

ACKNOWLEDGMENT

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NOMENCLATURE AND UNITS

B	= second virial coefficient, cu. ft. per mole
C	= third virial coefficient, (cu. ft.) ² per (mole) ²
D	= fourth virial coefficient, (cu. ft.) ³ per (mole) ³
c_p	= specific heat at constant pressure, B.t.u. per lb.-° F.
f	= free energy, B.t.u. per lb.
h	= enthalpy per unit mass, B.t.u. per lb.
Δh_v	= enthalpy change upon vaporization of a unit mass of equilibrium vapor, B.t.u. per lb.
ΔH_v	= enthalpy change upon vaporization of a mole of equilibrium vapor, B.t.u. per lb.-mole
J	= any unit conversion
M_1	= molecular weight of monomeric species
p	= absolute pressure, lb. per sq. inch
R	= gas constant
s	= entropy per unit mass, B.t.u. per lb.-° F.
Δs_v	= entropy change upon vaporization of a unit mass of equilibrium vapor, B.t.u. per lb.-° F.
T	= absolute temperature, ° R.
t	= temperature, ° F.
\bar{V}	= molal volume (per formula weight of monomer), cu. ft. per lb.-mole
v	= specific volume, cu. ft. per lb.
z	= compressibility factor, $p\bar{V}/RT$

Subscripts

0	= quantity at 0° R.
p	= constant pressure change
s	= quantity at saturation

Superscripts

g	= quantity in gas state
l	= quantity in liquid state
°	= standard state, 1 atm. for gas

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High-Temperature Properties of Cesium

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A virial equation of state for cesium with coefficients through the fifth virial is derived from compressibility data covering a temperature range from 1305° to 2550° F. and a pressure range from 16.5 to 487.9 p.s.i.a. Exact thermodynamic relationships involving the virial and vapor-pressure equations are used to generate saturation and superheat properties of the vapor. Values of the enthalpy, entropy, specific heat, and specific volume are tabulated for some 280 selected states in the temperature range from 1250° to 2550° F. and in the pressure range from 2.9 to 499.7 p.s.i.a. The thermal properties from 1050° to 2550° F. are also presented as a Mollier diagram.

DESIGN STUDIES of high-temperature power-converters employing alkali metal vapors as working fluids require an improved knowledge of their thermodynamic properties. In response to this need, experimental *PVT*, vapor-pressure, and liquid density data for three metals—sodium, potassium, and cesium—have been measured by Stone *et al.* (9, 10, 11). From the data for two of the metals, sodium and potassium, the authors have generated and published consistent sets of thermodynamic properties (5, 6). The present article extends this work to the third metal, cesium.

TREATMENT OF *PVT* DATA

Derivation of Virial Coefficients. Coefficients for the cesium equation of state in its volume expansion form (Equation 1) were derived from the experimental *PVT* data of Stone *et al.* (9), which cover a temperature range in the superheat region from 1305° to 2550° F. They were obtained graphically as described for potassium (5) by plotting experimental functions along constant temperature lines. However, two modifications of the method were required for cesium. First, the adjustment procedure (5), which was

helpful in defining virial coefficients at lower temperatures for sodium and potassium, was not required since data for cesium were more numerous and more precise than those reported for the other metals. In the second place, an additional virial coefficient—the fifth—was required for cesium in order to obtain a precise fit to all the data along any given isotherm. Only a rough value of the higher virial was required, and this was also obtained by a graphical procedure. The quantity $[(z-1)\bar{V}^3 - B\bar{V}^2 - C\bar{V}]$ was plotted vs. $1/\bar{V}$ at several higher temperatures using preliminary values for the lower virial coefficients; the average slope (600,000) at higher pressures was taken as the value of E .

The Virial Equation of State for Cesium. The equation of state, obtained as described above, is

$$\frac{p\bar{V}}{RT} = 1 + \frac{B}{\bar{V}} + \frac{C}{\bar{V}^2} + \frac{D}{\bar{V}^3} + \frac{E}{\bar{V}^4} \quad (1)$$

where,

$$\begin{aligned} \log|B| &= -3.6200 + 4000.0/T + \log T \\ B &< 0 \\ \log C &= 3.3551 - 5331.5/T + 10.825 \times 10^6/T^2 \\ C &> 0 \\ \log|D| &= 4.1856 + 880/T \\ D &< 0 \\ E &= 600,000 \end{aligned}$$

