# Vapor-Liquid Equilibria of Water-Diisopropylamine and Water-Di-n-Propylamine 

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#### Abstract

The vapor-liquid equilibria of diisopropylamine-water and di-n-propylamine-water solutions are reported. Data are given at $10^{\circ}, 15^{\circ}, 20^{\circ}, 25^{\circ}, 30^{\circ}$, and $39.95^{\circ} \mathrm{C}$. for diisopropylamine and at $10^{\circ}, 20^{\circ}, 30^{\circ}$, and $39.95^{\circ} \mathrm{C}$. for di-n-propylamine.


Water Vapor PRESSURES were taken from Lange's handbook (5). Between $10^{\circ}$ and $40^{\circ} \mathrm{C}$. the values given by Lange never differ by more than 0.01 mm . of Hg from values derived from the correlating equations used in calculating the 1967 ASME steam tables.

## EXPERIMENTAL

The vapor pressures of the pure amines and of aminewater mixtures were measured with a static apparatus described previously (3), which employs especially designed flasks that permit the liquid samples to be degassed without change in composition. These flasks are connected to mercury manometers, and the entire assembly is suspended in the constant temperature bath and may be agitated to hasten equilibrium.

Total pressure measurements were made using a cathetometer to read the manometers in the bath. If the pressure exceeded the range of the manometer, a back-pressure was provided by water in another flask immersed in the same bath. The manometer readings were corrected for temperature. The accuracy of the pressure measurements was about 0.1 mm . of Hg . Temperatures were measured with a calibrated glass thermometer and have an accuracy within $0.03^{\circ} \mathrm{C}$. A thermistor-operated Sargent temperature controller and an external circulating pump maintained the temperature within about $0.01^{\circ} \mathrm{C}$.
The liquid phase was analyzed by acid titration when the water concentration was greater than $10 \%$ by weight. For lower concentrations, a gas chromatograph was used with a 20 -foot by $1 / 4$-inch column packed with Fluoropac coated with $10 \%$ GE-SF- 96 Silicone (Aerograph). At $150^{\circ} \mathrm{F}$. and a helium flow of 30 ml . per minute the retention times were $21 / 2$ to 3 minutes for water and about 6 minutes for the amines. These methods gave amine concentration reproducibilities to 1 part in 250 to 500 , and water concentrations having a maximum deviation of 1 mole $\%$.

The liquid-phase compositions of the partially miscible solutions (Table I) were obtained by two methods. The diisopropylamine data and the water-rich phase data for di-n-propylamine were obtained by acid titration of the

| Table I. Liquid-Liquid Phase Equilibrium for Diisopropylamine and Di-n-propylamine |  |  |
| :---: | :---: | :---: |
|  | Mole Fraction Amine |  |
| Temp., ${ }^{\circ} \mathrm{C}$. | Water-rich phase <br> (Diiso | Amine-rich phase amine) |
| 30 | 0.0211 | 0.2155 |
| 40 | 0.0126 | 0.408 |
|  | (Di-n-p | amine) |
| 10 | 0.01389 | 0.269 |
| 20 | 0.00902 | 0.336 |
| 30 | 0.00677 | 0.425 |
| 40 | 0.00498 | 0.498 |

equilibrated phases. The amine-rich phase for di-npropylamine was determined from cloud points. Various solutions were made by weight and sealed in glass tubes. These tubes were immersed completely in a water bath, the temperature was slowly increased, and the cloud points were observed. These data were plotted and the values in Table I were interpolated from the curve.

## CALCULATION OF VAPOR COMPOSITION

The vapor composition cannot be measured directly with the apparatus used. However, since the vapor analysis is

