Table	П.	Measured	and	Estimated	Entropies	at	298.15°	Κ.
			in G	ibbs per <i>I</i>	Nole			

Compound	$S$ $^{\circ}_{\scriptscriptstyle 298}$ Measured	S ° <sub>298</sub> Estimated	% Diff.
КОН	18.85	14.2	-24.7
$K_2CO_3$	37.17	33.6	-9.6
$K_2SiO_3$	34.93	35.2	0.8
LiCl	14.19	13.5	-4.9
$Li_2SiO_3$	19.08	23.8	24.7
MgS	12.03	12.6	4.7
SnSO4	33.12	30.3	-8.5
$AlCl_3 \cdot 6H_2O$	76.03	82.5	8.5

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## Volumetric Behavior of a Polar-Nonpolar Gas Mixture: Trifluoromethane-Tetrafluoromethane System

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Compressibility factors were determined for the polar-nonpolar gaseous system,  $CF_3H-CF_4$ , for six compositions at temperatures between  $-30^\circ$  and  $95^\circ$  C., and at temperatures as low as  $-70^\circ$  C. for CF $_4$ . The data extend to pressures as high as 1238 p.s.i.a. and are accurate to better than one part per thousand. Values are given for the constants in the Martin-Hou equation of state which correlate the volumetric data with average deviations of less than 0.5% for the pure components and 1.1% for the mixtures. Second and third virial coefficients were determined from the data; the virial coefficients for the pure gases and the interaction virial coefficients are presented. Intermolecular-potential-function parameters are given for CF<sub>4</sub>, which correlate the second virial coefficient for this gas within its experimental uncertainty between -70° and 500° C.

 ${
m T}_{
m HE}$  APPARENT NEED for volumetric data on gas mixtures containing polar components, as well as a need for such data to interpret some vapor-liquid equilibrium data (12), prompted the determination of the volumetric behavior of the  $CF_3H$ — $CF_4$  system. Of the available methods of volumetric measurement, the Burnett method was selected as having the best combination of accuracy and efficiency.

The Burnett method is an isothermal experiment in which successive portions of a test gas, confined in a primary chamber, are allowed to expand into a secondary chamber, which is discharged and evacuated after each expansion. For each expansion, comparison of the ratio of the pressures before and after the expansion to the known volume ratio of the two chambers indi-

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cates the nonideality of the test gas. The Burnett method has been described in detail (1, 13).

#### EXPERIMENTAL

The experimental pressures were measured by balancing a hydraulic pressure generated by a dead-weight gage against the test-gas pressure, using a sensitive differential-pressure transducer (DP-cell). The latter employed a thin metallic diaphragm, the position of which was sensed by a magnetic-reluctance circuit. The pressure-measurement system had a sensitivity of 0.001 p.s.i. and an over-all accuracy of 0.01 p.s.i. The test-gas pressure was continually monitored on a recorder, except during and immediately following expansions, as a check for leaks and as evidence of the restoration of thermal equilibrium after each expansion.

The temperatures of the chamber walls, as well as that of the constant-temperature bath, were measured with an accuracy of  $0.01^{\circ}$  C. The calculations included corrections for small departures from run temperature. A very small portion of the test gas was at room temperature, because the DP-cell was located outside of the constant-temperature bath. A correction for this unthermostated test gas was included in the calculations in a fashion similar to that used by previous investigators (3).

The dead-weight gage was manufactured by the Ruska Instrument Corp. The weights were calibrated by Ruska against class S standards and were checked prior to this work for internal consistency. The piston area was certified by Ruska, based on light-wave micrometer readings accurate to  $5 \times 10^{-6}$  inch.

The temperature scale was established by calibration of the thermopiles  $in \ situ$  against a platinum-resistance thermometer previously calibrated by the National Bureau of Standards.

