

Table II. Properties of Sulfur Oxyfluorides

Formula	Name	M.P., ° C.	B.P., ° C.	F <sup>19</sup> NMR Chemical Shifts <sup>a</sup>
SO <sub>2</sub> F <sub>2</sub>	Sulfuryl fluoride	-120(?)	-52	-33.5
FSO <sub>2</sub> OF	Fluorine fluorosulfonate	-158.5	-31.3	-249, -37
FSO <sub>2</sub> OSO <sub>2</sub> F	Pyrosulfuryl fluoride	-58	51	-48.5
FSO <sub>2</sub> OOSO <sub>2</sub> F	Peroxydisulfuryl difluoride	-55.4	67.1	-40.4
FSO <sub>2</sub> OOF	Fluorosulfuryl hypofluorite	...	1	-291, -43.0

<sup>a</sup> Measured relative to CFCI<sub>3</sub>.

chloride to convert the fluorine to chlorine. The chlorine was condensed at liquid nitrogen temperature and measured volumetrically. The oxygen was transferred by means of a Toepler pump to a standard volume and measured. Small but various amounts (as much as 10% of the total gas) of SiF<sub>4</sub> were found, and corrections were made when the amounts were appreciable.

**Reactions of O<sub>2</sub>F<sub>2</sub>.** The material to be reacted with O<sub>2</sub>F<sub>2</sub> was placed in a small glass tube containing a magnetic stirrer. The O<sub>2</sub>F<sub>2</sub> was then distilled into the tube, which was cooled to -196°C. Distillation usually resulted in a slight decomposition of O<sub>2</sub>F<sub>2</sub>. If a solvent was used in the reaction, it was then distilled into a reactor and used to wash down the O<sub>2</sub>F<sub>2</sub> frozen on the side of the tube. The reaction was considered complete upon the disappearance of the red color of O<sub>2</sub>F<sub>2</sub>.

**Separation and Identification of Products.** Depending on the complexity of the products, trap-to-trap distillation or the Cady (1) codistillation method was used to separate the products. Both infrared and NMR analysis were used to characterize the products, but it was found that F<sup>19</sup> NMR analysis was usually the best method, since the chemical shifts for all the compounds in question had previously

been reported (4). The physical properties and NMR data for all of these compounds are given in Table II.

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## λ-Type Thermal Anomaly in Triuranium Octaoxide at 482.7° K.

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**Heat capacities of U<sub>3</sub>O<sub>8</sub> have been measured by adiabatic calorimetry from 300° to 555° K. A reversible λ-type thermal anomaly of uncertain origin, and with an enthalpy of transition of only 41 cal./[(mole ° K)], was observed at 482.7° K.**

APART from the potential importance of U<sub>3</sub>O<sub>8</sub> as a nuclear fuel, its thermophysical properties are of interest because of its polymorphism. The heat capacity of the α-phase was investigated from 5° to 350° K. by Westrum and Grønvdal (12). A λ-type thermal anomaly, which is presumed to be of magnetic origin, occurs at 25.3° K. Magnetic susceptibility data do indeed show a small anomaly at this temperature but do not reveal the primary nature of the anomaly (8). Heat capacity measurements have been made by other workers, covering the temperature regions 477° to 1254° K. (3) and 623° to 1273° K. (7). The temperature region 350° to 477° K., however, has remained unexplored. The present investigation endeavors to remedy this deficiency.

#### EXPERIMENTAL

The U<sub>3</sub>O<sub>8</sub> sample was prepared in this laboratory eight years ago by purifying uranyl nitrate hexahydrate, A. R. by the well-known procedure of Hönigschmidt (cf. 12). This U<sub>3</sub>O<sub>8</sub> was reduced to UO<sub>2</sub> by heating it in dry, purified

hydrogen gas at 500° until the formation of water ceased. The temperature was raised to 1200° and the sample kept at this temperature for 4 hours before cooling to room temperature. It then was oxidized in air at 800° to constant weight, transferred to a vitreous silica tube which was evacuated and sealed, heated at 800° for seven days, and gradually cooled to room temperature over a period of two months.