The degree to which the virial equation fits the measured data is shown graphically in Figure 1, in which compressibilities generated with Equation 1 are compared to experimental values at 100° intervals from 1350° to 2550° F. The

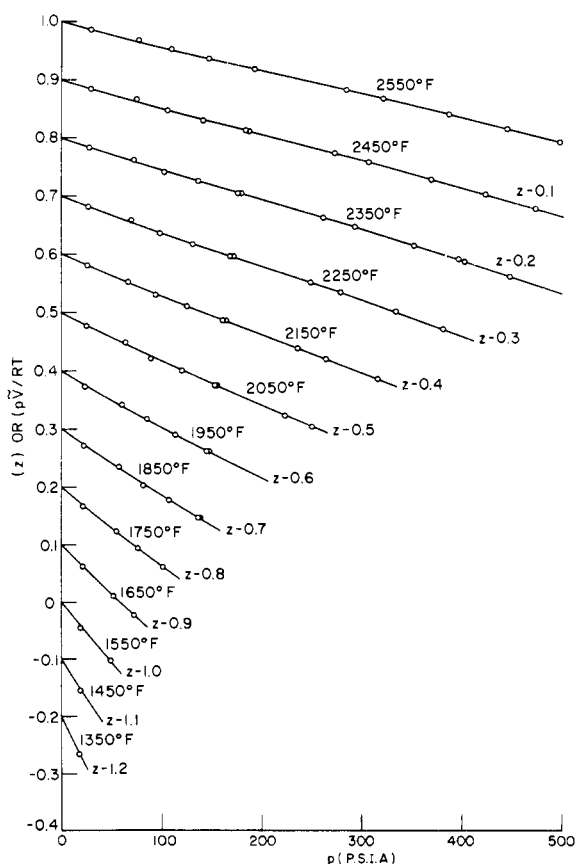


Figure 1. Compressibility of cesium vapor at several temperatures

degree of fit also can be shown mathematically. For example, all the observed compressibility data for cesium reported by Stone *et al.* (9) may be calculated from the virial equation with a standard per cent deviation of only ± 0.22 .

The equation of state was derived from experimental data over a temperature range from 1550° to 2500° F. Data were also obtained in a lower range from 1300° to 1550° F., but the number of experimental points along any given isotherm was insufficient to obtain reliable virial coefficients by the graphical method. Consequently, before the equation of state was acceptable for calculations below 1550° F., it was necessary to determine its fit to the observed low-temperature data. Compressibility factors were calculated and compared to the observed values in this region (Figure 1). The agreement was equivalent to that obtained at higher temperatures.

CALCULATION OF THERMODYNAMIC PROPERTIES AND FUNDAMENTAL DATA

For the design of high-temperature cycles, the important properties of the liquid and vapor are volume, enthalpy, entropy, and specific heat. These properties for cesium were computed from the virial equation of state, from exact thermodynamic relationships, or from the fundamental property equations included in this section.

Saturation Pressure of Liquid Cesium. Pressures at all temperatures were obtained from the vapor-pressure equation of Stone *et al.* (10). This equation,

$$\log p = 7.04023 - 7040.7/T - 0.53290 \log T \quad (2)$$

is based on experimentally determined data over the temperature range from the normal boiling point to 2350° F., and is in good agreement with the lower temperature work of Tepper and coworkers (12), Achener (1), and Bonilla, Sawhney, and Makansi (3).

Specific Volume of Liquid Cesium. Values were obtained from the correlating equation of Stone *et al.* (11). This equation,

$$\frac{1}{v^l} = 124.181 - 1.5970 \times 10^{-3}T - 1.6855 \times 10^{-6}T^2 \quad (3)$$

was derived from all available data and covers a temperature range from the melting point to 2304° F. The error limit given for the full range was $\pm 0.6\%$.

Specific Volume of Saturated and Superheated Cesium Vapor. The value of this property at each state was obtained from the virial equation of state.

Enthalpy, Entropy, and Specific Heat of Saturated and Superheated Cesium Vapor. These properties were computed along constant temperature lines from exact thermodynamic relationships with the value of the property for the monomeric gas as a base. The working relationships are the same as those used for potassium (5) except that terms for the fifth virial must be added.

