

plex manner. This may be due to the fact that the intrinsic intramolecular vibration frequencies of  $\text{Ni}(\text{NO}_2)_6^{4-}$  are lower than those of  $\text{Co}(\text{NO}_2)_6^{3-}$ .

In the spectrum of  $\text{Na}_2\text{Ba}[\text{Ni}(\text{NO}_2)_6]$  shown in Figure 2, the band around  $200\text{ cm}^{-1}$  shifts considerably with a lowering of temperature. Such a frequency shift was also seen for the band at  $214\text{ cm}^{-1}$  of  $\text{Na}_3[\text{Co}(\text{NO}_2)_6]$ . Therefore, the band around  $200\text{ cm}^{-1}$  of  $\text{Na}_2\text{Ba}[\text{Ni}(\text{NO}_2)_6]$  may be assigned to the lattice mode related to the  $\text{Na}^+$  ion. The lowest band around  $80\text{ cm}^{-1}$  which also exists in the spectrum of  $\text{K}_2\text{Ba}[\text{Ni}(\text{NO}_2)_6]$  is assigned to the lattice mode related to the  $\text{Ba}^{2+}$  ion.

### Conclusion

The far-infrared bands due to the lattice vibrations and the intramolecular vibrations of the hexanitro complex salts were both interpreted by the normal-coordinate analysis of the crystal as a whole, only the interaction between the outer cations and the oxygen atoms being taken into consideration. Low-temperature spectra yield useful information on the interpretation of the observed bands.

Lattice vibrations due to the displacements of the outer cations relative to the complex ion are observed below  $200\text{ cm}^{-1}$ . Those frequencies change depending upon the mass of the cations and the interionic force constants. In order to establish the characteristic frequencies of the lattice vibrations arising from the outer cations, further investigations on the far-infrared spectra of various types of complex salts are in progress.

### Appendix

For the intramolecular potential,  $V_{\text{intra}}$ , the modified Urey-Bradley force field (MUBFF) has been used

and the use of this type of potential in the vibrational analysis of the complex ions has been fully discussed in several previous papers.<sup>4b,12-14</sup> Force constants used in the present study are summarized in Table VI, where

TABLE VI  
INTRAMOLECULAR POTENTIAL CONSTANTS (MDYN/Å)  
FOR  $\text{Co}(\text{NO}_2)_6^{3-}$  AND  $\text{Ni}(\text{NO}_2)_6^{4-}$ <sup>a</sup>

	$\text{Co}(\text{NO}_2)_6^{3-}$	$\text{Ni}(\text{NO}_2)_6^{4-}$
$F_{\text{dia}}(\text{MN str})$	1.50	0.80
$F_{\text{dia}}(\text{NO str})$	9.30	9.30
$F_{\text{dia}}(\text{NMN def})$	1.10	0.60
$F_{\text{dia}}(\text{ONO def})$	1.78	2.00
$F_{\text{dia}}(\text{MNO def})$	0.50	0.40
	0.30 (Na salt)	
$F_{\text{dia}}(\text{NO}_2\text{ wag})$	0.55	0.32
$F_{\text{dia}}(\text{NO}_2\text{ twist})$	0.03 (assumed)	0.03
$F(\text{N}\cdots\text{N})$	0.05	0.20 <sup>b</sup>
$F(\text{O}\cdots\text{O})$	3.00	3.00
$F(\text{M}\cdots\text{O})$	0.20	0.10
$p(\text{NO},\text{NO})$	0.50	0.50

<sup>a</sup> In the Urey-Bradley approach, diagonal elements of the  $\mathbf{F}$  matrix corresponding to the bond stretching modes include  $F$  as well as  $K$  and those for angle deformation modes include  $F$  as well as  $H$ . The off-diagonal elements are expressed in terms of  $F$ .

<sup>b</sup> Taking into account the appreciable interaction between the Ni-N stretching and NMN deformation modes, the effective value of  $F(\text{N}\cdots\text{N})$  was assumed to be large.<sup>4b,12</sup>

$K$ ,  $H$ , and  $F$  denote the bond stretching, the angle deformation, and the repulsion between nonbonded atoms, respectively.  $F_{\text{dia}}$  means the diagonal element of the  $\mathbf{F}$  matrix (potential energy matrix) for the corresponding mode.  $p(\text{NO},\text{NO})$  is the resonance interaction constant. The value of  $p(\text{MN},\text{MN})$  cannot be determined unless the frequencies of the  $g$  species are available.

