being close-packed with between eight and twelve cations associated about it, then the entropy change is of the right order of magnitude.

The slightly negative entropy change in the case of molybdate and the large positive entropy change in the case of tungstate are more difficult to explain. One characteristic of the chemistries of Mo(VI) and W(VI) compounds which distinguishes them from Cr(VI) compounds is the increased tendency toward sixfold coordination. It is probably this propensity for sixfold coordination which leads to the extremely complex polyoxy anion structures observed in the aqueous chemistry of molybdenum and tungsten. Lindqvist[5] has demonstrated that solid $Na_2W_2O_7$ and $Na_2M_2O_7$ are isomorphous and do not contain discrete $M_2O_7{}^{2-}$ ions but are polymeric in nature, consisting of chains of MO_6 octahedra sharing corners, with additional MO_4 tetrahedra sharing corners with two adjacent octahedra. This suggests that in solution ditungstates and dimolybdates form polymeric anions of a similar configuration containing both four- and sixfold coordinated species. If solvation with the solvent occurs through cation interaction with oxygens primarily bonded to one metal species, then the effective number of sites for solvation is decreased from the

number available in the orthooxy anion. This decrease in solvation would cause an increase in entropy offsetting the decrease caused by formation of the oxide ion. The enthalpy of reaction for these three systems increases in the order Cr < Mo < W, which is the order which would be expected from comparison of heats of formation of the ortho salts. Whatever the explanation for the large entropy differences observed, it is clear that the basic properties of the ortho salts must be described in terms of the solvent in which they are placed.

Further studies on these systems are indicated. At present, work is underway to determine the effect of altering the cation ratio of the solvent. This should give information regarding the solvation of the oxide ion. If extensive polymers are formed in the case of dimolybdates and ditungstates, one could expect a significant difference in solubility of the ditungstates and dimolybdates as compared to the dichromates. Such studies are being initiated.

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Structure of $IrIO_2(CO)(P(C_6H_5)_3)_2 \cdot CH_2Cl_2$, the Oxygen Adduct of a Synthetic Irreversible Molecular Oxygen Carrier

BY J. A. McGINNETY, ROBERT J. DOEDENS, AND JAMES A. IBERS

The system $IrX(CO)(P(C_6H_5)_3)_2$ in benzene solution adds molecular oxygen reversibly when X is Cl and irreversibly when X is I. The crystal and molecular structure of $IrIO_2(CO)(P(C_6H_5)_3)_2 \cdot CH_2Cl_2$ has been determined from three-dimensional X-ray data collected by counter methods. The structure has been refined by least-squares techniques to a final R factor on F of 5.9% for 1255 reflections above background. The material crystallizes in space group $C_{2h}{}^5$-$P2_1/n$ of the monoclinic system, with four molecules in a cell of dimensions $a = 10.69$, $b = 22.86$, $c = 15.00$ A, $\beta = 93.05°$. The observed density is 1.80 g/cm³, the same as that calculated for the above cell. The crystal structure consists of the packing of the discrete monomeric units $IrIO_2(CO)(P(C_6H_5)_3)_2$ and CH_2Cl_2. The molecular structure of $IrIO_2(CO)(P(C_6H_5)_3)_2$ closely resembles that reported earlier for $IrClO_2(CO)(P(C_6H_5)_3)_2$; the iridium atom, the two oxygen atoms, the carbonyl group, and the iodine atom are situated in the equatorial plane of a trigonal bipyramid, with phosphorus atoms above and below this plane. The two oxygen atoms are equidistant from the iridium atom, with an average Ir–O distance of 2.06 A. The O–O distance is 1.509 ± 0.026 A in this irreversibly oxygenated iodo compound, whereas it is 1.30 ± 0.03 A in the reversibly oxygenated chloro compound. This very significant change in the O–O bond length is consistent with the π-bonding scheme that Griffith applied to the bonding of molecular oxygen to the heme group in oxyhemoglobin.

Introduction

Previously[1,2] we reported the crystal and molecular structure of the oxygen adduct of the synthetic reversible molecular oxygen carrier $IrCl(CO)P(C_6H_5)_3)_2$.[3] As part of a continuing study of the structures of oxygenated adducts of various synthetic oxygen carriers,

we have investigated the properties of the system $IrX(CO)(P(C_6H_5)_3)_2$, where X = Br, I, SCN, or N_3. This series of compounds, with the exception of the thiocyanato compound, reacts in solution with gaseous oxygen to form oxygenated species. For the halogen compounds the rate of reaction has been found[4] to increase in the order Cl < Br < I. We have found[5]

(1) J. A. Ibers and S. J. La Placa, *Science*, **145**, 920 (1964).
(2) S. J. La Placa and J. A. Ibers, *J. Am. Chem. Soc.*, **87**, 2581 (1965).
(3) L. Vaska, *Science*, **140**, 809 (1963).
(4) P. B. Chock and J. Halpern, *J. Am. Chem. Soc.*, **88**, 3511 (1966).

that whereas the addition is reversible for the chloro compound, it is irreversible in solution for the iodo compound. The question of whether this difference in reversibility in solution manifests itself in structural differences in the oxygen adducts in the solid state is one that has been answered by the determination of the crystal and molecular structure of the oxygen adduct of the iodo compound and comparison of it with the structure of the oxygen adduct of the chloro compound. A preliminary account of this work has been published,[5] and in this paper we present our results in full.

