1)	0°	. Oc	0°	1.75	2.30	2.13	-0.17	0.35	-0.13
$BC(2)^d$	0.1224	-0.1774	0.0476	2.5	3.3	2 , 9	0.3	0.3	0.0
$BC(3)^d$	0.1415	0.0369	0.0743	2.5	3.3	3.0	-0.4	-0.1	0.0
B(4)	0.0121	0.1172	0.0904	3.0	2.7	2.2	-0.2	0.4	-0.2
$BC(5)^d$	0.4146	0.4396	0.0744	2.9	2.8	2.6	-0.2	1.0	0.0
$BC(6)^d$	0.4823	0.2579	0.0470	2.8	2.7	2.7	-0.3	0.5	0.2
B(7)	0.1959	-0.1381	0.1256	2 . 4	4.0	3.3	-0.1	-0.1	0.5
B(8)	0.1267	0,4838	0.1538	3.5	3.5	2.5	-0.8	-0.1	-0.3
B(9)	0.4810	0.4845	0.1537	3.4	3.2	2 . 4	-0.4	0.8	-0.2
$\mathbf{B}(10)$	0.4613	0.2594	0.1260	2.8	3.3	2.5	-0.6	0.4	0.2
B(11)	0.0949	-0.3173	0.1087	3.1	3.0	3.4	0.0	0.2	0.3
B(12)	0.0962	-0.1757	0.1761	3.3	3.3	2.8	-0.2	0.0	0.2
H(2)	0.173	-0.233	0.015	2 . $2^{ m e}$					
H(3)	0.197	0.111	0.062	2.5					
H(4)	0.490	-0.262	0.085	1.9					
H(5)	0.319	0.456	0.061	3.0					
H(6)	0.436	0.152	0.019	2.0					
H(7)	0.281	-0.165	0.139	2.1					
H(8)	0.177	0.139	0.186	0.6					
H(9)	0.432	-0.458	0.191	0.4					
H(10)	0.397	0.170	0.142	2.6					
H(11)	0.128	-0.455	0.113	1.1					
H(12)	0.128	-0.227	0.221	0.9					

TABLE II

^{*a*} Estimated standard deviations on the atomic positions for Co, B or BC, and H are ± 0.0005 , ± 0.005 , and ± 0.05 A, respectively. The thermal parameters have the units A². The estimated standard deviations on the thermal parameters of Co or Cs, B or BC, and H are ± 0.04 , ± 0.2 , and ± 1.4 A², respectively. ^{*a*} Parameter fixed by the space group symmetry. ^{*d*} Half boron and half carbon. ^{*e*} Isotropic B in the case for hydrogen.



(B₉C₂H₁₁)₂Co⁻

Figure 1.—Framework of the $(B_9C_2H_{11})_2C_0^-$ ion. Hydrogen atoms, which are on each of the atoms with the exception of cobalt, are not shown.

TABLE III

Hydrogen	Distances	(A) in the $(B_9C_2H_{11})_2$	$Co^{-} Ion^{a}$
H(2)	1.08	H(8)	1.05
H(3)	0.93	H(9)	1.15
H(4)	0.93	H(10)	1.11
H(5)	1.10	H(11)	1.09
H(6)	1.06	H(12)	1.03
H(7)	1.00		

^a Estimated standard deviations of these distances are ± 0.06 A. These distances are shorter than expected for the internuclear B-H distance; the X-ray method gives values for such bonds which are systematically shortened by about 0.1 A. See: W. Smith and W. N. Lipscomb, J. Chem. Phys., 43, 1060 (1965); W. N. Lipscomb, *ibid.*, 22, 985 (1954). The standard deviations given by least squares, and stated in Table III, do not include this systematic effect.

scattering factors is expected to receive an abnormally low thermal parameter. No two atoms stand out as carbon atoms on this basis, but the above results are consistent with the carbon atoms being distributed in the pentagon adjacent to the cobalt. A better guide to locating the carbons are the interatomic distances. Figure 2 shows the icosahedra spread out with all of the interatomic distances in the icosahedra. In our previous study on $B_9C_2H_{11}Re(CO)_3Cs,^6$ which has a boron–carbon network quite similar to this structure, the bond distances were found to be: C–C, 1.61 A; C–B, 1.72 A; B–B 1.78 A. In Figure 2, the top half of the icosahedra, atoms 7–12, appear to have standard B–B distances. In the lower pentagon, atoms 2–6, the distances are all more or less B–C-type distances.

