# *An* **X-ray and Neutron Diffraction Structure Analysis of a Triply-bridged Binuclear Iridium Complex,**  $[(C_5(CH_3)_5Ir)_2(\mu-H)_3]^+ [ClO_4]^- \cdot 2C_6H_6]$

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### **Abstract**

The structure of  $[(C_5(CH_3)_5]r)_2(\mu-H)_3]^+[ClO_4]^{-1}$  $2C_6H_6$ ] has been studied with X-ray and neutron diffraction techniques at room temperature and 21 K, respectively. Crystal data from the X-ray diffraction analysis of the title compound at room temperature: space group *Pnnm; a =* 13.358(4), *b =* 14.083(4),  $c = 8.846(3)$  Å;  $V = 1664.3(9)$  Å<sup>3</sup>;  $Z = 2$ . Final agreement factors are  $R(F) = 0.026$  and  $R(wF) = 0.034$ for 1007 independent X-ray reflections with  $I > 3\sigma(I)$ and 130 parameters varied. Crystal data for the neutron diffraction analysis at 21 K: space group  $P2_12_12$ ;  $a = 13.261(3)$ ,  $b = 13.625(3)$ ,  $c = 8.612(2)$  Å;  $V =$ 1556.0(8)  $A^3$ ; Z = 2. Final agreement factors are  $R(F) = 0.038$  and  $R(wF) = 0.042$  for 2315 independent neutron reflections and 393 parameters varied. The neutron diffraction study represents the first accurate structure determination of a symmetric dimer with a metal-metal bond bridged by three hydride ligands. Average molecular parameters are:  $Ir-H = 1.78(1)$ ,  $Ir-Ir = 2.465(3)$ ,  $H\cdots H = 2.22(2)$ A, Ir-H-Ir =  $87.8(4)^\circ$ , H-Ir-H =  $77.2(7)^\circ$ , H-H-H  $= 60.0(7)^{\circ}.$ 

# **Introduction**

Over the past years, our group has been studying the structures of metal complexes of the type  $M-(H)<sub>n</sub>-M$  by neutron diffraction, in order to obtain basic structural information on hydrogen bridged metal-metal bonds. Thus far, symmetric bridged hydrides with  $n = 1,2,4$  have been accurately characterized [l-5] by ourselves and others; but the analysis of a compound having three hydrogen bridges  $[(C_5(CH_3)_5Ir)_2(\mu-H)_3]^+[BF_4]^-,$   $(1^+[BF_4]^-,$  was complicated by the presence of disorder, and was of relatively low precision  $[6,7]$ . In the present paper we report the neutron diffraction study of  $[ [C_5(CH_3)_5Ir)_2(\mu-H)_3]^+$ [ClO<sub>4</sub>]<sup>-</sup>·2C<sub>6</sub>H<sub>6</sub>], which provides more precise parameters for the  $M-(\mu-H)_{3}-M$ linkage, thereby completing the sequence of measurements on symmetrical  $M-(\mu-H)_{3}-M$  (n =  $1,2,3,4$ <sup>\*\*</sup>. From a catalytic viewpoint, 1<sup>+</sup> is of additional interest because it is the final hydrogenation product of the olefin hydrogenation and isomerization catalyst 2 [9], which also is a bridged hydride species. **1<sup>+</sup>** is formed starting from  $[C_5(CH_3)_5]rCl_2]_2$ by the stepwise replacement of bridging halides with hydrides which are themselves derived from the heterolytic splitting of hydrogen gas.

$$
\frac{C_0}{C_0 r} \sum_{C_1}^{H} \sum_{C_2}^{C_1 r} \frac{+2H_2}{C_1} + \left[ C_1 r - \frac{H}{H} \right] tr - C_1 r
$$
\n
$$
C_2 r^* = C_5 (CH_3)_5
$$

**As** mentioned above, the structure of **l+** had previously been studied by us in a neutron diffraction analysis at 20 K on a plate-like crystal of  $1^+$ [BF<sub>4</sub>]<sup>-</sup> [6,7]. Unfortunately, that earlier analysis was somewhat unsatisfactory, resulting in an incompletely refined structure  $(R(F) = 0.18)$ , due to the poor quality of the crystals and the disorder of the hydride ligands (which had thermal parameters corresponding to unusually elongated cigar-shaped ellipsoids). These factors prevented the structure analysis from attaining the desired precision, and therefore the study was abandoned at that stage, remaining unresolved and only partial results were published [6]. It was decided to reexamine **l',** and hopefully to grow a large crystal of better quality which would be more suitable for neutron diffraction. Preliminary experimentation had shown some promise with the  $[PF_6]$ <sup>-</sup> salt (J.C.), and an X-ray structural analysis

