# Thermodynamic Properties of $n$-Propyl Alcohol 

J. L. COSNER ${ }^{1}$, J. E. GAGLIARDO ${ }^{2}$, and T. S. STORVICK<br>University of Missouri, Columbia, Mo.

V ARIOUS USES of $n$-propyl alcohol as a solvent and chemical intermediate prompted a study of its derived thermodynamic properties. Sufficient volumetric, thermal, and vapor pressure data are available in the literature to evaluate these properties from $77^{\circ}$ to $540^{\circ} \mathrm{F}$. and pressures up to 700 p.s.i.a.
Extensive measurements of the volumetric behavior of $n$-propyl alcohol were made by Ramsey and Young (6) over the range from $30^{\circ}$ to $263.6^{\circ} \mathrm{C}$. and pressures from 800 to $53,480 \mathrm{~mm}$. of mercury. Foz Gazulla, Morcillo, and Mendes (2) measured vapor densities in the low pressure range from 100 to 760 mm . of mercury and the temperature range from $75^{\circ}$ to $130^{\circ} \mathrm{C}$. Saturated vapor and liquid densities were reported by Young (11).

The vapor pressure-temperature data for $n$-propyl alcohol have been summarized by Stull (9).

The heat capacity data at zero pressure, calculated from spectroscopic measurements and reported by Kobe, Harrison, and Pennington (4), compare favorably with the calorimetric data of Sinke and De Vries (8).

Calorimetric measurements of the heat of vaporization of $n$-propyl alcohol were reported by Bennewitz and Rossner (1). The normal boiling point data and the heat of vaporization at the normal boiling point were summarized in International Critical Tables (3).

The critical constants $T_{c}=537.3^{\circ} \mathrm{K} ., P_{c}=50.2 \mathrm{~atm}$., and $V_{c}=220 \mathrm{cc} . / \mathrm{gram}$ mole are reported by Lyderson, Greenkorn, and Hougen (5) and these values were used in this work. The molecular weight of $n$-propyl alcohol used in all the calculation in this work was 60.09 . The calculations were based on 1 pound of $n$-propyl alcohol, and the enthalpy and entropy were set equal to zero at $77^{\circ} \mathrm{F}$. for the saturated liquid at its vapor pressure.

## VAPOR PRESSURE AND HEATS OF VAPORIZATION

The vapor pressure data reported by Stull (9) were fitted to the following empirical equation:

$$
\begin{equation*}
\log P_{v}=(B / T-43)+C+D(T-43)+E(T-43)^{2} \tag{1}
\end{equation*}
$$

To improve the fit, Equation 1 was applied over a low and a high temperature range. The constants obtained were:

|  | Temperature Range, ${ }^{\circ} \mathrm{K}$ |  |
| :---: | :---: | :---: |
| Constant | $273-373$ | $373-537$ |
| $B$ | -1645.194 | -2173.147 |
| $C$ | 6.660707 |  |
| $D$ | $0.590031 \times 10^{-2}$ | 11.16614 |
| $E$ | $-6.46252 \times 10^{-6}$ | $3.520663 \times 10^{-2}$ |
|  |  | $3.5849 \times 10^{-6}$ |

The latent heats of vaporization were calculated with the thermodynamically rigorous Clapeyron equation

$$
\begin{equation*}
\frac{d P_{v}}{d T}=\frac{\Delta H_{0}}{T\left(V_{0}-V_{L}\right)} \tag{2}
\end{equation*}
$$

The necessary temperature derivatives of the vapor pressure were obtained from Equation 1. The saturated vapor $\left(V_{v}\right)$ and liquid $\left(V_{L}\right)$ volumes of Young allowed evaluation

[^0]of $\Delta H_{v}$. The heats of vaporization used in preparing Table I are smoothed values calculated from Equation 2. At atmospheric pressure the value given in Table I is 295.0 B.t.u. per pound and compares favorably with 295.9 B.t.u. per pound reported in International Critical Tables (3), less favorably with 299.0 B.t.u. per pound reported by Bennewitz and Rossner (1).

## VAPOR PHASE THERMODYNAMIC PROPERTIES

Residual volumes were calculated from the volumetric data of Ramsey and Young (6). Low pressure data on the high temperature isotherms were not reported. An extrapolating equation was developed for this low pressure vapor region based on dimerization of the $n$-propyl alcohol in the vapor phase. The data of Foz Gazulla, Morcillo, and Mendes (2) were used to confirm the extrapolated results. These two sources of volumetric data had so few common values that no direct comparison was possible.

