

The Structure of $K_2BaCo(NO_2)_6$ at 233°K.: A Static Jahn-Teller Distortion

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The structure of $K_2BaCo(NO_2)_6$ has been determined at 233°K; the crystals, which are cubic at room temperature, are orthorhombic, space group $Fmmm$, at 233°K with $a=10.51(3)$ Å, $b=10.68(3)$ Å, and $c=10.51(3)$ Å. The hexanitrocobaltate(II) ion, which appears to be octahedral with a cobalt-nitrogen distance of 1.98(2) Å at room temperature, is elongated along one axis to give cobalt-nitrogen distances of 1.90(2) Å, 1.94(6) Å, and 2.10(2) Å. The results are indicative of a dynamic Jahn-Teller distortion at room temperature and a static Jahn-Teller distortion at 233°K; the elongation along one axis is the type of distortion observed for most complexes of copper(II).

Introduction

In a recent paper,¹ we reported the structure of $K_2BaCo(NO_2)_6$; although a Jahn-Teller distortion² was expected for the low-spin d^7 complex, $Co(NO_2)_6^{4-}$, the structure study gave no evidence of such a distortion. The nitrite groups along the x, y, and z axes of the complex were related by a three-fold axis of the cubic space group ($Fm\bar{3}$) and, thus, no static distortion was possible. Although a dynamic Jahn-Teller effect could not be ruled out, the temperature parameters did not seem unusual when compared to values from other structure; however, an adequate comparison could not be made since values were not available for similar compounds which would not be subject to a Jahn-Teller distortion. Similar results from a neutron diffraction study³ have been reported for $K_2PbCu(NO_2)_6$, which should also show a Jahn-Teller distortion.

ESR and X-ray diffraction studies⁴ of powdered samples of $K_2PbCu(NO_2)_6$ have provided evidence for a static Jahn-Teller effect at 235°K. The ESR spectrum changed from isotropic at room temperature to anisotropic at the lower temperature; the x-ray data indicated a transition from cubic to tetragonal. ESR data⁵ have indicated a similar change in $K_2BaCo(NO_2)_6$.

In order to obtain more information about the Jahn-Teller distortion in $K_2BaCo(NO_2)_6$, we have determined the crystal structure at 233°K.

Experimental Section

An octahedral-shaped crystal of $K_2BaCo(NO_2)_6$, ca. 0.1 mm from apex-to-apex, was mounted on a precession camera equipped with a Cryo-Tip refrigeration system.⁶ An x-ray diffraction photograph of the $hk0$ zone was obtained at room temperature; the crystal was then cooled and additional photographs were obtained at 10° intervals. At 263°K, a transition occurred; the $hk0$ zone still showed mm symmetry and the same systematic absences ($h+k=2n+1$ absent) but the two axes were no longer equal ($a=10.51(3)$ Å, $b=10.86(3)$ Å). The $0kl$ zone at 263°K also showed mm symmetry with $c=10.51(3)$ Å. Although two axes were found to be of equal length and the powder diffraction pattern of the crystal could thus be indexed as tetragonal, the intensities of reflections in the two zones were not consistent with the four-fold symmetry necessary for a tetragonal cell. The diffraction symmetry and the systematic extinction of hkl reflections with odd values of $(h+k)$, $(k+l)$, or $(h+l)$ indicated a face-centered orthorhombic lattice. Since removal of the three-fold axis of the cubic space group, $Fm\bar{3}$, found at room temperature would give orthorhombic space group $Fmmm$, that space group was assumed; the successful refinement of the structure confirms that choice.

Photographs taken below the transition temperature always showed a second lattice of lower intensity; the second lattice was identical to the first but was elongated along a rather than along b and was rotated ca 5° with respect to the first. On warming the crystal to room temperature, the diffraction photographs again showed only one cubic lattice; the same result was obtained repeatedly with several different crystals.

Intensity data were collected with the precession camera, using Zr-filtered Mo $K\alpha$ radiation, and estimated visually; although there were two lattices, the reflections of the two were sufficiently separated to allow estimation of the intensities of the more intense lattice. A total of 244 independent, non-zero reflections from the hkl ($l=0-3$) and hkl ($h=0-3$) levels were collected; the intensities were corrected for Lorentz and polarization effects but no absorption corrections were made.

Computations were carried out on a Burroughs B-5500 computer and on a Univac 1108 computer; programs employed included modified versions of the

(1) J. A. Bertrand and D. A. Carpenter, *Inorg. Chem.*, 5, 415 (1966).

(2) H. A. Jahn and E. Teller, *Proc. Roy. Soc. (London)* A161, 220 (1937).

(3) N. W. Isaacs and C. H. L. Kennard, *J. Chem. Soc. (A)*, 386 (1969).

(4) H. Elliott, B. J. Hathaway, and R. C. Slade, *Inorg. Chem.*, 5, 669 (1966).

(5) B. J. Hathaway, Private Communication.