The mixtures were analyzed with a Beckman Instruments, Inc., GC-2A gas chromatograph which used an electronic integrator for the determination of peak areas. Each mixture was compared on the chromatograph to a standard of the same nominal composition. The standards were prepared by mixing the pure components to a total pressure of approximately 2 atm. The compositions of the standards, as determined from the mixing pressures, were then corrected for slight departures from ideality using second virial coefficients. Sufficient time was allowed for the homogenization of all gas mixtures. The estimated uncertainty in the reported compositions is 0.5% of the lesser mole fraction. The CF<sub>4</sub> used in this work was 0.996 pure; the 0.004 air impurity was determined chromatographically as described above. The CF<sub>3</sub>H was quoted as 0.9998 pure by the supplier, E. I. du Pont de Nemours & Co. The compressibility factors and second virial coefficients presented below were corrected for this small air impurity.

Measurement of a vapor pressure in this apparatus provides a convenient cross check among the independently established pressure, temperature, and purity of material. Accordingly, the vapor pressure of  $CF_3H$  was measured at 0° C.; the value found was 362.47 p.s.i.a. Martin and Hou (11) report an experimental value of 363.3 at 0° C. and a smoothed value of 362.61 p.s.i.a.

#### EXPERIMENTAL DATA

Two or three series of expansions were carried out at each condition of temperature and composition, the experimental pressures being staggered to define the isotherm better. The accuracy of the compressibility factors given in Table I varies between 3 and 7 parts

		Table I. Experi	mental Compre	ssibility Factors		
		Mole	Fraction CF.	= 1.0		
Pressure.			Compress	ibility Factor		
P.S.I.A.	95.00° C.	40.00° C.	0.00° C.	-30.00° C.	-50.00° C.	-70.00° C
1200	0.9075	0.7839	0.5825	а		
1000	0.9186	0.8111	0.6309	b		
800	0.9303	0.8433	0.7019	с		
600	0.9445	0.8802	0.7802	0.6000		
500	0.9528	0.8993	0.8191	0.6915	d	
400	0.9612	0.9189	0.8573	0.7668	0.6389	
300	0.9707	0.9391	0.8945	0.8333	0.7600	
250	0.9753	0.9490	0.9126	0.8640	0.8087	0.7058
200	0.9797	0.9589	0.9306	0.8933	0.8524	0.7843
150	0.9849	0.9692	0.9484	0.9213	0.8930	0.8480
100			0.9653	0.9483	0.9303	0.9040
	Mole Fraction	$CF_{3}H = 0.0505$		Mole Fraction	$CF_{3}H = 0.0872$	
	-50.00° C.	-70.00° C.		-50.00° C.	-70.00° C.	
400	0.6311			0.6192		
350	0.6993			0.6913		
300	0.7550			0 7515		
250	0.8056			0.8025		
200	0.8493	0.7803		0.8482	0 7764	
150	0.8911	0.8462		0.8899	0.8/39	
100	0.0011	0.9030		0.0000	0.9019	
		Mole F	Fraction $CF_{3}H =$	= 0.193		
	95.00° C.	40.00° C.	0.00° C.	-30.00° C.	-50.00° C.	
1200	0.8934	0.7557	0.5230			
1000	0.9055	0.7888	0.5775			
800	0.9210	0.8267	0.6643	e		
600	0.9386	0.8686	0.7571	0.5294		
500	0.9479	0.8903	0.8018	0.6506		
400	0.9577	0.9122	0.8449	0.7423	ſ	
300	0.9676	0.9343	0.8860	0.8184	0.7333	
250		0.9452	0.9059	0.8527	0.7900	
200		0.9560		0.8850	0.8397	
150				0.9153	0.8845	

(Continued)