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Angles (in Degrees) in $Co(B_9C_2H_{11})_2$ a								
BC(2)-Co-BC(3)	48.3	Co-BC(5)-B(4)	65.5	B(4)-B(9)-BC(5)	59.2			
BC(3)-Co-B(4)	49.2	B(4)-BC(5)-B(9)	61.1	BC(5)-B(9)-B(10)	59.6			
B(4)-Co-BC(5)	49.0	B(9)-BC(5)-B(10)	60.9	B(10)-B(9)-B(12)	60.2			
BC(5)-Co-BC(6)	48.8	B(10)-BC(5)-BC(6)	60.6	B(12)-B(9)-B(8)	60.0			
BC(6)-Co-B(2)	48.8	BC(6)-BC(5)-Co	65.5	B(8)-B(9)-B(4)	60.2			
Co-BC(2)-BC(6)	66.5	Co-BC(6)-BC(2)	64.8	BC(5)-B(10)-BC(6)	59.2			
BC(6)-BC(2)-B(11)	61.1	BC(2)-BC(6)-B(11)	60.4	BC(6)-B(10)-B(11)	59.7			
B(11)-BC(2)-B(7)	61.1	B(11)-BC(6)-B(10)	60.9	B(11)-B(10)-B(12)	59.8			
B(7)-BC(2)-BC(3)	60.9	B(10)-BC(6)-BC(5)	60.2	B(12)-B(10)-B(9)	60.5			
BC(3)-BC(2)-Co	66.1	BC(5)-BC(6)-Co	65.7	B(9)-B(10)-BC(5)	59.4			
Co-BC(3)-BC(2)	65.6	BC(2)-B(7)-BC(3)	58.3	BC(2)-B(11)-B(7)	59.1			
BC(2)-BC(3)-B(7)	60.9	BC(3)-B(7)-B(8)	59.6	B(7)-B(11)-B(12)	59.8			
B(7)-BC(3)-B(8)	61.3	B(8)-B(7)-B(12)	60.6	B(12)-B(11)-B(10)	60.4			
B(8)-BC(3)-B(4)	61.6	B(12)-B(7)-B(11)	60.3	B(10)-B(11)-BC(6)	59.4			
B(4)-BC(3)-Co	66.1	B(11)-B(7)-BC(2)	59.7	BC(6)-B(11)-BC(2)	58.5			
Co-B(4)-BC(3)	64.7	BC(3)-B(8)-B(4)	59.0	B(7)-B(12)-B(8)	59.8			
BC(3)-B(4)-B(8)	59.4	B(4)-B(8)-B(9)	59.6	B(8)-B(12)-B(9)	59.6			
B(8)-B(4)-B(9)	60.2	B(9)-B(8)-B(12)	60.4	B(9)-B(12)-B(10)	59.3			
B(9)-B(4)-B(5)	59.7	B(12)-B(8)-B(7)	59.6	B(10)-B(12)-B(11)	59.9			
B(5)-B(4)-Co	65.5	B(7)-B(8)-BC(3)	59.2	B(11)-B(12)-B(7)	59.9			

TABLE IV

^a Estimated standard deviations for all angles are $\pm 0.3^{\circ}$.



Figure 2.—Interatomic distances in the $(B_0C_2H_{11})_2C_0$ icosahedra. Multiple circles indicate the disordered positions containing half boron-half carbon atoms; single circles represent borons. The estimated standard deviations on these bond lengths are ± 0.01 A.

These results as well as the temperature factors indicate the carbons to be present in the lower pentagon in a disordered manner. Since atom 4 had distances to atoms 8 and 9 which are B-B type, we finally chose a model in which atoms 2, 3, 5, and 6 are half borons and half carbons and atoms 4, 7, 8, 9, 10, 11, and 12 are borons. Admittedly this is not unique but is one of the many ways of "playing the game." Various models using different ratios of B and C in the disordered positions were tried; temperature factors were affected, but no significant changes in the distances in the icosahedral structure resulted.

A difference function was calculated, and it indicated the 11 hydrogen atoms among the top 17 peaks. Six of these peaks were apparent anisotropies near Cs, Co, and B(8).

The final least squares, performed with all 24 atoms of the structure included, resulted in R = 0.042. The hydrogen atoms were refined with isotropic temperature factors, and all of the other atoms were refined with anisotropic temperature factors. The real and imaginary components of the anomalous dispersion correction were applied to the Cs and Co atoms. Atoms 2, 3, 5, and 6 in the icosahedron were treated as half boron and half carbon atoms. The B_{12} and B_{23} temperature factors of Cs were restrained to zero as required by the symmetry of the special position. The last shift in the parameter in each case was less than 0.001 of the estimated standard deviation of that parameter. The list of the observed and calculated structure factors is shown in Table I. The final set of parameters is shown in Table II. The sets of positions are: Cs in 4(e): $\pm (0, y, \frac{1}{4}; \frac{1}{2}, \frac{1}{2} + y, \frac{1}{4});$ Co in 4(a): $(0, 0, 0; 0, 0, \frac{1}{2}; \frac{1}{2}, \frac{1}{2}, 0; \frac{1}{2}, \frac{1}{2}, \frac{1}{2});$ B and C in 8(f): $\pm (x, y, z; x, \overline{y}, \frac{1}{2} + z; \frac{1}{2} + x, \frac{1}{2} + y,$ $z; \frac{1}{2} + x, \frac{1}{2} - y, \frac{1}{2} + z).$

The anisotropic temperature factor refinements were carried out primarily for the purpose of looking for hydrogen atoms. It has been our experience that it is often necessary to remove the anisotropies of the bonding atom in order to see the hydrogens in the difference Fouriers. With the exception of Cs, all of the temperature factors appear to be fairly isotropic. We offer no explanation for the sizable anisotropy of Cs.