Preparation of the Sample

IrCl(CO)(P(C₆H₅)₃)₂.—IrCl₃·3H₂O (1.0 g) was dissolved in warm diethylene glycol monomethyl ether (30 ml). Triphenylphosphine (5.0 g) was dissolved in the same solvent (20 ml) by warming and was added, with stirring, to the iridium chloride solution; a yellow precipitate formed. The mixture was refluxed until a clear, deep yellow solution was obtained. (If necessary, more solvent was added after 0.5 hr.) The mixture was then refluxed a further 2 hr. Solvent was then distilled off until about one-third of the original volume remained. The mixture was cooled, and methanol (100 ml) added. The product was filtered off and recrystallized from benzene under nitrogen using a Soxhlet apparatus. The over-all yield is 70–80%.

IrI(CO)(P(C₆H₅)₃)₂.—IrCl(CO)(P(C₆H₅)₃)₂ (1.0 g) was dissolved in refluxing benzene (30 ml) under nitrogen. To this was added a solution of lithium iodide (14 g) in hot ethanol (25 ml). A two-layer mixture was formed. More ethanol was added until a fully miscible, clear, refluxing solution was formed. The mixture was refluxed for 15 min and then methanol (100 ml) was added. After cooling, the product (0.8 g) was filtered off.

Exactly analogous procedures are used to prepare the series IrX(CO)(P(C₆H₅)₃)₂, where X = Br, SCN, or N₃ (from the chloro compound).

IrIO₂(CO)(P(C₆H₅)₃)₂.—A benzene solution of IrI(CO)(P(C₆H₅)₃)₂ reacts quantitatively with oxygen from the air when left exposed for 1 hr. Evaporation yields a brown powder which analyzes correctly for the O₂ adduct. *Anal.* Calcd for IrIO₂(CO)P(C₆H₅)₃)₂: Ir, 21.3; C, 49.2; H, 3.3; P, 6.9; O, 5.3; I, 14.0. Found: Ir, 23.6; C, 49.1; H, 3.7; P, 6.6; O, 4.7; I, 14.5.

Suitable crystals of this material could not be obtained from benzene solution. Those good crystals which formed in benzene disintegrated when not in contact with the mother liquor. It is possible that the crystals contain benzene of crystallization and that, once removed from solution, this benzene sets up a sufficient vapor pressure to destroy the crystal.

The powder obtained from the benzene solution was dissolved in dichloromethane; slow evaporation of the solution yielded good, reasonably stable crystals of the 1:1 aggregate of IrIO₂(CO)(P(C₆H₅)₃)₂ and CH₂Cl₂.

Anal. Calcd for IrIO₂(CO)(P(C₆H₅)₃)₂·CH₂Cl₂: I, 13.3; Cl, 7.4. Found: I, 13.0; Cl, 7.1. Additional verification of this formulation is provided by the X-ray molecular weight and, of course, by the structure determination itself.

The oxygenated compound was not prepared in dichloromethane solution from the parent compound because a dichloromethane adduct of low solubility is formed, and oxygenation of this proceeds very slowly.

Collection and Reduction of the X-Ray Data

Crystals, prepared as described above, are dark brown, thick needles. On the basis of optical examination and preliminary X-ray photography, the crystals were assigned to the monoclinic system. The lattice constants at 27°, obtained by least-squares refinement of the setting angles of 17 reflections which had been carefully centered on a Picker four-circle automatic X-ray diffractometer,[6] are: $a = 10.693 \pm 0.008$, $b = 22.864 \pm 0.015$, $c = 14.997 \pm 0.011$ A, $\beta = 93.05 \pm 0.02°$ (λ(Cu Kα_1) 1.54056 A). The systematic absences that were observed on Weissenberg and precession photographs are $h0l$ for $h + l$ odd and $0k0$ for k odd. These absences are consistent with the space group C_{2h}^5-P2₁/n. On this basis no piezoelectric effect is expected, and none could be observed. For four formula units in the cell the calculated density is 1.80 g/cm³; the observed density is also 1.80 g/cm³. No crystallographic symmetry conditions need be imposed on the molecules.

The procedures used in data collection and processing parallel those described previously.[6,7] The data were collected from a nearly spherical crystal of average diameter 0.204 mm. (Actually a more accurate description of the crystal shape is an oblate spheroid of maximum diameter 0.208 mm and minimum diameter 0.187 mm.) The linear absorption coefficient of the compound for Cu Kα radiation is 162.6 cm⁻¹, whereas it is 52.3 cm⁻¹ for Mo Kα. This difference is not especially large, and, for the reasons detailed earlier, Cu Kα radiation was chosen. The diffracted beams were filtered through 1.0-mil Ni foil.

The mosaicity of the crystal was checked by a narrow-source, open-counter scan[8] through several strong reflections and the average half-width was 0.07°, a satisfyingly low value.

The crystal was initially aligned along the 30Ī direction, but this orientation was changed before commencement of data collection by arbitrarily resetting the arcs of the goniometer head to zero. The effects of multiple scattering should therefore be minimal.[9] Data were collected by the θ–2θ scan technique at a takeoff angle of 1.8°. At this angle the peak intensity of a typical reflection was about 80% of its maximum value as a function of takeoff angle. A symmetric scan range of 2.2° in 2θ was used for all reflections. Stationary-

(5) J. A. McGinnety, R. J. Doedens, and J. A. Ibers, *Science*, **155**, 709 (1967).

(6) P. W. R. Corfield, R. J. Doedens, and J. A. Ibers, *Inorg. Chem.*, **6**, 197 (1967).

(7) R. J. Doedens and J. A. Ibers, *ibid.*, **6**, 204 (1967).

(8) T. C. Furnas, "Single Crystal Orienter Instruction Manual," General Electric Co., Milwaukee, Wis., 1957.