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<sup>\*\*</sup>It should be pointed out, however, that a neutron diffraction study of an asymmetric  $M-(\mu-H)_{3}-M$  compound,  $[Et_{4}$ -N]<sup>+</sup>[(triphos)Re<sub>2</sub>( $\mu$ -H)<sub>3</sub>(H)<sub>6</sub>]<sup>-</sup>·CH<sub>3</sub>CN, has been reported 181.

was carried out on this second derivative (T.W.). However, the  $[PF_6]$ <sup>-</sup> salt tended to grow as thin plates, still too small for a neutron analysis. In the present paper, we report that larger, high quality crystals can be obtained using the  $[ClO_4]$ <sup>-</sup> anion, and ultimately it was this third derivative which produced the crystal used in the neutron study reported here.

### Experimental

#### *Materials*

IrCl<sub>3</sub> 3H<sub>2</sub>O was purchased from Strem and used as received; H<sub>2</sub> gas (ultra-high purity,  $>99.999\%)$ and argon gas (pre-purified, >99.998%) were purchased from M.G. Industries. HCl gas (reagent grade) was purchased from Speciality Products and Equipment. All solvents were purchased from J. T. Baker and used as received, unless otherwise noted. The  $Na<sup>+</sup>$  salts of  $[BF<sub>4</sub>]<sup>-</sup>$ ,  $[ClO<sub>4</sub>]<sup>-</sup>$  and  $[B(Ph)<sub>4</sub>]<sup>-</sup>$  were purchased from Aldrich and used as received. Hexamethyldewarbenzene (HMBD-Aldrich) was stored in a freezer in sealed vials until required, and then opened and used immediately. Methanol was refluxed with, and then distilled from calcium hydride, under an atmosphere of argon. The synthetic steps are modification of published procedures [9, lo].

Proton NMR spectra were recorded in CDCl<sub>3</sub> solutions on a Jeol FX-90Q instrument at 90 MHz field, using residual  $CHCl<sub>3</sub>$  in the deuterated solvent as an internal standard. Although the final product is not particularily air-sensitive, some of the synthetic steps involved the handling of air sensitive precursors. Standard Schlenkware and vacuum line techniques were utilized for such steps, as noted in the following text.

# $[(C_5/CH_3)_5 Ir]_2 (\mu\text{-}H)_3$  +  $[ClO_4]^-$

This derivative was produced by dissolving the very soluble chloride salt in a minimum amount of distilled water and filtering the solution into a slight excess of a concentrated aqueous solution of Na+-  $\lceil$ ClO<sub>4</sub> $\rceil$ <sup>-</sup>. Precipitation ensued immediately, and the resulting perchlorate salt was collected by filtration (the filtrate was colorless), and washed with a minimum amount of cold water, and then dried exhaustively under vacuum. IR (KBr disk):  $1100 \text{ cm}^{-1}$  $\nu$ (Cl-O)(s). [Warning: perchlorate salts are notoriously shock-sensitive and prone to detonate. Although the explosive properties of this derivative were never investigated, it is recommended that no more than 25 mg be isolated (i.e.. in the dry state) at any one time.]

# *Crystal Growth*

*No* precautions for air sensitivity were needed in the crystal growth procedures. A portion of perchlorate was dissolved in a minimum amount  $(5 \text{ ml})$  of dry (stored over Linde type  $4(A)$  molecular sieves)  $CH_2Cl_2$ , and then benzene was slowly added to this solution until it became turbid, followed by a slow dropwise addition of  $CH<sub>2</sub>Cl<sub>2</sub>$  until the precipitate redissolved (total volume approximately 8 ml). The solution was then filtered into a suitable vapor diffusion crystallization setup with an excess volume (10 ml) of pentane as the precipitant and placed in the freezer  $(-22 \degree C)$  for a week. This technique slows crystal growth and thus enhances the chance of obtaining one large, good-quality crystal, as opposed to many small ones. Such a technique routinely resulted in the deposition of large, lemon-yellow, transparent crystals of  $1^+$   $[ClO_4]^-$ .