Enthalpy, Entropy, and Specific Heat of Monomeric Cesium Vapor. The monomeric gas properties, together with the

Table I. Summary of Heat Capacity Measurements of Liquid Cesium

Investigator and Reference	Temperature Range, ° F.	Heat Capacity Equation, B.t.u./Lb.-° F.
Achener (1)	152-1656	$c_p^l = 0.08543 - 9.605 \times 10^{-5}t + 5.985 \times 10^{-8}t^2$
Tepper (12)	620-1770	$c_p^l = 0.0545$
Lemmon (8)	570-2100	$c_p^l = 0.0600$

value selected for the enthalpy of sublimation at absolute zero, largely determine the magnitudes of the superheat properties. The equations for the three properties (relative to the solid crystal at 0° R.) are

$$(h^s)^\circ = 252.18 + 0.037361 T + 2480 e^{-31,290/T} \quad (4)$$

$$(s^s)^\circ = 0.080604 + 0.037361 \ln T + 0.371 e^{-28,598/T} \quad (5)$$

$$(c_p^s)^\circ = 0.037361 + 1.3099 e^{-25,663/T} \quad (6)$$

These were derived directly from the work of Evans and coworkers (4) and are based on their standard properties over the temperature range from 0° to 2800° F. and on the enthalpy of sublimation at 0° R. (18.62 mean kcal. per gram atom) as derived in the next section.

Enthalpy of Sublimation of Cesium at 0° R. This quantity is needed to relate the enthalpy of the ideal gas to the base state—solid at 0° R. Since the virial coefficients of cesium are known, a reliable value of the sublimation constant can be obtained from saturation pressure data with the third-law equation,

$$(\Delta H_\delta)_s = -RT \left[\frac{2B}{V_s} + \frac{3C}{2V_s^2} + \frac{4D}{3V_s^3} + \frac{5E}{4V_s^4} + \ln p_s - \ln \frac{p_s V_s}{RT} \right] - M_s T_s \left[\frac{f^\circ - h_\delta^\circ}{T} \right]_1 \quad (7)$$

Before this can be done, all the thermal functions must be known. The free-energy functions for monomeric cesium gas can be obtained from Evans *et al.* (4), and corresponding functions for the liquid may be derived from heat-capacity results. Three measurements of liquid heat capacity were found in the literature and pertinent data related to each are summarized in Table I. There is significant disagreement both in the magnitude and the temperature dependency of the specific heat.

Since there was no apparent reason to choose one set of specific-heat data over another, the selected value of both c_p^l and $(\Delta H_\delta)_s$ were based on a third-law analysis of the vapor-pressure measurements of Stone and coworkers (10). Normally, a third-law analysis is used to check the internal consistency of vapor-pressure measurements. In this case, however, the analysis was used to obtain the most consistent values for the specific heat of the liquid and the sublimation constant.

In two preliminary analyses, the value of $(\Delta H_\delta)_s$ was computed from the vapor-pressure data over the temperature range from 1215° to 2346° F. using, in one case a constant liquid specific heat of 0.0545 and, in the other, 0.0600 B.t.u. per lb.-° F. These are the heat capacities reported by Tepper *et al.* (12) and Lemmon *et al.* (8), respectively. Thermal data for the solid at 77° F. were taken from Hultgren *et al.* (7) and at higher temperatures from the work of Lemmon

et al. (8). When the sublimation quantities, so calculated, were plotted against temperature, those based on a liquid specific heat of 0.0600 B.t.u. per lb.-° F. exhibited a constant negative slope while those computed with the other value, a positive slope. It followed that a value of specific heat intermediate between the two, would lead to a generally constant value of $(\Delta H_\delta)_s$ for all the saturation pressures. By a trial and error procedure this value for the liquid specific heat was found to be 0.05683 B.t.u. per lb.-° F. The value of $(\Delta H_\delta)_s$ obtained by this procedure is 18.62 mean kcal. per gram atom. Although the agreement is perhaps fortuitous, this is practically the same as the value of 18.66 published by Hultgren and coworkers (7), which was obtained by a third-law treatment of several vapor-pressure measurements at lower temperatures.

Enthalpy and Entropy of Vaporization of Cesium. Latent heats of vaporization were calculated with

$$\Delta h_v = J p_s (1103.1/T - 0.036262) (v_s^g - v_s^l) \quad (8)$$

which was derived by a differentiation of Equation 2 and subsequent substitution into the Clapeyron equation. A value of v_s^l at each temperature was obtained from Equation 3 and a value of v_s^g from the virial equation of state (Equation 1).

