# Vapor Pressure and Latent Heat of Vaporization of Parahydrogen from the Triple to the Critical Point 

JOSEPH V. De PALMA and GEORGE THODOS

Northwestern University, Evanston, III.


#### Abstract

The existing vapor pressure measurements reported in the literature for parahydrogen have been utilized to establish the constants of the Frost-Kalkwarf equation applicable between the triple point and the critical point. For refinement, this equation has been applied specifically to vapor pressures above 1 atm . For vapor pressures below 1 atm , a truncated form of this equation has been used. These refined equations, along with the saturated vapor and liquid density relationships of Roder, Diller, Weber, and Goodwin were used with the Clapeyron equation to calculate the latent heats of vaporization. A plot of latent heats of vaporization is presented along with a table of comparison with the values of others.


Current interest in the use of parahydrogen as a fuel for rocket propulsion has stimulated interest for the accurate establishment of several of its physical properties. Two such interdependent properties that have practical usage are the vapor pressure and the latent heat of vaporization. Therefore, this study has been concerned with the procurement from the literature of experimental vapor pressure values to be used for the establishment of an analytical relationship capable of defining accurately the vapor pressure behavior of parahydrogen from its triple point to its critical point. Furthermore, such a relationship perinits the calculation of the latent heat of
vaporization of this substance through the use of the Clapeyron equation,

$$
\begin{equation*}
\left(\frac{\partial P}{\partial T}\right)_{v}=\frac{\lambda}{T\left(v_{0}-v_{l}\right)} \tag{1}
\end{equation*}
$$

Equation 1 is thermodynamically exact and consequently the accuracy of the latent heat of vaporization, $\lambda$, depends directly on the ability to obtain accurate values for $(\partial P / \partial T)_{e}$ and the quantity, $v_{0}-v_{1}$.


Figure 1. Relationship between $X$ and $Y$ for para hydrogen

## VAPOR-PRESSURE BEHAVIOR

A number of vapor pressure equations are presented in the literature. In a recent critical review, Miller (4) points out the superiority of the Frost-Kalkwarf equation (2) to represent accurately the vapor-pressure behavior of a substance over the complete range included between its triple and critical points. Consequently, this equation was applied to the existing vapor pressure meassurements presented in the literature. A critical literature survey reveals the extensive experimental vapor pressure measurements of Hoge and Arnold (3), which cover the range from the triple point to the near critical point vicinity. These data were adjusted to the N.B.S.- 1955 temperature scale by subtracting $0.01^{\circ} \mathrm{K}$. from the temperatures reported. These adjusted temperatures are used throughout this study. The recent contributions of the National Bureau of Standards (11) cover the range between the normal boiling point and the critical point of parahydrogen. Hoge and Arnold (3) report a triple point value of $T_{t}=13.803^{\circ} \mathrm{K} ., P_{t}=52.82 \mathrm{~mm}$. (adjusted to N.B.S.- 1955 temperature scale). In 1963, Barber and Hors-
ford (1) reported experimental measurements from the triple point to the normal boiling point. The triple point obtained by Barber and Horsford is $13.816^{\circ} \mathrm{K}$. and 52.95 mm . The temperatures of the investigation by Barber and Horsford (1) are consistent with the N.B.S.- 1955 temperature scale. These three references constitute the basic core of experimental data used to develop the constants of the Frost-Kalkwarf equation,

$$
\begin{equation*}
\ln P=A+\frac{B}{T}+C \ln T+D \frac{P}{T^{2}} \tag{2}
\end{equation*}
$$

In their original development, Frost and Kalkwarf point out the $D=a / R^{2}$, where $a$ is the pressure van der Waals constant and $R$ is the universal gas constant. The critical constants of parahydrogen reported by Roder, Diller, Weber, and Goodwin (7) are as follows: $T_{c}=32.976 \pm 0.015^{\circ} \mathrm{K}, P_{c}=12.759 \pm$ 0.028 atm. , and $\rho_{c}=0.01559 \pm 0.00005 \mathrm{gram}-\mathrm{moles} / \mathrm{cc}$. Thus, the critical volume of parahydrogen becomes $v_{c}=1 / 0.01559$ $=64.144 \mathrm{cc} / \mathrm{gram}-\mathrm{mole}$. The molecular weight of parahydro-

Table I. Comparison of Experimental and Calculated Vapor Pressures of Parahydrogen Using Equation 7

