

CONTRIBUTION FROM THE DEPARTMENT OF CHEMISTRY,
MASSACHUSETTS INSTITUTE OF TECHNOLOGY, CAMBRIDGE, MASSACHUSETTS 02139The Structure of Oxopentachloropropionatobis(triphenylphosphine)dirhenium(IV)^{1a}By F. A. COTTON AND BRUCE M. FOXMAN^{2b}

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The crystal and molecular structure of the compound $\text{Re}_2\text{OCl}_5(\text{C}_2\text{H}_3\text{CO}_2)[\text{P}(\text{C}_6\text{H}_5)_3]_2$ reported recently by Rouschias and Wilkinson has been determined by means of single-crystal X-ray diffractometry. The molecular structure is essentially, though not in all respects, that postulated by Rouschias and Wilkinson. There are two octahedrally coordinated rhenium atoms with the octahedra joined on a common edge; the bridging atoms on this common edge are Cl and O. The propionato group fills one apical coordination position in each octahedron and the remaining positions in each are filled by two Cl atoms and a phosphorus atom. The phosphorus atoms are *trans* to the bridging chlorine atom. Neglecting the orientations of the phenyl rings and the β -carbon atom of $\text{C}_2\text{H}_3\text{CO}_2$, the structure has a virtual mirror plane containing the bridging O and Cl atoms and the carboxyl carbon atom. The quadrilateral formed by the two bridging atoms and the two rhenium atoms is practically planar, with an Re-Re distance of $2.522 \pm 0.001 \text{ \AA}$. The crystals, which were black, with tetragonal prismatic morphology ($\rho = 1.94 \text{ g cm}^{-3}$ at $\sim 23^\circ$) belong to the tetragonal system with $a = 10.669 \pm 0.003 \text{ \AA}$ and $c = 34.838 \pm 0.007 \text{ \AA}$. For $Z = 4$, $\rho_{\text{calc}} = 1.945 \text{ g cm}^{-3}$. The structure was solved by Patterson and Fourier methods and refined by least squares. Systematic absences indicated the space group to be P4_1 or P4_3 . The use of 3156 hkl reflections and 446 $\bar{h}\bar{k}\bar{l}$ reflections permitted an unambiguous choice (for this particular crystal) of P4_3 , for which final weighted and unweighted residuals of 0.050 and 0.044, respectively, were obtained.

Introduction

It has now been well established that the chemistry of rhenium in the oxidation state III is largely that of dinuclear²⁻⁹ and trinuclear¹⁰⁻²⁰ metal atom cluster compounds, although a few mononuclear Re(III) complexes do appear to be adequately characterized.^{21,22} Compounds containing rhenium in mean oxidation states between II and III also seem to involve strong Re-Re bonds.^{23,24}

Very recently²⁵ it has been shown that the compound

$\text{La}_4\text{Re}_6\text{O}_{19}$, which has rhenium in the mean oxidation state $+4\frac{1}{3}$, contains Re_2O_{10} groups, which are pairs of octahedra sharing an edge, with the Re-Re distance equal to 2.42 \AA . The interpretation of this structure in terms of localized Re-Re bonding is not entirely straightforward, however, because the Re_2O_{10} groups are far from discrete. Instead they are extensively linked together by shared oxygen atoms and the substance has semimetallic electrical properties, indicative of considerable electron delocalization.

With the exception of the ternary oxide just mentioned, no Re-Re bonds are known in compounds in which the Re has an oxidation state greater than IV and compounds containing Re(IV) seem to have only a moderate tendency to form Re-Re bonds. Thus the best known Re(IV) species is the mononuclear ReCl_6^{2-} ion. However, rhenium(IV) chloride²⁶ has a structure⁸ in which there are chains of Re_2Cl_9 bioctahedra, with rhenium atoms occurring in pairs, separated by about 2.73 \AA , and the Re_2Cl_9^- ion²⁷ also has a bioctahedral structure²⁸ with an Re-Re distance of about 2.71 \AA . Thus, the limited information available suggests that the borderline, in terms of oxidation state, between predominant and negligible tendency to metal-metal bond formation occurs around oxidation number IV for rhenium.

In view of the borderline nature of Re(IV) with respect to Re-Re bond formation, it was considered of interest to investigate the structure of the compound $\text{Re}_2\text{OCl}_5(\text{C}_2\text{H}_5\text{CO}_2)[\text{P}(\text{C}_6\text{H}_5)_3]_2$ reported recently by Rouschias and Wilkinson.²⁹ These authors proposed structure I for the compound. This structure, which was quite speculative, would seem to require the

(1) (a) Supported by the U. S. Atomic Energy Commission; (b) N. I. H. Predoctoral Fellow, 1965-1966.

(2) (a) F. A. Cotton, N. F. Curtis, B. F. G. Johnson, and W. R. Robinson, *Inorg. Chem.*, **4**, 326 (1965); (b) F. A. Cotton and C. B. Harris, *ibid.*, **4**, 330 (1965).

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(6) F. A. Cotton, C. Oldham, and R. A. Walton, *ibid.*, **6**, 214 (1967).

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(8) M. J. Bennett, F. A. Cotton, B. M. Foxman, and P. F. Stokely, *J. Am. Chem. Soc.*, **89**, 2759 (1967).

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(11) W. T. Robinson, J. B. Fergusson, and B. R. Penfold, *Proc. Chem. Soc.*, 116 (1963).

(12) F. A. Cotton and J. T. Mague, *Inorg. Chem.*, **3**, 1094 (1964).

(13) F. A. Cotton and J. T. Mague, *Proc. Chem. Soc.*, 233 (1964); *Inorg. Chem.*, **3**, 1402 (1964).

(14) B. H. Robinson and J. E. Fergusson, *J. Chem. Soc.*, 5683 (1964).

(15) F. A. Cotton, S. J. Lippard, and J. T. Mague, *Inorg. Chem.*, **4**, 508 (1965).

(16) F. A. Cotton and R. A. Walton, *ibid.*, **5**, 1802 (1966).

(17) B. R. Penfold and W. T. Robinson, *ibid.*, **5**, 1758 (1966).

(18) M. Elder and B. R. Penfold, *ibid.*, **5**, 1763 (1966).

(19) J. H. Hickford and J. E. Fergusson, *J. Chem. Soc., A*, 113 (1967).

(20) M. J. Bennett, F. A. Cotton, and B. M. Foxman, *Inorg. Chem.*, **7**, 1563 (1968).

(21) Some examples will be found in: (a) J. Chatt and G. A. Rowe, *J. Chem. Soc.*, 4019 (1962); (b) N. F. Curtis, J. E. Fergusson, and R. S. Nyholm, *Chem. Ind. (London)*, 625 (1958).