CONTRIBUTION FROM BELL TELEPHONE LABORATORIES, INC.  
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## Rare Earth Orthogallates

BY M. MAREZIO, J. P. REMEIK, AND P. D. DERNIER

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Single crystals of the rare earth orthogallates,  $\text{REGaO}_3$  with  $\text{RE} = \text{Sm-Lu}$ , have been synthesized by the decomposition of the respective garnets at high pressure in the presence of a flux. Previously reported crystals of  $\text{LaGaO}_3$ ,  $\text{PrGaO}_3$ , and  $\text{NdGaO}_3$  have been prepared by the flux method at atmospheric pressure. The rare earth orthogallates are isostructural with the rare earth orthoferrites; namely, they have the orthorhombic perovskite-like structure. Their lattice parameters have been calculated by the least-squares method. The parameters  $a$  and  $c$  increase smoothly in going from Lu to La. However, the lattice parameter  $b$  goes through a maximum. It seems that this anomalous variation of the  $b$  parameter can be explained in terms of a variation of the coordination number of the rare earth cation.

### Introduction

In a recent note<sup>1</sup> we reported the synthesis and the crystal structure of  $\text{GdGaO}_3$ ,  $\text{YbGaO}_3$ , and  $\text{YGaO}_3$ . They were synthesized by decomposition at high pressure and high temperature of the respective garnets.

The new orthogallates are isostructural with the respective orthoferrites; namely, they have the orthorhombic perovskite-like structure. An interesting feature is that the rare earth orthoferrites can be synthesized at atmospheric pressure, whereas a pressure greater than 45 kbars is needed to synthesize  $\text{GdGaO}_3$ ,  $\text{YbGaO}_3$ , and  $\text{YGaO}_3$ . This was explained in terms of

(1) M. Marezio, J. P. Remeika, and P. D. Dernier, *Mater. Res. Bull.*, **1**, 247 (1966).

the stronger preference of the  $\text{Ga}^{3+}$  ions for the tetrahedral sites in oxide compounds. If a mixture of  $\text{RE}_2\text{O}_3$  (RE = Sm–Lu) and  $\text{Ga}_2\text{O}_3$  is fired at atmospheric pressure, the corresponding garnet  $\text{RE}_3\text{Ga}_5\text{O}_{12}$  is obtained. In the garnet structure three-fifths of the galliums are tetrahedrally coordinated and two-fifths octahedrally coordinated. All of the gallium ions in the  $\text{REGaO}_3$  compounds are in octahedral sites.

The present paper, which is part of a continuing program of study of the  $\text{Ga}^{3+}$  crystal chemistry in our laboratory, reports the synthesis and the crystal data of all the rare earth orthogallates.

### Experimental Section

With the exception of  $\text{LaGaO}_3$ ,  $\text{PrGaO}_3$ , and  $\text{NdGaO}_3$ , single crystals of all the other rare earth orthogallates were prepared by subjecting the respective garnets to high pressure and high temperature in the presence of a NaOH flux. The high-pressure runs were made by following the same procedure that has been described in detail earlier.<sup>1</sup> This time the temperature was monitored with a Pt–Pt–10% Rh thermocouple during each experiment. At high pressure the temperature fluctuation was less than  $\pm 20^\circ$  at  $1000^\circ$ . No correction was made for the effect of pressure on the thermocouple. A temperature of  $1000^\circ$  and a pressure of 45 kbars proved to be sufficient for the decomposition of all the rare earth gallium garnets. It is interesting to note that under the above conditions complete decomposition was obtained from  $\text{Sm}_3\text{Ga}_5\text{O}_{12}$  to  $\text{Tb}_3\text{Ga}_5\text{O}_{12}$ , whereas a partial decomposition was obtained from  $\text{Dy}_3\text{Ga}_5\text{O}_{12}$  to  $\text{Lu}_3\text{Ga}_5\text{O}_{12}$ , as evidenced by the occurrence of lines belonging to the garnet structure in the X-ray photographs of the quenched materials.