$$
\ln P=10.4791-\frac{105.651}{T}+0.425173 \ln T+0.047310 \frac{P}{T^{2}}
$$

|  | $T,{ }^{\circ} \mathrm{K}$. | Millimeters of Mercury |  | $\Delta P$ | Dev. \% |  | $T,{ }^{\circ} \mathrm{K}$. | Millimeters of Mercury |  | $\Delta P$ | Dev. \% |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $P_{\text {exptl }}$ | $P_{\text {ealod }}$ |  |  |  |  | $P_{\text {exptl }}$ | $P_{\text {onled }}$ |  |  |
| W | $32.976^{\text {a }}$ | 9696.8 | 9762.3 | 65.5 | 0.68 | H | 25.037 | 2488.5 | 2479.2 | -9.3 | -0.37 |
| W | 32.900 | 9589.6 | 9654.0 | 64.4 | 0.67 | W | 25.000 | 2467.1 | 2459.9 | -7.2 | -0.29 |
| H | 32.883 | 9566.2 | 9630.4 | 64.2 | 0.67 | H | 24.890 | 2414.1 | 2403.6 | -10.5 | -0.44 |
| H | 32.884 | 9564.4 | 9630.8 | 66.4 | 0.69 | V | 24.680 | 2301.5 | 2298.2 | -3.3 | -0.15 |
| H | 32.883 | 9559.3 | 9629.4 | 70.1 | 0.73 | H | 24.440 | 2189.4 | 2182.1 | -7.3 | -0.33 |
| W | 32.800 | 9449.5 | 9513.1 | 63.6 | 0.67 | H | 23.634 | 1827.1 | 1822.4 | -4.7 | -0.26 |
| W | 32.700 | 9312.0 | 9373.8 | 61.8 | 0.66 | W | 23.000 | 1572.3 | 1570.6 | -1.7 | -0.11 |
| H | 32.636 | 9219.5 | 9285.2 | 65.7 | 0.71 | H | 22.896 | 1534.9 | 1531.7 | -3.2 | -0.21 |
| W | 32.600 | 9176.8 | 9236.2 | 59.4 | 0.65 | V | 22.772 | 1488.7 | 1486.4 | -2.4 | -0.16 |
| W | 32.500 | 9042.8 | 9100.2 | 57.4 | 0.63 | H | 22.569 | 1416.1 | 1414.3 | -1.8 | -0.13 |
| H | 32.375 | 8875.1 | 8932.8 | 57.7 | 0.65 | H | 22.270 | 1313.8 | 1312.4 | -1.4 | -0.10 |
| H | 32.129 | 8557.7 | 8609.6 | 51.9 | 0.61 | H | 22.250 | 1308.4 | 1306.0 | -2.4 | -0.19 |
| W | 32.000 | 8399.0 | 8443.6 | 44.6 | 0.53 | W | 22.000 | 1225.4 | 1225.1 | -0.3 | -0.02 |
| H | 31.881 | 8255.1 | 8293.0 | 37.9 | 0.46 | H | 21.677 | 1127.2 | 1126.3 | -0.9 | -0.08 |
| W | 31.500 | 7792.5 | 7824.7 | 32.2 | 0.41 | H | 21.328 | 1026.8 | 1025.8 | -1.0 | -0.10 |
| H | 31.392 | 7660.2 | 7695.8 | 35.6 | 0.46 | H | 21.195 | 989.8 | 989.2 | -0.6 | -0.06 |
| V | 31.238 | 7503.2 | 7514.4 | 11.2 | 0.15 | H | 20.943 | 922.5 | 922.7 | 0.2 | 0.03 |
| H | 31.072 | 7302.5 | 7323.1 | 20.6 | 0.28 | H | 20.941 | 922.1 | 922.2 | 0.1 | 0.01 |
| W | 31.000 | 7221.1 | 7241.2 | 20.1 | 0.28 | H | 20.855 | 900.7 | 900.3 | -0.4 | -0.05 |
| V | 30.971 | 7200.2 | 7208.2 | 8.0 | 0.11 | H | 20.507 | 815.2 | 815.0 | -0.2 | -0.03 |
| H | 30.892 | 7102.8 | 7119.7 | 16.9 | 0.24 | H | 20.502 | 813.7 | 813.8 | 0.1 | 0.01 |
| W | 30.500 | 6679.2 | 6691.5 | 12.3 | 0.18 | H | 20.397 | 789.6 | 789.3 | -0.3 | -0.03 |
| H | 30.368 | 6544.6 | 6551.3 | 6.7 | 0.10 | B | 20.270 | 760.0 | 760.6 | 0.6 | 0.07 |
| V | 30.172 | 6357.3 | 6348.0 | -9.3 | -0.15 | H | 20.268 | 760.0 | 760.0 | 0.0 | 0.00 |
| W | 30.000 | 6169.0 | 6173.6 | 4.6 | 0.07 | W | 20.268 | 760.0 | 760.0 | 0.0 | 0.00 |
| H | 29.907 | 6080.2 | 6081.0 | 0.8 | 0.01 | H | 20.255 | 757.1 | 757.0 | -0.1 | -0.01 |
| H | 29.386 | 5583.9 | 5578.9 | -5.0 | -0.09 | H | 20.030 | 707.7 | 707.9 | 0.2 | 0.03 |
| V | 29.238 | 5453.3 | 5442.3 | -11.0 | -0.02 | H | 19.114 | 530.7 | 531.2 | 0.5 | 0.09 |
| W | 29.000 | 5233.6 | 5227.9 | -5.7 | -0.11 | B | 19.050 | 519.5 | 520.1 | 0.6 | 0.11 |
| H | 28.870 | 5121.5 | 5113.0 | -8.5 | -0.17 | H | 18.571 | 442.1 | 442.8 | 0.7 | 0.16 |
| H | 28.376 | 4705.7 | 4694.2 | -11.5 | -0.24 | B | 18.447 | 423.3 | 424.2 | 0.9 | 0.23 |
| W | 28.000 | 4401.9 | 4392.8 | -9.1 | -0.21 | H | 17.819 | 338.4 | 338.7 | 0.3 | 0.10 |
| H | 27.864 | 4299.6 | 4287.5 | -12.1 | -0.28 | B | 17.429 | 291.3 | 292.3 | 1.0 | 0.36 |
| V | 27.540 | 4049.9 | 4043.0 | -6.9 | -0.17 | H | 16.965 | 243.1 | 243.6 | 0.5 | 0.19 |
| H | 27.398 | 3952.7 | 3939.5 | -13.2 | -0.33 | H | 16.945 | 241.4 | 241.6 | 0.2 | 0.07 |
| V | 27.072 | 3718.2 | 3708.2 | $-10.0$ | -0.27 | H | 15.831 | 149.7 | 149.6 | -0.1 | -0.04 |
| W | 27.000 | 3669.7 | 3658.7 | -11.0 | -0.30 | B | 15.348 | 119.4 | 119.2 | -0.2 | -0.16 |
| H | 26.771 | 3517.5 | 3504.0 | -13.5 | -0.38 | B | 15.005 | 100.8 | 100.6 | -0.3 | -0.27 |
| H | 26.188 | 3142.6 | 3131.0 | -11.6 | -0.37 | B | 14.524 | 78.6 | 78.2 | -0.4 | -0.47 |
| W | 26.000 | 3026.7 | 3017.0 | -9.7 | -0.32 | B | 13.977 | 58.3 | 57.7 | -0.6 | -0.97 |
| H | 25.885 | 2960.3 | 2949.1 | -11.2 | -0.38 | B | $13.816^{6}$ | 52.95 | 52.5 | -0.4 | -0.77 |
| H | 25.561 | 2773.5 | 2762.5 | -11.0 | -0.40 | W | $13.803^{\text {b }}$ | 52.82 | 52.1 | -0.7 | -1.29 |
| V | 25.209 | 2578.1 | 2569.6 | -8.5 | -0.33 | H | $13.803^{\text {b }}$ | 52.80 | 52.1 | -0.7 | -1.25 |

gen reported by Roder et al. (7) is $M=2.01572$ grams/grammole. The evaluation of constant $D$ follows directly from the van der Waals constant, $a=27 R^{2} T_{c}^{2} / 64 P_{c}$. Thus,

$$
D=\frac{a}{R^{2}}=\frac{27 T_{c_{c}}{ }^{2}}{64 P_{c}}=\frac{27}{64} \frac{(32.976)^{2}}{(12.759 \times 760)}=0.047310\left({ }^{\circ} \mathrm{K} .\right)^{2} / \mathrm{mm} .
$$