|  | Table II. Experimental Vapor Pressures of Diisopropylamine-Water Mixtures |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Mole \% Amine | Temperature, ${ }^{\circ} \mathrm{C}$. |  |  |  |  |  |
|  | 10 | 15 | 20 | 25 | 30 | 39.95 |
|  | Pressure, Mm. Hg |  |  |  |  |  |
| 0.379 | 11.95 | 17.7 | 25.85 | 37.45 | 53.5 | 103.65 |
| 0.632 | 14.0 | 21.25 | 30.95 | 46.5 | 67.2 | 131.5 |
| 1.40 | 19.85 | 30.4 | 46.0 | 67.55 | 96.7 |  |
| 2.49 | 25.4 | 37.2 | 53.55 | 75.25 | $103.2{ }^{\text {a }}$ |  |
| 4.30 | 27.2 | 38.85 | 54.95 | 76.0 | $103.1{ }^{\text {a }}$ |  |
| 9.99 | 28.55 | 40.15 | 55.9 | 76.4 | $103.1{ }^{\text {a }}$ |  |
| 24.3 | 31.15 | 42.8 | 58.4 | 78.05 | 103.65 | $175.9^{\text {a }}$ |
| 36.5 | 33.4 | 45.2 | 60.9 | 80.55 | 105.8 | $175.9{ }^{\text {a }}$ |
| 57.4 | 36.3 | 48.5 | 64.35 | 84.45 | 109.85 | 178.85 |
| 70.3 | 37.8 | 50.15 | 66.15 | 86.3 | 111.55 | 180.55 |
| 79.8 | 38.5 | 50.75 | 66.8 | 86.8 | 112.2 | 180.8 |
| 85.2 | 38.7 | 51.0 | 66.85 | 86.7 | 111.9 | 179.45 |
| 87.2 | 38.8 | 51.15 | 67.15 | 86.75 | 111.45 | 178.2 |
| 94.6 | 38.65 | 50.65 | 65.95 | 84.75 | 108.0 | 169.75 |
| 100.0 | 37.9 | 49.25 | 63.5 | 80.75 | 102.35 | 157.75 |
| ${ }^{\circ} 2$ liquid phases. |  |  |  |  |  |  |

Table III. Experimental Vapor Pressures of Di-n-Propylamine-Water Mixtures

| Mole \% Amine | Temperature, ${ }^{\circ} \mathrm{C}$. |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | 10 | 20 | 30 | 39.95 |
|  | Pressure, Mm. Hg |  |  |  |
| 0.308 | 10.45 | 21.8 | 43.45 | 81.8 |
| 0.489 | 11.2 | 24.2 | 49.75 | 95.2 |
| ${ }^{\circ}$ | $15.15{ }^{\text {c }}$ | $29.9{ }^{\text {a }}$ | $55.45^{\text {a }}$ | $96.95{ }^{\text {a }}$ |
| 40.2 | 15.4 | 29.95 | $55.5{ }^{\text {a }}$ | $97.0{ }^{\text {a }}$ |
| 47.6 | 15.35 | 29.95 | 53.35 | $97.0{ }^{\text {a }}$ |
| 63.3 | 15.25 | 29.6 | 54.55 | 95.55 |
| 78.3 | 14.65 | 28.2 | 51.7 | 90.5 |
| 85.2 | 13.7 | 26.4 | 48.75 | 85.05 |
| 85.85 | 13.95 | 26.5 | 48.5 | 84.45 |
| 93.1 | 13.1 | 24.65 | 44.55 | 76.15 |
| 94.6 | 12.45 | 23.6 | 42.25 | 72.3 |
| 100.0 | 11.0 | 20.1 | 34.75 | 57.6 |

