# Refractive Indices of Ethane, Carbon Dioxide, and Isobutane 

George J. Besserer and Donald B. Robinson ${ }^{1}$<br>Department of Chemical and Petroleum Engineering, University of Alberta, Edmonton, Alta., Canada

Refractive index measurements on ethane, carbon dioxide, and isobutane are reported for temperatures of $100^{\circ}, 160^{\circ}, 220^{\circ}$, and $250^{\circ} \mathrm{F}$ at pressures up to about 1500 psia. Compressibility factors were calculated for ethane and carbon dioxide using densities obtained from the Lorentz-Lorenz relationship. These compared favorably with published values for the same materials.

A number of experimental studies have been made to determine the relationship between density and refractive index. The data have frequently been used to test the validity of the Lorentz-Lorenz theory and to examine by experimental measurement the temperature and density dependence of the Lorentz-Lorenz refractivity. These studies have often involved a lighter gaseous substance which is confined and then compressed so that the refractive index can be obtained over a range of pressures at each of several different temperatures.

Measurements of this kind have been made for carbon dioxide by Michels and Hamers (3) at temperatures of $25^{\circ}, 50^{\circ}$, and $100^{\circ} \mathrm{C}$ up to pressures of 2400 atm and by Phillips (4) at $34^{\circ} \mathrm{C}$ to pressures of 85 atm . Schmidt and Traube (7) and Vukalovich et al. (9) have made accurate measurements in the critical region. Measurements of refractive indices and the corresponding densities along the coexisting vapor-liquid curve have been made for ethane and isobutane by Sliwinski (8).

The object of the work reported here was to determine the refractive index of carbon dioxide, ethane, and isobutane along isotherms at $100^{\circ}, 160^{\circ}, 220^{\circ}$, and $250^{\circ} \mathrm{F}$ and at pressures up to about 1500 psia. These data were used with the Lorentz-Lorenz refractivities to calculate compressibility factors for carbon dioxide and ethane which existed in the gaseous region at all conditions.

## Experimental

Method. The experimental equipment and method have been described in detail by Besserer and Robinson (1). The sample under investigation was confined in a variable volume temperature-controlled stainless steel cell. The refractive index was measured by a minimum deviation angle method. The measured angles for the refractive index detemination had a repeatability of $\pm 4 \mathrm{sec}$ and a table error of $\pm 2 \mathrm{sec}$ over the measured range giving an overall possible error in the refractive index of $\pm 6 \times$ $10^{-5}$ in the range $1,0-1.4$. The temperature of the cell contents was measured with an iron constantan thermocouple pair with the reference junction placed in an ice bath. The thermocouple pair was calibrated at the triple point and steam point of triple distilled water. The temperatures are believed known to $\pm 0.1^{\circ} \mathrm{F}$. The pressure of the system was measured with a strain gage pressure transducer calibrated against the vapor pressure of carbon dioxide at $68^{\circ} \mathrm{F}$ and the pressure readings are believed known to $\pm 3$ psia.

Materials. The isobutane, carbon dioxide, and ethane used in the investigation were obtained from the Matheson Co., Canadian Liquid Air Ltd., and the Phillips Petroleum Co., respectively. In each case the purity was greater than $99.9 \mathrm{~mol} \%$.
${ }^{1}$ To whom correspondence should be addressed.

## Results and Calculations

The experimentally determined refractive indices for isobutane, carbon dioxide, and ethane are shown in Figures $1-3$, respectively. The curves for isobutane were obtained on the liquid-like region and consequently are terminated along the saturated liquid line. Smoothed values of the refractive indices at uniform increments of pressure up to 1500 psia are given for each material at each of the four temperatures in the range $100-250^{\circ} \mathrm{F}$ in Ta bles I-III. All refractive indices are given relative to vacuum using a monochromatic light source from a heliumneon gas laser at a wavelength of $6328 \AA$.

The Lorentz-Lorenz formulation for the relationship between refractive index and density is given by

$$
\begin{equation*}
\frac{M}{\rho} \cdot \frac{n^{2}-1}{n^{2}+2}=\mathrm{const}=R_{L L} \tag{1}
\end{equation*}
$$

Experimental data have shown that the refractivity, $R_{L L}$, is not exactly constant but increases slightly with increasing density, passing through a maximum and declining thereafter. Buckingham and Pople (2) have suggested that this effect can be approximated by

$$
\begin{equation*}
R_{L L}=R_{L L}{ }^{0}+\frac{B}{V}+\frac{C}{V^{2}} \tag{2}
\end{equation*}
$$

where $R_{L L}{ }^{0}$ is the limit of the molar refractivity at zero density, $B$ and $C$ are the second and third virial refractivity coefficients respectively, and $V$ is the molar volume.