When a reference point ( $T_{1}, P_{1}$ ), selected from a set of data, is substituted into Equation 2, the following expression results:

$$
\begin{equation*}
\ln P_{1}=A+\frac{B}{T_{1}}+C \ln T_{1}+D \frac{P_{1}}{T_{1}^{2}} \tag{3}
\end{equation*}
$$

By subtracting Equation 3 from Equation 2, and rearranging, it follows that:

$$
\begin{equation*}
\frac{\ln \frac{P}{P_{1}}-D\left[\frac{P}{T^{2}}-\frac{P_{1}}{T_{1}^{2}}\right]}{\ln \frac{T}{T_{1}}}=B \frac{\frac{1}{T}-\frac{1}{T_{1}}}{\ln \frac{T}{T_{1}}}+C \tag{4}
\end{equation*}
$$

If Equation 2 truly represents the behavior of the vapor pressure function and the reference is reliable, a plot of the vapor pressure modulus,

$$
\begin{equation*}
Y=\frac{\ln \frac{P}{P_{1}}-D\left[\frac{P}{T^{2}}-\frac{P_{1}}{T_{1}{ }^{2}}\right]}{\ln \frac{T}{T_{1}}} \tag{5}
\end{equation*}
$$

against the temperature modulus,

$$
X=\frac{\frac{1}{T}-\frac{1}{T_{1}}}{\ln \frac{T}{T_{1}}}
$$

should produce a straight line of slope $B$ and an intercept $C$. This premise has been substantiated with several substances (6), including the aliphatic hydrocarbons $(8,9)$.

The above procedure was followed in this study using the normal boiling point reported by each experimenter as the reference point for the corresponding set of data. The $X$ and $Y$ values resulting from the data of Hoge and Arnold (3), Weber, Diller, Roder, and Goodwin (11), and Barber and Horsford (1) were plotted as shown in Figure 1. The recent data of van Itterbeek et al. (10) have been used to calculate $X$ and $Y$ values which are also included for comparison in Figure 1. This was done to screen out any vapor pressure points that were in significant disagreement with the general trend of the data. Furthermore, this approach would indicate how well the behavior of parahydrogen conforms to Equation 2, as shown by the linearity of the $Y v s$. $X$ relationship. The data of the three sources produce a relationship that is linear from the critical point to nearly the triple point. As $T$ approaches $T_{1}$, Equation 6 becomes indeterminate. Application of L'Hospital's rule yields the limiting value

$$
\underset{T \rightarrow T_{1}}{\mathcal{L}} X=-1 / T_{1}=-1 / 20.268=-0.04934 .
$$

Because of the high sensitivity of the vapor pressure modulus, $Y$, near the reference point, some scatter can be encountered in its vicinity.

The method of least squares was applied to the values of Figure 1 to obtain a slope of $B=-105.651$ and an intercept of $C=0.425173$ to represent best the data of the three references. The value of $A=10.4791$ was established by substituting the normal boiling point reported by Hoge and Arnold (3), adjusted to the N.B.S.- 1955 temperature scale, $T_{b}=20.268^{\circ}$ K . in to Equation 2 along with the already established values for $B, C$, and $D$. The final equation for parahydrogen over the complete range included between the triple and critical points becomes:

$$
\begin{equation*}
\ln P=10.4791-\frac{105.651}{T}+0.425173 \ln T+0.047310 \frac{P}{T^{2}} \tag{7}
\end{equation*}
$$

where $P$ is in millimeters of mercury and $T$ is in degrees Kelvin. Vapor pressures calculated with Equation 7 are presented in


Figure 2. Relationship between $Q$ and $S$ for para hydrogen

Table II. Comparison of Experimental and Calculated Vapor Pressures of Parahydrogen Using Equation 8