The heats of vaporization so obtained from the Clapeyron equation are presented graphically in Figure 2 and are compared with values reported by Achener (1). The four results by Achener were measured directly by noting the heat required to vaporize a given mass of the liquid. The

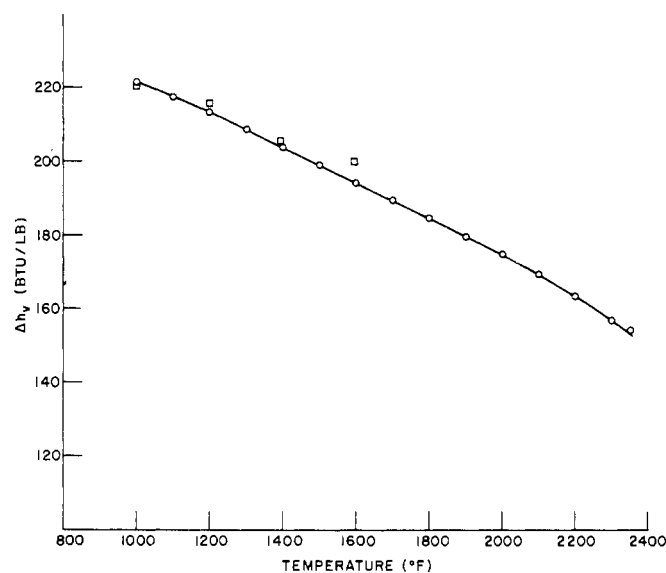


Figure 2. Enthalpy of vaporization of cesium

□ Achener
○ NRL

Table II. Saturation Properties of Cesium^a

Basis: $h = 0$ and $s = 0$ for solid cesium at 0° R.

t	p_s	v^l	v_s^g	h^l	Δh_v	h_s^g	s^l	Δs_v	s_s^g
1300.00	20.40	0.01101	6.3778	97.72	208.70	306.41	0.2310	0.1186	0.3496
1400.00	32.51	0.01128	4.1522	103.81	203.95	307.76	0.2344	0.1097	0.3441
1500.00	49.33	0.01157	2.8289	109.91	199.10	309.01	0.2376	0.1016	0.3392
1600.00	71.78	0.01189	2.0041	115.98	194.24	310.22	0.2406	0.0943	0.3349
1700.00	100.76	0.01222	1.4687	122.01	189.44	311.45	0.2434	0.0877	0.3311
1800.00	137.12	0.01258	1.1083	128.02	184.71	312.73	0.2461	0.0817	0.3278
1900.00	181.60	0.01297	0.8576	134.11	179.92	314.04	0.2487	0.0762	0.3249
2000.00	234.86	0.01339	0.6774	140.43	174.88	315.31	0.2512	0.0711	0.3223
2100.00	297.45	0.01384	0.5440	147.10	169.32	316.43	0.2538	0.0661	0.3200
2200.00	369.80	0.01433	0.4427	154.12	163.11	317.23	0.2564	0.0613	0.3178
2300.00	452.21	0.01487	0.3655	160.91	156.77	317.68	0.2589	0.0568	0.3157

^a Values are given to more significant figures than the data warrant in order to retain precision in differences.