The first three members of the orthogallate series, namely,  $\text{LaGaO}_3$ ,  $\text{PrGaO}_3$ , and  $\text{NdGaO}_3$ , can be synthesized at normal pressure. Their synthesis and crystal data have been reported previously.<sup>2</sup> Since it was our interest to have a complete set of lattice parameters of the orthogallates, we have prepared these materials and redetermined their lattice parameters. Single crystals were grown by self-nucleation from a melt starting at  $1300^\circ$  and consisting of the stoichiometric proportions of the constituent oxides, which were of 99.99% purity by emission spectroscopic analysis. The  $\text{PbO}(\text{B}_2\text{O}_3)_x$  flux used was of the same purity level. Cooling rates were of the order of 1–5°/hr. Complete details of the crystal growth of the  $\text{ABO}_3$  perovskite-like compounds from the flux system  $\text{PbO}(\text{B}_2\text{O}_3)_x$  will be published elsewhere.

Three compounds were selected for compositional analysis. Atomic absorption analysis gave 51.8% La and 26.0% Ga for  $\text{LaGaO}_3$  and 54.9% Nd and 25.0% Ga for  $\text{NdGaO}_3$ , compared with the calculated values 54.13% La, 27.17% Ga and 55.07% Nd, 26.61% Ga. X-Ray fluorescence analysis for  $\text{SmGaO}_3$  gave 56.7% Sm and 24.5% Ga compared with the calculated values 56.09% Sm and 26.01% Ga. In the case of the crystals grown from the NaOH flux at high pressure, an emission spectroscopic analysis was performed for Na and it indicated that Na is present as a negligible impurity in the range of 0.000X%.

The X-ray powder photographs were taken at room temperature ( $\sim 24^\circ$ ) with a Norelco camera of 114.6-mm diameter and Zr-filtered Cr radiation ( $K\alpha_1$  2.28962 Å and  $K\alpha_2$  2.29352 Å). The photographs were indexed on the basis of the orthorhombic cell with the lattice constants given in Table I. There was some doubt about the pattern of  $\text{LaGaO}_3$ . It could have been indexed on the trigonal cell isostructural with the one of  $\text{LaAlO}_3$ . After taking long exposure photographs with Cr  $K\alpha$  and one with Cu  $K\alpha$ , lines that could have only been explained with the orthorhombic cell appeared in the photographs. The final lattice parameters were obtained by accurate measurements of the  $2\theta$  values for at least ten reflections in the back reflection region.

(2) S. Geller, *Acta Cryst.*, **10**, 243 (1957).

TABLE I

	Orthorhombic cell				Pseudo-cell		
	$a$ , Å	$b$ , Å	$c$ , Å	Vol Å <sup>3</sup>	$a = c$ , Å	$b$ , Å	$\beta$ , deg
$\text{LaGaO}_3$	5.526	5.473	7.767	234.9	3.888	3.884	90.6
$\text{PrGaO}_3$	5.458	5.490	7.733	231.7	3.871	3.866	90.3
$\text{NdGaO}_3$	5.431	5.499	7.710	230.3	3.865	3.855	90.7
$\text{SmGaO}_3$	5.369	5.520	7.650	226.7	3.850	3.825	91.6
$\text{EuGaO}_3$	5.351	5.528	7.628	225.6	3.847	3.814	91.9
$\text{GdGaO}_3$	5.322	5.537	7.606	224.1	3.840	3.803	92.3
$\text{TbGaO}_3$	5.307	5.531	7.578	222.4	3.833	3.789	92.4
$\text{DyGaO}_3$	5.282	5.534	7.556	220.9	3.825	3.778	92.7
$\text{HoGaO}_3$	5.251	5.531	7.536	218.9	3.813	3.768	93.0
$\text{ErGaO}_3$	5.239	5.527	7.522	217.8	3.808	3.761	93.1
$\text{TmGaO}_3$	5.224	5.515	7.505	216.2	3.798	3.753	93.1
$\text{YbGaO}_3$	5.208	5.510	7.490	214.9	3.791	3.745	93.2
$\text{LuGaO}_3$	5.188	5.505	7.484	213.7	3.782	3.742	93.4

<sup>a</sup> The observed standard deviation of the lattice parameters determined in the present work was in all cases less than 1/12,000. As a conservative estimate of limits of error we consider the values of all parameters to be accurate within a precision of 1/5000.

