- v = molar volume, cc./gram-mole
- critical volume, cc./gram-mole  $v_{c} =$
- $z_c =$ critical compressibility factor,  $P_{c}v_{c}/RT_{c}$

#### Greek

- D =self-diffusivity, sq. cm./sec.
- = viscosity, g./cm. sec. ш
- μ\* viscosity at atmospheric pressure, g./cm. sec. =
- constant, 3.14159  $\pi =$
- density, g./cc. ρ =
- $\rho_c$  = critical density, g./cc.
- $\rho_R$  = reduced density,  $\rho/\rho_c$
- $\sigma$  = collision diameter, cm.
- $\chi$  = probability of nearness

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# **Thermodynamic Properties of Uranium Carbides** via the U-C-O System

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Equilibrium data of Piazza and Sinnott in the systems UO2-"UC2"-C-CO and UO2-"UC2"-UC-CO have been recalculated using both second and third law treatments. Internal consistency is excellent for the first system and acceptable for the second. Third law treatments predict the heats of formation,  $\Delta H_{286}^{o}$ , "UC<sub>2</sub>" and UC to be -20.5 and -23.2 kcal. per mole, respectively, in excellent agreement with combustion measurements. The "best" value of  $\Delta H_{128}^{\circ}$  "UC<sub>2</sub>" is -20.8° 0.7 kcal. per mole.

 $\mathbf{I}$ 'HE influence of oxygen on uranium dicarbide could be important in reactor technology. A panel assembled in Vienna in 1962 by the International Atomic Energy Agency (7) assumed the existence of an "oxygen-stabilized uranium dicarbide,"  $UC_xO_y$ . It has been estimated that the free energy of formation of oxygen-saturated "UC2" is about 5 kcal. more negative than that of the pure dicarbide (7, 15), from which arises the term "oxygen-stabilized."

Piazza and Sinnott studied equilibria in the system UO<sub>2</sub>-"UC $_2$ "-C-CO by measuring CO partial pressures (14). For their stated reaction

$$UO_2 + 4C = UC_2 + 2CO$$
 (1)

Piazza and Sinnott reported

$$\Delta F_{\rm F} = 164,500 - 74.23 \ T \pm 1200 \ \text{cal. per gram-mole}$$
 (2)

over the range 1714° to 1922° K. The Vienna Panel stated that "the derived heats and entropies are rather unreasonable." This note points out that auxiliary thermal data are now available to treat these measurements of Piazza and Sinnott in the usual fashion; that the internal consistency of these data is excellent; that reasonable heats and entropies are obtained; and that the data, combined with other information, predict that any stabilization by oxygen must be less than a few hundred calories. We also treat data on the  $UO_2$ -"UC<sub>2</sub>"-UC-CO system by the same authors (14).

Two methods are customarily used to treat data of the type obtained by Piazza and Sinnott: the second law slopeintercept treatment and the third law method which involves the use of free energy functions (9). We have used these standard methods and the CO partial pressures of Piazza and Sinnott to treat the reaction

$$UO_2 + 3.86C = UC_{1.86} + 2CO$$
(3)

Following Piazza and Sinnott, we estimated the activity of UC1.86 to be 0.95 in all experiments they performed and  $a_{\rm UO}$ , and  $a_{\rm C}$  to be unity. Auxiliary data have been obtained as follows. Uranium oxide thermal functions were calculated from the data of Conway and Hein (1) combined with the data of Moore and Kelly (13). The heat of formation of  $UO_2$ ,  $\Delta H_{1,298}^{\circ}$ , was taken to be -259.0 kcal. per mole from the compilation of Rand and Kubaschewski (16). All data on CO were taken from the JANAF tables (2). Data for C (graphite) were taken from the compilation of Stull and Sinke (17). Thermal functions for  $UC_{1.86}$  were calculated from the data of Levinson (12). The low-temperature  $C_{\nu}$ data of Jones, Gordon, and Long (8) were used for  $UO_2$ . The data of Westrum *et al.* (3, 18) were used for UC<sub>1.86</sub> with a small (and insignificant) correction for composition plus a configuration entropy term.

For the second law treatment of Reaction 3 an average  $\Delta C_p$  of 1.0 cal. ° K.<sup>-1</sup> was computed over the temperature range 1700° to 2000° K. The resulting equation calculated was

$$\Delta F^{\circ} = 182,300 - 1.0 \ T \ln T - 77.1 \ T \text{ cal.}$$
<sup>(4)</sup>

an equation which is cast in the form

$$\Delta F^{\circ} = \Delta H^{\circ}_{1800} - 1800(\overline{\Delta C_p}) - \overline{\Delta C_p} T \ln T + IT$$
(5)

where I is the computed intercept (9). We suggest that Piazza and Sinnott made a numerical error in computing the least squares slope and intercept; our calculated line fits their data better. The internal consistency of the data for Reaction 3 is excellent, much better than might be expected considering the narrow temperature range and small number of measurements, as demonstrated by the following:

Table I. Calculated Thermodynamic Functions

Temp	Free Energy Functions, $(F^{\circ}_{T} - H^{\circ}_{299})/T$ , Cal. ° K. <sup>-1</sup> GMole <sup>-1</sup>			Heat Capacities, $C_{p}^{\circ}$ , Cal. ° K. $^{-1}$ GMole $^{-1}$			Entropies, $S$ ?, Cal. ° K. <sup>-1</sup> GMole <sup>-1</sup>			Heat Contents, $H_{\mathcal{P}} - H_{\mathfrak{M}}^{\mathfrak{s}}$ , Kcal. GMole <sup>-1</sup>		
° K.		UC <sub>1.86</sub> <sup>a</sup>	UC	$UO_2$	UC <sub>1.86</sub>	UC	$UO_2$	UC 1.86 <sup>a</sup>	UC	UO <sub>2</sub>	UC <sub>1.86</sub>	UC
1700	-35.62	32.98	-27.20	21.95	29.1	15.42	52.14	49.52	39.40	28.090	28.127	20.739
1800	-36.57	33.94	-27.90	22.15	29.1	15.45	53.40	51.19	40.28	30.295	31.038	22.282
1900	-37.49	34.89	-28.58	22.35	29.1	15.49	54.60	52.76	41.12	32.519	33.948	23.829
2000	-38.37	35.82	-29.23	22.54	29.1	19.57	55.75	54.26	42.06	34.764	36.858	25.674
° Contribu	ition to en	tropy of 0.6	01 e.u. adde	d to take i	nto accour	nt random	orientatio	on of carbor	n vacancies	s.		

(6)

 $\Delta H_{288}^{\circ}$  (2nd) = 187.9 ± 8.0 kcal.

 $\Delta H_{298}^{\circ}$  (3rd) = 185.5 ± 0.24 kcal. (7)

 $\Delta S_{1800}^{\circ}$  (2nd) = 85.6 ± 4.3 cal. deg.<sup>-1</sup> (8)

 $\Delta S_{1800}^{\circ} (3rd) = 84.3 \text{ cal. deg.}^{-1}$  (9)

The errors quoted are the standard deviations for Equations 6 and 8 and the standard deviation of the mean for Equation 7. The (2nd) or (3rd) notation refers to the second law or third law method of calculation.

A C/U ratio of 1.86 in Equation 3 appears to be more accurate (10) than the value of 2 used by Piazza and Sinnott. No difference results in the calculated second law heat or entropy. A difference of 1.3 kcal. in the third law heat and of 0.7 cal. deg.<sup>-1</sup> in the calculated third law entropy results. The assumption of an activity of 0.95 for "UC<sub>2</sub>" contributes of the order of 0.2 kcal. to the results quoted in Equation 7. When the auxiliary data are available, the third law treatment is likely to give much more accurate values than those calculated by the second law.

Using the third law data cited above, we compute the heat of formation of UC<sub>1.86</sub>(s),  $\Delta H_{7.298}^{2}$ , to be -20.5 kcal. per mole. Huber, Head, and Holley (5) obtained -21.1  $\pm$  1.4 kcal. per mole by combustion calorimetry.

Agreement within 2.4 kcal. between the second and third law calculations is excellent but is also probably fortuitous. Agreement between the calculated heat of formation of UC<sub>1.86</sub> with the measurements of Huber, Head, and Holly (5) within 0.6 kcal. is also probably somewhat fortuitous. Nevertheless, the agreement is good and comparison between the two denies stabilization of "UC<sub>2</sub>" by oxygen, at least to the extent assumed (7,15).

Combining the standard deviation reported in Equation 7 with the uncertainty in the heat of formation of UO<sub>2</sub> yields 0.8 kcal. uncertainty in the heat of formation of UC<sub>1.86</sub>. Combination of this result with the uncertainty of Huber, Head, and Holley (5) leads to a "best" value of  $\pm 0.7$  kcal. The "best" value of the heat of formation,  $\Delta H_{296}$ , of "UC<sub>2</sub>" is thus -20.8  $\pm$  0.7 kcal. per mole.

### THE UO2-"UC2"-UC-CO SYSTEM

Piazza and Sinnott also studied the experimentally more difficult system  $UO_2$ -" $UC_2$ "-UC in equilibrium with CO from 1761° to 1953° K. by measuring the CO partial pressures. The reaction they studied is assumed to be

$$O_2 + 3.488 \text{ UC}_{1.86} = 4.488 \text{ UC} + 2\text{CO}$$
 (10)

The thermal data for UC were calculated from the data of Levinson (11) and of Harrington and Rowe (4) combined with the necessary low temperature data from Westrum *et al.* (18). An average  $\Delta C_p$  of -37.1 cal. °K.<sup>-1</sup> was calculated over the range 1700° to 2000°K., but for theis calculation we adjusted  $C_p$  of UC to 15.53 cal. deg.<sup>-1</sup> at 2000°K. instead of using the value reported in Table I.