Description of Structure.—The structure consists of a Cs^+ cation and a $(B_9C_2H_{11})_2Co^-$ anion. The cesium ion

is surrounded by ten hydrogen atoms at distances from 3.00 to 3.34 A; the next nearest neighbors to the cesium atom are two boron atoms, B(10), at 3.67 A.

The framework of the $(B_{9}C_{2}H_{11})_{2}C_{0}^{-}$ group is shown in Figure 1. The structure consists of two icosahedra in a staggered configuration joined at a common apex, which is the cobalt atom. The two carbon atoms are shown in disordered positions in the pentagon of atoms adjacent to the cobalt, occupying mainly positions 2, 3, 5, and 6. The interatomic distances in the framework are shown in Figure 2.

Each atom with the exception of cobalt in the framework shown in Figure 1 has bonded to it a hydrogen atom. The hydrogen bond lengths are shown in Table III; the differences among them are not experimentally significant. The above icosahedral structure is very similar to that found in $(B_9C_2H_{11})Fe(C_5H_5)^7$ and in $(B_9C_2H_{11})-Re(CO)_{3,6}$ where there is no disorder of the carbon atoms. Because of the disorder in $(B_9C_2H_{11})_2CoCs$, we cannot say if the carbon atoms are adjacent or not in each molecule, nor indeed if the crystal contains a single isomer or more than one isomer.

Table IV shows a list of angles in the $(B_9C_2H_{11})_2C_0$ -group.

Acknowledgment.—We wish to thank Professor M. F. Hawthorne and Dr. T. D. Andrews for providing us with the interest and the crystals that made this work possible.

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Notes

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The Preparation and Some Reactions of Phosphoramidic Difluoride¹

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The reaction of ammonia with pyrophosphoryl tetrachloride² was previously reported to give pyrophosphoryl tetramide, $P_2O_3(NH_2)_4$. In this reaction the preferential attack by ammonia on the P–Cl bond rather than the P–O–P bond was noted

$$\begin{array}{ccc} O & O & O \\ \parallel & \parallel \\ ClPOPC1 + 8NH_3 \longrightarrow H_2NPOPNH_2 + 4NH_4C1 \\ \mid & \parallel \\ Cl & Cl & H_2N & NH_2 \end{array}$$

The recently published facile preparation of pyrophosphoryl tetrafluoride³ prompted us to investigate its reaction with ammonia to determine whether the substitution of the phosphorus-bonded fluorine atoms by amido groups or the heterolytic cleavage of the P–O–P bond by ammonia occurred predominantly.

When pyrophosphoryl tetrafluoride was treated with ammonia, ammonium difluorophosphate^{4,5} and phosphoramidic difluoride were obtained



⁽¹⁾ Presented in part at the 148th National Meeting of the American Chemical Society, Chicago, Ill., Sept 1964; Abstracts, p 3K.

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The novel phosphoramidic diffuoride was characterized by infrared and ³¹P and ¹⁹F nuclear magnetic resonance spectrometry. The reaction of ammonia with pyrophosphoryl tetrafluoride is indeed different from that with pyrophosphoryl tetrachloride. The highly electronegative fluorine atoms in $P_2O_3F_4$ probably withdraw electrons toward themselves, thus weakening the POP bond. Since phosphoryl trifluoride⁶ in ethyl ether reacted rapidly with ammonia, even at -110° , to form phosphoryl triamide, phosphoramidic diffuoride would, thus, be expected to react further with excess of ammonia.

The reactions of phosphoramidic difluoride with phosphorus pentachloride and with triphenylphosphorus dichloride were investigated.

Phosphoramidic difluoride reacted with an equivalent amount of phosphorus pentachloride to yield trichlorophosphazophosphoryl difluoride

$$\begin{array}{c} O \\ \parallel \\ H_{2}NPF + PCl_{5} \longrightarrow Cl_{3}P = NPF \\ \parallel \\ F \end{array} \begin{array}{c} O \\ \parallel \\ PF \end{array} + 2HCl_{5} \end{array}$$

Its infrared spectrum is quite similar to that of Cl_3P == NPOCl₂⁷ (mp 35.5°; bp 102° (1 mm)), except for the additional P-F stretching frequency in Cl_3P ==N-POF₂.

Trichlorophosphazophosphoryl difluoride decomposed in an atmosphere of nitrogen at 200° to liberate gaseous products which were mainly OPF_2Cl and $OPCl_8$. The residue was an opaque elastomeric material with an empirical formula of $(NPFCl)_n$. It was not determined if the elastomer was a mixture of

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