Davage			Compressib	ility Factor		
Pressure, P.S.I.A.	95.00° C.	40.00° C.	0.00° C.	-30.00° C.	-50.00° C.	-70.00° C.
		Mole I	$Fraction CF_{3}H = 0.$	410		
1200	0.8610	0.6909	9			
1000	0.8795	0.7351	0.4413			
800	0.9003	0.7862	0.5616			
600	0.9233	0.8404	0.6979			
500	0.9354	0.8676	0.7579			
400	0.9478	0.8947	0.8130			
300	0.9606	0.9216	0.8643	0.7790		
250	0.9671	0.9348	0.8886	0.8224		
200	0.9734	0.9479	0.9121	0.8629		
150			0.9350	0.9001		
		Mole Fraction (	$CF_{3}H = 0.607$			
1200	0.8232	0.6049				
1000	0.8474	0.6643				
800	0.8750	0.7341				
600	0.9046	0.8048	0.6062			
500	0.9199	0.8395	0.6957			
400	0.9358	0.8736	0.7705			
300	0.9516	0.9063	0.8361	,		
250	0.9599	0.9225	0.8661	n 0.0000		
200	0.9679	0.9385	0.8950	0.8328		
150				0.8798		
100				0.9226		
		Mole Fraction C	$F_{3}H = 0.8407$			
1200	0.7597	0.4322				
1000	0.7921	0.5182				
800	0.8311	0.6324				
600	0.8732	0.7402				
500	0.8947	0.7898	i			
400	0.9161	0.8361	0.6885			
300	0.9370	0.8799	0.7862			
250	0.9476	0.9013	0.8279			
200		0.9218	0.8665			
150			0.9025	0.8434		
100			0.9365	0.9008		
		Mole Fraction	$\mathrm{CF}_{3}\mathrm{H} = 1.0$			
650	0.8340	0.6365				
600	0.8467	0.6743				
500	0.8728	0.7418				
400	0.8987	0.8017				
350	0.9115	0.8299	0.6692			
300	0.9245	0.8566	0.7348			
250	0.9373	0.8824	0.7901			
200	0.9498	0.9073	0.8392	,		
150		0.9315	0.8839	1		
100		0.9551	0.9251	0.8814		
1216.070, Z =	= 0.3851.	$^{d}P = 4$	66.481, Z = 0.4989	).	${}^{h}P = 224.920$	S, Z = 0.8074.
1025.151, Z	= 0.3569.	${}^{e}P = 6$	65.690, $Z = 0.4119$		${}^{i}P = 447.160$	S, Z = 0.6305.
745.116, Z =	= 0.4456.	$^{f}P = 3$	31.716, Z = 0.6914		P = 142.69	$\Theta, Z = 0.8199.$
		$^{\sigma}P = 12$	10.740, Z = 0.4056	i.		

Table I. (Continued)

in 10,000, tending to be poorer at the higher pressures and lower temperatures. The absence of high-pressure data for pure CF<sub>3</sub>H at 40° and 95° C. is due to the fact that pure CF<sub>3</sub>H would have condensed at room temperature in the DP-cell. The values presented in Table I are interpolations to round values of pressure. The original data at the experimentally measured pressures are available as a NAPS document.

## CORRELATION OF VOLUMETRIC DATA

 ${}^{a}P =$  ${}^{b}P =$  ${}^{c}P =$ 

Constants for the Martin-Hou equation of state (10) have been determined previously for pure CF<sub>4</sub> and for pure CF<sub>3</sub>H based on experimental data (11) for these

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gases. The numerical values, taken from technical bulletins available from Freon Products Division, E. I. du Pont de Nemours & Co., Inc., are shown in columns 2 and 4 of Table II. These constants, when used in the Martin-Hou equation,

$$P = \frac{RT}{(V-b)} + \frac{A_2 + B_2T + C_2 e^{-kT/T_c}}{(V-b)^2} + \frac{A_3 + B_3T + C_3 e^{-kT/T_c}}{(V-b)^3} + \frac{A_4 + B_4T}{(V-b)^4} + \frac{A_5 + B_3T + C_5 e^{-kT/T_c}}{(V-b)^5} + \frac{A_6 + B_6T}{e^{av}} \quad (1)$$

