

CONTRIBUTION FROM THE CHEMISTRY DEPARTMENT,
BROOKHAVEN NATIONAL LABORATORY, UPTON, NEW YORK 11973Neutron Diffraction Study of Decaborane¹BY ARMIN TIPPE² AND WALTER C. HAMILTON

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A three-dimensional neutron diffraction study on an ordered crystal of deuterated decaborane (¹¹B₁₀²H₁₄) at -160° has been completed. This is the first investigation of a boron hydride by neutron diffraction, and the accuracy of the hydrogen atom positions is greater than that of any X-ray or electron diffraction study. The molecular structure was found to be consistent with the cap model introduced by Kasper, Lucht, and Harker³ (space group C2/a-C_{2h}⁴ with $a = 14.23 \text{ \AA}$, $b = 20.26 \text{ \AA}$, $c = 5.62 \text{ \AA}$, and $\beta = 90.10^\circ$ at -160°; $Z = 8$). The bridging deuterium atoms are in asymmetric positions with respect to the bonding boron atoms in agreement with earlier X-ray investigations. The average values of the B-H bond lengths are 1.180 Å and 1.298, 1.347 Å for the unbridged and bridged hydrogen atoms, respectively. The standard deviations of the bond lengths are 0.004-0.007 Å, and the B-B bond distances and angles are in good agreement with those determined by X-rays. The data indicate that hydrogen-deuterium atom exchange takes place preferentially at the bridged positions. The coherent neutron scattering length for ¹¹B has been refined in an earlier neutron diffraction experiment at room temperature using a crystal exhibiting some disorder. The refined value is $(0.64 \pm 0.01) \times 10^{-12} \text{ cm}$.

Introduction

In 1950, Kasper, Lucht, and Harker (KLH)³ determined the structure of decaborane, B₁₀H₁₄, from X-ray diffraction data. The space group was found to be C2/a-C_{2h}⁴ with eight molecules in a unit cell with dimensions $a = 14.45 \text{ \AA}$, $b = 20.88 \text{ \AA}$, $c = 5.68 \text{ \AA}$, and $\alpha = \beta = \gamma = 90^\circ$. The reflections with h and k both odd were diffuse, suggesting a lack of long-range order in the 001 plane. KLH obtained a satisfactory structure on the basis of a disordered model with a smaller cell ($a' = a/2$, $b' = b/2$, $c' = c$). The space group assumed was Pnnm-D_{2h}¹². More recent X-ray investigations of decaborane derivatives^{4,5} and reanalysis of the original data⁶ on decaborane are in reasonably good agreement regarding B-B bond distances (see Table I); the positions of the hydrogen atoms were not determined to high accuracy in any of these studies.

We have undertaken neutron diffraction studies of both disordered and ordered crystals of decaborane to obtain more accurate hydrogen atom positions and B-H bond lengths. This is the first report of a neutron diffraction study of a boron hydride.

Because of the high absorption cross section of ¹⁰B for thermal neutrons, it seemed advisable to prepare decaborane using the isotope ¹¹B. Furthermore, since the deuterium scattering amplitude is greater than that for ¹H, it was felt that substitution of ¹H by ²H would be helpful in the attainment of high accuracy. Material with the nominal composition ¹¹B₁₀²H₁₄ was used in the experiments. ¹¹B₁₀²H₁₄ crystals, prepared by Dr. M. Dettke and Dr. J. Kurzidim, Fritz-Haber-Institut der Max Planck-Gesellschaft, Berlin, were obtained by sublimation (resulting in disordered crystals)

or crystallization from ether solution (resulting in ordered crystals). In the case of the disordered crystals a mass spectrometric isotopic analysis of the crystal used for the diffraction measurements gave an isotopic composition of 91.2% ²H ($\pm 0.5\%$) and 98.4% ¹¹B; in the case of the ordered crystals the isotopic composition was 80.1% ²H ($\pm 0.5\%$) and 98.4% ¹¹B. The starting material was identical, so there was apparently some H atom exchange in the solution-grown crystals.

A. The Disordered Crystal and Refinement of the Coherent Neutron Scattering Length for ¹¹B

The coherent neutron scattering length for ¹¹B is not well known; values varying from 0.60×10^{-12} to $0.66 \times 10^{-12} \text{ cm}$ have been reported.^{7,8} Therefore the boron scattering length was included as a parameter in a least-squares refinement using three-dimensional neutron diffraction data taken from a decaborane single crystal with the disordered structure. This crystal was used rather than the ordered one because of the smaller percentage of ¹H.

Data Collection and Processing.—A well-formed colorless crystal 2.5 mm long and 1 mm² in cross section was sealed in a quartz capillary and mounted with the c^* axis parallel to the ϕ axis on a four-circle goniometer at the Brookhaven National Laboratory high-flux beam reactor. Using a neutron wavelength of 1.047 Å and a θ - 2θ scan technique, the intensities of 421 independent "sharp" reflections (h and k both even) with a signal-to-background ratio greater than 1:1 were obtained at room temperature. Of these, 165 were measured more than once. For these reflections the agreement factor is $R = 0.06$ ($R = \sum |F - \bar{F}| / \sum |F|$). No correction was made for absorption. The absorption coefficient based upon the known isotopic composition was $\mu = 3.0 \text{ cm}^{-1}$. The maximum error introduced in a structure amplitude by neglect of absorption was 7%.