Those reflections for which overlapping could occur were discarded. The least-squares refinement program of Mueller, *et al.*,<sup>3</sup> was used. The powder diffraction data for all of the compounds are listed in Table II, where the reflections that were unobserved throughout the whole series have been omitted.

Details of the structural arrangement of the orthorhombic distortion of the perovskite structure are given by Coppens and Eibschütz,<sup>4</sup> who refined the structure of  $\text{YFeO}_3$ .

### Discussion

In Figure 1 the lattice parameters of the rare earth orthogallates are plotted against the atomic number of

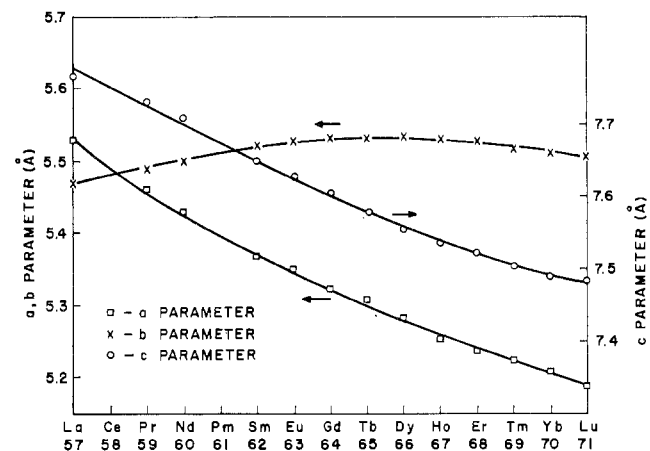


Figure 1.—Lattice parameters  $a$ ,  $b$ , and  $c$  vs. atomic number for the rare earth orthogallates.

the rare earth. Within experimental error the  $a$  and  $c$  parameters increase smoothly in going from Lu to La. Instead, the  $b$  parameter has an unexpected behavior. It goes through a maximum at about Dy–Tb–Gd. As shown by Eibschütz,<sup>5</sup> the lattice parameters of the corresponding rare earth orthoferrites have the same behavior. In this case the  $b$  parameter goes through a maximum at Gd. No explanation was given to account for this anomalous variation.

In the "ideal perovskite"  $\text{ABO}_3$  compounds the  $\text{A}^{3+}$

(3) M. H. Mueller, L. Heaton, and R. T. Miller, *ibid.*, **13**, 828 (1960).

(4) P. Coppens and M. Eibschütz, *ibid.*, **19**, 524 (1965).

(5) M. Eibschütz, *ibid.*, **19**, 337 (1965).

TABLE II  
POWDER DIFFRACTION DATA FOR SOME RARE EARTH ORTHOGALLATES

Table with columns for hkl, d (Å), and various rare earth elements (La, Pr, Nd, Sm, Eu, Tb, Dy, Ho, Er, Tm, Lu) with sub-columns for d\_c and I\_0.

ions are surrounded by 12 equidistant oxygen ions. In the orthorhombic distortion these 12 distances vary over a very large range; in fact, it is difficult to decide the coordination number of the cation A3+. Of all the orthorhombic perovskite-like compounds, only two have been refined, GdFeO3 and YFeO3. Although yttrium is not a rare earth, YFeO3 can be considered an orthoferrite whose A cation falls between Dy and Ho in size. From the rare earth-oxygen distances in these compounds it can be seen that there are eight first nearest oxygens and four second nearest oxygens. The averages of the distances of the first nearest oxygens in GdFeO3 and YFeO3 are 2.48 and 2.43 Å, respectively. The averages of the distances of the other four oxygens are 3.39 and 3.45 Å, while the averages of the distances of all 12 oxygens are 2.78 and 2.77 Å, respectively. The four second nearest oxygens are closer to Gd in GdFeO3 than to Y in YFeO3. Therefore, the anomalous behavior of the b parameter in the AFeO3 perovskite-like compounds, as A varies from Lu to La, can be explained if one considers the effect of the variation of the size of the A cation on the first nearest oxygens together with the effect of the same variation on the second nearest oxygens. As the size of A3+ increases, the average A-O distance of the first nearest oxygens increases, whereas the average A-O distance of the second nearest oxygens decreases. Therefore the unit

cell increases and decreases, respectively. Since the a and c parameters vary smoothly across the series, only the b parameter is appreciably affected by this double effect. The b parameter increases on going from Lu to Gd because in this region the first effect is the predominant one, whereas it decreases on going from Gd to La because in this region the second effect is the predominant one.