#### *X-ray Diffraction*

A preliminary room temperature X-ray diffraction structure analysis on a small yellow crystal of  $[(C_5 (CH_3)_5$ Ir)<sub>2</sub>( $\mu$ -H)<sub>3</sub>]<sup>+</sup>[ClO<sub>4</sub>]<sup>-</sup>·2C<sub>6</sub>H<sub>6</sub>] revealed the expected  $(C_5(CH_3)_5Ir)_2(\mu-H)_3$  cationic core (no hydrogen atoms were located in the X-ray study), one  $\left[\text{ClO}_4\right]^-$  counter-ion, and the two benzene molecules of solvation. The presence of benzene solvate was anticipated, from previous experience in handling the crystals, which had shown a strong tendency to desolvation. As a result, the crystal used for the intensity measurements was mounted in air on a glass fiber with Super Glue, and then rapidly coated with more Super Glue to minimize decomposition via desolvation. This technique proved successful, as evidenced by the fact that the faces of the crystal could still clearly be seen through the translucent adhesive during the crystal centering operations. Data were collected on a four-circle Nicolet/Syntex P2<sub>1</sub> automated diffractometer employing Mo K $\alpha$ radiation. The unit cell constants and orientation matrix were obtained from 15 centered reflections  $(20^{\circ} < 2\theta < 25^{\circ})$ . The unit cell parameters and other relevant data are presented in Table 1. Monitoring three check reflections  $[(640)(503)(631)]$  at 50-reflection intervals revealed no significant variation in intensity. The raw data were reduced to yield integrated intensities of 1007 independent reflections satisfying the criterion  $I > 3\sigma(I)$ . The intensities were corrected for Lorentz, polarization and absorption effects. The space group was determined to be Pnnm, consistent with the extinctions of *Okl,*   $k + l = 2n$ ; h0l,  $h + l = 2n$ . The coordinates of the unique iridium atom were obtained from a Patterson map, and the other non-hydrogen atoms were located from successive difference-Fourier maps. The cation has crystallographic  $2/m$  symmetry. All atoms were refined anisotropically. SHELX76 was used in the structure solution [l **11.** See also 'Supplementary Material'.

TABLE 1. Summary of crystal data and refinement results for  $[[(C_5(CH_3)_5Ir)_2(\mu-H)_3]^+[ClO_4]^{-}$ -2C<sub>6</sub>H<sub>6</sub>]

	$X$ -ray (298 K)	Neutron $(21 K)$
Space group	Pnnm	$P2_12_12$
q(A)	13.358(4)	13.261(3)
b(A)	14.083(4)	13.625(3)
c(A)	8.846(3)	8.612(2)
$V(A^3)$	1664.3(9)	1556.0(8)
Density (calc.) $(g \text{ cm}^{-3})$	1.667	1.783
Crystal dimensions (mm) $0.6 \times 0.3 \times 0.5$		$3.0 \times 1.5 \times 1.0$
Wavelength (A)	0.71069	1.15882(7)
Absorption coefficient		
$(\mu)$ (cm <sup>-1</sup> )	159.91	2.76
Sin $\theta/\lambda$ limit $(A^{-1})$	0.5385	0.6846
Total no. reflections		
measured	1374	4060
No. independent		
reflections	1320	2315
No. reflections used in		
structural analysis	1007 [ $I > 3σ(I)$ ]	2315
No. variable parameters	130	393
Final agreement factors		
R(F)	0.026	0.038
R(wF)	0.034	0.042
Goodness of fit $(F)$		1.342
$R(F^2)$		0.058
$R(wF^2)$		0.067
Goodness of fit $(F^2)$		1.358

#### *Neutron Diffraction*

Neutron diffraction data were collected at the Brookhaven High Flux Beam Reactor [12, 13] at 21 K on a large, transparent, lemon-yellow crystal which had been mounted on an aluminum pin approximately along the  $a^*$  direction and placed in a specially adapted closed-cycle helium refrigerator [ $14$ ]. Monitoring of one reflection (0.60) during cooling (298 to 21 K) revealed no significant variation in intensity. However, the crystal was found to undergo a reversible, non-destructive phase transition associated with ordering of the  $\left[ClO_4\right]^-$  counter-ions, with a change of space group from *Pnnm* at room temperature to  $P2<sub>1</sub>2<sub>1</sub>2$  at the temperature of neutron data collection  $(21 K)$ . Actually, this phase transition was not detected at first, since no change in the intensity of the monitored reflection (0 6 0) was discernible. It was only after difficulties were encountered during the structure analysis that it was suspected that a change in space group may have taken place. To test this hypothesis, the cooling experiment was repeated, on the same crystal used in the data collection, and five reflections were monitored in place of the previous one. Closely following the variation in intensity of these five reflections (allowed in  $P2_12_12$  and having strong intensity at 21 K but extinct in Pnnm) with temperature revealed the presence of the reversible phase transition. All five reflections rapidly diminished in intensity as the temperature increased towards the transition point, and disappeared completely at the transition temperature. The same situation occurred in reverse upon recooling through the transition. This left no doubt that we were observing a non-destructive, reversible phase change; the transition temperature was determined to be  $225(1)$  K.