(9) W. H. Zachariasen, *Acta Cryst.*, **18**, 705 (1965).

crystal, stationary-counter background counts of 10 sec were taken at each end of the scan. The scan rate was 2°/min. A counter aperture 4.0 mm high by 6.0 mm wide was placed 21 cm from the crystal. The pulse-height analyzer was set to accept approximately a 95% window centered on the Cu $K\alpha_1$ peak.

Intensity data were collected for all four forms of this monoclinic crystal (hkl, $h\bar{k}l$, $\bar{h}kl$, and $\bar{h}\bar{k}l$) out to a 2θ value for Cu $K\alpha_1$ of 70°. At higher scattering angles there were very few reflections above background. A total of 5987 intensities were measured.

As a general check on electronic and crystal stability, the intensities of four standard reflections were measured periodically during the collection of the intensity data. Over the entire period of data collection, these standard reflections slowly decreased in intensity, the total decrease being from 17 to 35% of the original intensity. Further checks at the completion of the run indicated unambiguously that these intensity changes were the result of crystal, rather than electronic, instability. Unfortunately, the decrease in intensity differed from standard to standard, and hence no simple correction to bring all of the intensity data to a reliable, common scale is possible. All that could be done was to take the average decrease in the standards as a function of time as the appropriate correction factor. Since each of the four equivalent forms of a given reflection typically is collected at a different time during the experiment and since the weighting scheme depends upon agreement among these forms, the averaged data and their weights should represent in a reasonable way the scattering from the crystal at the midpoint of the experiment. Attempts to find the nature of this decomposition of the crystal were not successful. The decrease in intensity is not due to loss of either oxygen or methylene chloride, since the contributions of these moieties to the structure factors of the standard reflections are in no way related to the relative loss in intensity that was observed. Thus the following are the percentage loss in intensity and the percentage contributions to F_o of CH_2Cl_2 and O_2 for each of the standard reflections: 023, 21.5, −8.0, 7.3; $\bar{1}03$, 16.6, 14.2, 20.4; $\bar{4}10$, 35.0, −6.2, 0.8; 006, 28.2, 4.0, 1.9.

The intensity data were corrected for background, and then for Lorentz-polarization effects and for absorption. The absorption correction was based on a spherical model with $\mu R = 1.66$, where μ is the linear absorption coefficient and R is the radius of the sphere. The individual values of F_o^2 from given equivalent sets were then averaged to yield 1573 independent F_o^2 values. Standard deviations on F^2, $\sigma(F^2)$, were estimated both from the range of equivalent forms and from the individual standard deviations and the larger of the two estimates was chosen.[6] For these individual standard deviations a value of $p = 0.05$[6] was used. Of the 1573 reflections, the range estimate was larger than the individual estimate for 865 reflections. The predicted value of the weighted R factor on F^2, $R_2(F^2)$, is 10.0%. Of the 1573 reflections, 318 had $F^2 < \sigma(F^2)$ and we describe these as being below background.

Only the 1255 reflections with $F^2 \geq \sigma(F^2)$ were used in refinement.

Solution and Refinement of the Structure

Initial values of the atomic coordinates of the iridium atom were obtained from a three-dimensional Patterson function.[10] These parameters were refined by usual least-squares procedures, with minimization of the function $\Sigma w(F_o - F_c)^2$, where F_o and F_c are the observed and calculated structure amplitudes and where the weights w are taken as $4F_o^2/\sigma^2(F_o^2)$. The usual tabulation of atomic scattering factors for P, Cl, O, C, and H were used;[11] the values obtained by Cromer and Waber[12] were used for Ir and I. The effects of anomalous dispersion were included in F_c;[13] the values of f' and f'' for Ir, I, P, and Cl were those given by Cromer.[14] This refinement of the scale factor, positional parameters, and an isotropic thermal parameter for the Ir atom led to a conventional R factor, $R_1 = (\Sigma||F_o| - |F_c||/\Sigma|F_o|)$, of 49% and a weighted R factor, $R_2 = (\Sigma w(F_o - F_c)^2/\Sigma w F_o^2)^{1/2}$, of 49%. The first difference Fourier map gave a clear indication of disorder in the position of the iodine atom. A similar disorder had been found in $IrClO_2(CO)(P(C_6H_5)_3)_2$.[1,2] In subsequent calculations the coordinates of both "half-iodine" atoms were included in the calculations. Succeeding least-squares and difference Fourier methods yielded, on the third difference Fourier map, the coordinates of all of the remaining nonhydrogen atoms, with the exception of those of the carbonyl group. Subsequent least-squares calculations were effected with the program NUGLS[6] in which the phenyl groups were refined as rigid groups (D_{6h} symmetry, C–C = 1.392 A). A difference Fourier map was computed at this point with no contributions from either half-iodine atom but with contributions from the rest of the atoms that had been located. This map showed two peaks of very nearly the same height (16.6 and 16.1 e/A³) at the positions corresponding to the half-iodine atoms. This we take as justification for the subsequent procedure of assigning an occupancy factor of 0.5 to each of the iodine positions. This map also provided an indication of anisotropy in the Ir and P motions. In a subsequent calculation these atoms were allowed to vibrate anisotropically and the value of R_2 decreased from 11.9 to 9.4%. A difference Fourier map at this stage showed the presence of two "half-carbonyl" groups slightly separated from the two half-iodine atoms. Owing to the slight separation, the atomic coordinates of the half-carbonyl groups were fixed at the positions indicated by the difference Fourier map (Ir–C ≈ 1.85 A; C–O ≈ 1.05 A; ∢Ir–C–O ≈ 174°) and these positions were not refined. A subsequent calculation, in which the isotropic thermal parameters

(10) In addition to various local programs for the CDC 3400, the programs used in this work were local modifications of Zalkin's FORDAP Fourier program, the Busing-Levy ORFFE error function program, and Johnson's ORTEP thermal ellipsoid plotting program.