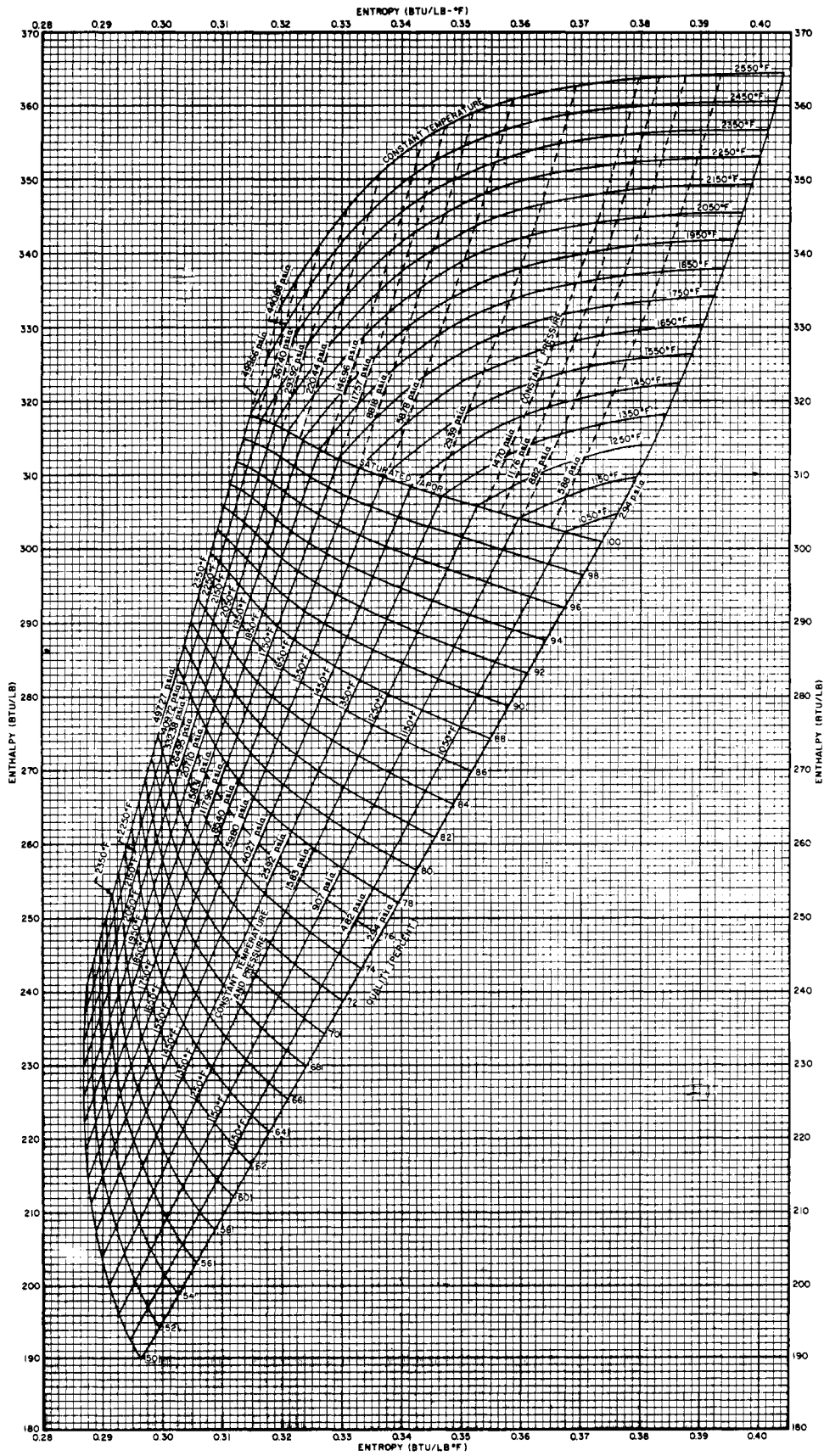


Figure 3. Mollier diagram for cesium

agreement is good, as three of his points fall within 1.0% of the corresponding NRL values and the fourth within 3.0%.

The entropy of vaporization of each saturation point was obtained by dividing the appropriate enthalpy change by the absolute temperature.

Enthalpy and Entropy of Liquid Cesium. These properties of the saturated liquid at each temperature were obtained by subtracting the enthalpy or entropy of vaporization from the corresponding property of the saturated vapor.

THERMODYNAMIC PROPERTIES OF CESIUM AND DISCUSSION

The saturation and superheat properties (1250° to 2550° F.) are presented in Tables II and III, and the thermal properties (1050° to 2550° F.) for both wet and dry vapors are presented graphically in a Mollier chart (Figure 3). The properties were computed from exact thermodynamic relationships involving the virial and vapor-pressure equations and are based on the properties of the ideal monomeric gas. The important virial equation was reduced from experimental *PVT* data covering a pressure range from 16.5 to 487.9 p.s.i.a. and a temperature range from 1305° to 2550° F. The vapor-pressure equation represents saturation data covering a range from 14.7 to 485.0 p.s.i.a. The experimental data, therefore, covers all states for which properties are reported except those with pressures below 16.5 p.s.i.a. Properties for these low-pressure states required short extrapolations of the virial and vapor-pressure equations and are of reduced accuracy.

Compressibility data and derived thermodynamic properties for cesium vapor at temperatures to 2060° F. were recently published by Tepper *et al.* (13). These properties may be compared directly to those generated in this study. If the NRL data are arbitrarily taken as reference and the property changes from p_s to 7.3 p.s.i.a. are compared at temperatures covering the overlapping range, the enthalpy changes reported by Tepper are 8% low, 32% high, and 350% high at temperatures of 1340°, 1700°, and 2060° F., respectively, and his entropy changes at the same temperatures are 8% low, 6% high, and 62% high. Although the observed discrepancies at lower temperatures in the two sets of properties can be explained by the random errors in Tepper's data (13), those at higher temperatures cannot be accounted for in this way. However, his values at higher temperatures appear to be completely out of line with existing spectroscopic information.