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(2) Guest from Fritz-Haber-Institut der Max-Planck-Gesellschaft, Berlin, Germany.

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TABLE I
COMPARISON OF BOND DISTANCES IN $B_{10}H_{14}$, $B_{10}H_{13}(C_2H_5)$, $B_{10}H_{13}I$, AND $^{11}B_{10}D_{14}$ (Å)

Distance	$B_{10}H_{14}^a$	$B_{10}H_{13}^b$	$B_{10}H_{13}(C_2H_5)^c$	$B_{10}H_{13}I^d$	$^{11}B_{10}^2H_{14}$
B1-B2	1.74	1.80	1.79	1.79	1.785
B1-B3	1.76	1.77	1.77	1.78	1.788
B1-B5'	1.74	1.77	1.74	1.77	1.754
B2-B3	1.73	1.72	1.72	1.73	1.715
B2-B4	1.76	1.76	1.77	1.80	1.787
B2-B5	1.81	1.80	1.77	1.77	1.782
B2-B5'	1.79	1.78	1.78	1.78	1.774
B3-B4	1.78	1.77	1.79	1.81	1.762
B4-B5	1.74	1.78	1.76	1.74	1.758
B5-B5'	1.78	1.71	1.78	1.75	1.772
B1-B4'	2.01	2.01	1.96	2.01	1.973
B1-D1(H1)	1.26	1.22	0.97		1.173
B2-D2(H2)	1.29	1.27	1.15		1.177
B3-D3(H3)	1.25	1.28	1.12	1.20	1.182
B4-D4(H4)	1.29	1.35	1.13		1.178
B5-D5(H5)	1.28	1.16	1.00		1.192
B1-D6(H6)	1.34	1.39	1.13		1.298
B3-D6(H6)	1.40	1.43	1.34	1.26	1.355
B4-D7(H7)	1.34	1.25	1.18		1.297
B3-D7(H7)	1.40	1.50	1.25		1.339

^a Reference 3. ^b Reference 6. ^c Reference 4. ^d Reference 5.

Least-Squares Refinements.—The least-squares refinements were carried out on a CDC 6600 computer using the BNL version of the Busing–Martin–Levy program.⁹ The function $\sum w(|F_o| - |F_c|)^2$ was minimized, where F_o and F_c are the observed and calculated structure factors. The weights w were assigned in the following way:¹⁰ $w^{1/2} = 1/|F_o|$ if $|F_o| > 4|F_{min}|$; $w^{1/2} = |F|/(16|F_{min}|^2)$ if $|F_o| \leq 4|F_{min}|$, where F_{min} is the smallest observed intensity. The usefulness of this weighting scheme was checked using the program NANOVA (available on request from W. C. Hamilton). The program divided the data into intensity classes and calculated any correlation between the values of $\Delta F/\sigma$ and intensity. In this case $\Delta F/\sigma$ was independent of F , as expected for a proper weighting function. The weighted and unweighted agreement factors

$$R_w = \left[\frac{\sum w ||F_o| - |F_c||^2}{\sum w |F_o|^2} \right]^{1/2}$$

$$R = \frac{\sum ||F_o| - |F_c||}{\sum |F_o|}$$

were used to assess the progress of the refinement.

The least-squares refinement took the structure of KLH as a point of departure. Three cycles of refinement with isotropic temperature factors reduced the values of R_w and R from 0.461 and 0.326 to 0.219 and 0.166.

Anisotropic thermal parameters were now added in the form

$$T_i = \exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$$

Refinement of all parameters except scattering lengths

led to $R_w = 0.085$ and $R = 0.068$. Next, the coherent scattering length for B was included as a parameter. In a full-matrix least-squares refinement (with fixed deuterium scattering factors) the refined value for the boron scattering amplitude is $(0.63 \pm 0.01) \times 10^{-12}$ cm. (Any difference in the distribution of 1H and 2H on the different sites had no significant effect on the scattering amplitudes for the different hydrogen positions. A refinement of all deuterium scattering amplitudes showed these amplitudes to be equal within the limits of $\pm 0.01 \times 10^{-12}$ cm.) This is an average value, calculated from the refined scattering amplitudes for each boron atom. (The maximum difference between two of these scattering amplitudes was 0.01×10^{-12} cm.) The value is relative only, being based on a deuterium scattering amplitude of 0.66×10^{-12} cm,¹¹ a hydrogen scattering amplitude of -0.378×10^{-12} cm, and the isotopic composition. Taking into account the small coherent scattering amplitude of ^{10}B (about 0.1×10^{-12} cm) and the isotopic composition of boron, the refined boron scattering amplitude of $(0.63 \pm 0.01) \times 10^{-12}$ cm corresponds to a value of $(0.64 \pm 0.01) \times 10^{-12}$ cm for ^{11}B . [The value of 0.1×10^{-12} cm for the ^{10}B scattering amplitude is derived from the total B coherent scattering length of 0.60 (S. W. Peterson and H. G. Smith, *J. Phys. Soc. Japan*, **17**, 335 (1962)), the value for ^{11}B determined here, and the known isotopic composition.] The structural parameters from this refinement did not differ significantly from those for the ordered crystal. Since they were somewhat less precisely determined, they are not quoted here. The structure factors for the common reflections agreed well.