(11) J. A. Ibers, "International Tables for X-Ray Crystallography," Vol. 3, The Kynoch Press, Birmingham, England, 1962.

(12) D. T. Cromer and J. T. Waber, *Acta Cryst.*, **18**, 104 (1965).

(13) J. A. Ibers and W. C. Hamilton, *ibid.*, **17**, 781 (1964).

(14) D. T. Cromer, *ibid.*, **18**, 17 (1965).

TABLE I

POSITIONAL, THERMAL, AND GROUP PARAMETERS FOR $IrIO_2(CO)P(C_6H_5)_3)_2 \cdot CH_2Cl_2$

	x	y	z	$\beta_{11}{}^a$ or B, Å²	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Ir	0.18567 (13)[b]	0.15625 (6)	0.25456 (9)	0.0184 (2)	0.00309 (4)	0.0075 (1)	0.00007 (8)	−0.0016 (1)	−0.00010 (6)
I''	−0.0683[f](4)	0.1482 (2)	0.2673 (3)	0.0157 (7)	0.0034 (1)	0.0090 (4)	−0.0002 (2)	0.0021 (4)	0.0002 (2)
I'	0.2442 (5)	0.1450 (2)	0.0778 (3)	0.0162 (8)	0.0025 (1)	0.0071 (3)	0.0000 (2)	−0.0002 (4)	0.0002 (2)
P₁	0.1928 (8)	0.2605 (3)	0.2426 (5)	0.0134 (13)	0.0030 (3)	0.0058 (6)	−0.0001 (5)	−0.0025 (7)	−0.0000 (3)
P₂	0.2138 (7)	0.0532 (3)	0.2707 (5)	0.0122 (12)	0.0024 (2)	0.0065 (6)	−0.0003 (4)	−0.0003 (7)	−0.0004 (3)
Cl₁	−0.3522 (17)	0.1813 (8)	0.3755 (11)	22.9 (6)					
Cl₂	−0.2437 (14)	0.2895 (10)	0.3344 (10)	18.7 (5)					
O₁	0.2260 (23)	0.1656 (9)	0.3879 (14)	0.0284 (41)	0.0027 (6)	0.0085 (17)	−0.0015 (15)	−0.0027 (21)	0.0006 (9)
O₂	0.3454 (23)	0.1664 (10)	0.3387 (16)	0.0199 (36)	0.0045 (8)	0.0147 (23)	−0.0001 (15)	−0.0015 (23)	0.0008 (11)
C	−0.3178 (43)	0.2477 (21)	0.4089 (30)	15.4 (1.5)					
CO'C	0.020	0.147	0.22	3.0 (1.3)					
CO'O	−0.075	0.147	0.20	14.3 (2.5)					
CO''C	0.165	0.147	0.13	5.6 (1.7)					
CO''O	0.155	0.147	0.06	12.8 (2.3)					

Rings		$x_c{}^c$	y_c	z_c	δ	ϵ	η
P₁	A	−0.0362 (12)	0.3256 (6)	0.1420 (8)	3.000 (19)	2.472 (14)	−0.680 (16)
P₁	B	0.4214 (13)	0.2985 (5)	0.1256 (8)	0.753 (15)	2.655 (13)	−0.752 (13)
P₁	C	0.2310 (12)	0.3245 (4)	0.4310 (8)	−1.282 (12)	2.880 (15)	1.990 (10)
P₂	L	0.4605 (12)	0.0154 (5)	0.1760 (8)	2.542 (16)	2.614 (12)	−2.530 (13)
P₂	M	−0.0104 (11)	−0.0272 (5)	0.1881 (8)	0.202 (21)	2.274 (11)	−2.563 (19)
P₂	N	0.2452 (11)	0.0062 (5)	0.4716 (8)	−2.143 (11)	2.891 (15)	1.303 (11)

[a] The form of the anisotropic thermal ellipsoid is $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}hl)]$.   [b] Numbers in parentheses here and in succeeding tables are estimated standard deviations in the least significant digits.   [c] These group parameters have been defined in previous papers.   See, for example, R. Eisenberg and J. A. Ibers, *Inorg. Chem.*, **4**, 773 (1965).

of the C and O atoms of the half-carbonyl groups and the anisotropic thermal parameters of the half-iodine were added to the variables, led to a value of $R_2 = 6.9\%$. Inclusion of the fixed contributions to F_c of the phenyl hydrogen atoms (C–H = 1.08 Å) reduced R_2 to 6.4%. In a final round of calculations the carbon atoms of the phenyl groups were assigned individually variable isotropic thermal parameters and the oxygen atoms were allowed to vibrate anisotropically. The final values of R_1 and R_2 for the 1255 reflections above background are 5.9 and 6.2%, respectively. The error in an observation of unit weight is 1.15; this is sufficiently close to the expected value to suggest that the absolute values of the weights are reasonable. The values of R_1 and R_2 based on F^2 are 7.7 and 11.8%, respectively. The latter value may be compared with that of 10.0% predicted from the agreement among the intensities of equivalent forms. The discrepancy may indicate some minor deficiencies in the model, for example, slight deviation from statistical disorder of the iodine atoms, failure to refine the half-carbonyl positions, or failure to allow all of the atoms (including the chlorine atoms of the methylene chloride) to vibrate anisotropically. Another possibility is that the crystal is deficient in methylene chloride. To check on this a least-squares calculation was carried out in which the occupancy factor of the three nonhydrogen atoms of the methylene chloride (as one variable) and the thermal parameters of these atoms were varied. The correlation coefficient between the occupancy factor and any of the thermal parameters was less than 0.6. The occupancy factor for methylene chloride converged to a value of 0.93 ± 0.02. We conclude that any possible deficiency in methylene chloride is very small.