predict pressures at the densities measured in this work which are in close agreement with the experimental

artin-Hou Constant	$\mathrm{CF}_4$	Interaction	$CF_{3}H$
$A_2$	$-1.675376  imes 10^4$	$-1.936576  imes 10^4$	$-2.294265  imes 10^4$
$A_3$	$3.002263 \times 10^{3}$	$-7.452064~ imes~10^{-1}$	$-4.282594 \times 10^{3}$
$A_4$	$1.152582 \times 10^{4}$	$2.614219  imes 10^4$	$4.971046 imes10^4$
$A_5$	$-2.366137 \times 10^{4}$	$-4.438648  imes 10^4$	$-6.511159 \times 10^{4}$
$A_{6}$	$5.838823  imes 10^7$	$6.635878 imes10^7$	$7.502357  imes 10^7$
$B_2$	$1.653808 \times 10^{11}$	$1.678043  imes 10^{1}$	$1.702634 \times 10^{12}$
$B_3$	$8.745026 imes10^{\circ}$	$1.764664 imes10^{1}$	$2.654826  imes 10^{1}$
$B_4$	$-2.350832  imes 10^{1}$	$-4.562773  imes 10^{1}$	$-8.855969 \times 10^{1}$
$B_5$	$4.785193  imes 10^{1}$	$6.644246  imes 10^{1}$	$1.086551  imes 10^{2}$
$B_6$	$-9.263923 \times 10^{4}$	$-1.015967~ imes~10^{5}$	$-1.114202 \times 10^{5}$
$C_2$	$-1.467135 \times 10^{5}$	$-2.471400  imes 10^{5}$	$-7.833461  imes 10^{5}$
$C_3$	$3.679677  imes 10^{3}$	$9.773246 imes10^{\scriptscriptstyle 5}$	$2.039583 imes10^6$
$C_5$	$-2.553920  imes 10^{5}$	$-6.251347  imes 10^{5}$	$-1.244522~ imes~10^{6}$
k	4.00	4.69	5.50
$T_{\iota}$	409.50	469.52	538.33
а	7.51278	7.46949	7.42645
Ъ	0.13200	0.10977	0.08753
R	10.73168	10.73168	10.73168

Table II. Recommended Values of Martin-Hou Constants for CF<sub>4</sub>-CF<sub>3</sub>H System

pressures (average deviations of 0.24% for CF<sub>4</sub>, 0.43% for CF<sub>3</sub>H, and maximum deviations of about 1% for both gases). The largest deviations occur at the highest pressures and in the vicinity of the critical point.

A quadratic composition dependency is usually assumed for the Martin-Hou constants,

$$K = y_1^2 K_{11} + 2 y_1 y_2 K_{12} + y_2^2 K_{22}$$
(2)

where K is any Martin-Hou constant, subscript 1 refers to  $CF_4$ , and 2 refers to  $CF_3H$ . The interaction constants,  $K_{12}$  in Equation 2, which best correlate the mixture data are listed in Table II. Only interaction constants bearing some rational relationship to the pure-gas constants were considered (9). The values given predict the experimental pressures for the mixtures with an average deviation of 1.1% and a maximum deviation of 2.8%.

Table	III.	Second	Virial	Coefficients	for	CF <sub>4</sub> -CF <sub>3</sub> H	System
			(Cc	./gram-mole	)		

° C.	$B_{11}$ (CF <sub>4</sub> )	$B_{12}$	$B_{22}$ (CF <sub>3</sub> H)
$95 \\ 40 \\ 0 \\ -30 \\ -50 \\ -70$	$\begin{array}{r} -45.75 \pm 0.2 \\ -77.26 \pm 0.2 \\ -110.87 \pm 0.2 \\ -146.36 \pm 0.2 \\ -177.71 \pm 0.5 \\ -216.75 \pm 0.7 \end{array}$	$\begin{array}{c} -51.13 \pm 1.0 \\ -85.28 \pm 1.4 \\ -122.75 \pm 1.4 \\ -159.33 \pm 1.8 \end{array}$	$\begin{array}{r} -109.50 \pm 0.2 \\ -165.50 \pm 0.2 \\ -233.60 \pm 0.4 \\ -311.60 \pm 1.0 \end{array}$