(22) There are other reports of mononuclear Re(III) complexes, but in many the structural conclusions lack adequate proof.

(23) F. A. Cotton, W. R. Robinson, and R. A. Walton, *Inorg. Chem.*, **6**, 1257 (1967).

(24) M. J. Bennett, F. A. Cotton, and R. A. Walton, *J. Am. Chem. Soc.*, **88**, 2866 (1966); *Proc. Roy. Soc. (London)*, **A303**, 175 (1968).

(25) J. M. Longo and A. W. Sleight, *Inorg. Chem.*, **7**, 108 (1968); N. Morrow and L. Katz, American Crystallographic Association Meeting, Minneapolis, Minn., Aug 1967, Paper K-6.

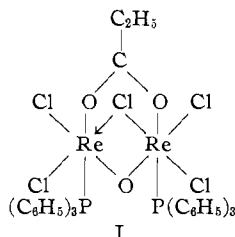
(26) F. A. Cotton, W. R. Robinson, and R. A. Walton, *Inorg. Chem.*, **6**, 223 (1967).

(27) F. Bonati and F. A. Cotton, *ibid.*, **6**, 1353 (1967).

(28) F. A. Cotton and P. F. Stokely, unpublished observations.

(29) G. Rouschias and G. Wilkinson, *J. Chem. Soc., A*, 465 (1966).

existence of a fairly short Re–Re separation and hence a high—possibly unprecedented—amount of Re(IV)–Re(IV) bonding. It seemed that a knowledge of the Re–Re distance and other structural features of this molecule might enable us significantly to refine our views concerning the influence of oxidation number on the tendency to metal–metal bond formation. This paper reports the results of the X-ray crystallographic study.



Experimental Section

Crystals of $\text{Re}_2\text{OCl}_5(\text{C}_2\text{H}_5\text{CO}_2)[\text{P}(\text{C}_6\text{H}_5)_3]_2$ were prepared by a slight modification of the method of Rouschias and Wilkinson.²⁹ Propionic acid (25 ml) and $\text{ReOCl}_3[\text{P}(\text{C}_6\text{H}_5)_3]_2$ ³⁰ (2.0 g) were refluxed for 1 hr. The mixture was allowed to cool and the supernatant liquid was decanted. Black crystals, 0.9 g (56%), were collected, washed with propionic acid and pentane, and dried under vacuum. No $\text{ReCl}_4[\text{P}(\text{C}_6\text{H}_5)_3]_2$ was observed. *Anal.* Calcd for $\text{Re}_2\text{Cl}_5\text{P}_2\text{O}_3\text{C}_30\text{H}_36$: C, 40.26; H, 3.03; Cl, 15.24. Found: C, 40.8; H, 3.22; Cl, 16.0.

A preliminary optical examination showed the crystals to be tetragonal prisms. Single crystals were examined by Weissenberg and precession photography and found to belong to the tetragonal system. The unit cell parameters were determined at 23° as $a = 10.669 \pm 0.003 \text{ \AA}$ and $c = 34.898 \pm 0.007 \text{ \AA}$ from measurements on a General Electric manual diffractometer using Cu $K\alpha$ radiation ($\lambda_{\alpha 1} 1.5405 \text{ \AA}$, $\lambda_{\alpha 2} 1.5443 \text{ \AA}$). The values quoted for a and c and corresponding uncertainties were obtained in a manner previously described by us.²⁰ The observed density, $1.94 \pm 0.01 \text{ g cm}^{-3}$, measured by flotation in aqueous zinc bromide, agrees well with the calculated value of 1.945 g cm^{-3} for $Z = 4$, $V = 3972.4 \text{ \AA}^3$, and mol wt = 1163.4.

The observed Laue symmetry of 4/m and the systematic absences, $00l$ for $l \neq 4n$, observed on Weissenberg photographs of levels $hk0$ and $hk1$, and precession photographs of $h0l$, $h1l$, $h2l$, $0kl$, $1kl$, and $2kl$ suggested the enantiomorphic space groups $P4_1$ (C_4^2 , no. 76) and $P4_3$ (C_4^4 , no. 78).

The crystal used in the data collection was a tetragonal prism with dimensions of $0.06 \times 0.06 \times 0.11 \text{ mm}$. The crystal was aligned with its c axis coincident with the ϕ axis of the goniometer. Intensity data were collected on a General Electric quarter-circle automated XRD-6 diffractometer equipped with a Datex electronic system. Using nickel-filtered Cu $K\alpha$ radiation, the intensities of 3938 independent reflections of the complete index set hkl and 554 related reflections of the index set $\bar{h}\bar{k}l$ (where the indices are ≥ 0), accessible within the sphere bounded by $(\sin \theta)/\lambda \leq 0.617$, were measured. The intensities were measured using a scintillation counter with the pulse height discriminator set to receive 95% of the Cu $K\alpha$ radiation with the window centered on the Cu $K\alpha$ peak. The crystal was placed 14.6 cm from the source and 17.9 cm from the 2° circular screening aperture; a moving-crystal, moving-counter (θ – 2θ scan at 4°/min in 2θ) technique was employed. The integrated intensities ($I = P - B_1 - B_2$) were obtained from the total counts (P) of scan³¹ from $2\theta(\text{calcd}) - 1.33^\circ$ to $2\theta(\text{calcd}) + 1.33^\circ$ and stationary background measurements (B_1 , B_2) at the limits of each scan. A take-off angle of 2° was used, and the scan range of 2.66° conformed well with the mosaic spread of the crystal. Manual

scans of the strongest reflections showed that the maximum counting rates never exceeded the linear response range of the scintillation counter. Every 6 hr during data collection the crystal orientation was checked, and three standard reflections were monitored. A random variation with time ($\pm 2\%$ from the mean) of the standard reflections was consistent with the random instability of the X-ray equipment rather than decomposition. Of the 3938 independent reflections of the hkl set, 782 were rejected using two criteria: (1) $I \leq 0$; (2) $I/(P + B_1 + B_2)^{1/2} < 3$. The remaining 3156 reflections were corrected for Lorentz and polarization effects. Similarly, the 554 related reflections ($\bar{h}\bar{k}l$) were processed and corrected as above; 108 of these reflections were rejected using the stated criteria.

All nonrejected data were corrected for absorption effects³² ($\mu_{\text{Cu}} = 155.1 \text{ cm}^{-1}$). An experimental absorption curve for the $00l$ reflections (obtained at $\chi = 90^\circ$ upon rotation about the ϕ axis) and the calculated curve for the same reflections agreed well. The calculated transmission factors fell in the range 0.495–0.600. Only the 3156 independent reflections (complete index set hkl) were used in the structure solution and refinement.