B. The Ordered Crystal

Using the principles of derivative structures (Buer-

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ger¹²) KLH determined the space group for the ordered structure to be $C2/a-C_{2h}^4$ (twofold axis in the c direction). There are two possible molecular structures for the ordered state derivable from the disordered structure and consistent with this space group: the cap and the crown model (see KLH). According to the structures of $B_{10}H_{13}I^5$ and $B_{10}H_{13}(C_2H_5)^4$ and consistent with the structures of other boron hydrides, the cap model should be correct for decaborane. However, no direct test using X-ray or neutron diffraction methods has previously been made. Electron diffraction data have been shown to be consistent with the cap model.¹³

Data Collection and Processing.—Because rotating-crystal and precession photographs showed no diffuseness of reflections for h and k both odd (the indication for the disordered structure), the crystal used in this experiment was considered to be well ordered. The crystal (volume 5.7 mm³) was cut from a bigger one and sealed in a quartz capillary in the presence of dry argon gas. The chunky crystal may be described by the Miller indices (hkl) of its bounding planes (not necessarily faces of natural growth) and the perpendicular distances (d) of these planes from an arbitrary origin on the surface of the crystal: h, k, l, d (cm): $-3, 13, 0, 0.13$; $2, 1, 0, 0.00$; $1, -2, 0, 0.00$; $-4, -3, 0, 0.118$; $0, 0, 1, 0.15$; $0, 0, -1, 0.15$.

Using the same technique as described in section A (c^* parallel to ϕ , λ 1.073 Å), the intensities of 1333 independent reflections with a signal-to-background ratio greater than 1:1 were obtained at $-160 \pm 1^\circ$. For 729 reflections, measured more than once, the agreement factor on F^2 was $R = 0.06$ ($R = \Sigma|F^2 - \bar{F}^2| / \Sigma|F^2|$). The low temperature was achieved by blowing a cold nitrogen gas stream at the crystal using the apparatus ALTA described by Rudman and Godel.¹⁴ The temperature was automatically controlled using a thermocouple as sensor. Using the 2θ values of 87 strong reflections with $2\theta > 60^\circ$ the cell parameters at -160° were determined to be $a = 14.23 \pm 0.02$ Å, $b = 20.26 \pm 0.02$ Å, $c = 5.62 \pm 0.01$ Å, $\alpha = \gamma = 90^\circ$ (fixed), and $\beta = 90.10 \pm 0.1$. The least-squares calculation was carried out using the program CELDIM (available from W. C. H.). The slight difference from 90° of the monoclinic angle is real, as Dierks and Dietrich¹⁵ have found a value of 90.07° at room temperature for a single crystal of $^{11}B_{10}^2H_{14}$ from the same preparation using X-ray data.

After applying the usual Lorentz factor, the observed intensities were corrected for absorption. The absorption coefficient based upon the known isotopic composition was $\mu = 3.4$ cm⁻¹. For the extinction correction (applied during the least-square refinement) the approximation introduced by Zachariasen¹⁶ was used. The parameter g defined by Zachariasen was refined as described below.

Least-Squares Refinement.—In the least-squares refinement the function $\Sigma w(F_o^2 - F_c^2)^2$ was minimized. F_o and F_c are the observed and calculated structure factors, respectively. To get a proper weighting scheme, the F_o^2 values were divided into 28 intensity classes and the weight for each class was calculated using the program NANOVA as described in section A. Starting with the positional parameters from the disordered structure, these parameters were refined along with isotropic thermal parameters, a scale factor, and the extinction parameter g for both the cap and the crown models. Two cycles of refinement led to weighted and unweighted R values, R_w and R , of 0.099 and 0.116 for the cap model and 0.261 and 0.436 for the crown model, where $R_w = [\Sigma w(F_o^2 - F_c^2)^2 / \Sigma w F_o^4]^{1/2}$ and $R = \Sigma|F_o^2 - F_c^2| / \Sigma F_o^2$. According to these results the cap model as expected is the right structure model for decaborane. Therefore only this model was used in the final refinement. Taking into account the lack of knowledge about the hydrogen-deuterium distribution, the atomic scattering amplitudes for all seven hydrogen positions were allowed to vary in two cycles of refinement along with the positional parameters, anisotropic thermal parameters, and the extinction parameter g . The final R values (based on F^2) are $R_w = 0.063$ and $R = 0.095$. The refined scattering factors for the hydrogen-deuterium atoms are listed in Table II. The observed and calculated structure factors (F^2) are presented in Table III and the final positional and thermal parameters are in Table IV. The refined value for the extinction parameter is $g = (1.23 \pm 0.16) \times 10^{-4}$. The extinction correction was in general very small; only a few strong reflections gave rise to a correction factor between 0.83 and 0.98. The refined value for the extinction parameter g may be interpreted for the two limiting cases introduced by Zachariasen.¹⁶ In the first case (type I crystals) the extinction is dominated by the mosaic spread. The refined value $g = (1.23 \pm 0.16) \times 10^{-4}$ corresponds to a mosaic spread parameter of 47.4 sec. In the second case (type II crystals), the extinction is dominated by the domain size. Here the calculation led to an average domain size of 1.32×10^{-5} cm.