A final difference Fourier map showed no peaks higher than 0.9 e/Å³, approximately 35% of the height of a phenyl carbon atom in this structure.

A final statistical analysis of the F_o and F_c values as a function of the indices, the scattering angle, and the magnitude of F_o showed no unusual trends and suggests that the relative weighting scheme is also a reasonable one.

In Table I we give the final values of the positional and thermal parameters, along with their standard deviations as derived from the inverse matrix. In Table II we give the fractional coordinates of the carbon atoms of the phenyl rings that can be derived from the data of Table I. In Table III we give the values in electrons of F_o and F_c (based on the parameters of Table I) for the 1255 reflections above background. The values of F_c for $F_o < \sigma(F_o)$ in no case exceeded 1.5 $\sigma(F_o)$. In Table IV we give the root-mean-square amplitudes of vibration for those atoms that were refined anisotropically. An indication of the directions of vibration is provided by Figures 1–3.
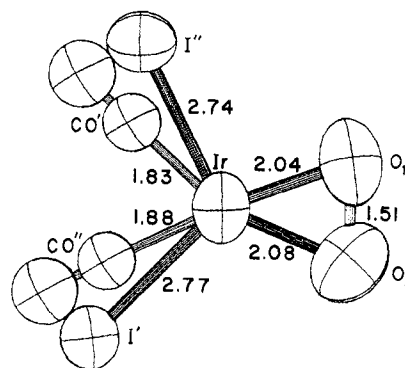
Figure 1.—A view of the disorder around iridium in the equatorial plane of the trigonal bipyramid.

Description of the Structure

The crystal structure consists of the packing of two monomeric molecular units, $IrO_2(CO)(P(C_6H_5)_3)_2$ and CH_2Cl_2. In the iridium moiety, disorder occurs between iodine and the carbonyl group; in analogous examples such disordered groups are directly superimposed, but in this case the disorder was resolved.

TABLE II

DERIVED PARAMETERS FOR GROUP CARBON ATOMS

	x	y	z	B, A^2
		Group 1 (P, A)		
1	0.057(2)	0.295(1)	0.191(1)	6.2(4)
2	0.023(2)	0.351(1)	0.217(1)	7.2(4)
3	−0.070(3)	0.382(1)	0.168(2)	8.5(4)
4	−0.129(2)	0.356(1)	0.093(1)	7.7(4)
5	−0.095(2)	0.300(1)	0.067(1)	8.5(4)
6	−0.002(3)	0.270(1)	0.116(1)	7.2(4)
		Group 2 (P, B)		
1	0.323(2)	0.282(1)	0.178(1)	6.6(4)
2	0.441(2)	0.256(1)	0.192(1)	6.6(4)
3	0.539(2)	0.273(1)	0.140(1)	7.4(4)
4	0.519(2)	0.315(1)	0.073(1)	6.6(4)
5	0.402(2)	0.341(1)	0.059(1)	7.4(4)
6	0.304(2)	0.324(1)	0.112(1)	6.6(4)
		Group 3 (P, C)		
1	0.216(2)	0.297(1)	0.349(1)	5.1(4)
2	0.119(1)	0.296(1)	0.407(2)	7.0(4)
3	0.134(2)	0.324(1)	0.489(1)	6.8(4)
4	0.245(2)	0.352(1)	0.513(1)	6.4(4)
5	0.343(1)	0.353(1)	0.455(2)	6.8(4)
6	0.328(2)	0.325(1)	0.373(1)	7.0(4)
		Group 4 (P, L)		
1	0.351(2)	0.029(1)	0.218(1)	5.4(4)
2	0.461(2)	0.060(1)	0.240(1)	6.9(4)
3	0.570(2)	0.046(1)	0.198(1)	6.7(4)
4	0.570(1)	0.002(1)	0.134(1)	6.5(4)
5	0.460(2)	−0.029(1)	0.112(1)	6.7(4)
6	0.351(2)	−0.015(1)	0.154(1)	6.9(4)
		Group 5 (P, M)		
1	0.086(3)	0.008(1)	0.225(1)	5.2(4)
2	0.052(2)	−0.043(1)	0.268(1)	6.6(4)
3	−0.044(3)	−0.078(1)	0.232(1)	6.8(4)
4	−0.106(3)	−0.062(1)	0.152(1)	7.0(4)
5	−0.073(2)	−0.011(1)	0.108(1)	6.8(4)
6	0.023(3)	0.024(1)	0.144(1)	6.6(4)
		Group 6 (P₂ N)		
1	0.238(2)	0.028(1)	0.385(1)	5.8(4)
2	0.149(2)	0.045(1)	0.445(2)	6.6(4)
3	0.157(2)	0.023(1)	0.532(1)	7.2(4)
4	0.253(2)	−0.015(1)	0.559(1)	6.6(4)
5	0.341(2)	−0.032(1)	0.499(2)	7.2(4)
6	0.333(2)	−0.011(1)	0.412(1)	6.6(4)

Figure 2.—A perspective drawing of the molecular structure of IrIO₂(CO)(P(C₆H₅)₃)₂·CH₂Cl₂. For pictorial purposes the carbonyl group and the dichloromethane molecule have artificially low thermal parameters. Only one set of the disordered carbonyl and iodine positions is shown.
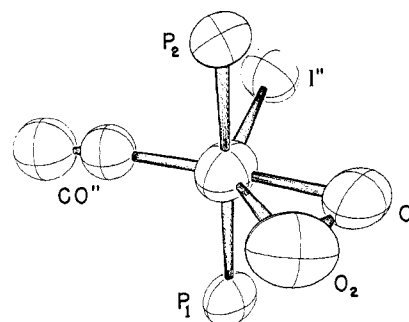
Figure 3.—A perspective drawing of the inner coordination sphere around iridium. Only one set of the disordered carbonyl and iodine positions is shown. With the exception of the carbonyl group, which has been assigned artificially low thermal parameters, the sizes of the various ellipsoids are dictated by the perspective view and by the amplitudes of vibration of the individual atoms.