Spectrographic analysis of the U<sub>3</sub>O<sub>8</sub> sample showed the presence of these impurities (in parts per million): Al 60, B < 0.08, Cu 3, Fe 10, Mg 30, Ni 20, and Si 200. The weight increase upon oxidation of UO<sub>2</sub> to U<sub>3</sub>O<sub>8</sub> corresponded to the theoretical value within 0.01 O/U atom, and x-ray powder photographs showed only the presence of the U<sub>3</sub>O<sub>8</sub> phase with cell dimensions in accord with the literature (12) both before and after the heat capacity determinations.

Heat capacity was determined by the quasi-adiabatic technique using a silver calorimeter (laboratory designation W-22) and platinum resistance thermometer (A-8) in the previously described Mark IV (11) cryostat. The calorimeter was sealed after admitting 110 torr of helium gas

at 300° K. The total mass of the loaded sample was 164.892 grams in vacuo. A sample density of 8.41 grams per cc. (10) was used in making the buoyancy correction. The heat capacity of the sample was greater than 63% of the total at all temperatures. All determinations of mass, temperature, resistance, voltage, and time were referred to calibrations or standardizations made by the National Bureau of Standards.

## RESULTS AND DISCUSSION

The experimental heat capacities at the mean temperatures of the determinations are presented in chronological sequence in Table I. These data are given in terms of the defined thermochemical calorie of 4.184 Joules and a gram formula mass (g.f.m.) for  $U_3O_8$  of 842.085. The results have been adjusted for curvature. The probable error of the experimental heat capacities and derived thermodynamic functions is 0.1%.

A  $\lambda$ -type thermal anomaly with a peak temperature of 482.7° K. is evident in the  $C_p$  vs.  $T$  plot of Figure 1. The measurements were repeated six times through the anomalous region, and the results of different series are given in Table II.

The chronological sequence of the  $\Delta Ht$  runs are given in Table I. Each involves a sequence of two or more energy increments which together span the approximate temperature interval 430° to 535° K. Slight adjustments of the actual lower and higher temperatures are made from the smoothed curve but electrical energies added to the calorimeter are summed directly as the major contribution. A lattice contribution is interpolated across the anomalous region and the heat capacity in excess of this is integrated to give the excess enthalpy—i.e., the  $\Delta Ht$ . The average value of  $\Delta Ht$  between 430° and 535° K. is  $41 \pm 1$  calories per g.f.m., and the corresponding entropy increment is 0.084 calorie per (g.f.m. °K.). A sharp break in the slope of the cooling curve occurs at the maximum temperature found in the heat-capacity measurements. This further confirms the reversibility of the transformation (and the absence of undercooling of the high-temperature phase).

$\Delta Ht$  Runs A, C, E, and G individually encompass so large a temperature region that they yield directly enthalpy increments rather than heat capacities. The  $\Delta H$ 's thus obtained in regions of essentially normal behavior were in excellent accord with those integrated from the heat-capacity determinations.

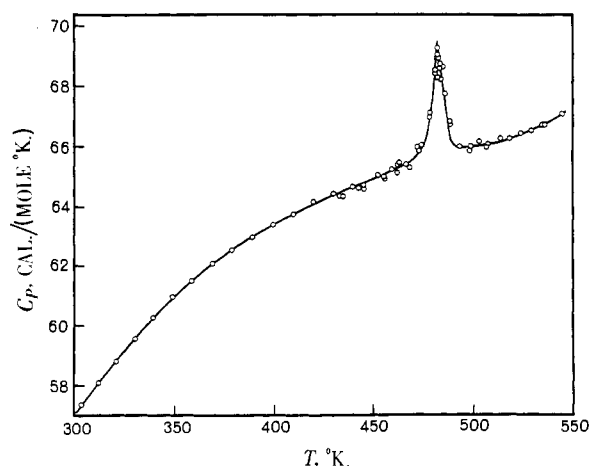


Figure 1. The heat capacity of  $U_3O_8$  from 300° to 550° K.