The curves of residual volume $v s$. pressure were prepared and used to calculate the fugacity coefficients. The relationship used was

$$
\begin{equation*}
\ln \left(\frac{f}{P}\right)=\frac{-1}{R T} \int_{0}^{P} \alpha d P \tag{3}
\end{equation*}
$$

Equation 3 was evaluated graphically.
The enthalpies of the superheated vapor were calculated as differences between the enthalpy of the gas at pressure $P$ and the value in the zero pressure or ideal state for each isotherm, The thermodynamic relationship used was

$$
\begin{equation*}
\left(H-H^{\circ}\right)_{T}=\int_{0}^{P}\left[T\left(\frac{\partial \alpha}{\partial T}\right)_{P}-\alpha\right] d P \tag{4}
\end{equation*}
$$

To integrate Equation 4, a curve of $\alpha$ vs. temperature was prepared by cross-plotting the $\alpha$-pressure diagram. The derivatives, $T\left(\frac{\partial \alpha}{\partial T}\right)_{P}$, were evaluated numerically using the Gregory-Newton backward and forward interpolation formulas (10) at the ends of the curves and Stirling's formula (7) in the central position. A plot of $T\left(\frac{\partial \alpha}{\partial T}\right)_{P}$ was prepared and Equation 4 evaluated graphically.

The entropies of the superheated vapor were calculated as differences between the entropy of the gas at pressure $P$ and the ideal gas state at zero pressure by the expression

$$
\begin{equation*}
\left(S-S^{\circ}\right)_{T}=\int_{0}^{P}-\left(\frac{\partial V}{\partial T}\right)_{P} d P \tag{5}
\end{equation*}
$$

When the residual volume is substituted in Equation 5, the integral becomes indeterminate at the lower limit. Equation 5 can be evaluated from $P=0$ to $P=P^{*}$, an ideal gas state at unit fugacity. When this is substracted from Equation 5, the lower limit $P=0$ is eliminated and the result is

$$
\begin{equation*}
\left(S-S^{*}\right)_{\tau}=\left(S-S^{\circ}\right)_{r}=\int_{P^{*}}^{P}\left(\frac{\partial \alpha}{\partial T}\right)_{P} d P-\int_{P^{*}}^{P}\left(\frac{R}{P}\right) d P \tag{6}
\end{equation*}
$$

Equation 6 was evaluated using the results of Equations

Table I. Thermodynamic Properties of Saturated n-Propyl Alcohol

| $\begin{aligned} & \text { Temp., } \\ & \circ \text { F. } \end{aligned}$ | Pressure, P.S.I.A. | Volume, Cu. Ft./Lb. |  | Enthalpy, B.t.u./Lb. |  |  | Entropy, B.t.u./Lb. ${ }^{\circ}$ R. |  |  | Fugacity Ratio, f/P |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  | Vaporization | Vapor |  | Vapori- <br> zation | Vapor |  |
|  |  | Liquid | Vapor | Liquid |  |  | Liquid |  |  |  |
| 77 | 0.39 | 0.0200 | 244.7 | 0.0 | 337.5 | 337.5 | 0.0 | 0.629 | 0.629 | 0.997 |
| 100 | 0.85 | 0.0203 | 116.4 | 12.60 | 333.0 | 345.6 | 0.103 | 0.595 | 0.698 | 0.994 |
| 140 | 2.83 | 0.0208 | 37.26 | 35.74 | 324.5 | 360.0 | 0.243 | 0.541 | 0.784 | 0.986 |
| 180 | 7.94 | 0.0214 | 13.98 | 60.24 | 315.6 | 375.9 | 0.358 | 0.493 | 0.851 | 0.973 |
| 280 | 14.696 | 0.0217 | 7.86 | 90.02 | 295.0 | 385.0 | 0.455 | 0.442 | 0.897 | 0.959 |
| 220 | 19.17 | 0.0219 | 6.21 | 99.81 | 291.8 | 390.0 | 0.486 | 0.429 | 0.915 | 0.952 |
| 260 | 40.18 | 0.0231 | 2.94 | 139.4 | 268.4 | 402.7 | 0.580 | 0.373 | 0.953 | 0.924 |
| 300 | 76.17 | 0.0237 | 1.55 | 165.4 | 249.6 | 415.0 | 0.659 | 0.328 | 0.987 | 0.886 |
| 340 | 132.9 | 0.0249 | 0.880 | 208.9 | 218.8 | 427.7 | 0.752 | 0.274 | 1.026 | 0.842 |
| 380 | 216.5 | 0.0265 | 0.530 | 243.6 | 194.1 | 437.7 | 0.823 | 0.231 | 1.054 | 0.788 |
| 420 | 333.1 | 0.0289 | 0.321 | 284.0 | 160.6 | 445.0 | 0.905 | 0.182 | 1.087 | 0.747 |
| 460 | 488.2 | 0.0320 | 0.186 | 329.7 | 116.1 | 445.8 | 0.977 | 0.126 | 1.103 | 0.687 |
| 500 | 688.5 | 0.0410 | 0.0995 | 380.6 | 58.1 | 440.0 | 1.053 | 0.060 | 1.113 | 0.652 |
| 507 | 737.9 | 0.0587 | 0.0587 | 414.5 | 0.0 | 414.5 | 1.115 | 0.0 | 1.115 |  |