The relative positions of all of the atoms in the equatorial plane of the iridium moiety are shown in Figure 1. In the asymmetric environment caused by the presence of methylene chloride, the two molecules which are disordered are best considered separately. Molecule 1 (containing I′ and CO′) has a fairly symmetrical arrangement of the carbonyl and iodine ligands relative to the oxygen group. Molecule 2 (containing I″ and CO″) has a much more unsymmetrical arrangement of ligands. There are several close contacts between methylene chloride and these ligands (see Table V), and we propose it is these steric requirements that are the primary cause of the resolving of the disordered groups. In addition to the steric repulsions, weak hydrogen bonds may be formed (e.g., between H₁ and I′); it has been reported[15] that a hydrogen bond can be

formed between an iodine ligand and a hydrogen atom on chloroform.

In Figure 2 the over-all molecular structure is displayed, and in Figure 3 the inner coordination sphere around the iridium atom is shown. (For the sake of clarity, these two diagrams do not show the disordered array; instead molecule 2 is drawn.) The inner coordination sphere is trigonal bipyramidal when O₂ is considered as a monodentate ligand. A best weighted least-squares plane was defined through the two half-iodine atoms, the two oxygen atoms, and the iridium atom. The displacements in angstroms from this plane (with their standard deviations) are: Ir, 0.0014 (14); I′, 0.0022 (38); I″, 0.0017 (43); O₁, 0.043 (20); O₂, 0.036 (23). Thus there are no significant deviations from this plane. However, the P₁–Ir–P₂ angle of 170.9 (0.3)° differs significantly from the expected value of 180° for a perfect trigonal bipyramid. The two sets of atoms I′, Ir, I″ (*i.e.*, approximately the equatorial plane) and P₁, Ir, P₂ define two planes which are in-

(15) M. F. Rettig and R. S. Drago, *J. Am. Chem. Soc.*, **88**, 2966 (1966).

Table III

Observed and Calculated Structure Amplitudes ($\times 10$) (in electrons) for $IrIO_2(CO)(P(C_6H_5)_3)_2 \cdot CH_2Cl_2$

Table IV

Root-Mean-Square Amplitudes of Vibration (Å)

Atom	1	2	3
Ir	0.275 (2)	0.286 (2)	0.343 (2)
I′	0.256 (7)	0.282 (6)	0.311 (7)
I″	0.284 (7)	0.303 (7)	0.331 (7)
P₁	0.220 (14)	0.280 (12)	0.312 (12)
P₂	0.240 (14)	0.268 (13)	0.284 (12)
O₁	0.259 (32)	0.297 (31)	0.425 (30)
O₂	0.330 (31)	0.341 (33)	0.422 (33)

clined at 91.2 (1)°. The two phosphorus atoms and the iridium atom lie in a plane nearly at right angles to the equatorial plane, and the phosphorus atoms are bent away from the two half-iodine atoms toward the oxygen atoms.

Principal intramolecular distances are given in Table V. The Ir–P distances do not differ significantly from each other and the average Ir–P distance of 2.39 Å is very close to the distance 2.37 Å found in IrCl(O₂)-(CO)(P(C₆H₅)₃)₂. The Ir–I distances differ by four standard deviations; however, the true error in the positional parameters of the half-iodine atoms may be larger than the values quoted because of the proximity of the half-carbonyl groups, the positions of which

were not refined. The average Ir–I distance of 2.75 Å is about the value expected. The Ir–O distances are very similar and the average Ir–O distance of 2.06 Å is the same as that found in IrCl(O₂)(CO)(P(C₆H₅)₃)₂. The O₁–O₂ distance of 1.509 (26) Å is significantly longer than the value of 1.30 (3) Å observed in IrCl(O₂)-(CO)(P(C₆H₅)₃)₂. The O–O distance in the iodo compound is that of a typical peroxide; the O–O distance in the chloro compound is that of a typical superoxide. Hence, these two compounds might best be named chlorosuperoxocarbonylbis(triphenylphosphine)iridium-(II) and iodoperoxocarbonylbis(triphenylphosphine)-iridium(III).

The iridium, half-iodine, phosphorus, and oxygen atoms were allowed to vibrate anisotropically. For these atoms the root-mean-square amplitudes of vibration along the principal axes of vibration are given in Table IV. While there can be some discussion as to the reality of motion derived in this way, it can be seen that these results are not completely random. None of the vibrations is extremely eccentric. For the iridium atom, the largest principal axis of vibration is in the equatorial plane, parallel to the O₁–O₂ direction. The vibrations of the two phosphorus atoms are ap-