Solution and Refinement of Structure

The three-dimensional Patterson function³³ gave a set of ten unique vectors attributable to Re–Re interactions. A straightforward solution was not readily obtained; however, when the triclinic vector set was generated from the unique 4/m set, the solution became immediately obvious, with Re(1) placed at $x = 0.14$, $y = 0.32$; Re(2) at $x = 0.16$, $y = 0.36$; and $\Delta z(\text{Re–Re}) = 0.07$. A structure factor calculation (assuming space group $P4_1$) at this point, with phases based on the two rhenium atoms, gave $R_1 = \Sigma ||F_o| - |F_c|| / \Sigma |F_o| = 0.369$, where F_o is the observed and F_c the calculated structure factor; the weighted residual, $R_2 = \{ \Sigma w \cdot [|F_o| - |F_c|]^2 / \Sigma w |F_o|^2 \}^{1/2}$, using equal (unit) weights for all reflections was 0.412. For this and subsequent calculations, Re(1) was assigned a z coordinate of zero, and this parameter was held fixed throughout the course of the refinement.

One cycle of full-matrix, least-squares³⁴ refinement in which scale factor and positional and isotropic thermal parameters for the two rhenium atoms were varied reduced R_1 to 0.244 and R_2 to 0.311. A difference Fourier map at this stage revealed the positions of the five chlorine and two phosphorus atoms. It was noted that one of the seven “chlorine atoms,” as they were temporarily designated, bridged the Re–Re bond, while the other six were terminal. Two more cycles of refinement reduced R_1 to 0.124 and R_2 to 0.165. A difference Fourier map at this point revealed the positions of the 36 phenyl carbon atoms, a bridging propionate group, and a bridging oxygen atom. At this

(31) A modified version of D. P. Shoemaker's MIXO-2 MIT X-ray goniometer package (1962) was used to compute settings and prepare input for the automated diffractometer. It should be noted that for the above experiment, the conditions for multiple diffraction are maximized: W. A. Zachariasen, *Acta Cryst.*, **18**, 705 (1965). Since agreement between observed and calculated structure factors was later found to be very good, the effects are probably small in this case.

(32) Calculated using a local modification of the GONO9 program of W. C. Hamilton, goniostat absorption corrections (1960).

(33) Calculated using GINPUT-GENFOR, a Fourier series analysis program by R. B. Roof, Jr., D. R. Cromer, and A. C. Larson, Los Alamos Scientific Laboratory, 1965.

(34) C. T. Prewitt, a Fortran IV full-matrix crystallographic least-squares program, SFLS-5, 1966. The function minimized during refinement is $\Sigma w ||F_c| - |F_o||^2$.

time, real and imaginary corrections³⁵ for anomalous scattering by Re, Cl, and P were included in the calculated structure amplitudes,³⁶ and refinement was continued in each of the enantiomorphous space groups $P4_1$ and $P4_3$. In the absence of anomalous scattering effects, there are two equally likely (and indistinguishable) solutions of the structure for this particular crystal; that is, the solution (obtained from the Patterson function) for space group $P4_1$ (enantiomer A) is equivalent to its inverse (enantiomer B) for space group $P4_3$. The set of coordinates for enantiomer B in space group $P4_3$ could be obtained, for example, by reflection of the enantiomer A structure in (001). Statistically one would expect to find an equal number of each of the above enantiomer-space group combinations. (It is important to note that structures involving the enantiomorphous pair—enantiomer B in space group $P4_1$ or enantiomer A in $P4_3$ —are not under consideration here. In these the atoms are quite differently related to the unit cell, and it seems unlikely, though not impossible, that such crystals occur in the reaction product.) If anomalous scattering occurs, the two enantiomer-space group combinations are no longer equivalent,³⁷ and can be distinguished in several ways, the chief ones being: (a) by statistical tests on the crystallographic R factor; (b) by comparison of the observed F_{hkl} and $F_{\bar{h}\bar{k}\bar{l}}$ with the corresponding calculated structure amplitudes for both the enantiomer A and enantiomer B solutions; and (c) by an assessment of the relative stereochemical plausibilities of the two structures.

The essential features of this part of the refinement are summarized in Table I. Atomic scattering factors for Re are due to Cromer and Waber,³⁸ while, for the lighter atoms, those compiled by Ibers³⁹ were used.

From the data of Table I, it is readily seen that a significance test on the crystallographic R factors⁴⁰ allows rejection, with better than 99.5% confidence, of space group $P4_1$ (enantiomer A) as the correct solution. Further support for $P4_3$ (enantiomer B) as the correct choice is derived from a comparison of the calculated F_{hkl} and $F_{\bar{h}\bar{k}\bar{l}}$ with the corresponding observed structure amplitudes for each model at various stages in the refinement. The agreement for the cases in which there are significant differences between the calculated

(35) D. T. Cromer, *Acta Cryst.*, **18**, 17 (1965).

(36) C. T. Prewitt, Ph.D. Thesis, Massachusetts Institute of Technology, 1962, p 163. The formulas used to calculate the dispersion-corrected real and imaginary parts of the structure factor are (neglecting temperature factors)

$$A = \sum_{\text{all atoms}} \{ (f_0 + \Delta f') \cos \phi - \Delta f'' \sin \phi \}$$

$$B = \sum_{\text{all atoms}} \{ (f_0 + \Delta f') \sin \phi + \Delta f'' \cos \phi \}$$

where $\phi = 2\pi(hx + ky + lz)$.

(37) Indeed, the "enantiomers" are no longer that but are dimensionally different structures.

(38) D. T. Cromer and J. T. Waber, *Acta Cryst.*, **18**, 104 (1965).

(39) J. A. Ibers in "International Tables for X-Ray Crystallography," Vol. III, The Kynoch Press, Birmingham, England, 1962, p 202.

(40) W. C. Hamilton, *Acta Cryst.*, **18**, 502 (1965).

TABLE I

SUMMARY OF THE STEPS IN LEAST-SQUARES REFINEMENTS^a

Remarks	Space group $P4_1$ (enantiomer A)		Space group $P4_3$ (enantiomer B)		Ratio ^c $R_2(P4_1)/R_2(P4_3)$
	R_1	R_2	R_1	R_2	
(1) All atom parameters varied (isotropic temperature factors); equal (unit) weights	0.0556	0.0642	0.0527	0.0609	1.054
(2) All atom parameters varied (anisotropic temperature factors for Re, Cl and P; isotropic temperature factors for other atoms); empirical weighting scheme ^b	0.0456	0.0525	0.0438	0.0503	1.025
(3) Structure factor calculation only; parameters of step 2 used to calculate R_1 and R_2 for index set hkl and partial index set $\bar{h}\bar{k}\bar{l}$	0.0477	0.0576	0.0445	0.0512	1.125

^a Each set of residuals, R_1 and R_2 , listed are those obtained when the indicated type of refinement was cycled to convergence (individual parameter change less than ~ 0.3 esd). ^b The empirical weighting scheme was chosen according to Cruickshank's criterion that $w|F_o| - |F_c|^2$ should be constant. The details are as follows: for $|F_o| < 74.0$, $w^{-1} = 53.2 - 0.488|F_o|$; for $74.0 \leq |F_o| < 235.0$, $w^{-1} = 0.0977|F_o| + 9.77$; for $|F_o| \geq 235.0$, $w^{-1} = 1.07|F_o| - 218.2$. Anisotropic temperature factors were of the form $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. Cf. D. W. J. Cruickshank in "Computing Methods in Crystallography," J. S. Rollett, Ed., Pergamon Press Inc., New York, N. Y., 1965, p 114. ^c For rejection of the hypothesis that enantiomer A is the correct solution a value for this ratio of > 1.001 is necessary.