The space group was determined to be  $P2<sub>1</sub>2<sub>1</sub>2$ in the low temperature neutron case, consistent with the extinctions of  $h00$ ,  $h = 2n$  and  $0k0$ ,  $k = 2n$ . In this space group, the anion lies across the twofold axis. The non-hydrogen atoms were located with MULTAN [15], and the hydrogen atoms were then located from successive difference-Fourier maps. Several cycles of refinement of the non-bridging atoms, in which all positional and anisotropic thermal parameters were varied, were performed by means of differential synthesis [ 161, until all non-bridging hydrogen atoms had been located. Full-matrix least-squares refinement [17] was then carried out, yielding an  $R(F)$  value of 0.060. The bridging hydrogen atoms were located from a difference-Fourier map (see Fig. 2) and found to be disordered about the two-fold axis. The disordered model (with six half-hydrogens) was refined with full-matrix leastsquares until convergence was obtained at an  $R(F)$ value of 0.038. Due to the disorder of the bridging hydrogens, the thermal parameters of these atoms were refined isotropically in the final cycles. An isotropic type I, extinction correction with Lorentzian mosiac was applied [18]. The quantity minimized in the least-squares refinement was  $\sum w (F_{\text{obs}}^2 F_{\text{coh}}^2$ <sup>2</sup> with  $w = (\sigma_{\text{count}}^2 + 0.02F_{\text{e}}^2)^2$ <sup>-1</sup>. A final difference map was essentially featureless. Atomic positional and-thermal parameters are given in Tables 2 and 3 respectively. See also 'Supplementary Material'.

Neutron scattering lengths were taken from a compilation of Koester [19]. Computer programs employed include MULTAN [15] (direct methods), DIFSYN [16] (differential synthesis program), a modified version of UPALS [17] (full-matrix leastsquares program), the absorption correction procedure of Coppens et al. [20], Johnson's ORTEP2 [21] and locally written programs.

#### **Results and Discussion**

Crystals of the title complex diffracted neutrons well, and refinement of the structure proceeded smoothly, once the change in space group on cooling was recognized. Aside from the ordering of the  $[C1O_4]^-$  ions, no substantial difference in the crystal packing was observed between the room temperature

TABLE 2. Atomic coordinates for  $[(C_5(CH_3)_5Ir)_2(\mu\text{-}H)_3]^+$ .  $[CIO<sub>4</sub>]<sup>-</sup>$  2C<sub>6</sub>H<sub>6</sub>] from the neutron diffraction study at 21 K

Atom	Occupancy $x$ .		у	z
Ir 1	1.000000	0.5900(1)	0.0225(1)	0.2140(1)
H B1	0.500000	0.5117(9)	$-0.036(1)$	0.076(1)
HB2	0.500000	0.5161(9)	$-0.057(1)$	0.330(2)
HB3	0.500000	0.4766(9)	$-0.9103(9)$	0.229(2)
Centroid <sup>a</sup>		0.5000	0.0000	0.2115
C1	1.000000	0.7409(2)	$-0.0233(1)$	0.2839(2)
C <sub>2</sub>	1.000000	0.7144(2)	0.0661(1)	0.3630(2)
C <sub>3</sub>	1.000000	0.6979(2)	0.1407(1)	0.2473(2)
C <sub>4</sub>	1.000000	0.7142(2)	0.0976(1)	0.0967(2)
C <sub>5</sub>	1.000000	0.7412(2)	$-0.0038(2)$	0.1195(2)
Centroidb		0.7217	0.0555	0.2221
C11	1.000000	0.7685(2)	$-0.1186(2)$	0.3580(2)
C <sub>21</sub>	1.000000	0.7104(2)	0.0807(2)	0.5354(2)
C <sub>31</sub>	1.000000	0.6726(2)	0.2461(2)	0.2782(2)
C <sub>41</sub>	1.000000	0.7105(2)	0.1490(2)	$-0.0563(2)$
C <sub>51</sub>	1.000000	0.7712(2)	$-0.0742(2)$	$-0.0051(2)$
<b>H11A</b>	1.000000	0.7509(5)	0.3194(4)	$-0.2844(7)$
H11C	1.000000	0.6510(4)	0.3784(5)	$-0.3814(9)$
H21A	1.000000	0.7842(4)	0.1008(5)	$-0.4220(6)$
H21B	1.000000	0.6896(6)	0.0121(4)	$-0.4046(6)$
H21C	1.000000	0.6566(5)	0.1376(4)	$-0.4328(6)$
H31A	1.000000	0.7430(4)	0.2886(4)	$-0.7181(8)$