$$
\ln P=11.2586-\frac{110.157}{T}+0.240055 \ln T+0.047310 \frac{P}{T^{2}}
$$

Millimeters of Mercury

|  | $T{ }^{\circ} \mathrm{K}$. | $P_{\text {exptl }}$ | $P_{\text {calcd }}$ | $\Delta P$ | Dev. \% |  | $T,{ }^{\circ} \mathrm{K}$. | $P_{\text {expt }}$ | $P_{\text {calcd }}$ | $\Delta P$ | Dev. \% |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| W | $32.976^{\text {a }}$ | 9696.8 | 9690.2 | -6.6 | -0.07 | H | 25.037 | 2488.5 | 2489.6 | 1.1 | 0.04 |
| W | 32.900 | 9589.6 | 9585.0 | -4.6 | -0.05 | W | 25.000 | 2467.1 | 2470.1 | 3.0 | 0.12 |
| H | 32.883 | 9566.2 | 9561.9 | -4.3 | -0.04 | H | 24.890 | 2414.1 | 2413.6 | -0.5 | -0.02 |
| H | 32.884 | 9564.4 | 9562.4 | -2.0 | -0.02 | V | 24.680 | 2301.5 | 2307.8 | 6.3 | 0.27 |
| H | 32.883 | 9559.3 | 9561.0 | 1.7 | 0.02 | H | 24.440 | 2189.4 | 2191.2 | 1.8 | 0.08 |
| W | 32.800 | 9449.5 | 9447.8 | $-1.7$ | -0.02 | H | 23.634 | 1827.1 | 1829.6 | 2.5 | 0.14 |
| W | 32.700 | 9312.0 | 9312.2 | 0.2 | 0.00 | W | 23.000 | 1572.3 | 1576.3 | 4.0 | 0.26 |
| H | 32.636 | 9219.5 | 9225.8 | 6.3 | 0.07 | H | 22.896 | 1534.9 | 1537.2 | 2.3 | 0.15 |
| W | 32.600 | 9176.8 | 9178.1 | 1.3 | 0.01 | V | 22.772 | 1488.7 | 1491.6 | 2.9 | 0.19 |
| W | 32.500 | 9042.8 | 9045.4 | 2.6 | 0.03 | H | 22.569 | 1416.1 | 1419.0 | 2.9 | 0.20 |
| H | 32.375 | 8875.1 | 8882.1 | 7.0 | 0.08 | H | 22.270 | 1313.8 | 1316.5 | 2.7 | 0.20 |
| H | 32.129 | 8557.7 | 8566.3 | 8.6 | 0.10 | H | 22.250 | 1308.4 | 1310.0 | 1.6 | 0.12 |
| W | 32.000 | 8399.0 | 8403.9 | 4.9 | 0.06 | W | 22.000 | 1225.4 | 1228.6 | 3.2 | 0.26 |
| H | 31.881 | 8255.1 | 8256.5 | 1.4 | 0.02 | H | 21.677 | 1127.2 | 1129.0 | 1.8 | 0.16 |
| W | 31.500 | 7792.5 | 7797.3 | 4.8 | 0.06 | H | 21.328 | 1026.8 | 1027.8 | 1.0 | 0.10 |
| H | 31.392 | 7660.2 | 7670.8 | 10.6 | 0.14 | H | 21.195 | 989.8 | 990.9 | 1.1 | 0.11 |
| V | 31.238 | 7503.2 | 7492.5 | -10.7 | -0.14 | H | 20.943 | 922.5 | 924.0 | 1.5 | 0.16 |
| H | 31.074 | 7302.5 | 7304.5 | 2.0 | 0.03 | H | 20.941 | 922.1 | 923.5 | 1.4 | 0.15 |
| W | 31.000 | 7221.1 | 7223.9 | 2.8 | 0.04 | H | 20.855 | 900.7 | 901.4 | 0.7 | 0.08 |
| V | 30.971 | 7200.2 | 7191.3 | -8.9 | -0.12 | H | 20.507 | 815.2 | 815.5 | 0.3 | 0.03 |
| H | 30.892 | 7102.8 | 7104.2 | 1.4 | 0.02 | H | 20.502 | 813.7 | 814.3 | 0.6 | 0.07 |
| W | 30.500 | 6679.2 | 6682.2 | 3.0 | 0.04 | H | 20.397 | 789.6 | 789.7 | 0.1 | 0.01 |
| H | 30.368 | 6544.6 | 6543.8 | -0.8 | -0.01 | B | 20.270 | 760.0 | 760.7 | 0.7 | 0.09 |
| V | 30.172 | 6357.3 | 6343.1 | -14.2 | -0.22 | H | 20.268 | 760.0 | 760.1 | 0.1 | 0.01 |
| W | 30.000 | 6169.0 | 6170.7 | 1.7 | 0.03 | W | 20.268 | 760.0 | 760.1 | 0.1 | 0.01 |
| H | 29.907 | 6080.2 | 6079.1 | -1.1 | -0.02 | H | 20.255 | 757.1 | 757.1 | 0.0 | 0.00 |
| V | 29.238 | 5453.3 | 5446.4 | -6.9 | -0.13 | H | 20.030 | 707.7 | 707.7 | 0.0 | 0.00 |
| H | 29.386 | 5583.9 | 5581.8 | -2.1 | -0.04 | H | 19.114 | 530.7 | 529.8 | -0.9 | -0.17 |
| W | 29.000 | 5233.6 | 5233.6 | 0.0 | 0.00 | B | 19.050 | 519.5 | 518.7 | -0.8 | -0.16 |
| H | 28.870 | 5121.5 | 5119.5 | -2.0 | -0.04 | H | 18.571 | 442.1 | 440.9 | -1.2 | -0.27 |
| H | 28.376 | 4705.7 | 4703.2 | -2.5 | -0.05 | B | 18.447 | 423.3 | 422.3 | -1.0 | -0.24 |
| W | 28.000 | 4401.9 | 4403.1 | 1.2 | 0.03 | H | 17.819 | 338.4 | 336.4 | -2.0 | -0.59 |
| H | 27.864 | 4299.6 | 4298.2 | -1.4 | -0.03 | B | 17.429 | 291.3 | 289.9 | -1.4 | -0.49 |
| V | 27.540 | 4049.9 | 4054.3 | 4.4 | 0.11 | H | 16.965 | 243.1 | 241.0 | -2.1 | -0.87 |
| H | 27.398 | 3952.7 | 3951.1 | -1.6 | -0.04 | H | 16.945 | 241.4 | 239.0 | -2.4 | -1.00 |
| V | 27.072 | 3718.2 | 3720.2 | 2.0 | 0.05 | H | 15.831 | 149.7 | 147.1 | -2.6 | -1.72 |
| W | 27.000 | 3669.7 | 3670.7 | 1.0 | 0.03 | B | 15.348 | 119.4 | 116.8 | -2.6 | -2.14 |
| H | 26.771 | 3517.5 | 3516.0 | -1.5 | -0.04 | B | 15.005 | 100.8 | 98.3 | -2.5 | -2.50 |
| H | 26.188 | 3142.6 | 3142.9 | 0.3 | 0.01 | B | 14.524 | 78.6 | 76.2 | -2.4 | -3.08 |
| W | 26.000 | 3026.7 | 3028.7 | 2.0 | 0.07 | B | 13.977 | 58.3 | 55.9 | -2.4 | -4.05 |
| H | 25.885 | 2960.3 | 2960.7 | 0.4 | 0.01 | B | $13.816^{\text {b }}$ | 52.95 | 50.8 | -2.1 | -4.01 |
| H | 25.561 | 2773.5 | 2773.7 | 0.2 | 0.01 | W | $13.803^{\text {b }}$ | 52.82 | 50.4 | -2.4 | -4.52 |
| V | 25.209 | 2578.1 | 2580.3 | 2.2 | 0.08 | H | $13.803^{\text {b }}$ | 52.80 | 50.4 | -2.4 | -4.49 |

Millimeters of Mercury
${ }^{a}$ Reported critical point.
${ }^{b}$ Reported triple point.

Table I along with the corresponding experimental values and resulting deviations. Each point is identified in this table by B for Barber and Hosford (1), H for Hoge and Arnold (3), V for van Itterbeek et al. (10) and W for Weber et al. (11). These calculations were carried out by an iterative scheme on a digital computer. These iterations were carried out until the correction between successive values of $P$ was less than 0.01 mm . The deviations reported in Table I between experimental and calculated values are small, but deviate, systematically from the triple point to the critical point. A maximum deviation of $1.29 \%$ is encountered at the triple point.

Despite the fact that Equation 7 is capable of producing vapor pressures well within ordinary limits of accuracy, improvements could be made to represent more precisely the vapor pressure behavior of parahydrogen.