F_{hkl} 's and $F_{\bar{h}\bar{k}\bar{l}}$'s is excellent for only the $P4_3$ solution. In addition, from a chemical standpoint, the choice of enantiomer B is certainly the more likely one (see Table V and Discussion below).

A final electron density difference map showed no features higher than $0.6 \text{ e}^-/\text{\AA}^3$ (near the rhenium atoms) and $0.3\text{--}0.4 \text{ e}^-/\text{\AA}^3$ elsewhere. Many of the smaller peaks ($\sim 0.3 \text{ e}^-/\text{\AA}^3$) were near positions where hydrogen atoms attached to the phenyl rings would be expected to lie; however, a structure factor calculation including phenyl hydrogen atoms and one cycle of least-squares refinement varying carbon atom parameters only showed no change in R_1 or R_2 . Since a full-matrix refinement varying positional and thermal parameters for all 81 atoms was impossible owing to storage limitations, this phase of the structure solution was abandoned. It is worthwhile to note, however, that the data are probably of good enough quality for a complete refinement if computer limitations were not a factor. Calculation of the structure factors for the rejected reflections gave no $|F_c|$ greater than twice the minimum observable.

Table II lists the observed and calculated structure amplitudes, $10|F_o|$ and $10|F_c|$, respectively, each in electrons; data are included for both the hkl and $\bar{h}\bar{k}\bar{l}$ sets. Final atomic fractional coordinates and isotropic temperature factors are shown in Table III. Anisotropic temperature factors are given in Table IV. Estimated standard deviations quoted in Tables III and IV were obtained from the inverse matrix of the final least-squares refinement cycle.

Results

The structure (enantiomer B, space group $P4_3$) is depicted in Figures 1 and 2, which also show the num-

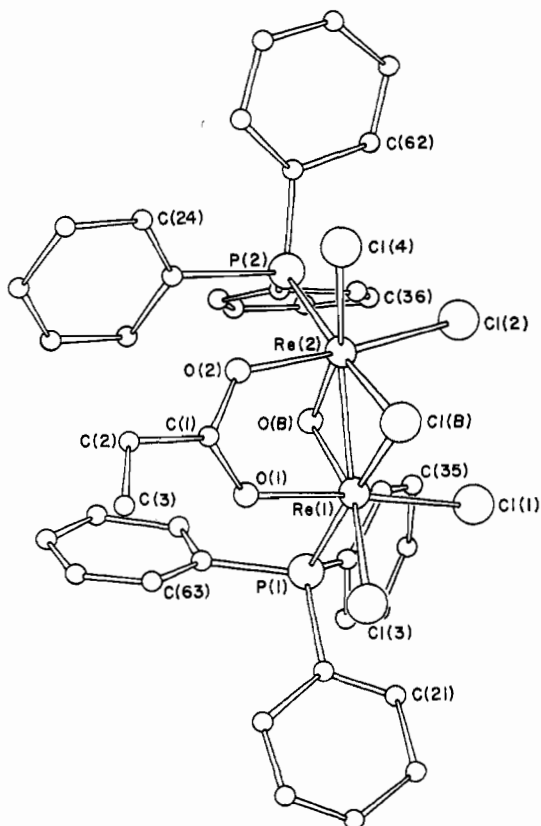


Figure 1.—A perspective drawing of the molecular structure of $\text{Re}_2\text{OCl}_5(\text{C}_2\text{H}_5\text{CO}_2)[\text{P}(\text{C}_6\text{H}_5)_3]_2$. The atom numbering system is also illustrated.

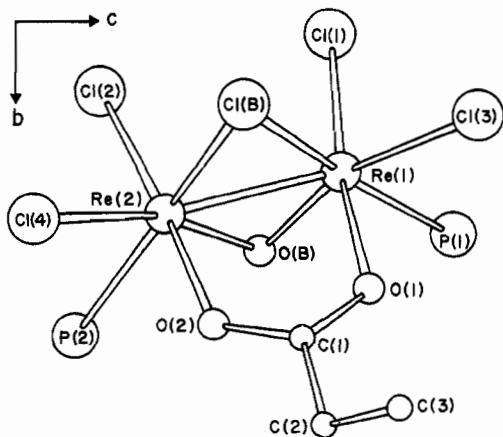


Figure 2.—Projection down $[100]$ of the inner coordination sphere of $\text{Re}_2\text{OCl}_5(\text{C}_2\text{H}_5\text{CO}_2)[\text{P}(\text{C}_6\text{H}_5)_3]_2$.

bering scheme used in Tables V–VII, where the intramolecular dimensions are presented. In the figures, all atoms associated solely with $\text{Re}(1)$ have been assigned odd serial numbers, and all atoms associated solely with $\text{Re}(2)$ have been assigned even serial numbers. Analogously, phenyl rings 1, 3, and 5 are attached to $\text{P}(1)$, while rings 2, 4, and 6 are attached to $\text{P}(2)$. The numbering scheme for the phenyl carbon

atoms is arranged so that each carbon atom has a label $\text{C}(mn)$, in which m is the atom sequence number and n is the ring number. For example, the carbon atoms attached directly to $\text{P}(1)$ are $\text{C}(11)$, $\text{C}(13)$, and $\text{C}(15)$. In Figure 1, only one carbon atom on each ring has been labeled; the other labels may be determined using this label and the foregoing rule.

Table VIII gives the equations for two important planes in the inner coordination sphere and the distances of the atoms from the planes. Root-mean-square vibrational amplitudes and their direction cosines have been omitted from this report, since the observed complex thermal motion of the heavy atoms afforded no useful chemical or structural insight.