TABLE II

REFINED SCATTERING AMPLITUDE FOR THE HYDROGEN ATOMS

Atom	Mean scattering factor $\times 10^{12}$, cm	Atom	Mean scattering factor $\times 10^{12}$, cm
D1	0.421 \pm 0.01	D5	0.575 \pm 0.01
D2	0.502 \pm 0.01	D6	0.336 \pm 0.01
D3	0.394 \pm 0.01	D7	0.354 \pm 0.01
D4	0.401 \pm 0.01		

Results

The intramolecular distances and bond lengths are listed in Tables V and VI, respectively. The designation of the atoms follows that of KLH and is shown in Figure 1. The B-B bond distances are in good agreement with those found in earlier X-ray studies (see Table I). A comparison of the B-H distances between the neutron

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TABLE III

OBSERVED AND CALCULATED STRUCTURE FACTORS (F^2)

h	k	l	Observed F^2	Calculated F^2	h	k	l	Observed F^2	Calculated F^2	h	k	l	Observed F^2	Calculated F^2
0	0	0	100	100	1	0	0	100	100	2	0	0	100	100
0	1	0	100	100	1	1	0	100	100	2	1	0	100	100
0	2	0	100	100	1	2	0	100	100	2	2	0	100	100
0	3	0	100	100	1	3	0	100	100	2	3	0	100	100
0	4	0	100	100	1	4	0	100	100	2	4	0	100	100
0	5	0	100	100	1	5	0	100	100	2	5	0	100	100
0	6	0	100	100	1	6	0	100	100	2	6	0	100	100
0	7	0	100	100	1	7	0	100	100	2	7	0	100	100
0	8	0	100	100	1	8	0	100	100	2	8	0	100	100
0	9	0	100	100	1	9	0	100	100	2	9	0	100	100
0	10	0	100	100	1	10	0	100	100	2	10	0	100	100
0	11	0	100	100	1	11	0	100	100	2	11	0	100	100
0	12	0	100	100	1	12	0	100	100	2	12	0	100	100
0	13	0	100	100	1	13	0	100	100	2	13	0	100	100
0	14	0	100	100	1	14	0	100	100	2	14	0	100	100
0	15	0	100	100	1	15	0	100	100	2	15	0	100	100
0	16	0	100	100	1	16	0	100	100	2	16	0	100	100
0	17	0	100	100	1	17	0	100	100	2	17	0	100	100
0	18	0	100	100	1	18	0	100	100	2	18	0	100	100
0	19	0	100	100	1	19	0	100	100	2	19	0	100	100
0	20	0	100	100	1	20	0	100	100	2	20	0	100	100
0	21	0	100	100	1	21	0	100	100	2	21	0	100	100
0	22	0	100	100	1	22	0	100	100	2	22	0	100	100
0	23	0	100	100	1	23	0	100	100	2	23	0	100	100
0	24	0	100	100	1	24	0	100	100	2	24	0	100	100
0	25	0	100	100	1	25	0	100	100	2	25	0	100	100
0	26	0	100	100	1	26	0	100	100	2	26	0	100	100
0	27	0	100	100	1	27	0	100	100	2	27	0	100	100
0	28	0	100	100	1	28	0	100	100	2	28	0	100	100
0	29	0	100	100	1	29	0	100	100	2	29	0	100	100
0	30	0	100	100	1	30	0	100	100	2	30	0	100	100
0	31	0	100	100	1	31	0	100	100	2	31	0	100	100
0	32	0	100	100	1	32	0	100	100	2	32	0	100	100
0	33	0	100	100	1	33	0	100	100	2	33	0	100	100
0	34	0	100	100	1	34	0	100	100	2	34	0	100	100
0	35	0	100	100	1	35	0	100	100	2	35	0	100	100
0	36	0	100	100	1	36	0	100	100	2	36	0	100	100
0	37	0	100	100	1	37	0	100	100	2	37	0	100	100
0	38	0	100	100	1	38	0	100	100	2	38	0	100	100
0	39	0	100	100	1	39	0	100	100	2	39	0	100	100
0	40	0	100	100	1	40	0	100	100	2	40	0	100	100
0	41	0	100	100	1	41	0	100	100	2	41	0	100	100
0	42	0	100	100	1	42	0	100	100	2	42	0	100	100
0	43	0	100	100	1	43	0	100	100	2	43	0	100	100
0	44	0	100	100	1	44	0	100	100	2	44	0	100	100
0	45	0	100	100	1	45	0	100	100	2	45	0	100	100
0	46	0	100	100	1	46	0	100	100	2	46	0	100	100
0	47	0	100	100	1	47	0	100	100	2	47	0	100	100
0	48	0	100	100	1	48	0	100	100	2	48	0	100	100
0	49	0	100	100	1	49	0	100	100	2	49	0	100	100
0	50	0	100	100	1	50	0	100	100	2	50	0	100	100
0	51	0	100	100	1	51	0	100	100	2	51	0	100	100