TABLE 2. *(continued)* 

Atom	Occupancy $x$		y	z
43 1 B	1.000000	0.6237(5)	0.2773(4)	0.1902(6)
131 C	1.000000	0.6354(4)	0.2554(4)	$-0.6110(6)$
14 1 A	1.000000	0.6631(5)	0.2134(4)	$-0.0522(6)$
14 1 B	1.000000	0.6797(5)	0.1006(4)	$-0.1459(6)$
14 I C	1.000000	0.7867(5)	0.1691(4)	$-0.0910(6)$
151 A	1.000000	0.7523(5)	0.4500(4)	0.1192(6)
151B	1.000000	0.7621(5)	0.3533(4)	$-0.0136(7)$
151C	1.000000	0.6470(4)	0.4186(4)	0.0070(8)
11	0.500000	0.000000	0.000000	$-0.2688(2)$
$\overline{\mathbf{1}}$	1.000000	0.9182(2)	0.0336(2)	$-0.3636(3)$
)2	1.000000	0.9661(3)	0.0797(4)	$-0.1746(6)$
IS 1	1.000000	0.9534(2)	0.1641(2)	0.1819(3)
<b>S2</b>	1.000000	0.9507(2)	0.1553(2)	0.3436(3)
<b>S3</b>	1.000000	0.9501(2)	0.2379(2)	0.4381(3)
IS4	1.000000	0.9528(2)	0.3304(2)	0.3697(4)
<b>S5</b>	1.000000	0.9554(2)	0.3401(2)	0.2079(4)
CS6	1.000000	0.9560(2)	0.2571(2)	0.1150(3)
IS 1	1.000000	0.9543(6)	0.0985(5)	0.1085(7)
IS2	1.000000	0.9496(5)	0.0816(6)	$-0.6034(7)$
IS3	1.000000	0.9485(6)	0.2284(5)	$-0.4365(6)$
IS4	1.000000	0.9545(6)	0.3962(5)	$-0.5575(9)$
IS5	1.000000	0.9572(6)	0.4137(5)	0.156(1)
iS6	1.000000	0.9586(6)	0.2649(6)	$-0.0106(7)$

aBridging hydride centroid. bRing centroid.

TABLE 3. Anisotropic temperature factors for [ [(Cs(CH\$&),(lr-H),1+[C104]-\*2CaH6] from the neutron diffraction study at  $\ddot{\cdot}$ 

Atom	$\beta_{11}$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$
Ir1	0.0004(1)	0.0005(1)	0.0004(1)	0.0000(1)	0.0000(1)	0.0001(1)
HB1	0.0013(6)	0.009(1)	0.007(1)	$-0.0006(7)$	$-0.0003(8)$	$-0.005(1)$
HB2	0.0015(7)	0.007(1)	0.014(2)	$-0.0016(7)$	$-0.0021(1)$	0.008(1)
HB3	0.0029(7)	0.0021(5)	0.019(2)	$-0.0002(5)$	0.001(1)	$-0.001(1)$
C1	0.0005(1)	0.0007(1)	0.0016(2)	$-0.0000(1)$	0.0000(1)	0.0001(1)
C <sub>2</sub>	0.0003(1)	0.0009(1)	0.0014(2)	0.0000(1)	$-0.0000(1)$	$-0.0001(1)$
C <sub>3</sub>	0.0003(1)	0.0006(1)	0.0017(2)	$-0.0001(1)$	0.0003(1)	0.0000(1)
C <sub>4</sub>	0.0004(1)	0.0004(1)	0.0013(2)	0.0000(1)	0.0000(1)	0.0000(1)
C <sub>5</sub>	0.0004(1)	0.0007(1)	0.0015(2)	$-0.0001(1)$	0.0003(1)	$-0.0002(1)$
C11	0.0009(1)	0.0006(1)	0.0027(2)	0.0002(1)	0.0001(1)	0.0004(1)
C <sub>21</sub>	0.0006(1)	0.0013(1)	0.0016(2)	0.0000(1)	$-0.0000(1)$	$-0.0001(1)$
C <sub>31</sub>	0.0011(1)	0.0004(1)	0.0024(2)	$-0.0002(1)$	0.0003(1)	$-0.0002(1)$
C <sub>41</sub>	0.0012(1)	0.0008(1)	0.0011(2)	$-0.0000(1)$	0.0001(1)	0.0001(1)
C <sub>51</sub>	0.0011(1)	0.0010(1)	0.0017(1)	0.0001(1)	0.0002(1)	$-0.0005(1)$
<b>H11A</b>	0.0057(4)	0.0015(2)	0.0090(7)	0.0002(3)	$-0.0020(5)$	0.0002(3)
H11B	0.0048(4)	0.0036(3)	0.0075(7)	$-0.0010(3)$	0.0024(5)	$-0.0025(4)$
<b>H11C</b>	0.0019(3)	0.0037(3)	0.015(1)	$-0.0001(2)$	$-0.0016(5)$	$-0.0024(5)$
H21A	0.0015(3)	0.0067(4)	0.0057(6)	$-0.0011(3)$	$-0.0007(4)$	$-0.0017(4)$
H21B	0.0070(5)	0.0024(2)	0.0046(5)	$-0.0006(3)$	0.0008(4)	0.0006(3)
H21C	0.0041(4)	0.0030(3)	0.0059(7)	0.0016(3)	$-0.0004(4)$	$-0.0010(4)$
<b>H31A</b>	0.0025(3)	0.0019(2)	0.0131(8)	$-0.0008(2)$	0.0003(5)	$-0.0010(4)$
H31B	0.0042(4)	0.0021(2)	0.0066(6)	0.0007(2)	$-0.0019(4)$	0.0006(3)
<b>H31C</b>	0.0035(3)	0.0028(2)	0.0056(6)	0.0008(3)	0.0021(4)	$-0.0007(4)$
<b>H41A</b>	0.0043(4)	0.0019(2)	0.0066(6)	0.0013(3)	0.0005(4)	0.0011(3)
H41B	0.0052(4)	0.0031(3)	0.0040(6)	$-0.0008(3)$	$-0.0021(4)$	$-0.0005(4)$
						(continued)