Discussion

Effect of Anomalous Scattering.—Ueki, Zalkin, and Templeton⁴¹ have recently discussed the serious coordinate errors that can result from the neglect of the imaginary component, $\Delta f''$, of the anomalous scattering in structure solutions involving polar space groups. As these authors have pointed out, unless the $\Delta f''$ contribution (which is always positive) is considered, the anomalous scatterer will appear to be closer to the X-ray source and detector than it actually is. Consequently, if the origin is fixed by an anomalous scatterer in a polar space group, the other atoms will undergo shifts to compensate for this. In the sense of the crystallographic point symmetry (4, as opposed to Laue symmetry $4/m$) we have collected an incomplete set of data; the point group symmetry demands collection of both index sets hkl and $h\bar{k}l$ (where h , k , and l are all ≥ 0). The coordinate error introduced by neglect of $\Delta f''$ in the refinement of polar structures with this type of incomplete data set has been termed *polar dispersion error*;⁴² further, similar errors, of smaller magnitude, will be introduced if incorrect or inaccurate $\Delta f''$ corrections are used in the refinement.

Cruickshank and McDonald⁴³ have recently discussed polar dispersion errors in detail. In particular, they found that an estimate of Δz , the coordinate error introduced by including $\Delta f''$ in the calculations and choosing the *incorrect* structure model (for example, enantiomer A in our case), is given by

$$\Delta z = \frac{2}{\pi S_{\max}} \left(\frac{\Delta f''}{|f|} \right)_{1/2 S_{\max}}$$

where $S = (2 \sin \theta) / \lambda$, $|f|$ is the modulus of the complex atomic scattering factor, and the quantity $(\Delta f'' / |f|)$ is taken as the phase shift due to anomalous scattering, evaluated at $S = 1/2 S_{\max}$. Use of this formula predicts shifts in the z coordinate (with reference to $\text{Re}(1)$) of 0.081 Å for light atoms, 0.027 Å for chlorine atoms, and 0.041 Å for phosphorus atoms. The average observed shifts were 0.084, 0.028, and 0.036 Å, respectively. A rather surprising fact is that a similar formula predicts

(41) T. Ueki, A. Zalkin, and D. H. Templeton, *Acta Cryst.*, **20**, 836 (1966).

(42) A. Zalkin, T. E. Hopkins, and D. H. Templeton, *Inorg. Chem.*, **5**, 1767 (1966).

(43) D. W. J. Cruickshank and W. S. McDonald, *Acta Cryst.*, **23**, 9 (1967).

TABLE II

OBSERVED AND CALCULATED STRUCTURE AMPLITUDES (X10) (IN ELECTRONS) FOR Re2OCl6(C2H5CO2)[P(C6H5)3]2

Table with multiple columns and rows of numerical data representing observed and calculated structure amplitudes. The table is organized into several sections with varying column widths and signs (+, -, *) indicating phase information.

in this case an error in the z coordinate of an oxygen atom of 0.014 Å if the correction, Δf'' = 0.1 electron, is neglected.

respectively) in the inner coordination sphere, grouped into what one would normally expect to be chemically equivalent pairs. For enantiomer A, the lighter atoms are shifted so that (expected) chemically equivalent bond lengths are different by as much as 0.12 Å; for

TABLE II (Continued)

Table with multiple columns of numerical data, likely representing bond lengths or angles for the complex. The columns are organized in pairs, with each pair corresponding to a different structural parameter. The data is presented in a grid-like format with varying column widths.

enantiomer B, the lighter atoms are shifted in the opposite direction, so that bond lengths of (expected) chemically equivalent pairs are, in general, more nearly equal. It would appear that a greater Δf'' for Re (Cromer's value = 5.37) would bring the bond lengths nearer to equality. It is estimated that a Δf'' in the range 7-8 electrons would reduce the

discrepancies, assuming inherent mirror symmetry for the molecule, to within acceptable limits. Certainly one must exercise caution in making statements of this sort; Figure 1 shows that the chemically equivalent "halves" of the molecule are not without significant structural distortions. This will be discussed further below; however, it seems possible at this point that the

TABLE III

Atom	FINAL ATOMIC COORDINATES ^a FOR $\text{Re}_2\text{OCl}_3(\text{C}_2\text{H}_5\text{CO}_2)_2[\text{P}(\text{C}_6\text{H}_5)_3]_2$			$B,^b \text{ \AA}^2$
	x	y	z	
Re(1)	0.14245 (7)	0.31929 (7)	0.0 ^c	
Re(2)	0.16479 (7)	0.37300 (7)	-0.07005 (3)	
Cl(1)	0.02508 (44)	0.13706 (37)	-0.00442 (16)	
Cl(2)	0.05397 (43)	0.20920 (42)	-0.09641 (14)	
Cl(3)	0.24923 (45)	0.24541 (53)	0.05404 (14)	
Cl(4)	0.30421 (44)	0.38031 (53)	-0.12291 (14)	
Cl(B)	0.30723 (39)	0.23102 (44)	-0.03933 (14)	
P(1)	-0.01730 (40)	0.40266 (40)	0.04460 (12)	
P(2)	0.04027 (41)	0.53269 (42)	-0.10641 (12)	
C(11)	-0.0025 (18)	0.3401 (17)	0.0936 (6)	3.32 (33)
C(21)	-0.0283 (21)	0.2185 (21)	0.0990 (7)	4.52 (42)
C(31)	-0.0128 (25)	0.1606 (25)	0.1358 (8)	5.97 (56)
C(41)	0.0264 (30)	0.2437 (32)	0.1646 (10)	7.60 (73)
C(51)	0.0528 (27)	0.3592 (28)	0.1595 (9)	6.75 (63)
C(61)	0.0408 (20)	0.4214 (21)	0.1228 (6)	4.51 (42)
C(12)	0.0356 (17)	0.5256 (17)	-0.1584 (5)	3.26 (32)
C(22)	0.0047 (19)	0.6340 (19)	-0.1787 (6)	4.08 (39)
C(32)	-0.0067 (25)	0.6206 (25)	-0.2206 (8)	5.81 (55)
C(42)	0.0082 (25)	0.5065 (25)	-0.2370 (7)	5.66 (53)
C(52)	0.0389 (26)	0.3967 (26)	-0.2176 (8)	5.95 (58)
C(62)	0.0486 (21)	0.4063 (20)	-0.1778 (7)	4.44 (42)
C(13)	-0.0111 (17)	0.5711 (17)	0.0465 (5)	3.27 (32)
C(23)	-0.1122 (21)	0.6419 (20)	0.0317 (6)	4.45 (42)
C(33)	-0.1072 (26)	0.7761 (26)	0.0327 (8)	6.13 (58)
C(43)	-0.0009 (26)	0.8348 (26)	0.0482 (9)	6.20 (58)
C(53)	0.1006 (26)	0.7608 (26)	0.0645 (8)	6.22 (57)
C(63)	0.0955 (21)	0.6306 (21)	0.0610 (7)	4.65 (43)
C(14)	0.1110 (17)	0.6793 (17)	-0.0956 (6)	3.30 (34)
C(24)	0.2007 (21)	0.7307 (21)	-0.1205 (7)	4.67 (43)
C(34)	0.2675 (25)	0.8424 (25)	-0.1105 (8)	5.93 (55)
C(44)	0.2302 (28)	0.9049 (27)	-0.0759 (9)	7.03 (65)
C(54)	0.1425 (28)	0.8561 (29)	-0.0511 (9)	6.95 (67)
C(64)	0.0828 (19)	0.7425 (20)	-0.0612 (6)	4.28 (41)
C(15)	-0.1777 (17)	0.3646 (17)	0.0347 (5)	3.07 (32)
C(25)	-0.2202 (16)	0.3359 (15)	-0.0032 (5)	3.06 (30)
C(35)	-0.3441 (18)	0.3112 (19)	-0.0088 (6)	3.85 (38)
C(45)	-0.4315 (19)	0.3161 (18)	0.0211 (6)	3.90 (37)
C(55)	-0.3899 (20)	0.3424 (20)	0.0563 (6)	4.35 (40)
C(65)	-0.2661 (19)	0.3681 (19)	0.0642 (6)	3.96 (38)
C(16)	-0.1300 (18)	0.5383 (19)	-0.0968 (6)	3.77 (36)
C(26)	-0.1962 (20)	0.4278 (20)	-0.1025 (6)	4.25 (40)
C(36)	-0.3337 (21)	0.4337 (22)	-0.1007 (7)	4.70 (43)
C(46)	-0.3869 (23)	0.5432 (23)	-0.0907 (7)	5.27 (49)
C(56)	-0.3199 (25)	0.6477 (25)	-0.0833 (8)	5.87 (54)
C(66)	-0.1882 (20)	0.6528 (20)	-0.0871 (6)	4.29 (41)
C(1)	0.3044 (18)	0.5400 (18)	-0.0159 (6)	3.59 (36)
C(2)	0.3836 (20)	0.6578 (20)	-0.0069 (7)	4.77 (44)
C(3)	0.4746 (23)	0.6333 (23)	0.0237 (7)	5.53 (51)
O(1)	0.2580 (11)	0.4746 (11)	0.0104 (3)	3.06 (22)
O(2)	0.2811 (11)	0.5187 (11)	-0.0514 (4)	3.17 (22)
O(B)	0.0473 (10)	0.4238 (11)	-0.0329 (3)	2.52 (20)