VIRIAL COEFFICIENTS

The coefficients of the second and third terms in the virial equation

$$Z = 1 + \frac{B(T)}{V} + \frac{C(T)}{V^2} + \frac{D(T)}{V^3} + \cdots$$
(3)

were determined from plots of (Z - 1)V vs. 1/V. They are given for the pure gases in Tables III and IV. Also given in Tables III and IV are the second and third interaction virial coefficients, which were calculated from the composition dependencies of the virial coefficients of the mixtures and which refer to Equations 4 and 5 (6). The virial coefficients of CF<sub>4</sub> at 0° C. reported here are in good agreement with values reported elsewhere (5).

$$B(T) = y_1^2 B_{11}(T) + 2y_1 y_2 B_{12}(T) + y_2^2 B_{22}(T)$$
(4)

$$C(T) = y_1^3 C_{111}(T) + 3y_1^2 y_2 C_{112}(T) + 3y_1 y_2^2 C_{122}(T) + y_2^3 C_{222}(T)$$
(5)

The estimated accuracies of the virial coefficients are given in the tables. The temperature dependencies of the virial coefficients are shown in Figures 1 and 2. In Figure 2, the expected low-temperature maximum in the third virial coefficient (2) is in evidence, but is not fully defined by these data.

The linearity of the (Z - 1)V vs. 1/V plots indicated that Equation 3. truncated after the third term, is

## Table IV. Third Virial Coefficients for CF<sub>4</sub>-CF<sub>3</sub>H System (Cc. per gram-mole)<sup>2</sup>

° C.	$C_{111}$ (CF4)	$C_{112}$	$C_{122}$	$C_{222}$ (CF <sub>3</sub> H)
95	$4890 \pm 150$	$4700 \pm 500$	$4200 \pm 300$	$8410~\pm~250$
40 0	$6030 \pm 125$ $7070 \pm 100$	$4950 \pm 200 \\ 6400 \pm 200$	$7100 \pm 200$ $10.250 \pm 300$	$11,630 \pm 100$ $15,700 \pm 175$
$-30_{50}$	$8660 \pm 100$	$7200 \pm 500$	$11,800 \pm 600$	$15,000 \pm 1000$
-70	$9320 \pm 500$			



Figure 1. Second virial coefficients



Figure 2. Third virial coefficients

#### Table V. Optimum Values of Parameters for CF<sub>4</sub> for Spherical-Shell and Spherical-Kihara Functions

#### Spherical-Shell Function

Equilibrium separation, $r_o$ , A. Core diameter, $d$ , A. Energy parameter, $\epsilon/k$ , °K.	$4.60 \\ 2.585 \\ 302.75$	H H H	0.01 0.005 0.1
Spherical-Kihara Function	n		
	4 0 0		o o 1

Collision diameter, $\circ$ , A.	$4.33 \pm 0.01$
Core radius, a, A.	$0.715 \pm 0.005$
Energy parameter, $\varepsilon/k$ , ° K.	$287.625 \pm 0.1$

valid, within the limits of accuracy of the compressibilities measured here, to pressures as high as approximately 0.9 times the critical pressure at temperatures below approximately 1.1 times the critical temperature and to still higher pressures at higher temperatures.

#### CORRELATION OF SECOND VIRIAL COEFFICIENTS OF CF4

The second virial coefficients of  $CF_4$  reported here, together with values at temperatures as high as 500° C. reported by other investigators (5, 7), were correlated within their experimental accuracies using either the spherical-shell (4) or a spherical Kihara (8) intermolecular potential-energy function. The generalized second virial coefficient is given in tabulated form for each of the functions in the references cited. The optimum values of the function parameters for  $CF_4$  are given in Table V. The second-virial-coefficient data used at temperatures above the range of this work are considered accurate to 0.3 cc. per gram-mole, based on the accuracy of the volumetric data from which they were obtained.