TABLE 3. *(continued)* 

Atom	$\beta_{11}$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$
<b>H41C</b>	0.0024(3)	0.0049(3)	0.0064(6)	$-0.0008(3)$	0.0009(3)	0.0026(4)
<b>H51A</b>	0.0057(5)	0.0031(3)	0.0040(5)	$-0.0017(3)$	$-0.0007(4)$	$-0.0000(3)$
H51B	0.0039(4)	0.0014(2)	0.0106(8)	0.0005(2)	0.0027(4)	0.0013(3)
<b>H51C</b>	0.0012(3)	0.0039(3)	0.0131(9)	$-0.0003(3)$	0.0018(4)	0.0036(5)
<b>C11</b>	0.0005(1)	0.0014(1)	0.0015(2)	$-0.0004(1)$	0.000000	0.000000
01	0.0011(1)	0.0013(1)	0.0031(2)	$-0.0004(1)$	$-0.0007(1)$	0.0004(2)
O <sub>2</sub>	0.0024(2)	0.0070(3)	0.0177(7)	$-0.0031(2)$	0.0052(3)	$-0.0100(4)$
CS <sub>1</sub>	0.0010(1)	0.0023(1)	0.0041(3)	$-0.0002(1)$	$-0.0002(2)$	0.0010(2)
CS <sub>2</sub>	0.0008(1)	0.0020(1)	0.0036(3)	$-0.0003(1)$	$-0.0000(2)$	0.0010(2)
CS <sub>3</sub>	0.0010(1)	0.0021(1)	0.0052(3)	$-0.0002(1)$	$-0.0001(2)$	0.0006(2)
CS <sub>4</sub>	0.0008(1)	0.0017(1)	0.0102(4)	$-0.0000(1)$	$-0.0002(2)$	0.0006(2)
CS <sub>5</sub>	0.0010(1)	0.0021(1)	0.0109(4)	$-0.0003(1)$	$-0.0008(2)$	0.0030(2)
CS6	0.0012(1)	0.0029(1)	0.0060(3)	$-0.0004(1)$	$-0.0002(2)$	0.0024(2)
HS1	0.0037(4)	0.0045(4)	0.0090(7)	$-0.0006(3)$	0.0003(5)	$-0.0017(5)$
HS2	0.0038(4)	0.0025(3)	0.0092(7)	0.0000(3)	$-0.0004(5)$	0.0015(4)
HS3	0.0055(4)	0.0048(4)	0.0058(7)	$-0.0000(4)$	0.0003(5)	$-0.0011(4)$
HS4	0.0036(4)	0.0033(3)	0.017(1)	$-0.0003(3)$	0.0009(6)	$-0.0014(5)$
HS5	0.0036(4)	0.0035(3)	0.022(1)	$-0.0003(3)$	$-0.0012(6)$	0.0055(6)
HS6	0.0039(4)	0.0073(5)	0.0065(7)	$-0.0006(4)$	$-0.0003(4)$	0.0035(5)

The form of the anisotropic thermal factor is:  $exp[-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2hl\beta_{13} + 2kl\beta_{23})].$ 

and low-temperature phases. Surprisingly, however, an examination of ORTEP plots revealed the same cigar-shaped hydride thermal ellipsoids (Fig. 1) which had been observed in the previously examined  $[BF<sub>4</sub>]$ <sup>-</sup> derivative [6,7]. This occurred in spite of the much improved agreement  $[R(F) = 0.038$  for this study versus 0.18 for the previous one] and the fact that the present data set had been collected at the low temperature of 21 K. In the model illustrated in Fig. 1, the hydride ligands have been refined anisotropically, and their ellipsoids can be seen to elongated in the plane normal to the Ir-Ir bond.