The values in the low pressure region of the $Y, X$ relationship are slightly offset from the otherwise straight line. On the basis of this observation, it was decided to fit this low pressure region with a separate vapor pressure-temperature relationship. Therefore, a new straight line was fitted through the ( $Y, X$ ) values corresponding to vapor pressures above 1 atm . and placing more
weight on the values in the high pressure region. Again the method of least squares was applied to produce a slope of $B=$ -110.157 and an intercept of $C=0.240055$. Constant $A$ was determined using the vapor pressure value of Weber et al. (11), $T=29.000^{\circ} \mathrm{K}$. and $P=6.8863 \mathrm{~atm}$. ( 5233.6 mm .). This procedure yielded $A=11.2586$. Therefore, the vapor pressure equation for the high pressure range becomes,

$$
\begin{equation*}
\ln P=11.2586-\frac{110.157}{T}+0.240055 \ln T+0.047310 \frac{P}{T^{2}} \tag{8}
\end{equation*}
$$

Vapor pressures calculated with Equation 8 are presented in Table II along with the corresponding experimental values and their deviations. These deviations remain small for pressures down to the normal boiling point and then increase to a maximum of $4.52 \%$ at the triple point. In view of this, the application of Equation 8 should be restricted to vapor pressures above 760 mm . of mercury.

Since the last term of Equation 2 accounts for the reversal in curvature of the $\ln P v s .1 / T$ relationship in the vicinity of the

Table III. Comparison of Experimental and Calculated Vapor Pressures of Parahydrogen Using Equation 11

$$
\ln P=5.83882-\frac{87.2596}{T}+1.69482 \ln T
$$

|  |  | Millimeters of Mercury |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | T, ${ }^{\circ} \mathrm{K}$ | $P_{\text {expt }}$ | $P_{\text {caled }}$ | $\Delta P$ | Dev. \% |
| H | $13.803^{\text {a }}$ | 52.80 | 52.76 | -0.04 | -0.08 |
| W | $13.803^{\text {a }}$ | 52.82 | 52.76 | -0.06 | -0.11 |
| B | $13.816^{\circ}$ | 52.95 | 53.16 | 0.21 | 0.39 |
| B | 13.977 | 52.28 | 58.30 | 0.02 | 0.03 |
| B | 14.524 | 78.62 | 78.70 | 0.08 | 0.10 |
| B | 15.005 | 100.84 | 100.87 | 0.02 | 0.02 |
| B | 15.348 | 119.38 | 119.36 | -0.01 | -0.01 |
| H | 15.831 | 149.70 | 149.62 | -0.08 | -0.05 |
| B | 16.288 | 183.03 | 183.28 | 0.25 | 0.14 |
| H | 16.945 | 241.40 | 241.17 | -0.23 | -0.10 |
| H | 16.965 | 243.10 | 243.15 | 0.05 | 0.02 |
| B | 17.429 | 291.29 | 291.81 | 0.52 | 0.18 |
| H | 17.819 | 338.40 | 338.15 | -0.25 | -0.07 |
| B | 18.447 | 423.28 | 423.65 | 0.37 | 0.08 |
| H | 18.571 | 442.10 | 442.21 | 0.11 | 0.02 |
| B | 19.050 | 519.53 | 519.63 | 0.10 | 0.02 |
| H | 19.114 | 530.70 | 530.70 | 0.00 | 0.00 |
| H | 20.030 | 707.70 | 707.81 | 0.11 | 0.02 |
| H | 20.255 | 757.10 | 757.03 | $-0.07$ | -0.01 |
| W | 20.268 | 760.00 | 760.00 | 0.00 | 0.00 |
| H | 20.268 | 760.00 | 760.00 | 0.00 | 0.00 |
| B | 20.270 | 760.00 | 760.56 | 0.56 | 0.07 |
| H | 20.397 | 789.60 | 789.40 | -0.20 | -0.02 |
| H | 20.502 | 813.70 | 813.91 | 0.21 | 0.02 |
| H | 20.507 | 815.20 | 815.07 | $-0.13$ | -0.02 |
| H | 20.855 | 900.70 | 900.57 | $-0.13$ | -0.01 |
| H | 20.941 | 922.10 | 922.54 | 0.44 | 0.05 |
| H | 20.943 | 922.50 | 923.08 | 0.58 | 0.06 |
| H | 21.195 | 989.80 | 989.60 | -0.20 | -0.52 |
| H | 21.328 | 1026.8 | 1026.2 | -0.6 | -0.05 |
| H | 21.677 | 1127.2 | 1126.8 | $-0.4$ | -0.03 |
| W | 22.000 | 1225.4 | 1225.7 | 0.3 | 0.02 |
| H | 22.250 | 1308.4 | 1306.4 | $-2.0$ | -0.15 |
| H | 22.270 | 1313.8 | 1312.9 | -0.9 | -0.07 |
| H | 22.569 | 1416.1 | 1414.5 | $-1.5$ | -0.11 |
| H | 22.896 | 1534.9 | 1531.6 | $-3.3$ | -0.22 |

${ }^{a}$ Reported triple point.
critical point (9), it becomes unnecessary to carry it along for the establishment of a vapor pressure relationship in the low pressure region. Therefore, the truncated form of Equation 2 becomes,

$$
\begin{equation*}
\ln P=A+\frac{B}{T}+C \ln T \tag{9}
\end{equation*}
$$

Again, a reference point ( $P_{1}, T_{1}$ ) is chosen and substituted into Equation 9. The resulting equation is then subtracted from Equation 9 and rearranged to yield the expression,

$$
\begin{equation*}
\frac{\ln \frac{P}{P_{1}}}{\frac{1}{T}-\frac{1}{T_{1}}}=B+C \frac{\ln \frac{T}{T_{1}}}{\frac{1}{T}-\frac{1}{T_{1}}} \tag{10}
\end{equation*}
$$

The reference point selected for this procedure was again the normal boiling point.

The vapor pressure modulus, $Q=\left(\ln P / P_{1}\right) /\left(1 / T-1 / T_{1}\right)$ was plotted against the temperature modulus $S=\left(\ln T / T_{1}\right) /$ $\left(1 / T-1 / T_{1}\right)$ as shown in Figure 2. The fact that the $Q, S$ relationship is linear in the low pressure region and curved in the high pressure region supports the assumption that the term $D P / T^{2}$ is negligible at low pressure. Application of L'Hos-


Figure 3. Latent heat of vaporization for parahydrogen
pital's rule to the temperature modulus yie'ds $\underset{T \rightarrow T_{1}}{\mathfrak{L}}\left(\frac{\mathrm{n}}{} T / T_{1}\right) /$ $\left(1 / T-1 / T_{1}\right)=-T=-20.268$. The method of least squares was applied to the values from the triple point to the normal boiling point. The results of this analysis produced an intercept $B=-87.2596$ and a slope of $C=1.69482$. The normal boiling point was then used to determine constant $A=5.83882$. Thus the vapor pressure function for parahydrogen covering the range from the triple point to the normal boiling point becomes,

$$
\begin{equation*}
\ln P=5.83882-\frac{87.2596}{T}+1.69482 \ln T \tag{11}
\end{equation*}
$$

Yalues calculated with Equation 11, along with corresponding experimental values and deviations are presented in Table III. A review of these comparisons indicates that Equation 11 is capable of predicting accurate vapor pressures up to almost 1400 mm . Since the objective of Equation 11 is to define the vapor pressure behavior of parahydrogen up to the boiling point, this equation more than fulfills its purpose. Therefore, the combined results of Equations 11 and 8 enable the calculation of accurate vapor pressures from the triple point to the critical point. At the normal boiling point, these equations properly blend into each other as indicated by their derivatives at this point. For Equation $8, \mathrm{~d} P / \mathrm{d} T=226.0 \mathrm{~mm} . /^{\circ} \mathrm{K}$, and for Equation 11, $\mathrm{d} P / \mathrm{d} T=225.0 \mathrm{~mm} . /^{\circ} \mathrm{K}$.