Table I. Heat Capacity of Triuranium Octaoxide<sup>a</sup>

T	C <sub>p</sub>	T	C <sub>p</sub>	T	C <sub>p</sub>
SERIES I		SERIES III		SERIES V	
303.02	57.39	434.70	64.25	483.20	68.87
311.54	58.06	445.02	64.69	483.97	68.97
320.48	58.82	455.28	64.98	484.74	68.74
329.92	59.56	462.97	65.44	485.51	68.60
339.28	60.27	468.14	65.27	486.28	67.72
348.54	60.98	473.27	65.82	488.72	66.70
358.38	61.49	478.37	66.94	493.35	65.98
368.77	62.08	483.40	68.68	499.06	65.97
379.09	62.48	488.42	66.80	507.67	66.06
389.32	62.94	493.51	65.96	518.62	66.23
399.49	63.36	498.61	65.81	529.50	66.49
409.58	63.73	506.72	65.93		
419.64	64.15	517.75	66.15	SERIES VI	
429.63	64.39	528.73	66.51	$\Delta H$ Run C	
439.58	64.63			$\Delta Ht$ Run D	
		SERIES IV		SERIES VII	
SERIES II		$\Delta H$ Run A		$\Delta H$ Run E	
432.69	64.34	$\Delta Ht$ Run B		458.93	65.35
442.49	64.60	536.38	66.68	461.95	65.01
452.22	65.03	SERIES V		$\Delta Ht$ Run F	
461.90	65.39	434.75	64.33	481.25	68.49
472.02	65.95	445.35	64.49	482.02	68.88
482.49	68.26	455.88	64.81	482.51	68.87
492.91	66.27	466.34	65.38	482.99	69.00
503.40	66.12	474.12	66.01	483.48	68.56
513.88	66.24	478.68	67.10	484.25	68.20
524.33	66.40	481.21	68.38	$\Delta H$ Run G	
534.74	66.67	482.28	69.23		
545.10	67.03				
555.41	67.49				

<sup>a</sup> Units: calories, g.f.m., °K.

Table II. Enthalpy of Transformation of Triuranium Octaoxide<sup>a</sup>

Designation	Number of Runs <sup>b</sup>	$H_{535} - H_{430}$
Series II	11	6917
Series III	14	6918
$\Delta Ht$ Run B	2	6917
Series V	19	6916
$\Delta Ht$ Run D	2	6917
$\Delta Ht$ Run F	11	6916
Average value	=	$6917 \pm 1$
Lattice, $H_{535} - H_{430}$	=	6876
$\Delta Ht$	=	$41 \pm 1$

<sup>a</sup> Units: calories, g.f.m., °K. <sup>b</sup> This represents the number of energy increments used to span the approximate temperature interval, 430° to 535° K.

The smoothed heat capacity and derived thermodynamic functions are given at selected temperatures in Table III. The integrations were performed by a digital computer except in the anomalous region where it was done graphically. These values of thermodynamic functions are predicated on the 273.15° K. values of Westrum and Grønvold (12). Near 300° K., both sets of data overlap smoothly and agreement is very good; the accord is everywhere within the estimated precision limits of both sets of data. The contributions due to isotopic mixing and nuclear spin are not included in the entropy and Gibbs energy functions; hence these practical values are suitable for use in chemical thermodynamic calculations. Known transformations of  $U_3O_8$  are summarized in Table IV. Recently George and Karkhanavala (5) prepared  $\delta$ - $U_3O_8$  by heating  $\alpha$ - $U_3O_8$  to 1350° C. for about 15 days and cooling it at a rate of 100° per day. They reported studies of the electrical conductivity and differential thermal analysis of the  $\delta$  phase and

Table III. Thermodynamic Properties of Triuranium Octaoxide<sup>a</sup>

$T$	$C_p$	$S^\circ$	$H^\circ - H^\circ_0$	$-(G^\circ - H^\circ_0)/T$
298.15	56.97	67.54	10216	33.27
300	57.13	67.89	10322	33.48
350	60.97	77.00	13279	39.06
400	63.36	85.31	16393	44.33
450	64.83	92.87	19603	49.31
482.7 <sup>b</sup>	69.21	97.49	21755	52.42
500	66.03	99.82	22906	54.01
550	67.23	106.15	26231	58.46

<sup>a</sup>Units: calories, g.f.m., °K. <sup>b</sup>At the peak of the lambda-type transformation.