There are a number of publications in which thermodynamic properties of cesium vapor have been calculated by estimating gas imperfections from spectroscopic data. The properties generated in this article may be compared to those computed in recent publications by Agapova, Paskar, and Fokin (2) and Weatherford, Tyler, and Ku (14). If the NRL data are again taken as the reference and property changes from p_s to 2.9 p.s.i.a. are compared, the enthalpy changes reported by Agapova are 23 to 44% low and the entropy changes are 9% low. By a similar comparison, the enthalpy changes reported by Weatherford are 8 to 41% low and the entropy changes 4 to 8% low.

With the appearance of this article, thermodynamic properties based on experimental *PVT* data have now been published for three alkali metals—sodium, potassium, and cesium. The reported properties for each metal have been evaluated successfully for internal consistency and it is believed that they represent the best values available. They should be suitable for many calculations currently required for the design of systems using the metals as heat-transfer or working fluids. Even so, it is recognized that the reported properties can be evaluated better as additional properties are measured directly for the three metals. Particularly important would be reliable determina-

tions of the heat of vaporization and the specific heat of the liquid and vapor. A good example of this type of evaluation is provided by the recent measurements at Aerojet-General Nucleonics (1) of the heats of vaporization of cesium. These were measured directly and have been shown to substantiate those computed from the Calpeyron equation.

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Superscripts

g	= quantity in gas state
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$^\circ$	= standard state, 1 atm. for gas

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Salt Effects in Liquid-Liquid Equilibria

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The effect of lithium chloride, sodium chloride, and potassium chloride on the distribution of acetic acid, propionic acid, and butyric acid between benzene and water at 25° C. has been investigated. The method of Hand, originally developed for ternary systems, was used to correlate the distribution data for quaternary systems containing salt. A regression in two variables was used to express the slope and the intercept of the Eisen-Joffe Equation as linear functions of salt concentration. The order of decreasing salting out, on a weight basis, is



for all systems studied. On a mole basis, the effect of lithium chloride diminishes with increase in the number of carbon atoms in the carboxylic acid.

INVESTIGATIONS of salt effects in liquids have dealt primarily with the solubilities of nonelectrolytes, such as benzene, in water. Excellent summaries of this work appear in the literature (9, 11, 14). The effect of salts on the distribution of a solute between two immiscible solvents, a problem having engineering applications, has received very little attention.

There is general agreement on the physical phenomena responsible for the salt effects (1). The mathematical interpretation of these phenomena has resulted in several relationships which are, at best, semiquantitative. The limited applicability of these expressions emphasizes the complexity of the salt effects.

THEORIES OF THE SALT EFFECT

Three principal theories have been advanced to explain the salt effect. The hydration theory assumes that each salt ion binds a constant number of water molecules as a shell of oriented water dipoles surrounding the ion. This "bound" water is then unavailable as solvent for the nonelectrolyte. Considering the wide variation in hydration numbers reported in the literature, it appears that this concept permits only a qualitative estimate of the magnitude of the salt effect.

The electrostatic theory, which considers only the action of coulombic forces and omits other factors, was used by

Debye and McAulay (7) to compute the partial free energy of a molecule relative to the pure solvent from the reversible electrical work. Because of simplifications and approximations made in its derivation, the Debye-McAulay equation is a limiting equation and can be regarded as a first approximation only. Butler (4), using a similarly simplified model, obtained an equation virtually identical with that of Debye and McAulay. Later Debye (6), taking into account the heterogeneity of the mixture of water and neutral solute, expressed the total free energy of the system, including the contribution due to the field of the ion, as a function of distance from the ion.

Since the electrostatic interaction between an ion and a neutral molecule is short range in nature, additional interaction terms of the London or van der Waals type must be considered more fully. These terms involve the polarizability of salt ions, solvent molecules, and nonelectrolyte solute molecules, as well as the special force fields originating from any component dipoles that may be present.

McDevit and Long (12), applying the "internal pressure" concept of Tammann (17) to nonpolar nonelectrolytes, calculated the free energy of transfer of the latter from pure water to the salt solution.

PREVIOUS WORK

The Setschenow equation (15)

$$\ln \frac{C^0}{C} = k_s X, \quad (1)$$

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