Treating their data in the fashion described above, with  $a_{\rm UC}$  and  $a_{\rm UC_{1sc}}$  assumed to be 0.95 and  $a_{\rm UO_2}$  to be unity, we obtain the equation

$$\Delta F^{\circ} = 216,100 + 37.1 \ T \ln T - 380.1 \ T \text{ cal.}$$
(11)

The internal consistency of the data is shown by the following:

 $\Delta H_{298}^{\circ}$  (2nd) = 164.1 ± 7.5 kcal. (12)

 $\Delta H_{28}$  (3rd) = 174.8 ± 0.3 kcal. (13)

 $\Delta S_{1\infty0}^{\circ} (2\text{nd}) = 64.9 \pm 4.0 \text{ cal. } \text{deg}^{-1}$ (14)

$$\Delta S_{1500}^{\circ} (3rd) = 70.6 \text{ cal. deg.}^{-1}$$
(15)

The internal consistency is not particularly good; however, this is not unexpected.

Combination of Equations 3 and 10 in an appropriate fashion leads to the reaction

UC (s) + 0.86 C (graphite) = UC<sub>1.86</sub> (16)

 $A \Delta H_{296}^{\circ}$  of +2.37 kcal., using the third law results, is calculated for Reaction 16. An alternative way of expressing the information is to state that the heat of formation of UC,  $\Delta H_{7,296}^{\circ}$ , is computed to be -23.2 kcal. per mole. Huber, Head, and Holley (6) have recently remeasured this quantity indirectly by combustion calorimetry in an oxygen bomb and obtained a  $\Delta H_{7,296}^{\circ}$  of -23.3  $\pm$  0.9 kcal. for UC<sub>0.996  $\pm$  0.005 and -23.0  $\pm$  1.0 kcal. for UC<sub>1.033  $\pm$  0.005.</sub></sub>

The agreement of the measurements of Huber, Head, and Holley (6) with Piazza and Sinnott (14) appears to be exceedingly good. The discrepancy between the second law and third law calculations probably reflects the difficulty of obtaining equilibrium in the system studied by Piazza and Sinnott. It may also reflect a shift in activity of the UC phase with temperature. A "best" value for the heat of formation,  $\Delta H_{1,298}^{2}$ , of UC can probably be fixed by Huber, Head, and Holley when they have completed their calculations.

Thermodynamic functions calculated from the data cited are tabulated in Table I.

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# Surface Tension–Viscosity Relation for Liquids

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An empirical equation is proposed,  $\gamma = A_{exptl.}B/\eta$ , relating viscosity  $(\eta)$  to surface tension  $(\gamma)$  where A and B are dimensional constants. Experimental literature data for 33 liquid systems (organic and inorganic) are used to demonstrate the validity of the equation. Water, benzene, and toluene liquid systems are used to show that the constant B is a function of (Mk/R) where M is the molecular weight, k is the thermal conductivity, and R is the gas constant. The data suggest that this equation may be applicable over the entire liquid phase regime for any given substance.

(1)

(16)

(17)

(18)

Pressures,"

New York, 1965.

contract with the Union Carbide Corp.

SEVERAL relations have been postulated relating surface tension of a liquid in air to viscosity. In 1932, Silverman and Roseveare (9) postulated that

 $\gamma - \frac{1}{4} = \frac{A}{n} + B$ 

where K and n are empirical constants. In 1956, Sanyal and Mitra (8) postulated the following relations:

Rand, M.H., Jubaschewski, O., "Thermochemical Properties of Uranium Compounds," Interscience, New York, 1963.

Stull, D.R., Sinke, G.C., Advan. Chem. Ser., No. 18 (1956).

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Associated liquids: 
$$T\log \eta - a\gamma^{5/6} = b$$
 (3)

(4)

Nonassociated liquids: 
$$T \log \eta - a\gamma = b$$

where A and B are empirical constants for a given substance. In 1953, Murkerjee (5) postulated that viscosity is related to surface tension by

$$\gamma = \mathbf{K} \eta^{n1/3} \tag{2}$$

where a and b are empirical constants. In 1962, Pratap and Narsimhan (7), suggested that surface tensions and viscosity are related by the following equation:

$$\frac{u_L \cdot \eta \cdot V^{2/3}}{(M \cdot T)^{1/2}} = 9.27_{\text{exptl.}} \frac{1.091 \cdot N^{1/3} \cdot \gamma \cdot V_m^{2/3}}{R \cdot T}$$
(5)



Figure 1. Least squares fit of experimental data for some members of the paraffin series