3 and 4. Dividing Equation 4 by $T$ and substracting from this, $R$ times Equation 3 gives the first integral in Equation 6. Thus:

$$
\begin{equation*}
\left(S-S^{\circ}\right)_{r}=\left(\frac{H-H^{\circ}}{T}\right)-R \ln \left(\frac{f}{P}\right)-R \ln \left(\frac{P}{P^{*}}\right) \tag{7}
\end{equation*}
$$

Equation 7 was used to evaluate the superheated vapor entropy values.

The ideal gas, or zero pressure enthalpies were calculated from the heat capacity data of Kobe, Harrison, and Pennington (4) with the expression

$$
\begin{equation*}
d H=C_{p}^{\circ} d T \tag{8}
\end{equation*}
$$

The ideal gas entropy values were obtained from the expression

$$
\begin{equation*}
d S=C_{\rho}^{\circ} / T d T \tag{9}
\end{equation*}
$$

Saturated enthalpy values were obtained by plotting a temperature-enthalpy diagram and extrapolating the isobars to the saturated temperature. The saturated entropy values were obtained in a similar manner on a temperatureentropy diagram.
Enthalpy and entropy data for the saturated liquid were obtained by subtracting the enthalpy and entropy of vaporization from their saturated vapor values. The data were smoothed by plotting on enthalpy-entropy, tempera-


Figure 1. Pressure-enthalpy diagram for n-propyl alcohol
Reference state: $H=0, S=0$ at $77^{\circ} \mathrm{F}$., saturated liquid