^a Estimated standard deviations (in parentheses) in this and following tables occur in the last significant digit in each case.

^b Thermal parameters for Re, Cl, and P are given in Table IV.

^c Parameter fixed in polar space group.

apparent nonequivalence of the bonds is not completely due to packing considerations.

There appears to be some uncertainty as to whether the value of $\Delta f''$ depends upon the number of anomalous scatterers in the unit cell; however, current thought indicates that this is not an important factor.⁴⁴⁻⁴⁶ The need for more accurate values of $\Delta f''$ and its angular dependence and/or experimental verification of

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(45) J. M. Bijvoet, *ibid.*, **15**, 620 (1962).

(46) S. R. Hall and E. N. Maslen, *ibid.*, **20**, 383 (1966).

current literature values is apparent. Ueki, Zalkin, and Templeton⁴¹ compared their X-ray determination of thorium nitrate pentahydrate with the results of a neutron diffraction study of the same compound⁴⁷ and estimated that $\Delta f''$ should be greater than theoretical values^{35,48} by 1.0-1.6 electrons. On the other hand, there has been experimental confirmation of the $\Delta f''$ values of both light and heavy atoms.⁴⁹

Description of the Structure.—The crystal structure, which consists of normally packed, discrete dinuclear molecules, is defined by the unit cell dimensions, the space group symmetry operations, and the positional parameters given in Table II. Figure 1 shows the over-all molecular structure and numbering scheme; in Figure 2 only the inner coordination about the rhenium atoms is shown. The shorter intermolecular contacts appear to be normal and are not listed; there are no rhenium atom intermolecular contacts at distances of less than 4.5 Å. As in the structure of $\text{ReNCl}_2(\text{P}(\text{C}_6\text{H}_5)_3)_2$,⁵⁰ the crystal packing is mainly dependent upon the intermolecular interactions of the phenyl rings. Principal intramolecular distances and angles, as derived from the data of Table II, are given in Table VI.⁵¹ The effects of different packing environments for the chemically related "halves" of the molecule are seen in the marked discrepancies between related pairs of Cl-Re-P, Re-P-C, C-P-C, and P-C-C_{para} angles.

The molecular structure as a whole possesses neither crystallographic nor virtual symmetry higher than C_1 . However, the inner coordination sphere (Figure 2) has effective mirror symmetry if one ignores atom C(3) of the propionate group and keeps in mind the previous discussion of anomalous scattering. The dinuclear molecular unit consists of coordination octahedra which share an edge. In effect, the short Re-Re bond is triply bridged by a chlorine atom, an oxygen atom, and a propionate group. For the sake of simplicity, the weighted averages of (presumably) chemically equivalent bond lengths will be used in the following discussion of bond lengths. The major structural features may now be discussed.

The Re-Re distance, 2.522 Å, definitely indicates direct metal-to-metal bonding. It is close to the distances 2.43-2.51 Å found in trinuclear halorhenium-(III) structures,^{10,13,17,18,20} in which the rhenium-rhenium bond order is generally agreed^{52,53} to be 2, regardless of whether an MO or a VB analysis of the bonding is used. It is to be noted that this bond length is unaffected by the $\Delta f''$ corrections since both rhenium atoms experience the same phase shift due to anomalous scattering.

Thus, one of the principal questions prompting this

(47) J. C. Taylor, M. H. Mueller, and R. L. Hitterman, *ibid.*, **20**, 842 (1966).

(48) C. H. Dauben and D. H. Templeton, *ibid.*, **8**, 841 (1955).

(49) See ref 44-46 and references therein.

(50) R. J. Doedens and J. A. Ibers, *Inorg. Chem.*, **6**, 204 (1967).

(51) Distances and angles were calculated using updated versions of the programs MGEOM by J. S. Wood and DISTAN by D. P. Shoemaker.

(52) F. A. Cotton and T. E. Haas, *Inorg. Chem.*, **3**, 10 (1964).

(53) J. E. Fergusson, B. R. Penfold, M. Elder, and B. H. Robinson, *J. Chem. Soc.*, 5500 (1965).