Table IV. Calculated Vapor-Liquid Equilibria of Diisopropylamine-Water

| $x$ | v | $\gamma_{1}$ | $\gamma$, | $G^{*}$ | $P$ | $x$ | $y$ | $\gamma$ : | $\gamma_{2}$ | $G^{E}$ | $P$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | At $10^{\circ} \mathrm{C}$. |  |  |  |  | At $25^{\circ} \mathrm{C}$. |  |  |  |  |  |
| 0 |  |  |  |  | 9.21 | 0 |  |  |  |  | 23.76 |
| 0.05 | 0.672 | 9.709 | 1.029 | 79.1 | 27.40 | 0.05 | 0.696 | 13.094 | 1.026 | 90.7 | 76.06 |
| 0.10 | 0.686 | 5.169 | 1.081 | 131.6 | 28.55 | 0.10 | 0.697 | 6.589 | 1.083 | 154.4 | 76.40 |
| 0.15 | 0.698 | 3.612 | 1.137 | 169.5 | 29.43 | 0.15 | 0.699 | 4.428 | 1.146 | 200.5 | 76.76 |
| 0.20 | 0.710 | 2.842 | 1.195 | 197.8 | 30.35 | 0.20 | 0.701 | 3.359 | 1.215 | 235.6 | 77.33 |
| 0.25 | 0.722 | 2.386 | 1.257 | 218.9 | 31.29 | 0.25 | 0.706 | 2.734 | 1.290 | 261.9 | 78.17 |
| 0.30 | 0.735 | 2.084 | 1.323 | 234.2 | 32.23 | 0.30 | 0.712 | 2.326 | 1.372 | 280.9 | 79.16 |
| 0.35 | 0.748 | 1.869 | 1.395 | 244.7 | 33.14 | 0.35 | 0.719 | 2.041 | 1.462 | 293.9 | 80.24 |
| 0.40 | 0.760 | 1.703 | 1.474 | 250.8 | 33.97 | 0.40 | 0.726 | 1.828 | 1.562 | 301.3 | 81.32 |
| 0.45 | 0.771 | 1.569 | 1.567 | 252.9 | 34.69 | 0.45 | 0.734 | 1.662 | 1.677 | 303.7 | 82.30 |
| 0.50 | 0.782 | 1.460 | 1.672 | 251.0 | 35.37 | 0.50 | 0.741 | 1.528 | 1.811 | 301.4 | 83.20 |
| 0.55 | 0.794 | 1.370 | 1.794 | 245.3 | 36.00 | 0.55 | 0.750 | 1.418 | 1.968 | 294.4 | 84.03 |
| 0.60 | 0.806 | 1.298 | 1.930 | 236.0 | 36.63 | 0.60 | 0.759 | 1.330 | 2.150 | 282.7 | 84.87 |
| 0.65 | 0.819 | 1.238 | 2.089 | 223.1 | 37.24 | 0.65 | 0.770 | 1.256 | 2.365 | 266.3 | 85.63 |
| 0.70 | 0.832 | 1.185 | 2.291 | 206.6 | 37.76 | 0.70 | 0.782 | 1.193 | 2.641 | 245.6 | 86.26 |
| 0.75 | 0.846 | 1.136 | 2.560 | 186.0 | 38.19 | 0.75 | 0.792 | 1.132 | 3.039 | 219.8 | 86.63 |
| 0.80 | 0.860 | 1.092 | 2.937 | 160.8 | 38.52 | 0.80 | 0.802 | 1.078 | 3.611 | 187.8 | 86.80 |
| 0.85 | 0.875 | 1.052 | 3.514 | 130.0 | 38.73 | 0.85 | 0.840 | 1.061 | 3.910 | 150.7 | 86.71 |
| 0.90 | 0.894 | 1.017 | 4.466 | 92.6 | 38.80 | 0.90 | 0.868 | 1.029 | 4.769 | 107.8 | 86.15 |
| 0.95 | 0.942 | 1.010 | 4.870 | 50.0 | 38.62 | 0.95 | 0.914 | 1.007 | 6.106 | 57.2 | 84.48 |
| 1.00 |  |  |  |  | 37.90 | 1.00 |  |  |  |  | 80.75 |

At $15^{\circ} \mathrm{C}$.

| 0 |  |  |  |  | 12.79 |
| :--- | ---: | ---: | ---: | ---: | ---: |
| 0.05 | 0.681 | 10.801 | 1.025 | 81.4 | 39.02 |
| 0.10 | 0.691 | 5.627 | 1.078 | 137.7 | 40.15 |
| 0.15 | 0.699 | 3.880 | 1.137 | 178.9 | 41.03 |
| 0.20 | 0.708 | 3.016 | 1.199 | 209.7 | 41.98 |
| 0.25 | 0.718 | 2.505 | 1.266 | 232.6 | 42.98 |
| 0.30 | 0.729 | 2.166 | 1.337 | 249.2 | 43.98 |
| 0.35 | 0.738 | 1.926 | 1.416 | 260.8 | 44.97 |
| 0.40 | 0.748 | 1.741 | 1.506 | 267.5 | 45.84 |
| 0.45 | 0.758 | 1.597 | 1.607 | 269.9 | 46.68 |
| 0.50 | 0.768 | 1.479 | 1.722 | 267.5 | 47.44 |
| 0.55 | 0.778 | 1.383 | 1.856 | 261.4 | 48.16 |
| 0.60 | 0.789 | 1.305 | 2.010 | 251.2 | 48.86 |
| 0.65 | 0.802 | 1.240 | 2.194 | 237.3 | 49.51 |
| 0.70 | 0.814 | 1.182 | 2.427 | 219.2 | 50.07 |
| 0.75 | 0.826 | 1.129 | 2.744 | 196.6 | 50.50 |
| 0.80 | 0.839 | 1.082 | 3.195 | 169.1 | 50.81 |
| 0.85 | 0.855 | 1.042 | 3.844 | 135.6 | 51.00 |
| 0.90 | 0.895 | 1.030 | 4.188 | 97.0 | 51.02 |
| 0.95 | 0.934 | 1.010 | 5.203 | 52.5 | 50.58 |
| 1.00 |  |  |  |  | 49.25 |