0	52	0	100	100	1	52	0	100	100	2	52	0	100	100
0	53	0	100	100	1	53	0	100	100	2	53	0	100	100
0	54	0	100	100	1	54	0	100	100	2	54	0	100	100
0	55	0	100	100	1	55	0	100	100	2	55	0	100	100
0	56	0	100	100	1	56	0	100	100	2	56	0	100	100
0	57	0	100	100	1	57	0	100	100	2	57	0	100	100
0	58	0	100	100	1	58	0	100	100	2	58	0	100	100
0	59	0	100	100	1	59	0	100	100	2	59	0	100	100
0	60	0	100	100	1	60	0	100	100	2	60	0	100	100
0	61	0	100	100	1	61	0	100	100	2	61	0	100	100
0	62	0	100	100	1	62	0	100	100	2	62	0	100	100
0	63	0	100	100	1	63	0	100	100	2	63	0	100	100
0	64	0	100	100	1	64	0	100	100	2	64	0	100	100
0	65	0	100	100	1	65	0	100	100	2	65	0	100	100
0	66	0	100	100	1	66	0	100	100	2	66	0	100	100
0	67	0	100	100	1	67	0	100	100	2	67	0	100	100
0	68	0	100	100	1	68	0	100	100	2	68	0	100	100
0	69	0	100	100	1	69	0	100	100	2	69	0	100	100
0	70	0	100	100	1	70	0	100	100	2	70	0	100	100
0	71	0	100	100	1	71	0	100	100	2	71	0	100	100
0	72	0	100	100	1	72	0	100	100	2	72	0	100	100
0	73	0	100	100	1	73	0	100	100	2	73	0	100	100
0	74	0	100	100	1	74	0	100	100	2	74	0	100	100
0	75	0	100	100	1	75	0	100	100	2	75	0	100	100
0	76	0	100	100	1	76	0	100	100	2	76	0	100	100
0	77	0	100	100	1	77	0	100	100	2	77	0	100	100
0	78	0	100	100	1	78	0	100	100	2	78	0	100	100
0	79	0	100	100	1	79	0	100	100	2	79	0	100	100
0	80	0	100	100	1	80	0	100	100	2	80	0	100	100
0	81	0	100	100	1	81	0	100	100	2	81	0	100	100
0	82	0	100	100	1	82	0	100	100	2	82	0	100	100
0	83	0	100	100	1	83	0	100	100	2	83	0	100	100
0	84	0	100	100	1	84	0	100	100	2	84	0	100	100
0	85	0	100	100	1	85	0	100	100	2	85	0	100	100
0	86	0	100	100	1	86	0	100	100	2	86	0	100	100
0	87	0	100	100	1	87	0	100	100	2	87	0	100	100
0	88	0	100	100	1	88	0	100	100	2	88	0	100	100
0	89	0	100	100	1	89	0	100	100	2	89	0	100	100
0	90	0	100	100	1	90	0	100	100	2	90	0	100	100
0	91	0	100	100	1	91	0	100	100	2	91	0	100	100
0	92	0	100	100	1	92	0	100	100	2	92	0	100	100
0	93	0	100	100	1	93	0	100	100	2	93	0	100	100
0	94	0	100	100	1	94	0	100	100	2	94	0	100	100
0	95	0	100	100	1	95	0	100	100	2	95	0	100	100
0	96	0	100	100	1	96	0	100	100	2	96	0	100	100
0	97	0	100	100	1	97	0	100	100	2	97	0	100	100
0	98	0	100	100	1	98	0	100	100	2	98	0	100	100
0	99	0	100	100	1	99	0	100	100	2	99	0	100	100
0	100	0	100	100	1	100	0	100	100	2	100	0	100	100

TABLE IV
ATOMIC PARAMETERS^a

Atom	<i>x</i>	<i>y</i>	<i>z</i>	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
B1	0.0336 (2)	0.3317 (1)	0	15 (1)	8 (1)	90 (6)	-2 (1)	0	0
B2	0.0993 (2)	0.2738 (2)	0.1738 (5)	13 (1)	8 (1)	77 (8)	1 (1)	-5 (3)	-6 (2)
B3	0.1192 (2)	0.2778 (2)	-0.1277 (5)	11 (2)	12 (1)	95 (8)	-3 (1)	7 (3)	-1 (2)
B4	0.0976 (2)	0.1998 (1)	0	12 (1)	7 (1)	104 (7)	0 (1)	0	0
B5	0.0198 (2)	0.2086 (1)	0.2407 (5)	16 (1)	7 (1)	77 (7)	1 (1)	3 (2)	0 (1)
D1	0.0456 (3)	0.3889 (2)	0	28 (2)	9 (1)	251 (16)	0 (1)	0	0
D2	0.1591 (3)	0.2874 (2)	0.3093 (7)	18 (2)	17 (1)	178 (14)	4 (1)	-8 (4)	-10 (3)
D3	0.1853 (4)	0.2946 (3)	-0.2395 (6)	19 (3)	18 (2)	207 (19)	-6 (2)	17 (7)	10 (5)
D4	0.1565 (3)	0.1590 (2)	0	15 (2)	12 (1)	237 (15)	4 (1)	0	0
D5	0.0353 (3)	0.1747 (2)	0.4097 (6)	37 (2)	13 (1)	165 (11)	1 (1)	0 (4)	15 (3)
D6	0.0452 (4)	0.3117 (3)	-0.2176 (5)	21 (4)	12 (2)	64 (18)	-3 (2)	-13 (6)	0 (3)
D7	0.0883 (4)	0.2205 (2)	-0.2172 (5)	20 (3)	13 (2)	75 (16)	-4 (2)	-2 (5)	-8 (3)