TABLE IV
ANISOTROPIC TEMPERATURE FACTORS ($\times 10^5$)

	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}	Equiv isotropic B
Re(1)	540 (7)	554 (7)	52 (1)	80 (5)	-3 (2)	-4 (2)	2.50
Re(2)	511 (6)	598 (7)	50 (1)	-34 (6)	4 (2)	-19 (2)	2.50
Cl(1)	1020 (45)	479 (33)	96 (5)	-68 (32)	43 (13)	8 (11)	3.84
Cl(2)	810 (42)	692 (40)	74 (4)	-169 (33)	-19 (10)	-34 (10)	3.48
Cl(3)	797 (45)	1283 (58)	68 (4)	333 (41)	0 (11)	88 (12)	4.26
Cl(4)	757 (44)	1318 (60)	67 (4)	-73 (40)	64 (10)	-34 (12)	4.24
Cl(B)	637 (37)	794 (45)	72 (4)	151 (34)	-16 (10)	-27 (10)	3.34
P(1)	602 (38)	646 (39)	36 (3)	94 (31)	-12 (8)	-9 (9)	2.48
P(2)	658 (41)	677 (41)	48 (4)	-71 (32)	9 (9)	19 (10)	2.80

TABLE V
VARIATIONS IN PRINCIPAL BOND LENGTHS

	Space group $P4_1$ enantiomer A	Space group $P4_3$ enantiomer B	Wtd av ^a
Re(1)-Cl(1)	2.320	2.318	2.312 (3)
Re(2)-Cl(2)	2.312	2.302	2.312 (3)
Re(1)-Cl(3)	2.321	2.340	2.356 (4)
Re(2)-Cl(4)	2.389	2.371	2.356 (4)
Re(1)-Cl(B)	2.437	2.421	2.411 (3)
Re(2)-Cl(B)	2.388	2.399	2.411 (3)
Re(1)-P(1)	2.451	2.474	2.490 (3)
Re(2)-P(2)	2.523	2.506	2.490 (3)
Re(1)-O(B)	1.948	1.896	1.889 (8)
Re(2)-O(B)	1.827	1.882	1.889 (8)
Re(1)-O(1)	2.075	2.097	2.095 (8)
Re(2)-O(2)	2.066	2.092	2.095 (8)

^a Weighted averages used in this table, Table VII, and the text were computed according to "Tables of Interatomic Distances and Configuration in Molecules and Ions," The Chemical Society, London, 1958, p S1.

study, that is, whether Re-Re bonding would occur in this Re(IV) compound, is clearly answered. Such a bond does occur and has a strength comparable to that of a double bond according to the criterion of distance. However, this remark is not to be taken literally, nor should the metal-metal interaction be oversimplified in any other way. Each Re(IV) species individually has three 5d electrons left after metal-ligand bonding. Since the dinuclear molecule is reported to be diamagnetic,²⁹ a complete description of the metal-metal interaction must account for the pairing up of six electron spins. Besides direct metal-metal bonding, indirect (*i.e.*, superexchange) coupling of spins, intra-atomic coupling of spins, and direct but weak (as in Cu(II) and Cr(II) carboxylate dimers) interatomic coupling of spins are all possible contributors to an over-all description of the metal-metal interaction. It must also be borne in mind that the presence of bridging atoms here undoubtedly imposes constraints on the metal-metal distance so that a bond order-bond length relationship of the type which is valid for species with little or no bridging⁶⁴ cannot be applied uncritically in the present case.

With regard to the bridging ligands, it is notable that in this one molecule there is found each one of the types which occur separately in other polynuclear rhenium species, *viz.*, halogen, oxygen, and carboxyl. The single bridging atoms Cl(B) and O(B) lie in en-

TABLE VI
PRINCIPAL MOLECULAR DIMENSIONS

Bond distances, Å		Intramolecular nonbonded contacts, Å	
Re(1)-Re(2)	2.522 (1)	Re(1)-C(1)	2.973
Re(1)-Cl(1)	2.318 (4)	Re(2)-C(1)	2.994
Re(2)-Cl(2)	2.302 (5)	Re(1)-O(2)	3.151
Re(1)-Cl(3)	2.340 (5)	Re(2)-O(1)	3.169
Re(2)-Cl(4)	2.371 (5)	Re(1)-Cl(2)	3.687
Re(1)-Cl(B)	2.421 (4)	Re(2)-Cl(1)	3.715
Re(2)-Cl(B)	2.399 (4)	Cl(1)-Cl(3)	3.349
Re(1)-P(1)	2.474 (4)	Cl(2)-Cl(4)	3.364
Re(2)-P(2)	2.506 (4)	Cl(1)-P(1)	3.341
Re(1)-O(B)	1.896 (11)	Cl(3)-P(1)	3.318
Re(2)-O(B)	1.882 (11)	Cl(2)-P(2)	3.472
Re(1)-O(1)	2.097 (11)	Cl(4)-P(2)	3.302
Re(2)-O(2)	2.092 (12)	Cl(1)-Cl(2)	3.316
P(1)-C(11)	1.84 (2)	Cl(1)-Cl(B)	3.399
P(1)-C(13)	1.80 (2)	Cl(2)-Cl(B)	3.365
P(1)-C(15)	1.79 (2)	Cl(1)-O(B)	3.226
P(2)-C(12)	1.82 (2)	Cl(2)-O(B)	3.186
P(2)-C(14)	1.78 (2)		
P(2)-C(16)	1.85 (2)		
C(1)-O(1)	1.25 (2)		
C(2)-O(2)	1.28 (2)		
C(1)-C(2)	1.55 (3)		
C(2)-C(3)	1.47 (3)		
Interbond angles, deg			
Cl(1)-Re(1)-P(1)	88.4 (2)	Re(1)-P(1)-C(11)	113.3 (6)
Cl(2)-Re(2)-P(2)	92.4 (2)	Re(1)-P(1)-C(13)	110.9 (6)
Cl(1)-Re(1)-Cl(3)	92.0 (2)	Re(1)-P(1)-C(15)	117.0 (6)
Cl(2)-Re(2)-Cl(4)	92.1 (2)	Re(2)-P(2)-C(12)	119.5 (6)
Cl(3)-Re(1)-P(1)	87.1 (2)	Re(2)-P(2)-C(14)	105.4 (6)
Cl(4)-Re(2)-P(2)	85.2 (2)	Re(2)-P(2)-C(16)	116.8 (7)
P(1)-Re(1)-O(1)	90.7 (3)	C(11)-P(1)-C(13)	108.9 (8)
P(2)-Re(2)-O(2)	88.1 (3)	C(31)-P(1)-C(51)	105.6 (9)
Cl(3)-Re(1)-Cl(B)	88.4 (2)	C(51)-P(1)-C(11)	100.3 (8)
Cl(4)-Re(2)-Cl(B)	88.3 (2)	C(12)-P(2)-C(14)	105.9 (9)
Re(2)-Re(1)-Cl(1)	100.2 (1)	C(14)-P(2)-C(16)	110.5 (9)
Re(1)-Re(2)-Cl(2)	99.6 (1)	C(16)-P(2)-C(12)	98.9 (9)
Cl(B)-Re(2)-O(2)	87.4 (4)	C(1)-C(2)-C(3)	111 (2)
Cl(B)-Re(1)-O(1)	88.8 (3)	P(1)-C(11)-C(41)	178 (1)
Cl(1)-Re(1)-Cl(B)	91.6 (2)	P(1)-C(13)-C(43)	179 (1)
Cl(2)-Re(2)-Cl(B)	91.4 (2)	P(1)-C(15)-C(45)	177 (1)
Re(2)-Re(1)-O(1)	86.1 (3)	P(2)-C(12)-C(42)	175 (1)
Re(1)-Re(2)-O(2)	85.6 (3)	P(2)-C(14)-C(44)	177 (1)
Re(1)-O(B)-Re(2)	83.8 (5)	P(2)-C(16)-C(46)	174 (1)
Re(1)-Cl(B)-Re(2)	63.1 (1)		