## ACKNOWLEDGMENT

The donation of samples of  $CF_4$  and  $CF_3H$  by the Freon Products Division, E. I. du Pont de Nemours & Co., is sincerely appreciated. Fellowship support from NDEA is gratefully acknowledged and appreciated.

#### NOMENCLATURE

- B = second virial coefficient, cc./gram-mole (except in Equation 1)
- C =third virial coefficient, (cc./gram-mole)<sup>2</sup> (except in Equation 1)
- K = general Martin-Hou constant
- P = absolute pressure, p.s.i.a.
- T = absolute temperature, ° R, or ° K.
- V = molar volume
- y =mole-fraction
- Z = compressibility factor

#### Subscripts

- c = critical property
- $1 = \text{refers to } CF_4$
- $2 = \text{refers to } CF_3H$  (except in Equation 1)

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# Pressure-Volume-Temperature-Concentration Relation of Aqueous NaCl Solutions

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The derivatives  $\left(\frac{\partial \mathbf{v}}{\partial P}\right)_{T,x}$  of NaCl solutions have been experimentally determined

in the temperature range  $0^{\circ}$  to  $175^{\circ}$  C. for NaCl concentrations of 0 to 25 grams per 100 grams of solution and pressures up to 350 kg, per sq. cm. An interpolation formula which describes the pressure-volume-temperature-concentration (*P*-v-*T*-x) relation has been developed to fit these experimental results and the density data from the literature.

BECAUSE of the importance of converting salt water to potable water, there is considerable interest in the thermodynamic properties of aqueous NaCl solutions. The petroleum industry is also interested in these properties for predicting steam drive and aquifer performance. Pressure, specific volume, temperature, and concentration data are the basic data required for calculating thermodynamic properties of solutions. Although no equation can be found to represent the relations among the properties in an absolutely correct manner without understanding the fundamental nature of the substances and their interactions, an empirical equation which adequately represents the property relations can be used to interpolate the experimental data, facilitate calculations involving integration and differentiation, and provide a concise representation of a large mass of data. An efficient evaluation of the properties is nearly impossible without interpolation formulas.

A survey of literature indicated the lack of experimental data on compressibility for NaCl solutions at temperatures above 40° C.; measurements were therefore made to determine the values of the derivative  $\left(\frac{\partial v}{\partial P}\right)_{T.x}$  for 0, 5, 10, 15, 20, and 25% of NaCl in solutions at 21.7°, 50.9°, 81.0°, 119.4°, 146.9°, and 174.4° C.

up to pressures of 350 kg. per sq. cm. By using test results and low-pressure density data from the literature, a formula was then improvised to describe the P-v-T-x relations. The estimated maximum deviation of the observations from this formula was less than 1.5 p.p.t. over the entire region investigated.

#### EXPERIMENTAL

The compressibilities of a liquid are often experimentally determined by measuring the total change in specific volume caused by the change in pressure from 1 to P atm. However, this experiment was designed to obtain the differential change in specific volume associated with a small pressure change at pressure P, at the fixed temperature, by using a calibrated positive displacement pump. The apparatus used is shown in Figure 1. The vessel used for testing the solutions was made of stainless steel, surrounded by three electrical heating coils of 650 watts each, and insulated by a 1.5-inch layer of mineral wool. The electric current through the coil was controlled by a variable resistor that maintained the desired temperature of the solution in the vessel. To minimize temperature variation of the solution, 415 feet of 0.01-inch o.d. nickel-chromium wire was placed inside the vessel to distribute the heat uniformly through the solution. The pump piston, which was calibrated to 0.01 cc. of piston displacement, could be advanced or retreated at uniform speed of 0 to 100 cc. of

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