## LATENT HEATS OF VAPORIZATION

Equations 8 and 11 have been used to obtain values of the slope, $\mathrm{d} P / \mathrm{d} T$, along the vapor pressure curve. For pressures above the normal boiling point, the differentiated form of Equation 8,

$$
\begin{equation*}
\frac{\mathrm{d} P}{\mathrm{~d} T}=\frac{0.240055 / T+110.157 / T^{2}+0.094620 P / T^{3}}{1 / P-0.047310 / T^{2}} \tag{12}
\end{equation*}
$$

Table IV. Derived Quantities Used to Calculate Latent Heats of Vaporization

| $T,{ }^{\circ} \mathrm{K}$. | $P, \mathrm{~mm}$. | Cc./Gram-Mole |  | $\begin{gathered} \mathrm{d} P / \mathrm{d} T \\ \operatorname{Atm} . /^{\circ} \mathrm{K} . \end{gathered}$ | $\lambda$ <br> Cal/ <br> Gram- <br> Mole | T, ${ }^{\circ} \mathrm{K}$. | $P, \mathrm{~mm}$. | Cc./Gram-Mole |  | $\begin{gathered} \mathrm{d} P / \mathrm{d} T \\ \operatorname{Atm} . /^{\circ} \mathrm{K} . \end{gathered}$ | $\lambda$ <br> Cal./ <br> GramMole |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $v_{0}$ | $v{ }^{2}$ |  |  |  |  | $v_{0}$ | $v$ |  |  |
| $13.803^{\text {a }}$ | 52.8 | 16,025 | 26.173 | 0.0403 | 215.6 | 29.300 | 5503.0 | 213.64 | 36.076 | 1.2035 | 151.6 |
| 14.000 | 59.1 | 14,503 | 26.230 | 0.0440 | 216.0 | 29.400 | 5595.0 | 209.38 | 36.247 | 1.2179 | 150.1 |
| 14.500 | 77.7 | 11,378 | 26.376 | 0.0544 | 216.8 | 29.500 | 5688.1 | 205.19 | 36.423 | 1.2325 | 148.6 |
| 15.000 | 100.6 | 9055.3 | 26.526 | 0.0663 | 217.4 | 29.600 | 5782.3 | 201.06 | 36.604 | 1.2472 | 147.0 |
| 15.500 | 128.3 | 7301.9 | 26.682 | 0.0798 | 217.9 | 29.700 | 5877.7 | 197.00 | 36.789 | 1.2620 | 145.4 |
| 16.000 | 161.4 | 5959.5 | 26.842 | 0.0949 | 218.2 | 29.800 | 5974.1 | 193.00 | 36.981 | 1.2769 | 143.8 |
| 16.500 | 200.6 | 4917.7 | 27.008 | 0.1117 | 218.4 | 29.900 | 6071.8 | 189.05 | 37.178 | 1.2919 | 142.1 |
| 17.000 | 246.6 | 4098.8 | 27.179 | 0.1303 | 218.4 | 30.000 | 6170.5 | 185.16 | 37.380 | 1.3071 | 140.4 |
| 17.500 | 299.9 | 3447.5 | 27.357 | 0.1507 | 218.4 | 30.100 | 6270.4 | 181.33 | 37.590 | 1.3225 | 138.6 |
| 18.000 | 361.3 | 2923.7 | 27.542 | 0.1728 | 218.2 | 30.200 | 6371.5 | 177.55 | 37.806 | 1.3379 | 136.7 |
| 18.500 | 431.5 | 2498.1 | 27.734 | 0.1968 | 217.8 | 30.300 | 6473.8 | 173.82 | 38.030 | 1.3535 | 134.9 |
| 19.000 | 511.1 | 2149.0 | 27.933 | 0.2225 | 217.2 | 30.400 | 6577.3 | 170.13 | 38.261 | 1.3692 | 132.9 |
| 19.500 | 600.8 | 1860.1 | 28.141 | 0.2501 | 216.4 | 30.500 | 6681.9 | 166.49 | 38.501 | 1.3851 | 130.9 |
| 20.000 | 701.4 | 1619.1 | 28.358 | 0.2795 | 215.4 | 30.600 | 6787.8 | 162.89 | 38.749 | 1.4011 | 128.9 |
| 20.268 | 760.0 | 1506.8 | 28.479 | 0.2960 | 214.8 | 30.700 | 6894.9 | 159.34 | 39.008 | 1.4172 | 126.8 |
| 20.500 | 813.8 | 1417.1 | 28.586 | 0.3123 | 215.3 | 30.800 | 7003.2 | 155.82 | 39.277 | 1.4335 | 124.6 |
| 21.000 | 938.8 | 1245.3 | 28.824 | 0.3456 | 213.8 | 30.900 | 7112.8 | 152.33 | 39.557 | 1.4500 | 122.4 |
| 21.500 | 1076.8 | 1098.8 | 29.073 | 0.3809 | 212.2 | 31.000 | 7223.6 | 148.88 | 39.849 | 1.4666 | 120.0 |
| 22.000 | 1228.5 | 973.39 | 29.335 | 0.4181 | 210.3 | 31.100 | 7335.7 | 145.45 | 40.155 | 1.4833 | 117.6 |
| 22.500 | 1394.8 | 865.34 | 29.612 | 0.4573 | 208.2 | 31.200 | 7449.1 | 142.05 | 40.476 | 1.5002 | 115.1 |
| 23.000 | 1576.3 | 771.82 | 29.903 | 0.4984 | 206.0 | 31.300 | 7563.8 | 138.67 | 40.813 | 1.5173 | 112.6 |
| 23.500 | 1773.8 | 690.48 | 30.212 | 0.5415 | 203.5 | 31.400 | 7679.7 | 135.31 | 41.168 | 1.5345 | 109.9 |
| 24.000 | 1988.1 | 619.37 | 30.539 | 0.5866 | 200.8 | 31.500 | 7797.0 | 131.96 | 41.543 | 1.5519 | 107.0 |
| 24.500 | 2219.9 | 556.93 | 30.887 | 0.6338 | 197.8 | 31.600 | 7915.6 | 128.62 | 41.941 | 1.5695 | 104.1 |
| 25.000 | 2470.1 | 501.84 | 31.259 | 0.6832 | 194.6 | 31.700 | 8035.6 | 125.29 | 42.365 | 1.5872 | 101.0 |
| 25.500 | 2739.4 | 453.04 | 31.656 | 0.7346 | 191.2 | 31.800 | 8156.9 | 121.94 | 42.818 | 1.6051 | 97.81 |
| 26.000 | 3028.7 | 409.62 | 32.084 | 0.7883 | 187.4 | 31.900 | 8279.6 | 118.59 | 43.305 | 1.6232 | 94.40 |
| 26.500 | 3338.8 | 370.81 | 32.545 | 0.8442 | 183.3 | 32.000 | 8403.6 | 115.21 | 43.832 | 1.6414 | 90.79 |
| 27.000 | 3670.6 | 336.00 | 33.045 | 0.9025 | 178.8 | 32.100 | 8529.1 | 111.79 | 44.406 | 1.6599 | 86.95 |
| 27.250 | 3844.9 | 319.91 | 33.311 | 0.9325 | 176.4 | 32.200 | 8655.9 | 108.32 | 45.038 | 1.6785 | 82.83 |
| 27.500 | 4025.0 | 304.62 | 33.590 | 0.9632 | 173.9 | 32.300 | 8784.2 | 104.78 | 45.742 | 1.6973 | 78.38 |
| 27.750 | 4211.0 | 290.08 | 33.882 | 0.9945 | 171.2 | 32.400 | 8913.9 | 101.13 | 46.536 | 1.7163 | 73.52 |
| 28.000 | 4402.9 | 276.24 | 34.188 | 1.0264 | 168.5 | 32.500 | 9045.1 | 97.334 | 47.452 | 1.7355 | 68.14 |
| 28.250 | 4601.0 | 263.03 | 34.510 | 1.0590 | 165.6 | 32.500 32.600 | 9177.7 | 93.315 | 48.538 | 1.7549 | 62.04 |
| 28.500 28.750 | 4805.4 | 250.43 238.38 | 34.850 35.208 | 1.0923 1.1263 | 162.5 | 32.600 32.700 | 9311.8 | 88.949 | 49.884 | 1.7745 | 54.90 |
| 29.000 | 5233.4 | 226.85 | 35.208 35.588 | 1.1263 1.1609 | 155.9 | 32.800 | 9447.5 | 83.969 | 51.687 | 1.7944 | 46.01 |
| 29.100 | 5322.2 | 222.37 | 35.747 | 1.1750 | 154.5 | 32.900 | 9584.6 | 77.567 | 54.571 | 1.8144 | 33.24 |
| 29.200 | 5412.0 | 217.97 | 35.910 | 1.1892 | 153.1 | $32.976^{\text {b }}$ | 9689.8 | 64.144 | 64.144 | 1.8298 | 0.00 |