The fact that the average A-O distance of the second nearest oxygens becomes smaller on going from LuFeO3 to LaFeO3 seems to suggest that the coordination number of the cation A3+ varies across the series. Even though LaFeO3 is the least distorted among the rare earth orthoferrites, it can be seen from the intensities of the reflections in the powder pattern that the distortion from the ideal perovskite is still very appreciable which means that the coordination number should be far from 12.

It is reasonable to assume that the same explanation holds true in the case of the orthogallates.

The same behavior of the b parameter is also present in the rare earth orthochromites,2,7 orthovanadites,2,8 and orthorhodites.9 In the case of orthoalumi-

(6) S. Geller, J. Chem. Phys., 24, 1236 (1956).

(7) S. Quezel-Ambunaz and M. Mareshal, Bull. Soc. Franc. Mineral. Crist., 86, 204 (1963).

(8) B. Reuter, Bull. Soc. Chim. France, 1053 (1965).

(9) R. Chazalon, E. F. Bertaut, and T. Q. Duc, Bull. Soc. Franc. Mineral. Crist., 87, 98 (1964).

nates<sup>10,11</sup> the  $b$  parameter decreases slightly on going from  $\text{TmAlO}_3$  to  $\text{SmAlO}_3$ . Starting with  $\text{NdAlO}_3$ , the orthoaluminates have the hexagonal perovskite-like arrangement.<sup>10</sup> In the cases of orthoscandates<sup>2,11</sup> and orthoindates<sup>12</sup> the  $b$  parameter increases throughout the series from Lu to La. These results seem to indicate that when the B cation of the  $\text{ABO}_3$  perovskite-like compounds is a larger ion such as indium or scandium,

(10) S. Geller and V. B. Bala, *Acta Cryst.*, **9**, 1019 (1956).

(11) S. J. Schneider, R. S. Roth, and J. L. Waring, *J. Res. Natl. Bur. Std.*, **A66**, 345 (1961).

(12) R. D. Shannon, *Inorg. Chem.*, **6**, 1474 (1967).

the substitution of the A cation with a larger rare earth does not have an appreciable effect on the second nearest oxygens. Therefore, the coordination number of the rare earth probably varies across the series  $\text{REBO}_3$  when B = Al, Ga, Fe, Cr, or V, whereas it is almost constant when B = Sc or In.

We plan to do the refinements of a number of these perovskite-like structures.

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## The Crystal Structure of Lanthanum Carbonate Octahydrate

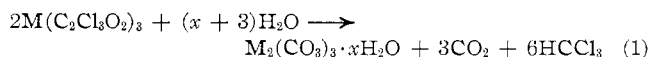
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The crystal structure of lanthanum carbonate octahydrate,  $\text{La}_2(\text{CO}_3)_3 \cdot 8\text{H}_2\text{O}$ , has been determined from an X-ray diffraction study of a single-crystal specimen. Four formula units are contained in the orthorhombic unit cell, space group  $Pccn$ , with lattice parameters  $a = 8.984 \pm 0.004$ ,  $b = 9.580 \pm 0.004$ , and  $c = 17.00 \pm 0.01 \text{ \AA}$  ( $\rho_{\text{obsd}} = 2.72 \pm 0.02 \text{ g cm}^{-3}$ ;  $\rho_{\text{calcd}} = 2.732 \text{ g cm}^{-3}$ ). The final  $R$  factor was 0.061 for three-dimensional counter data collected with Cu  $K\alpha$  radiation [ $(\sin \theta)/\lambda \leq 0.50$ ].  $\text{La}_2(\text{CO}_3)_3 \cdot 8\text{H}_2\text{O}$  crystallizes in an irregular layer type of structure in which the basic layer is formed by alternate rows of carbonate ions and metals. The structure contains two distinctive 10-coordinate metal polyhedra in which coordination sites are occupied both by water molecules and by bidentate and unidentate carbonates. The polyhedra symmetry may be related to a dodecahedron with two sites occupied by bidentate carbonate ions. One-fourth of the water molecules are not bound to the metal ions and are situated between the layers. The average La–O bond is 2.61  $\text{\AA}$ ; individual La–O distances range from 2.52 to 2.74  $\text{\AA}$ . Hydrogen bonding is apparent.