At $20^{\circ} \mathrm{C}$.

| 0 |  |  |  |  | 17.54 |
| :--- | ---: | ---: | ---: | ---: | ---: |
| 0.05 | 0.690 | 11.947 | 1.025 | 86.0 | 55.07 |
| 0.10 | 0.695 | 6.117 | 1.080 | 145.9 | 55.90 |
| 0.15 | 0.700 | 4.161 | 1.141 | 189.7 | 56.66 |
| 0.20 | 0.706 | 3.199 | 1.206 | 222.7 | 57.55 |
| 0.25 | 0.713 | 2.629 | 1.276 | 247.7 | 58.53 |
| 0.30 | 0.721 | 2.255 | 1.353 | 265.3 | 59.57 |
| 0.35 | 0.729 | 1.987 | 1.440 | 277.8 | 60.57 |
| 0.40 | 0.738 | 1.787 | 1.535 | 285.0 | 61.55 |
| 0.45 | 0.746 | 1.629 | 1.644 | 286.9 | 62.42 |
| 0.50 | 0.754 | 1.501 | 1.772 | 284.9 | 63.21 |
| 0.55 | 0.763 | 1.398 | 1.920 | 278.2 | 63.97 |
| 0.60 | 0.774 | 1.314 | 2.088 | 266.8 | 64.72 |
| 0.65 | 0.785 | 1.244 | 2.292 | 251.7 | 65.43 |
| 0.70 | 0.799 | 1.187 | 2.532 | 232.3 | 66.10 |
| 0.75 | 0.812 | 1.135 | 2.857 | 208.0 | 66.59 |
| 0.80 | 0.820 | 1.079 | 3.419 | 178.4 | 66.80 |
| 0.85 | 0.825 | 1.056 | 3.738 | 142.1 | 66.82 |
| 0.90 | 0.888 | 1.034 | 4.280 | 102.3 | 66.60 |
| 0.95 | 0.923 | 1.007 | 5.761 | 54.9 | 65.82 |
| 1.00 |  |  |  |  | 63.50 |

At $39.95^{\circ} \mathrm{C}$.

|  |  | 55.15 |
| :--- | :--- | :---: |
|  |  | 175.9 |
| Heterogeneous |  | 175.9 |
| Heterogeneous |  | 175.9 |
| Heterogeneous |  | 175.9 |
| Heterogeneous |  | 175.9 |
| Heterogeneous |  | 175.9 |
| Heterogeneous |  | 175.9 |
| Heterogeneous |  | 175.9 |
| Heterogeneous |  | 176.69 |
| 1.727 | 1.782 | 350.6 |
| 1.574 | 1.939 | 346.9 |
| 1.450 | 2.123 | 337.7 |
| 1.348 | 2.344 | 323.1 |
| 1.263 | 2.610 | 303.3 |
| 1.199 | 2.910 | 278.1 |
| 1.157 | 3.189 | 248.3 |
| 1.090 | 3.913 | 212.7 |
| 1.047 | 4.712 | 168.9 |
| 1.020 | 5.601 | 118.1 |
| 1.003 | 6.761 | 60.8 |
|  |  |  |
|  |  | 189.25 |

Table V. Calculated Vapor-Liquid Equilibria of Di-n-Propylamine-Water

usually the least reliable measurement in phase equilibrium studies, probably no accuracy is lost and time is gained by omitting this measurement, particularly at low pressures. Vapor phase compositions were calculated by the total pressure method (1, 2, 4), in which the Gibbs-Duhem equation is integrated at constant temperature using pressurecomposition data.

Two computer programs were used to carry out the numerical integration. One used the method of Boissonnas (2). Starting with $x_{2}=0, p_{2}$ for the first equally spaced increment of $X$ is obtained by assuming that the other component obeys Raoult's law, so that

$$
\begin{equation*}
\Delta p_{2}=\Delta P+p_{1}^{\circ} \Delta x_{2} \tag{1}
\end{equation*}
$$

in which $p_{2}$ and $x_{2}$ are the partial pressure and mole fraction of component 2. $P$ is the total pressure and $p_{i}^{\circ}$ is the vapor pressure of pure component 1 . Succeeding increments are obtained from the equation,

$$
\begin{equation*}
\Delta P_{2}=\frac{\Delta P}{1-\frac{p_{1} x_{2}}{p_{2} x_{1}}} \tag{2}
\end{equation*}
$$

in which $p_{1}, p_{2}, x_{1}$, and $x_{2}$ are the values at the beginning
of the increment. This is a rather crude approximation, but in many instances it gives results as good as theoretically better approximations. The second program integrated the Gibbs-Duhem equation in the form

$$
\begin{equation*}
\ln \gamma_{2}=-\frac{x_{1}}{x_{2}} \ln \gamma_{1}+\int_{1}^{x_{1}} \frac{\ln \gamma_{1}}{x_{2}} d x_{1} \tag{3}
\end{equation*}
$$