Table II. Thermodynamic Properties of Superheated n-Propyl Alcohol

| $\begin{aligned} & \text { Temp., } \\ & \circ \mathbf{F} . \end{aligned}$ | Volume, Cu . Ft./Lb. | Enthalpy, B.t.u./Lb. 0 P.S.I.A. | Entropy, <br> B.t.u. per <br> Lb. ${ }^{\circ}$ R. | Fugacity, Ratio, $f / P$ | $\begin{aligned} & \text { Temp., } \\ & \circ \text { F } . \end{aligned}$ | Volume, Cu. Ft./Lb. 100 | Enthalpy, B.t.u./Lb. | Entropy, <br> B.t.u. per <br> Lb. ${ }^{\circ}$ R. <br> . ${ }^{\text {a }}$ ) | Fugacity Ratio, $f / p$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 77 | Undefined | 337.7 | 0.509 | 1.000 | 340 | 1.243 | 436.3 | 1.042 | 0.8804 |
| 100 |  | 346.0 | 0.604 | 1.000 | 380 | 1.349 | 458.8 | 1.101 | 0.9008 |
| 140 |  | 361.3 | 0.732 | 1.000 | 420 | 1.439 | 481.5 | 1.155 | 0.9155 |
| 180 |  | 377.4 | 0.833 | 1.000 | 460 | 1.526 | 504.2 | 1.210 | 0.9274 |
| 220 |  | 394.4 | 0.918 | 1.000 | 500 | 1.610 | 527.5 | 1.252 | 0.9368 |
| 260 |  | 412.1 | 0.993 | 1.000 | 540 | 1.690 | 550.9 | 1.296 | 0.9440 |
| 300 |  | 430.6 | 1.058 | 1.000 |  |  |  |  |  |
| 340 |  | 450.5 | 1.119 | 1.000 | 150 P.S.I.A. (345 ${ }^{\circ}$ F. ${ }^{\text {a }}$ ) |  |  |  |  |
| 380 420 |  | 469.9 490.6 | 1.174 1.226 | 1.000 1.000 | 360 | 0.803 | 438.4 | 1.051 | 0.8423 |
| 460 |  | 512.0 | 1.275 | 1.000 | 380 | 0.848 | 452.2 | 1.082 | 0.8561 |
| 500 |  | 534.1 | 1.321 | 1.000 | 420 | 0.928 | 476.3 | 1.138 | 0.8772 |
| 540 |  | 556.5 | 1.362 | 1.000 | 460 | 0.984 | 499.8 | 1.189 | 0.8957 |
| 10 P.S.I.A. (190 ${ }^{\circ} \mathrm{F}^{\circ}$ ) |  |  |  |  | 500 540 | 1.044 | 524.0 | 1.236 | ${ }^{0.9098}$ |
| 220 | 11.414 | 392.5 | 0.915 | 0.9708 | 200 P.I.S.A. (374 ${ }^{\circ} \mathrm{F} .{ }^{\text {a }}$ ) |  |  |  |  |
| 260 | 11.803 | 410.6 | 0.990 | 0.9795 |  |  |  |  |  |
| 300 | 13.349 | 429.4 | 1.056 | 0.9838 | 380 | 0.584 | 442.5 | 1.062 | 0.8124 |
| 340 | 14.094 | 449.2 | 1.117 | 0.9870 | 420 | 0.656 | 471.3 | 1.123 | 0.8431 |
| 380 | 14.837 | 468.6 | 1.172 | 0.9892 | 460 | 0.711 | 496.0 | 1.176 | 0.8662 |
| 420 460 | 15.565 16.294 | 489.5 511.2 | ${ }_{1}^{1.225}$ | 0.9908 0.9920 | 500 540 | 0.761 0.807 | 520.6 544.9 | 1.224 1.269 | 0.8843 0.8979 |
| 500 | 17.011 | 533.4 | 1.321 | 0.9927 | 540 | 0.807 | 544.9 | 1.269 | 0.8979 |
| 540 | 17.743 | 556.0 | 1.361 | 0.9934 | 250 P.S.I.A. (393 ${ }^{\circ} \mathrm{F} .{ }^{\text {a }}$ ) |  |  |  |  |
| 14.696 P.S.I.A. ( $208{ }^{\circ} \mathrm{F} .{ }^{\text {a }}$ ) |  |  |  |  | 400 | 0.452 | 446.7 | 1.076 | 0.7895 |
| 220 | 7.932 | 392.5 | 0.913 | 0.9526 | 420 | 0.493 0.546 | 463.8 496.5 | 1.109 1.164 | ${ }_{0}^{0.8239}$ |
| 260 | 8.201 | 409.7 | 0.988 | 0.9709 | 500 | 0.590 | 517.1 | 1.214 | 0.8602 |
| 300 | 9.016 | 429.1 | 1.054 | 0.9770 | 540 | 0.632 | 542.0 | 1.260 | 0.8773 |
| 340 | 9.532 | 448.7 468.4 | 1.116 | ${ }_{0}^{0.9812}$ | 300 P.S.I.A. ( $410^{\circ} \mathrm{F} .{ }^{\circ}$ ) |  |  |  |  |
| 380 420 | 10.041 10.546 | 468.4 489.2 | 1.171 1.224 | 0.9843 0.9867 |  |  |  |  |  |
| 460 | 11.047 | 510.8 | 1.273 | 0.9883 | 420 | 0.373 | 453.4 | 1.093 | 0.7725 |
| 500 | 11.545 | 533.0 | 1.319 | 0.9899 | 460 | 0.433 | 486.1 | 1.154 | 0.8094 |
| 540 | 12.043 | 555.7 | 1.359 | 0.9913 | 500 540 | 0.475 0.510 | 513.5 539.2 | 1.205 1.252 | 0.8363 0.8564 |
| 25 P.S.I.A. (233 ${ }^{\circ}$ F. ${ }^{\text {a }}$ ) |  |  |  |  | 400 P.S.I.A. (439 ${ }^{\circ} \mathrm{F} .{ }^{\text {a }}$ ) |  |  |  |  |
| 240 260 | 4.712 4.883 | 395.6 407.5 | 0.934 0.971 | ${ }_{0}^{0.9451}$ | 460 | 0.284 | 471.3 | 1.131 | 0.7283 |
| 300 | 5.373 | 428.4 | 1.037 | 0.9610 | ${ }^{480}$ | 0.308 | 489.4 | 1.161 | 0.7524 |
| 340 | 5.530 | 447.5 | 1.098 | ${ }_{0}^{0.9681}$ | 500 540 | 0.329 0.363 | 504.8 533.0 | 1.188 1.238 | 0.7898 0.8170 |
| 380 | ${ }_{6}^{5.837}$ | 467.2 488.3 | ${ }_{1}^{1.155}$ | 0.9734 0.9774 | 540 | 0.363 | 533.0 | 1.238 | 0.8170 |
| 420 | 6.141 6.442 | 488.3 510.0 | 1.217 1.256 | 0.9774 0.9804 | 500 P.S.I.A. ( $459^{\circ} \mathrm{F} .{ }^{\text {. }}$ ) |  |  |  |  |
| 460 500 | 6.442 6.739 | 510.0 532.4 | 1.256 1.302 | 0.9884 0.988 |  |  |  |  |  |
| 540 | 7.032 | 554.9 | 1.346 | 0.9849 | 480 500 | ${ }_{0}^{0.236}$ | 492.8 | 1.171 | 0.7441 |
| 50 P.S.I.A. ( $273{ }^{\circ} \mathrm{F}^{\circ}{ }^{\circ}$ ) |  |  |  |  | 520 | 0.254 0.266 | 509.1 525.7 | 1.198 1.224 | 0.7633 0.7785 |
| 280 | 2.377 | 410.0 | 0.973 | 0.9167 | 540 | 0.266 | 525.7 | 1.224 |  |
| 300 | 2.441 | 422.1 | 1.009 | 0.9244 | 600 P.S.I.A. (479 ${ }^{\circ} \mathrm{F}^{\text {a }}$ ) |  |  |  |  |
| 340 | 2.677 | 444.2 | 1.072 | 0.9375 | 500 | 0.167 | 472.9 | 1.152 | 0.6966 |
| 380 | 2.842 | 464.5 | 1.129 | 0.9476 | 520 | 0.190 | 497.3 | 1.182 | 0.7211 |
| 420 | ${ }^{3.005}$ | 486.2 | 1.181 | 0.9555 0.9618 | 540 | 0.205 | 515.9 | 1.209 | 0.7402 |
| 460 | 3.160 | 505.7 | 1.231 | 0.9618 |  |  |  |  |  |
| 460 | 3.160 | 505.7 | 1.231 | 0.9618 | 700 P.S.I.A. ( $501{ }^{\circ} \mathrm{F} .{ }^{\circ}$ ) |  |  |  |  |
| 500 | ${ }_{3.315}$ | 530.8 | 1.277 | 0.9667 | 520 | 0.136 | 480.5 | 1.161 | 0.6774 |
| 540 | 3.464 | 553.8 | 1.322 | 0.9706 | 540 | 0.152 | 505.3 | 1.196 | 0.7025 |
| ${ }^{\text {a }}$ Saturated temperature. |  |  |  |  |  |  |  |  |  |