vironments similar to those observed in other rhenium cluster compounds. Thus, the Re-Cl(B) distance, 2.411 Å, and Re(1)-Cl(B)-Re(2) angle, 63.1°, are quite similar to the corresponding distances and angles in Re_3Cl_9 ,¹³ for example, which are 2.46 (4) Å and

TABLE VII
 C-C BOND DISTANCES (Å) IN THE TRIPHENYLPHOSPHINE GROUPS^a

C(11)-C(21)	1.34 (3)	C(14)-C(24)	1.40 (3)
C(21)-C(31)	1.43 (3)	C(24)-C(34)	1.43 (3)
C(31)-C(41)	1.40 (4)	C(34)-C(44)	1.44 (3)
C(41)-C(51)	1.28 (5)	C(44)-C(54)	1.38 (4)
C(51)-C(61)	1.45 (3)	C(54)-C(64)	1.41 (4)
C(61)-C(11)	1.41 (3)	C(64)-C(14)	1.41 (2)
C(12)-C(22)	1.40 (3)	C(15)-C(25)	1.43 (2)
C(22)-C(32)	1.47 (3)	C(25)-C(35)	1.36 (3)
C(32)-C(42)	1.36 (4)	C(35)-C(45)	1.40 (2)
C(42)-C(52)	1.39 (4)	C(45)-C(55)	1.33 (2)
C(52)-C(62)	1.40 (3)	C(55)-C(65)	1.38 (3)
C(62)-C(12)	1.45 (3)	C(65)-C(15)	1.40 (2)
C(13)-C(23)	1.41 (3)	C(16)-C(26)	1.39 (3)
C(23)-C(33)	1.43 (4)	C(26)-C(36)	1.47 (3)
C(33)-C(43)	1.40 (4)	C(36)-C(46)	1.35 (3)
C(43)-C(53)	1.46 (4)	C(46)-C(56)	1.35 (4)
C(53)-C(63)	1.40 (4)	C(56)-C(66)	1.41 (3)
C(63)-C(13)	1.40 (3)	C(66)-C(16)	1.41 (3)

^a See footnote *a* in Table V; weighted average = 1.403 (5) Å.

 TABLE VIII
 BEST WEIGHTED LEAST-SQUARES PLANES^a

Atoms	Plane	Eq.
Re(1)-Re(2)-O(1)-O(2)- C(1)-C(2)	1	-0.8030x + 0.5929y + 0.0604z = 0.8000
Re(1)-Re(2)-O(B)-Cl(B)	2	0.5949x + 0.7675y + 0.2388z = 3.5165

Distances of Atoms from Planes (Å)

Atoms	Plane 1	Atoms	Plane 2
Re(1)	-0.001	Re(1)	0.002
Re(2)	0.000	Re(2)	0.000
O(1)	0.014	O(B)	-0.021
O(2)	-0.035	Cl(B)	-0.002
C(1)	-0.025		
C(2)	0.061		

^a Calculated using the method outlined in "International Tables for X-Ray Crystallography," Vol. II, The Kynoch Press, Birmingham, England, 1962, p 93.

61 (2)°. Similarly, the Re-O(B) distance, 1.889 Å, is comparable to the Re-O distance of 1.95 (4) Å, observed in La₄Re₆O₁₉.²⁵ In K₄Re₂OCl₁₀,⁵⁵ where the (55) J. C. Morrow, *Acta Cryst.*, **15**, 851 (1962).

Re-O-Re linkage is linear, the distance is 1.86 (1) Å, and it has been suggested that there is considerable double-bond character.⁵⁶

The carboxylate bridge system differs slightly in some of its dimensions from that found in Re₂(O₂CC₆H₅)₄Cl₂.⁹ In that compound the Re-O distances are 2.018 (5) Å while here they are 2.095 Å. Also, the Re-Re-O angles are slightly acute (86°). Since the Re-Re distance here (2.522 Å) is appreciably greater than that in Re₂(O₂-CC₆H₅)₄Cl₂ (2.235 Å), the acute angles and at least part of the apparent weakening of the Re-O bonds are obviously attributable to the fact that the Re-Re distance somewhat exceeds the optimum "bite" of the carboxyl group. The C-O and C-C distances in the propionate group are normal.

Calculations of the weighted least-squares best planes (Table VIII) show that the two sets of atoms, *viz.*, O(1), O(2), Re(1), Re(2), C(1), C(2) and O(B), Cl(B), Re(1), Re(2), involved in the bridging, each form planar groups with no appreciable deviations. The angle between these planes is 94.7°.

The mean Re-P bond length, 2.490 Å, lies within the range of rhenium-phosphorus bonds in accurately determined structures such as ReNCl₂(P(C₂H₅)₂C₆H₅)₃ (2.490 (5), 2.442 (4), 2.469 (5) Å),⁵⁷ ReNCl₂(P(C₆H₅)₃)₂ (2.448 (2) Å),⁵⁰ and Re₂Cl₆(P(C₂H₅)₃)₂ (2.45 (1) Å).⁸ Within the (C₆H₅)₃P ligands the mean bond distances (P-C, 1.813 (8) Å, and C-C, 1.403 (5) Å) agree well with those in P(C₆H₅)₃ itself.⁵⁸ There are significant distortions in the C-P-C and P-C-C_{para} angles, however, presumably because of packing forces.

The mean Re-Cl distances are very similar to those in other rhenium compounds.^{8,55} The longer bond distance, 2.355 Å, occurs for the Re-Cl bond *trans* to the strong Re-O(B) bond, while the shorter distance, 2.310 Å, occurs for the Re-Cl bond *trans* to a carboxylate oxygen atom.

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(58) J. J. Daly, *J. Chem. Soc.*, 2799 (1964).