^a The coordinates are expressed as fractions of the cell parameters. The β_{ij} 's are multiplied by 10^4 . Estimated standard deviations in the last significant figure are in parentheses.

TABLE V
INTRAMOLECULAR DISTANCES AND THEIR ESD'S FOR
¹¹B₁₀D₁₄ (80.1% D)

Atoms	Distances, Å	σ , Å
Interatomic Distances		
B1-B2	1.785	0.005
B1-B3	1.788	0.005
B1-B5'	1.754	0.004
B2-B3	1.715	0.004
B2-B4	1.787	0.005
B2-B5	1.782	0.005
B2-B5'	1.774	0.005
B3-B4	1.762	0.005
B4-B5	1.758	0.004
B5-B5'	1.772	0.006
B4'-B1	1.973	0.004
Unbridged Hydrogen Bonds		
B1-D1	1.173	0.005
B2-D2	1.177	0.006
B3-D3	1.182	0.007
B4-D4	1.178	0.005
B5-D5	1.192	0.004
Bridged Hydrogen Bonds		
B1-D6	1.298	0.005
B3-D6	1.355	0.007
B4-D7	1.297	0.005
B3-D7	1.339	0.007

TABLE VI
INTRAMOLECULAR BOND ANGLES

Atoms	Angle, deg	σ , deg
B2-B1-B5'	60.2	0.2
B2-B1-D1	125.0	0.3
B3-B1-B2	57.4	0.2
B3-B1-B4'	116.6	0.2
B3-B1-D1	120.3	0.3
B3-B1-D6	49.0	0.3
B3-B1-B5'	108.8	0.2
B4'-B1-D1	117.2	0.3
B4'-B1-B5'	55.8	0.2
B5'-B1-D1	121.5	0.2
D1-B1-D6	106.8	0.3
B1-B2-B3	61.4	0.2
B1-B2-B5'	59.0	0.2
B1-B2-D2	125.3	0.4
B3-B2-B4	60.4	0.3
B3-B2-D2	120.7	0.4
B4-B2-B5	59.0	0.2
B4-B2-D2	124.0	0.4
B5-B2-B5'	59.8	0.2
B5-B2-D2	119.5	0.4
B5'-B2-D2	120.2	0.4
B1-B3-B2	61.2	0.3
B1-B3-D3	125.5	0.5
B1-B3-B4	105.3	0.2
B1-B3-D6	46.3	0.3
B2-B3-B4	61.8	0.3
B2-B3-D3	132.0	0.4
B4-B3-D3	127.9	0.5
B4-B3-D7	47.0	0.3
D3-B3-D6	106.0	0.5
D3-B3-D7	108.1	0.5
B1'-B4-B5	55.6	0.2
B1'-B4-D4	116.5	0.3
B1'-B4-D7	90.6	0.3
B2-B4-B3	57.8	0.2
B2-B4-B5	60.3	0.2
B2-B4-D4	125.5	0.3
B2-B4-D7	104.0	0.3
B3-B4-D4	120.3	0.3
B3-B4-D7	49.1	0.3
B5-B4-D4	121.3	0.2
D4-B4-D7	107.3	0.3
B1'-B5-B4	68.4	0.2
B1'-B5-D5	115.3	0.3
B2-B5-B4	60.6	0.2
B2-B5-B5'	59.9	0.2
B2-B5-D5	118.8	0.3
B4-B5-D5	116.1	0.3
B5'-B5-D5	127.2	0.2

and X-ray studies is not so useful because of the low accuracy of the X-ray results. Nevertheless, this study confirms the slight *but highly significant* asymmetry in the positions of the bridging hydrogen atoms H6 and H7. These atoms are further away from the terminal boron B3 than from B1 and B4. The average values of 1.180 Å for the unbridged B-H distance and 1.298, 1.347 Å for the bridged B-H distances are in good agreement with those found in pentaborane¹⁷ (1.20 and 1.35 Å for the unbridged and bridged B-H distances, respectively). Table VII gives the shortest nonbonded intramolecular and the shortest intermolecular distances, respectively. These values also agree with those found in pentaborane. The bond angles in the boron framework agree with those found in B₁₀H₁₃I.⁵ The folding angle

(17) W. F. Dulmage and W. N. Lipscomb, *Acta Cryst.*, **5**, 260 (1952).