ture-enthalpy, and temperature-entropy diagrams. Tabulated data are given in Tables I and II, and the data are plotted on a pressure-enthalpy diagram (Figure 1).

The calculated data were checked for internal consistency with the combined first and second law of thermodynamics

$$
\begin{equation*}
d H=T d S+V d P \tag{10}
\end{equation*}
$$

Equation 10 was integrated graphically at both constant pressure and constant entropy and compared with the enthalpy from the calculated data. These checks showed that the data presented are internally consistent.
The accuracy of the data cannot be checked by direct comparison, because the necessary data are lacking. The vapor phase enthalpy values are probably accurate to
$\pm 1.00$ B.t.u. per pound and the entropy values to $\pm 0.001$ B.t.u. per pound ${ }^{\circ} R$.

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## NOMENCLATURE

$B, C, D, E=$ constants in Equation 15
$. C_{p}=$ constant pressure heat capacity
$C_{p}=$ pure component fugacity

| $H$ | $=$ enthalpy |
| ---: | :--- |
| $P$ | $=$ pressure |
| $R$ | $=$ universal gas constant, 1.9872 B.t.u. $/$ lb.-mole $\circ R$. |
| $S$ | $=$ entropy |
| $T$ | $=$ temperature |
| $V$ | $=$ volume |
| $\alpha=$ residual volume |  |

Superscripts

- = zero pressure, ideal gas state
* = unit fugacity, ideal gas state


## Subseripts

$c=$ critical constants
$L=$ liquid phase
$v=$ vapor phase

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## CORRECTION

In "Prediction of Equilibrium Ratios from Nomograms of Improved Accuracy" [B.C. Cajander, H.G. Hipkin, and J.M. Lenoir, J. Chem. Eng. Data 5, 251 (1960)] the figures entitled "Equilibrium ratios of aliphatic hydrocarbons at 10 p.s.i.a." in the right column on page 254 are incorrect.

The line labeled butadiene-1,2 is low by about $30 \%$. The line labeled 2 or 3 -methylbutene- 1 is correct for 3 -methyl-butene-1, but 2 -methylbutene-1 should lie between pentene1 and isoprene. The corrected figures appear as shown below.




[^0]:    ${ }^{1}$ Present address, Columbia Southern Chemical Co., Barberton, Ohio.
    ${ }^{2}$ Present address, U.S.S. Glennon (DD840), F.P.O., New York, N. Y.