### Introduction

Methods of preparation, thermal decomposition, and structural properties of lanthanon carbonates have been the subject of numerous recent studies.<sup>1–10</sup> Unfortunately, many of these results and their interpretation are contradictory. The inconsistencies are due in part to preparatory difficulties and lack of structural data. A popular method for the preparation of the carbonates is homogeneous precipitation by hydrolysis of the trichloroacetate (eq 1) described first by Salutsky and Quill.<sup>1</sup>



By use of this and related methods, normal carbonates,  $\text{M}_2(\text{CO}_3)_3 \cdot x\text{H}_2\text{O}$ , have been prepared for all the lan-

(1) M. L. Salutsky and L. L. Quill, *J. Am. Chem. Soc.*, **72**, 3306 (1950).

(2) R. G. Charles, *J. Inorg. Nucl. Chem.*, **27**, 1498 (1965).

(3) E. L. Head, 6th Rare Earth Research Conference, Gatlinburg, Tenn., May 1967.

(4) E. L. Head and C. E. Holley, Jr., "Rare Earth Research III," Gordon and Breach, New York, N. Y., 1964, p 51.

(5) E. L. Head and C. E. Holley, Jr., "Rare Earth Research IV," Gordon and Breach, New York, N. Y., 1965, p 707.

(6) R. L. N. Sastry, S. R. Yoganarisimhan, P. N. Mehrota, and C. N. R. Rao, *J. Inorg. Nucl. Chem.*, **28**, 1165 (1966).

(7) M. Clarus Strouth, Ph.D. Thesis, Michigan State University, 1962.

(8) G. Pannetier, J. Nataf, and A. Desreigne, *Bull. Soc. Chim. France*, **5**, 318 (1965).

(9) P. Caro and L. Eyring, personal communication.

(10) S. D. Ross and J. Goldsmith, *Spectrochim. Acta*, **20**, 781 (1964).

thanons.<sup>2–7</sup> Values reported for  $x$  vary from 8 for the lighter metals to 2 or 3 for the heavier ones and also seem dependent on the preparatory conditions. In thermal decomposition studies of the hydrates, intermediate phases have not been observed consistently.<sup>2–4,6,8</sup> The crystal structures of the lanthanon carbonates are unknown, although many reported lattice parameters are based on values analogous to those obtained from a single crystal of lanthanite<sup>11</sup>  $\{(\text{La}, \text{Ce})_2(\text{CO}_3)_3 \cdot 8\text{H}_2\text{O}\}$ .

In order to determine the structural interrelationships in the lanthanon carbonates and to resolve some of the contradictory reports of their properties, a single-crystal X-ray diffraction study of the inherent structures has been initiated. The structure of  $\text{La}_2(\text{CO}_3)_3 \cdot 8\text{H}_2\text{O}$  is described below.

### Experimental Section

Crystals of lanthanum carbonate octahydrate,  $\text{La}_2(\text{CO}_3)_3 \cdot 8\text{H}_2\text{O}$ , prepared by slow hydrolysis of lanthanum trichloroacetate, were obtained from Dr. L. L. Quill of this department. Composition was determined from carbon and hydrogen analyses (Spang Microanalytical Laboratory, Ann Arbor, Mich.) and from metal analyses by ignition to the oxide. *Anal.* Calcd for  $\text{La}_2(\text{CO}_3)_3 \cdot 8\text{H}_2\text{O}$ : La, 46.15; C, 5.98; H, 2.69. Found:

(11) D. Palache, H. Berman, and C. Frondel, "Dana's Systems of Mineralogy," Vol. II, 7th ed, John Wiley and Sons, Inc., New York, N. Y., 1944.