The value of the activity coefficient, $\gamma_{1}$, at infinite dilution was obtained by assuming that the other component obeyed Raoult's law and extrapolating $p_{1} /\left(x_{1} p_{1}^{\circ}\right)$ to $x_{1}=0$.

By substituting an integration formula such as the trapezoidal rule or Simpson's rule into Equation 3, an expression of the following form may be obtained:

$$
\begin{equation*}
\ln \gamma_{2}=A+B \ln \gamma_{\gamma} \tag{4}
\end{equation*}
$$

in which $A$ and $B$ are constants. This is solved simultaneously with the relation

$$
\begin{equation*}
P=x_{1} p_{i}^{\circ} \gamma_{1}+x_{2} p_{i}^{i} \gamma_{2} \tag{5}
\end{equation*}
$$

to obtain $\gamma_{1}$ and $\gamma_{2}$ for each value of $x$. For the first increment the trapezoidal rule was used; after that, Simpson's rule. Both programs assume that the gas phase is ideal.


Figure 1. Pressure-composition data for diisopropylamine-water

Both Equations 2 and 3 become indeterminate at an azeotrope. For azeotropic systems the pressure curve must be integrated from each extreme to the azeotrope.

To obtain the pressure for equally spaced values of the mole fraction, the experimental data in Tables II and III were plotted on large scale graph paper and values of $P$ at intervals of 0.025 mole fraction were taken from these graphs. The corresponding values of $x$ and $P$ were used in the numerical integration.

In the second program (Equations 3, 4, and 5) the increment size was altered as a heterogeneous region was approached, so as to terminate exactly at the saturation concentration. The activity coefficient at the other side of the immiscible region was determined by equating $\gamma x$ for the equilibrium liquid phases. The interval was again altered to obtain a succeeding value of $x$ divisible by 0.025 . The trapezoidal rule was used in the immediate neighborhood of the heterogenous region.

The results in Tables IV and V were obtained from both programs. In general, for amine concentrations greater than the azeotropic concentration, the data were integrated by the program using Equations 1 and 2. The other side of the curve was integrated with the other program when it was necessary to span a heterogenous region or when a better match at the azeotrope was obtained.

The extreme flatness of these curves causes some problems. It makes it very difficult to pick the correct azeotropic composition, and it makes the calculations very sensitive to errors in the pressure curve. Pressure differences of a few hundredths of a millimeter can make a significant difference in the calculated vapor composition, and the accuracy of the data is only about 0.1 mm . Very minute changes in the total pressure can significantly improve agreement at the azeotrope. The mismatch at the azeotrope is always less than 1 mole $\%$, however, and the two programs agree everywhere within 1 mole $\%$.

The difficulty at the azeotrope could be avoided by using the method of Barker (1), except that it requires an analytical expression which can accurately represent the data. A 10 -constant Redlich-Kister equation will not fit the amine-water data within the accuracy of the measurements.

Some of the results are shown in Figures 1 and 2. The points are experimental pressure-composition data, except


Figure 2. Pressure-composition data for di-n-propylamine-water
for those on each side of the heterogeneous regions. These show the limits of the miscibility gap. The pressures within the partially miscible region were obtained from data omitted from the graphs. The broken lines are extrapolations. This portion of the vapor composition curve can be calculated from pressure-liquid composition data with good accuracy by assuming that water obeys Raoult's law.

## NOMENCLATURE

$G^{*}=$ excess Gibbs free energy, cal./g. mole
$P=$ total pressure, mm . of Hg
$p=$ partial pressure, mm. of Hg
$x=$ mole fraction in liquid
$y=$ mole fraction in vapor

## Greek Letter

$\gamma=$ activity coefficient

## Subscripts

$1=$ component 1 or amine
$2=$ component 2 or water

## Superscript

$=$ pure component

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