Table V. Comparison of Latent Heats of Vaporization Resulting from this Study with Experimental and Calculated Values of Others

|  | $\lambda$, Cal./Gram-Mole |  |  |  |  | $\lambda$, Cal./Gram-Mole |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $T,{ }^{\circ} \mathrm{K}$. | This study | White et al. (11) ${ }^{a}$ | Mullins et al. (5) | $\begin{gathered} \text { Roder } \\ \text { et al. (7) } \end{gathered}$ | $T,{ }^{\circ} \mathrm{K}$. | This study | White et al. (11) ${ }^{a}$ | Mullins et al. (5) | Roder et al. (7) |
| 13.803 | 215.6 |  | 216.87 | 216.8 | 24.000 | 200.8 |  |  | 200.8 |
| 13.990 | 216.0 |  | 217.07 | 217.1 | 24.41 | 198.4 | 204.7 | $\ldots$ | 198.5 |
| 14.000 | 216.0 |  |  | 217.1 | 25.000 | 194.6 | ... | $\ldots$ | 195.0 |
| 14.990 | 217.4 |  | 217.93 | 218.3 | 26.000 | 187.4 |  | $\ldots$ | 187.8 |
| 15.000 | 217.4 |  |  | 218.3 | 26.33 | 184.7 | 193.4 | $\ldots$ | 185.2 |
| 15.990 | 218.2 |  | 218.44 | 218.6 | 27.000 | 178.8 |  |  | 179.2 |
| 16.000 | 218.2 | $\cdots$ |  | 218.5 | 28.000 | 168.5 |  | $\ldots$ | 168.7 |
| 16.990 | 218.4 |  | 218.50 | 218.4 | 28.12 | 167.1 | 177.5 | $\ldots$ | 167.3 |
| 17.000 | 218.4 |  |  | 218.4 | 29.000 | 155.9 |  |  | 155.8 |
| 17.990 | 218.2 |  | 218.05 | 217.9 | 29.65 | 146.2 | 145.2 |  | 145.9 |
| 18.000 | 218.2 |  |  | 217.9 | 30.000 | 140.4 |  |  | 140.1 |
| 18.990 | 217.2 |  | 217.03 | 216.8 | 30.97 | 120.8 | 116.5 |  | 120.5 |
| 19.000 | 217.2 |  |  | 216.8 | 31.000 | 120.0 |  |  | 119.8 |
| 19.990 | 215.4 |  | 215.38 | 215.3 | 31.85 | 96.13 | 93.8 |  | 96.1 |
| 20.000 | 215.4 |  |  | 215.2 | 32.000 | 90.79 |  |  | 90.8 |
| 20.268 | 214.8 |  | $214.80^{\text {b }}$ | $214.8{ }^{\text {b }}$ | 32.400 | 73.52 |  |  | 73.9 |
| 21.000 | 213.8 |  | 214.03 | 212.5 | 32.69 | 55.67 | 51.8 |  | 56.4 |
| 22.000 | 210.3 |  | 209.94 | 209.5 | 32.700 | 54.90 | . $\cdot$ |  | 55.6 |
| 23.000 | 206.0 |  |  | 205.6 | 32.900 | 33.24 | . | $\ldots$ | 33.9 |