(6) Air Products and Chemicals, Inc., Advanced Products Dept., Allentown, Penn

Busing-Martin-Levy ORFLS and ORFFE least-squares and function and error programs and various locally written programs. In structure factor calculations, the scattering factors tabulated by Ibers⁷ were employed for all atoms.

All atoms were assigned coordinates from the room temperature cubic structure; because of the lower symmetry the nitrogen and oxygen atoms along each crystallographic axis were no longer equivalent and could be varied independently. Least-squares refinement using isotropic temperature factors and with all reflections weighted at unity converged to a conventional R value ($R = \sum ||F_o| - |F_c|| / \sum |F_o|$) of 0.09. Final coordinates and thermal parameters are listed in Table I; final F_o and F_c values are listed in Table II.

Table I. Final positional and thermal parameters for $K_2BaCo(NO_3)_6$

Atom	x	y	z	B(Å ²)
K	1/4	1/4	1/4	2.4(1)
Ba	1/2	0	0	1.1(1)
Co	0	0	0	0.8(1)
N1	0	0	0.184(6)	2.2(5)
N2	0	0.193(2)	0	1.1(4)
N3	0.181(2)	0	0	0.3(4)
O1	0.099(1)	0	0.249(3)	2.5(3)
O2	0	0.255(2)	0.107(3)	2.1(3)
O3	0.246(2)	0.098(1)	0	1.8(3)

Table II. Observed and calculated structure factors (electrons) for $K_2BaCo(NO_3)_6$

h	k	l	F _o	F _c	h	k	l	F _o	F _c	h	k	l	F _o	F _c
0	0	0	222 243	222 243	10	4	2	222 243	222 243	11	7	3	211 74	211 74
0	0	2	102 96	102 96	10	4	2	102 96	102 96	11	7	3	102 96	102 96
0	2	0	102 96	102 96	10	4	2	102 96	102 96	11	7	3	102 96	102 96
2	0	0	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
2	0	2	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
2	2	0	108 90	108 90	10	4	2	108 90	108 90	11	7	3	108 90	108 90
2	2	2	108 90	108 90	10	4	2	108 90	108 90	11	7	3	108 90	108 90
4	0	0	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
4	0	2	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
4	2	0	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
4	2	2	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
6	0	0	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
6	0	2	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
6	2	0	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
6	2	2	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
8	0	0	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
8	0	2	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
8	2	0	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
8	2	2	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
10	0	0	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
10	0	2	110 96	110 96	10	4	2	110 96	110 96	11	7	3	110 96	110 96
10	2	0	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301
10	2	2	286 301	286 301	10	4	2	286 301	286 301	11	7	3	286 301	286 301

Discussion

The only change in the hexanitrocobaltate(II) ion in going from room temperature to 233°K is in the lengths of the cobalt-nitrogen bonds. At room temperature, all of the cobalt-nitrogen bonds are equivalent and 1.98(2) Å in length; below the transition point the three independent cobalt-nitrogen bond lengths are 1.94(6) Å, 2.10(2) Å, and 1.90(2) Å along the x, y, and z axes, respectively. Although the standard deviation of the bond along x is extremely high, the complex appears to be elongated along one axis in the same way as most copper(II) complexes. The magnitude of the distortion agrees well with theoretical predictions.⁸

(7) J. A. Ibers in «International Tables for X-Ray Crystallography», Vol. 3, The Kynoch Press, Birmingham, England, 1962.

(8) A. D. Liehr and C. J. Ballhausen, *Annals of Physics*, 3, 304 (1958).

Table III.

Atoms	This Work	Reference 1
Co-N1	1.93(6)Å	
Co-N2	2.10(2)	
Co-N3	1.91(2)	
Co-N(Avg)	1.98(8)	1.98(2)Å
N1-O1	1.25(4)	
N2-O2	1.31(3)	
N3-O3	1.26(2)	
N-O(Avg)	1.27(3)	1.23(2)
Ba-O1	2.84(2)	
Ba-O2	2.89(2)	
Ba-O3	2.87(2)	
Ba-O(Avg)	2.87(2)	2.87(2)
O1-N1-O1'	113(5)°	
O2-N2-O2'	118(3)	
O3-N3-O3'	115(2)	
O-N-O'(Avg)	115(3)	114(2)°

Three possible explanations⁴ for the lack of distortion in the hexanitrocobaltate(II) ion at room temperature have been given: (1) the environment of each ion is tetragonal with the unique axis in one of three mutually perpendicular directions, the distribution of these directions throughout the lattice being random; (2) the ion environments are undergoing a dynamic tetragonal distortion in three perpendicular directions in the lattice; (3) the statically distorted hexanitrocobaltate(II) anions are undergoing free rotation in the

lattice. The third possibility is not consistent with the observation of discrete nitrogen and oxygen positions at room temperature and the first possibility is not consistent with a transition to an ordered structure as the temperature is lowered; thus, only the second possibility—a dynamic tetragonal distortion—is consistent with all of the observations.

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