TABLE V

SELECTED INTRA- AND INTERMOLECULAR DISTANCES

Ir–P₁	2.391 (8)	Cl₁–I'	5.070 (19)
Ir–P₂	2.386 (8)	Cl₂–I'	3.954 (15)
Ir–I'	2.767 (5)	Cl₁–I''	3.598 (18)
Ir–I''	2.738 (5)	Cl₂–I''	3.895 (15)
Ir–O₁	2.035 (20)	O₁–I'	4.688 (22)
Ir–O₂	2.082 (22)	O₂–I'	4.032 (24)
O₁–O₂	1.509 (26)	O₁–I''	3.568 (22)
Ir–CO'C	1.83	O₂–I''	4.515 (24)
Ir–CO''C	1.88	O₁–P₁	3.082 (21)
CO'C–CO'O	1.04	O₁–P₂	3.112 (21)
CO''C–CO''O	1.05	O₂–P₁	3.018 (24)
P₁–I'	3.676 (9)	O₂–P₂	3.093 (24)
P₁–I''	3.825 (9)	I'–CO'C	3.29
P₂–I'	3.603 (9)	I'–CO''C	1.18
P₂–I''	3.714 (9)	I''–CO'C	1.21
I'–CO'O	3.96	I''–CO''C	3.32
I'–CO''O	0.98	C–Cl₁	1.63 (4)
I''–CO'O	1.01	C–Cl₂	1.70 (4)
I''–CO''O	4.02	Cl₁–H₁	2.22
P₁–P₂	4.76 (11)	Cl₁–H₂	2.22
P₁–CO'C	3.19	Cl₂–H₁	2.34
P₁–CO'O	3.89	Cl₂–H₂	2.32
P₁–CO''C	3.10	H₁–I'	2.96
P₁–CO''O	3.78	H₁–I''	4.44
P₂–CO'C	3.05	H₂–I'	3.42
P₂–CO'O	3.86	H₂–I''	5.23
P₂–CO''C	3.03	H₁–CO'C	5.08
P₂–CO''O	3.84	H₁–CO''C	3.51
O₁–CO'C	3.28	H₁–CO'O	5.19
O₁–CO''C	3.91	H₁–CO''O	2.93
O₂–CO'C	3.85	H₂–CO'C	4.96
O₂–CO''C	3.61	H₂–CO''C	3.72
P₁–P₁A₁ᵃ	1.79	H₂–CO'O	4.99
P₁–P₁B₁	1.81	H₂–CO''O	2.88
P₁–P₁C₁	1.81		
P₂–P₂L₁	1.79		
P₂–P₂M₁	1.82		
P₂–P₂N₁	1.81		

ᵃ By P₁A₁, we denote the α-carbon atom on ring A, that is, the carbon atom attached to P₁.

TABLE VI

SELECTED INTERATOMIC ANGLES

I'–Ir–I''	109.6 (2)	P₁–Ir–I'	90.6 (2)
P₁–Ir–P₂	170.9 (3)	P₁–Ir–I''	96.2 (2)
O₁–Ir–O₂	43.0 (7)	P₂–Ir–I'	88.4 (2)
CO'C–Ir–CO''C	69.2	P₂–Ir–I''	92.7 (2)
O₁–Ir–I''	95.6 (7)	Ir–CO'C–CO'O	173.4
O₁–Ir–I'	154.7 (7)	Ir–CO''C–CO''O	173.5
O₂–Ir–I'	111.8 (8)	Cl₁–C–Cl₂	115.2 (29)
O₂–Ir–I''	138.6 (8)	P₁A₁–P₁–P₁B₁	106.2
I'–Ir–CO'C	89.0	P₁B₁–P₁–P₁C₁	105.9
I'–Ir–CO''C	19.8	P₁A₁–P₁–P₁C₁	104.3
I''–Ir–CO'C	20.6	P₂L₁–P₂–P₂M₁	105.9
I''–Ir–CO''C	89.8	P₂M₁–P₂–P₂N₁	104.0
O₁–Ir–P₁	87.9 (6)	P₂L₁–P₂–P₂N₁	103.9
O₂–Ir–P₁	84.6 (7)		
O₁–Ir–P₂	89.1 (6)		
O₂–Ir–P₂	87.3 (7)		

proximately similar in direction and magnitude to those of the iridium atom. The motions of the half-iodine atoms will not be as meaningful as those of the other atoms because of the inadequacy of the description of the half-carbonyl groups. The oxygen atoms have their maximum vibration approximately in the equatorial plane in a direction most clearly shown in Figure 1.

This motion of the oxygen atoms may be at least partially due to the superposition of the two separate disordered molecules; in the asymmetric molecule 2, the oxygen atoms could be displaced away from the iodine atom.

If one assumes that the "riding" model for correction of the effects of thermal motion on bond distances[16] is appropriate, then one finds that none of the bond distances within the inner coordination sphere, including the O–O distance, is lengthened appreciably.

The geometry of the triphenylphosphine groups does not differ significantly from that found in a number of previously studied structures.

The C–Cl bond lengths of 1.63 (4) and 1.70 (4) Å in the dichloromethane molecule are somewhat shorter than the distance of 1.77 Å found in a previous microwave study[17] of the gaseous molecule. In view of the very large thermal motions of the dichloromethane molecule in this structure, such shortening is not unexpected. The Cl–C–Cl bond angle of 115.2 (2.9)° does not differ significantly from the microwave value of 111.8°.

Discussion

The mode of attachment of oxygen in molecular oxygen carriers has long been a subject of discussion. A comparison can now be made between the structures of two closely related compounds, $IrCl(O_2)(CO)(P(C_6H_5)_3)_2$ and $IrI(O_2)(CO)(P(C_6H_5)_3)_2$, and this is of interest because the chloro compound is reversibly oxygenated whereas the iodo compound is irreversibly oxygenated. The results obtained here may perhaps be related to the natural molecular oxygen carriers, the most important of which is hemoglobin.