Fig. 1. ORTEP view down the Ir-Ir bond illustrating the  $\Gamma$ g. 1. ONTET view down the  $n-1$  bond mustating the eigar shaped thermal empsolus resulting from freating each pair of half-occupied bridging hydrogens as a single site. In this model,  $HB(1)$  lies on the  $C_2$  axis. Methyl hydrogens have been removed for clarity. Thermal ellipsoids are drawn at the 75% probability level [21].



rig.  $\lambda$ . Difference Fourier section  $(x - 1/2)$  in the plane normal to the Ir-Ir bond, with the two-fold axis horizontal and showing the six disordered bridging hydrogen peaks [17].

Further investigation into this phenomenon uncovered disorder in the structure, as evidenced by a six-peak, pseudohexagonal arrangement of the nuclear density in the plane of the hydrogens (Fig. 2). For clarity, the two sets of three half-hydrogens related by the two-fold axis have been linked to form two equilateral triangles. This disorder involving



Fig. 3. ORTEP view normal to the  $Ir-Ir$  bond  $[21]$ .

the bridging hydride ligands then was taken into account, and the refinement was continued with the six half-occupied hydrogens (refined isotropically), resulting in the final structure model shown in Fig. 3. (It should be pointed out that this Figure shows only one of the two sets of three half-occupied hydrogens.) The *R* factor remained unchanged (0.038) with the six half-hydrogen isotropic model.

Bond distances and angles involving the bridging hydrides are given in Table 4. The M-H bond distances are normal, with a mean value of  $1.78(1)$  Å. This distance can be compared with mean values in, for example, the quadruply-bridged  $Re-(\mu-H)_4-Re$ dimer [5],  $[(PEt_2Ph)_2]_2(\mu-H)_4(H)_4$  with Re-H = 1.878(8) Å; and the two doubly-bridged  $M-(\mu$ -

**TABLE 4. Bridging hydride bond distances (A) and angles (<sup>o</sup>) in the**  $[(C_5(CH_3)_5]r)_2(\mu-H)_3]^+$  **cation, from the neutron diffraction** study at 21 K

$I_{\rm I}-I_{\rm I}$	2.465(3)		
$Ir-HB1$	1.76(1)	$Ir-Centroida$	1.23(1)
$I_{r}-HB1$	1.81(1)	$H1-Centroida$	1.27(2)
$Ir-HB2$	1.77(2)	$H2$ –Centroid <sup>a</sup>	1.30(2)
$Ir-HB2'$	1.79(1)	$H3$ -Centroid <sup>a</sup>	1.27(2)
$Ir-HB3$	1.76(1)	$H1'$ -Centroid <sup>a</sup>	1.27(2)
$Ir-HB3'$	1.77(1)	$H2'$ -Centroid <sup>a</sup>	1.30(2)
Mean Ir-H	1.78(1)	H3'-Centroid <sup>a</sup>	1.27(2)
		Mean H-Centroid <sup>a</sup>	1.28(2)
$HB1-HB2$	2.20(2)		
$HB1-HB3$	2.20(2)	$H1$ –Centroid <sup>a</sup> –Ir	90(1)
$HB2-HB3$	2.24(2)	$H2$ -Centroid <sup>a</sup> -Ir	89(1)
$H$ B1-HB1'	1.02(3)	$H3$ -Centroid <sup>a</sup> -Ir	90(1)
$HBI-HB3'$	1.51(2)	$H1'$ -Centroid <sup>a</sup> -Ir	92(1)
$HB2-HB2$	1.61(3)	$H2'$ –Centroid <sup>a</sup> –1r	90(1)
$HB2-HB3'$	0.98(2)	$H3'$ -Centroid <sup>a</sup> -Ir	90(1)
		Mean $H$ –Centroid <sup>a</sup> –Ir	90(1)
$HB1-Ir-HB2$	77.2(6)		
$HB1-Ir-HB3$	77.2(7)		
$HB2-Ir-HB3$	78.7(7)	Centroid <sup>a</sup> -Centroid <sup>b</sup>	3.04(1)
$HB1' - Ir - HB2'$	75.7(6)	Centroid <sup>a</sup> -Ir-Centroid <sup>b</sup>	178.8(6)
$HB1' - Ir - HB3'$	76.1(7)	$Ir-Centroida-Ir$	178(1)
$HB2' - Ir - HB3'$	78.1(7)		
Mean H-Ir-H	77.2(7)		
$HB1-HB2-HB3$	59.4(7)	$Ir-HB1-Ir'$	87.3(4)
$HB2-HB1-HB3$	61.1(7)	$Ir-HB2-Ir'$	87.8(5)
$HB1-HB3-HB2$	59.5(6)	$Ir-BH3-Ir'$	88.4(4)
Mean $H-H-H$	60.0(7)	Mean $Ir-H-Ir'$	87.8(4)
HB1-Ir-Centroid <sup>b</sup>	135.0(4)		
$HB2-Ir-Centroidb$	131.9(4)		
$HB3-Ir-Centroidb$	133.9(4)		
$HB1' - Ir - Centroidb$	136.4(4)		
HB2'-Ir-Centroidb	132.5(5)		
$HB3' - Ir - Centroidb$	134.0(4)		
Mean H-Ir-Centroid <sup>b</sup>	134.0(4)		