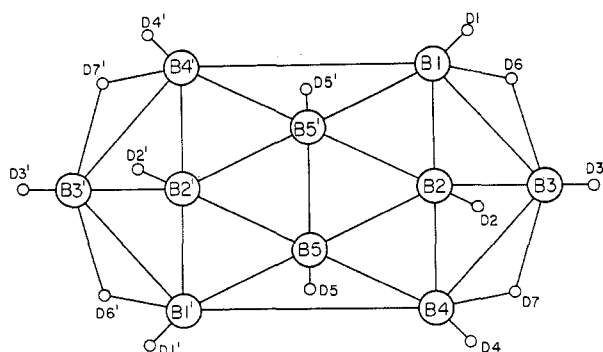


Figure 1.—Topological diagram of decaborane.

TABLE VII
SHORTEST NONBONDED INTRAMOLECULAR AND INTERMOLECULAR DISTANCES

Atoms	Nonbonded intramolecular distances, Å
D6-D1	1.985 (0.006)
D6-D3	2.029 (0.008)
D6-D7	1.949 (0.008)
D6-D7'	2.009 (0.009)
D7-D3	2.043 (0.008)
D7-D4	1.994 (0.006)

Atoms	Shortest intermolecular distances, Å
D7-D5 ^a	2.461 (0.006)
B3-D2 ^a	3.300 (0.006)

^a In each case the second atom is in the molecule related to the basic molecule by a translation of $-c$.

of the decaborane molecule between the two planes B5-B5'-B1 and B5-B5'-B1' is $71.8 \pm 0.3^\circ$, compared to the 76° found by KLH. This difference may be partly due to the fact that the data were collected at different temperatures. A calculation of this angle using the data from the crystal with disordered structure (at room temperature) gave $73.2 \pm 0.6^\circ$.

According to Table II, hydrogen atoms are preferentially at bridge positions. Using the neutron coherent scattering lengths for ^2H and ^1H , 0.66×10^{-12} and -0.378×10^{-12} cm, respectively, the average percentage is 30 and 19% ^1H at bridged and unbridged positions, respectively. (Nevertheless, in the tables and figures we refer to the hydrogen atoms as D.) The crystal was grown in ethereal solution from material with an average isotopic composition of 8.8% ^1H and 91.2% ^2H . Apparently during the crystallization a deuterium-hydrogen exchange took place. According to the isotopic composition calculated from the refined scattering lengths of the hydrogen-deuterium atoms, such exchange takes place preferentially at the bridge position. This is in agreement with nmr experiments¹⁸ and charge distribution calcula-

tions.¹⁹ The significant differences in the refined scattering factors for the unbridged hydrogen positions (see Table II) are also qualitatively understandable by considering the charge distributions for the boron atoms in decaborane as calculated by Lipscomb¹⁹ and listed in Table VIII. From these values one would expect that the strongest B-H bonds are B5-D5 and B2-D2. Therefore the ^2H - ^1H exchange should be least at these positions. The opposite is true for the B3-D3 bond. The strengths of the B1-D1 and B4-D4 bonds are equal and between those for B3-D3 and B2-D2. The refined scattering factors listed in Table II confirm these expectations.

TABLE VIII
CHARGE DISTRIBUTIONS IN DECABORANE¹⁹

Atom	Charge
B1	0.12
B2	-0.10
B3	0.33
B3	0.12
B5	-0.04

In Table IX the rms components of the thermal motion (in Å) along the principal axes of the thermal displacement ellipsoids are listed. These values have been calculated from the thermal parameters listed in Table II, using the BNL version of the ORFFE program.²⁰ For the unbridged deuterium atoms, the direction of maximum vibration is perpendicular to the B-H bonds as expected; for the bridged deuterium atoms the maximum thermal displacement is in a plane parallel to the

TABLE IX
ATOMIC THERMAL DISPLACEMENTS^a

Atom	r_1	σ	r_2	σ	r_3	σ
B1	11.5	0.6	12.0	0.4	13.9	0.5
B2	10.1	0.7	10.9	0.8	13.9	0.7
B3	9.6	0.9	12.7	0.6	16.2	0.6
B4	11.2	0.6	12.4	0.4	12.9	0.4
B5	10.9	0.5	11.7	0.6	13.2	0.6
D1	13.8	0.8	17.0	0.8	20.0	0.6
D2	11.9	1.1	16.8	0.8	19.7	0.8
D3	11.0	1.5	18.8	1.1	20.8	1.1
D4	11.3	1.1	16.4	0.8	19.5	0.6
D5	13.4	0.8	18.6	0.6	19.6	0.7
D6	8.9	1.8	14.0	1.3	16.9	1.3
D7	9.9	1.4	13.5	1.3	17.6	1.1

^a The rms amplitudes r_1 , r_2 , r_3 are parallel to the principal axes of the thermal displacement ellipsoid (units: Å $\times 10^{-2}$).

B-B bond belonging to the bridge. Normal to this plane the vibrations are essentially isotropic. One can see this clearly in Figure 2 which shows a stereoview of the decaborane molecule. Figure 3 shows a stereoview of the unit cell of decaborane. In order to give a clear view of the packing of the molecules, the hydrogen atoms are not included.