[^0]was used to calculate these slopes at convenient temperature intervals. From the triple point up to and including the normal boiling point, the differentiated form of Equation 11,
\[

$$
\begin{equation*}
\frac{\mathrm{d} P}{\mathrm{~d} T}=\frac{P}{T}\left[1.69482+\frac{87.2596}{T}\right] \tag{13}
\end{equation*}
$$

\]

was used to obtain the values of the slopes at convenient intervals. The saturated vapor and liquid molar volumes, $v_{0}$ and $v_{l}$, respectively, were obtained from the saturated density relationships of Roder et al. (7), which were developed from their experimental studies. Using this information, the latent heats of vaporization for parahydrogen were calculated from the rearranged Clapeyron equation,

$$
\begin{equation*}
\lambda=\left(v_{v}-v_{l}\right) T \frac{\mathrm{~d} P}{\mathrm{~d} T} \tag{14}
\end{equation*}
$$

The resulting latent heats of vaporization in cc. mm./grammole were divided by the conversion factors 760 mm ./atm. and $41.2929 \mathrm{cc} . \mathrm{atm}$./cal. to produce values in cal./gram-mole. Values of slopes, saturated vapor, and liquid molar volumes, and latent heats of vaporization are presented for convenient temperatures in Table IV. These calculated latent heats of vaporization have been plotted to produce Figure 3. It will be noted in this figure that the latent heat function of parahydrogen reaches a maximum at approximately $17.0^{\circ} \mathrm{K}$.

To compare these results with the experimental latent heats of vaporization of White, Hu , and Johnston (2), values were calculated at their corresponding experimental temperatures. These comparisons are presented in Table V. Also included in this table are the calculated values of Mullins, Ziegler, and Kirk ( $b$ ) and Roder et al. (7).
The latent heats of vaporization of this investigation presented in Table V show good agreement when compared with the experimental and calculated values reported by others (5,7, 12). Mullins, Ziegler, and Kirk ( 5 ) calculated latent heats of vaporization from the triple point up to and including $22.0^{\circ}$ K., using a virial-type equation of state. The only experimental values for temperatures above the normal boiling point are presented by White, Hu, and Johnston (12). The values reported by Roder, Diller, Weber, and Goodwin( $\gamma^{\prime}$ ) were calculated using the Clapeyron equation and three separate vapor pressure relationships applicable in different regions included between the triple point and the critical point. The values resulting from this study also cover the complete vapor-liquid range and were found to be in good agreement with those already reported in the literature. These findings add creditability to the work already cited ( $5,7,12$ ) in which the equations utilized to calculate the latent heat of vaporization of parahydrogen differed in each case.

## ACKNOWLEDGMENT

The authors express gratitude to the National Aeronautics and Space Administration for support through NASA Grant SC-3467. Also thanks are due the Northwestern University Computing Center for the allocation of computing time for this study.

## NOMENCLATURE

$a=$ van der Waals constant, (cc./gram-mole) ${ }^{2}$ atm.
$A, B, C, D=$ constants for Frost-Kalkwarf equation, Equation (2) $M=$ molecular weight
$P=$ vapor pressure, mm of mercury
$P_{1}=$ vapor pressure of reference point, mm . of mercury
$P_{c}=$ critical pressure, atm
$P_{t}=$ triple point pressure, mm . of mercury
$Q=$ vapor pressure modulus, $\left(\ln P / P_{1}\right) /\left(1 / T-1 / T_{1}\right)$
$R=$ gas constant
$S=$ temperature modulus, $\left(\ln T / T_{1}\right) /\left(1 / T-1 / T_{1}\right)$
$T=$ temperature, ${ }^{\circ} \mathrm{K}$.
$T_{1}=$ temperature of reference point, ${ }^{\circ} \mathrm{K}$.
$T_{b}=$ normal boiling point, ${ }^{\circ} \mathrm{K}$.
$T_{c}=$ critical temperature, ${ }^{\circ} \mathrm{K}$.
$T_{t}=$ triple point temperature, ${ }^{\circ} \mathrm{K}$.
$v_{c}=$ critical volume, cc./gram-mole
$v_{a}=$ molar volume of saturated vapor, cc./gram-mole
$v_{l}=$ molar volume of saturated liquid, cc./gram-mole
$X=$ temperature modulus, $\left(1 / T-1 / T_{1}\right) /\left(\ln T / T_{1}\right)$
$Y=$ vapor pressure modulus, $\left\{\ln P / P_{1}-D\left(P / T^{2}-\right.\right.$ $\left.\left.P_{1} / T_{1}{ }^{2}\right)\right\} /\left(\ln T / T_{1}\right)$
Greek
$\lambda=$ latent heat of vaporization, cal./gram-mole
$\rho_{c}=$ critical density, gram-moles/cc.

## LITERATURE CITED

(1) Barber, C. R., Horsford, A., Brit. J. Appl. Phys. 14, 920 (1963).
(2) Frost, A. A., Kalkwarf, D. R., J. Chem. Phys. 21, 264 (1953).
(3) Hoge, H. J., Arnold, R. D., J. Res. Natl. Bur. Stds. 47, 63 (1951).
(4) Miller, D. G., Ind. Eng. Chem. 56, No. 3, 46 (1964).
(5) Mullins, J. C., Ziegler, W. T., Kirk, B. S., Tech. Rept. 1, Eng. Expt. Station, Georgia Institute of Technology, Nov. 1, 1961.
(6) Reynes, E. G., Thodos, George, Ind. Eng. Chem. Fundamentals 1, 127 (1962).
(7) Roder, H. M., Diller, D. E., Weber, L. A., Goodwin, R. D., Cryogenics 3, 16 (1963).
(8) Smith, C. H., Thodos, George, A. I.Ch.E.J. 6, 569 (1960).
(9) Sondak, N. E., Thodos, George, Ibid., 2, 347 (1956).
(10) van Itterbeek, A., Verbeke, O., Theewes, F., Staes, K., De Boelpacp, J., Physica 30, 1238 (1964).
(11) Weber, L. A., Diller, D. E., Roder, H. M., Goodwin, R. D., Cryogenics 2, 236 (1962).
(12) White, D., Hu, J. H., Johnston, H. L., J. Phys. Chem. 63, 1181 (1959).

Received for review January 26, 1965. Accepted October 11, 1965.


[^0]:    a Recalculated by Roder et al. (7) as suggested by White et al. (11).
    ${ }^{b}$ Defined by Roder et al. (7).