The prime (') notation is used to denote the symmetry (two-fold) related sites. aBridging hydride central bCsMes ring  $\mathbf{r}$  centre  $\mathbf{r}$ .





TABLE 5. *(continued)* 



H)<sub>2</sub>-M complexes  $H_2Os_3(CO)_{10}$  with  $Os-H = 1.850$ -(5) A [4] and  $[W_2(\mu-H)_2(\text{CO})_{10}]^{2-}$  with W-H = 1.926(3) Å [2]. The slighty shorter mean Ir-H bond distance observed here is consistent with iridium's position, to the right of tungsten, rhenium and osmium in the periodic table. The mean Ir-H-Ir bond angle of  $87.8(4)$ <sup>o</sup> is at the low end of the recorded range  $(85^\circ - 159^\circ)$ ; this is consistent with the observed short Ir-Ir bond distance of 2.465(3) Å (electron counting, to satisfy the 18electron rule, suggests an M=M triple bond). This angle may be compared to that in  $[(PEt_2Ph)_2Re]_2(\mu-H)_4(H)_4$ which is also presumed to have an M=M triple bond (again, from electron counting considerations) with  $Re-Re = 2.538(4)$  Å, and which exibits an average acute M-H-M bond angle of  $85.0(3)^\circ$ .

The disordered hydride ligands in  $1^+$  are in an essentially equilateral triangular arrangement in a plane at right angles to the  $Ir-Ir$  axis, with  $H--H$ non-bonding contact distances of 2.20-2.24(2) A. Other bond lengths and angles (Table 5) are within the expected normal ranges.

#### **Conclusions**

In conclusion, the  $[(C_5(CH_3)_5Ir)_2(\mu-H)_3]^+$  cation has been reexamined by neutron diffraction, this time utilizing a large, high quality crystal of the  $[ClO<sub>4</sub>]$ derivative. The selected crystal diffracted strongly, resulting in a data set which refined to yield a low *R* factor  $(R(F) = 0.038)$ . The unusual shapes of the bridging hydride probability ellipsoids (see Fig. 2) have been shown to be due to a hitherto overlooked disorder in the structure, which results in a situation where the three hydrides are seen as six disordered half-occupied hydrogens, with each pair of halfhydrogens clearly resolved. When this disorder is taken into account and the six half-hydrogens are refined isotropically, the fmal structure model is obtained (see Fig. 3). In spite of the presence of disorder, the *R* factor and the standard deviations are sufficiently low to consider this a definitive structural analysis, thus filling in the gap in the symmetric  $M-(H)<sub>n</sub>-M$  series, with characterization of compounds now up to and including  $n = 4$ . The bond distance trend in the series for the third-row metals is  $W-H(br) = 1.926(3)$ , Re-H(br) = 1.878(8), Os-H- $(br) = 1.850(5)$  and Ir-H(br) = 1.78(1) Å. Given the short H-H non-bonding contact  $(1.870(8)$  Å) in the  $n = 4$  complex, it may not be possible for more than four hydrides to 'fit' around any M-M bond for transition metals (221. Thus, one may have to turn to larger atoms (the actinides and/or lanthanides) to further expand the  $M-(H)<sub>n</sub>$ —M series.

# **Supplementary Material**

Atomic positional and thermal parameters from the X-ray study (Tables Sl and S2); observed and calculated structure factors (Table S3, X-ray; Table S4, neutron) (30 pages) are available from author **R.B.** 

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