To us the most reasonable explanation of the bonding in these compounds is analogous to that proposed for oxygenated heme.[18] The processes now discussed are theoretical and have no semblance to the chemical intermediates of the oxygenation reaction. Prior to bonding, the O_2 group forms its prepared, or "valence," state; in this state the valence description is the same as in ethylene; *i.e.*, the oxygen atoms are sp² hybridized and the molecule has two lone electron pairs on each oxygen atom with one σ bond and one π bond, both occupied by one pair of electrons. In this state the O_2 molecule is diamagnetic. The O_2 molecule acts as a ligand by donating one pair of electrons from the π bond to iridium (and thus the O–O axis must be at right angles to the direction of coordination). Prior to oxygenation, the central atom, Ir(I), a d⁸ system, was dsp² hybridized. The prepared state of iridium can be considered to be dsp³ hybridized with one directional vacant orbital which, in oxygenation, receives the pair of electrons from the O_2 group. In conjunction with the formation of this bond, there is back-donation from one of the filled iridium d orbitals to the π* orbital of the O_2 group. As electrons begin to be donated to iridium from oxygen, a dipole is set up which tends to

(16) W. R. Busing and H. A. Levy, *Acta Cryst.*, **17**, 142 (1964).
(17) R. J. Myers and W. D. Gwinn, *J. Chem. Phys.*, **20**, 1420 (1952).
(18) J. S. Griffith, *Proc. Roy. Soc.* (London), **A235**, 23 (1956).

stop further donation; the back-donation reduces this dipole and thus the over-all bond strength depends very largely upon the degree of back-donation that is possible. The extent of back-donation depends upon the availability of electrons at iridium and the chloro complex would be expected to form a weaker bond with oxygen than the iodo complex. In the formation of the bond to iridium, electrons are removed from an O–O bonding orbital and the back-donation is to a molecular orbital very similar to the lowest antibonding orbital in the "prepared" valence state of oxygen. Thus the more strongly the O_2 group is bound to iridium, the longer is the expected O–O distance. In this way, we rationalize the observed effects.

Infrared absorptions due to the O_2 group attached to a metal have been reported for a wide range of compounds. In general, the only band observed is in the region 840–900 cm^{-1}; such bands have been assigned to the O–O stretch. The positions of this absorption for some compounds of the general formula $IrX(O_2)$-$(CO)(P(C_6H_5)_3)_2$ are 858 cm^{-1} (X = Cl), 862 cm^{-1} (X = Br, I), and 855 cm^{-1} (X = N_3). The facile assignment of this absorption as an O–O stretch does not explain why a difference of 0.2 A in O–O bond length between the oxygenated chloro and iodo compounds only results in a change of 4 cm^{-1} in the position of the absorption. Normally the position of an absorption should vary greatly with the bond length, *i.e.*, bond strength.

For the system Ir–O_2 (C_{2v} symmetry), three infrared-active modes are expected: two symmetric and one antisymmetric. In this system, the absorption due to the antisymmetric mode is expected to occur at a wavelength lower than has been investigated. Of the two symmetric modes, we believe that the observed band corresponds to that mode analogous to the ν_1 mode in H_2O. The displacement coordinates of the oxygen atoms in this mode can be resolved into components along and perpendicular to the O–O bond. Because of the greater strength of the bond in the chloro compound, the displacement along the O–O bond should be less there than in the iodo compound. Because the O_2 group is more tightly bonded to iridium in the iodo compound the displacement normal to the O–O bond should be less there than in the chloro compound. Thus it seems probable that these changes in the displacements counterbalance one another and lead to essentially the same over-all displacement and hence the same observed frequency. The usual relations between bond strength and stretching frequencies strictly apply only to diatomic gaseous molecules, and the present case is an example where extension of these relations to polyatomic molecules is not successful.

Stereochemistry of Organophosphorus Complexes of Transition Metals

BY J. M. JENKINS AND J. G. VERKADE

Syntheses for 20 new complexes of the type $[M(CO)_4L_2]$ and $[M'X_2L_2]$ where M = Cr, Mo, or W; M' = Pd or Pt; X = Cl or I; and L = $P(CH_3)_3$, $P(CH_2Cl)_3$, $P(CH_2O)_3CCH_3$, $P(N(CH_3)_2)_3$, $P(OCH_2)_3CC_2H_5$, and $P(OCH_2)_3C(CH_2)_4CH_3$ are described. An improved method is presented for the preparation of the dichloro Pt and Pd(II) complexes from the corresponding dichlorobis(benzonitrile)metal(II) compound in cases where the ligand or derived complex is easily solvolyzed. Both near- and far-infrared measurements are interpreted in terms of the *cis* and *trans* geometry of all of the complexes. Metal-phosphorus stretching frequencies from 220 to 214 cm^{-1} were assigned in three of the carbonyl complexes. The required number of metal–chlorine modes could be identified in all of the Pt and Pd(II) complexes: 329–362 cm^{-1} in the *trans* compounds and 289–334 cm^{-1} in the *cis*. Metal–iodine modes in the Pt and Pd(II) systems could be assigned with reasonable certainty in only two of the *trans* (151, 158 cm^{-1}) and one of the *cis* (157, 150 cm^{-1}) compounds because their intensities were generally low.

A number of examples of P^{31}–P^{31} spin–spin coupling in various complexes containing two phosphorus nuclei have been reported.[1-5] In order to gain some insight into the nature of this interaction and its relation to the

phosphorus–metal bond, a new experimental approach has been developed involving the study of P^{31}–P^{31} spin–spin coupling constants in disubstituted complexes as a function of ligand strength and coordination geometry.[6] To enlarge the scope of our study, the preparation of new disubstituted complexes containing various phosphorus ligands was undertaken.

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(6) F. B. Ogilvie, J. M. Jenkins, and J. G. Verkade, submitted for publication.