Table IV. Thermal Properties in U<sub>3</sub>O<sub>8</sub> Transformations

Transformation	Nature	$T_i$	$\Delta Ht$	$\Delta St$	Source
$\alpha''' \rightarrow \alpha''$	Magnetic	4.2			(8)
$\alpha'' \rightarrow \alpha'^b$	Reversible	25.3	12	0.56	(12)
$\alpha' \rightarrow \alpha$	Reversible	482.7	41.7	0.086	
$\alpha \rightarrow \beta$	Reversible	1043	265	0.254	(7)
$\beta \rightarrow \gamma$	Reversible	1213	1105	0.91	(7)
$\alpha \rightarrow \delta$	Monotropic	1350	(under optimum conditions)		(6)
$\delta \rightarrow \alpha$	Monotropic	393	-1050		(6)

<sup>a</sup>Units: calories, g.f.m., °K. <sup>b</sup>The primed quantities are not established as distinct phases.

found that the conductivity *vs.* temperature plot showed a discontinuity at 388° to 398° K. (6). Their results also confirmed the presence of the  $\delta \rightarrow \alpha$  transformation at 393° K. upon heating. This (endothermic) transformation has  $\Delta Ht = 1050 \pm 50$  calories per mole and is irreversible (monotropic). The sample of  $\alpha$ -U<sub>3</sub>O<sub>8</sub> on which the present investigation was conducted was not treated under conditions productive of  $\delta$ -U<sub>3</sub>O<sub>8</sub>. Moreover, the thermal anomaly found in this investigation occurs at 482.7° K. [about 90° higher than that reported by George and Karkhanavala (5, 6)]. It is reversible and reproducible with respect to the  $\Delta Ht$  and the peak temperature,  $T_i$ . However, scattering of the experimental heat-capacity points within a few degrees of the peak temperature suggests that the crystal lattice transformation is sensitive to the rate of heating. A knowledge of structural parameters, especially in the anomalous regions of temperature, is obviously needed. Ermischer *et al.* (13) have shown that the ratio of orthorhombic to hexagonal phase is dependent on the rate of heating and cooling and the atmosphere. Their x-ray diffraction studies on U<sub>3</sub>O<sub>8</sub> show a reversible and slow transition in the temperature range of 473° to 523° K. Cordfunke (14) has confirmed the existence of the thermal anomaly in U<sub>3</sub>O<sub>8</sub> at 383° K. using a high-temperature x-ray diffraction camera. He considers that the structure of U<sub>3</sub>O<sub>8</sub> remains orthorhombic and that the hexagonal pattern occasionally observed on the x-ray films is only

a consequence of the approach to coincidence of lattice constants  $a$  and  $b$ .

Chodura and Malý (4) proposed the presence of two different coexisting structures in  $\alpha$ -U<sub>3</sub>O<sub>8</sub>, one orthorhombic and the other monoclinic. More recently Blinova *et al.* (2) found that their x-ray diffraction data are inconsistent with the existence of such structures. Neutron diffraction investigations on powdered U<sub>3</sub>O<sub>8</sub> made by Andresen (1) show that two-thirds of the uranium atoms are surrounded by seven oxygen atoms at the corners of a pentagonal double pyramid and one-third by six oxygen atoms at the corners of an octahedron. To resolve the discrepancies between the x-ray and neutron diffraction studies of U<sub>3</sub>O<sub>8</sub>, Loopstra (9) investigated a polycrystalline sample of  $\alpha$ -U<sub>3</sub>O<sub>8</sub> by neutron diffraction with improved resolution and found that the uranium atoms are surrounded by six oxygen atoms in close contact at distances between 2.07 and 2.23 Å. with a seventh oxygen at 2.44 Å. for one type of U atom and at 2.71 Å. for the other type of U atom. The minute entropy increment observed may thus be due to slight, but abnormal expansion of the crystal lattice and changes in the distances of the neighboring oxygen atoms from the central U atom.

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