(19) W. N. Lipscomb, "Boron Hydrides," W. A. Benjamin, Inc., New York, N. Y., 1963.

(20) W. R. Busing, K. L. Martin, and H. A. Levy, Report ORNL-TM-306, Oak Ridge National Laboratory, 1964.

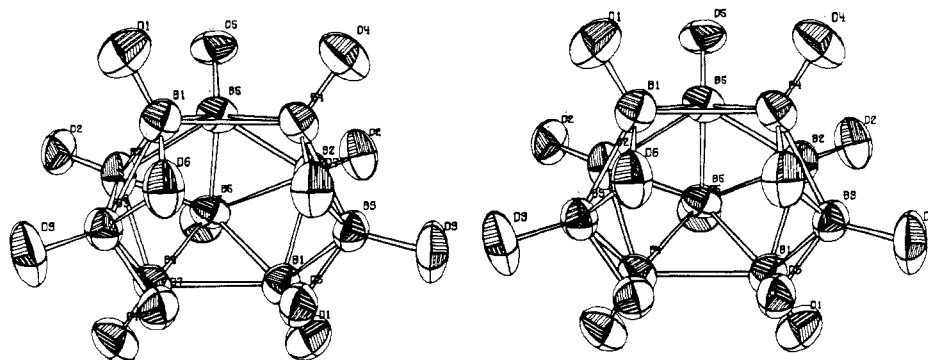


Figure 2.—Stereoview of the decaborane molecule. The ellipsoids indicate the magnitudes of thermal vibration.

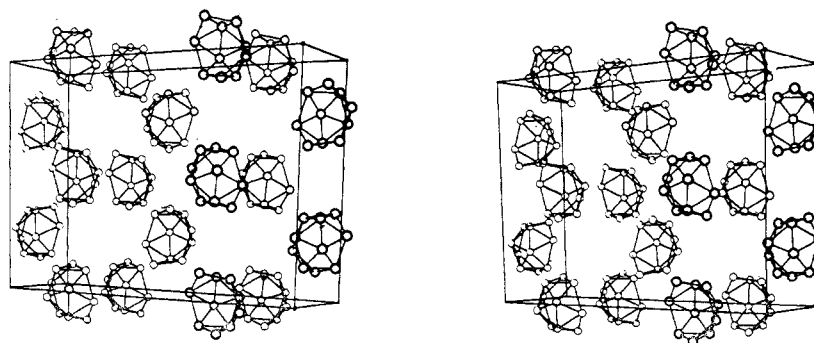


Figure 3.—Packing diagram of the decaborane unit cell (stereoview). The hydrogen atoms are not included.

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CONTRIBUTION FROM THE DEPARTMENTS OF CHEMISTRY, MUHLENBERG COLLEGE, ALLENTOWN, PENNSYLVANIA 18104, AND LEHIGH UNIVERSITY, BETHLEHEM, PENNSYLVANIA 18015

Metal Complexes of the Dibenzamido Anion

BY DAVID N. STEHLY AND CHARLES S. KRAIHANZEL

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Several new metal complexes of the dibenzamido anion, $M(C_6H_5CONCOC_6H_5)_n$, where $M = Be(II), Al(III), Fe(III), Cu(II),$ and $Hg(II)$, have been synthesized. The compounds were characterized by elemental analysis and from their infrared spectra. A comparison of the electronic properties of some isoelectronic and isostructural dibenzamido and dibenzoylmethano complexes is presented.

Introduction

Tris(dibenzamido)chromium(III), $Cr(dba)_3$, and tris(diacetamido)chromium(III), $Cr(daa)_3$, were reported recently as the first well-characterized complexes of diacylamido anions.¹ Further attempts to obtain other new diacetamido-metal complexes have been generally unsuccessful. However, several new complexes of the dibenzamido anion have been isolated and characterized and will be reported on at this time.²

(1) C. S. Kraihanzel and D. N. Stehly, *Inorg. Chem.*, **6**, 277 (1967).

(2) D. N. Stehly and C. S. Kraihanzel, presented in part at the 3rd Middle Atlantic Regional Meeting of the American Chemical Society, Philadelphia, Pa., Feb 1968.

Experimental Section

Reagents.—Anhydrous beryllium chloride was obtained from Alfa Inorganics, Inc. Anhydrous chromium(III) chloride³ and iron(III) chloride⁴ were prepared in the laboratory. Anhydrous aluminum chloride (Baker and Adamson) was sublimed before use. Copper(II) acetate monohydrate and mercury(II) acetate were purchased from the J. T. Baker Chemical Co. The preparation of dibenzamide has already been reported.¹ Dibenzoylmethane (Eastman Kodak Co.) was used as received.

Chloroform, *n*-hexane, and anhydrous diethyl ether were used as received from J. T. Baker Chemical Co. ("Baker

(3) G. B. Heisig, B. Fawkes, and R. Hedin, *Inorg. Syn.*, **2**, 193 (1946).

(4) A. R. Pray, *ibid.*, **5**, 154 (1957).