Figure 1. Effect of temperature and pressure on the refractive index of liquid isobutane

Table I. Effect of Temperature and Pressure on the Refractive Index of Liquid Isobutane

|  | Refractive index ${ }^{a}$Pressure <br> psia |  |  |  |
| :---: | :---: | :---: | :---: | :---: | | $100.0^{\circ} \mathrm{F}$ | $159.8^{\circ} \mathrm{F}$ | $220.0^{\circ} \mathrm{F}$ | $250.0^{\circ} \mathrm{F}$ |  |
| :---: | :---: | :---: | :---: | :---: |
| 100 | 1.3068 | $\ldots$ | $\ldots$ | $\ldots$ |
| 200 | 1.3078 | 1.2779 | $\ldots$ | $\ldots$ |
| 300 | 1.3088 | 1.2800 | $\ldots$ | $\ldots$ |
| 400 | 1.3097 | 1.2820 | 1.2406 | $\ldots$ |
| 500 | 1.3107 | 1.2838 | 1.2449 | 1.2166 |
| 600 | 1.3117 | 1.2852 | 1.2486 | 1.2239 |
| 700 | 1.3125 | 1.2866 | 1.2519 | 1.2300 |
| 800 | 1.3132 | 1.2878 | 1.2548 | 1.2349 |
| 900 | 1.3141 | 1.2889 | 1.2576 | 1.2389 |
| 1000 | 1.3149 | 1.2901 | 1.2599 | 1.2422 |
| 1100 | 1.3158 | 1.2913 | 1.2620 | 1.2452 |
| 1200 | 1.3165 | 1.2925 | 1.2641 | 1.2480 |
| 1300 | 1.3173 | 1.2937 | 1.2661 | 1.2506 |
| 1400 | 1.3181 | 1.2948 | 1.2680 | 1.2530 |
| 1500 | 1.3190 | 1.2960 | 1.2700 | 1.2553 |

${ }^{a}$ Relative to vacuum at 6328 A .


Figure 2. Effect of temperature and pressure on the refractive index of carbon dioxide

Plots of $\left(R_{L L}=R_{L L}{ }^{0}\right) V$ vs. $1 / V$ were made for Michels and Hamers' (3) high-density data for carbon dioxide and from this the following equation was derived:

$$
\begin{equation*}
R_{\mathrm{LL}}=\frac{n^{2}-1}{n^{2}+2} \cdot V=6.600+\frac{1.25}{V}-\frac{264}{V^{2}} \tag{3}
\end{equation*}
$$

where $V$ is in cc/g-mol.
A similar expression was used by Sliwinski (8) for ethane:

$$
\begin{equation*}
R_{L L}=11.26+\frac{23.2}{V}-\frac{1670}{V^{2}} \tag{4}
\end{equation*}
$$

The smoothed experimental values of the refractive indices were used in these expressions to find the molar

Table II. Effect of Temperature and Pressure on Refractive Index of Carbon Dioxide

|  | Refractive index ${ }^{a}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Pressure <br> psia | $100.2^{\circ} \mathrm{F}$ | $160.0^{\circ} \mathrm{F}$ | $219.8^{\circ} \mathrm{F}$ | $250.0^{\circ} \mathrm{F}$ |
| 100 | 1.0027 | 1.0025 | 1.0023 | 1.0021 |
| 200 | 1.0055 | 1.0050 | 1.0046 | 1.0042 |
| 300 | 1.0086 | 1.0077 | 1.0069 | 1.0064 |
| 400 | 1.0120 | 1.0104 | 1.0093 | 1.0087 |
| 500 | 1.0158 | 1.0134 | 1.0118 | 1.0110 |
| 600 | 1.0199 | 1.0165 | 1.0144 | 1.0135 |
| 700 | 1.0246 | 1.0199 | 1.0172 | 1.0161 |
| 800 | 1.0299 | 1.0235 | 1.0201 | 1.0187 |
| 900 | 1.0364 | 1.0274 | 1.0231 | 1.0214 |
| 1000 | 1.0454 | 1.0314 | 1.0261 | 1.0241 |
| 1050 | 1.0508 | 1.0336 | 1.0276 | 1.0255 |
| 1100 | 1.0570 | 1.0358 | 1.0292 | 1.0269 |
| 1150 | 1.0657 | 1.0382 | 1.0308 | 1.0283 |
| 1200 | 1.0805 | 1.0408 | 1.0324 | 1.0298 |
| 1250 | 1.1105 | 1.0434 | 1.0341 | 1.0312 |
| 1300 | 1.1333 | 1.0462 | 1.0358 | 1.0327 |
| 1350 | 1.1438 | 1.0492 | 1.0376 | 1.0342 |
| 1400 | 1.1499 | 1.0520 | 1.0395 | 1.0358 |
| 1500 | 1.1580 | 1.0587 | 1.0432 | 1.0390 |
|  |  |  |  |  |
| $a_{\text {Relative to vacuum at } 6328 \mathrm{~A} .}$ |  |  |  |  |

Table III. Effect of Temperature and Pressure on the Refractive Index of Ethane

|  | Refractive index ${ }^{a}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Pressure, <br> psia | $100.3^{\circ} \mathrm{F}$ | $160.4^{\circ} \mathrm{F}$ | $219.6^{\circ} \mathrm{F}$ | $249.6^{\circ} \mathrm{F}$ |
| 100 | 1.0047 | 1.0043 | 1.0039 | 1.0037 |
| 200 | 1.0100 | 1.0089 | 1.0080 | 1.0075 |
| 300 | 1.0159 | 1.0138 | 1.0122 | 1.0115 |
| 400 | 1.0230 | 1.0191 | 1.0166 | 1.0157 |
| 500 | 1.0315 | 1.0251 | 1.0212 | 1.0199 |
| 600 | 1.0438 | 1.0318 | 1.0264 | 1.0245 |
| 700 | 1.0611 | 1.0395 | 1.0319 | 1.0295 |
| 750 | 1.0788 | 1.0440 | 1.0348 | 1.0321 |
| 800 | 1.1365 | 1.0490 | 1.0378 | 1.0348 |
| 850 | 1.1620 | 1.0540 | 1.0410 | 1.0375 |
| 900 | 1.1712 | 1.0591 | 1.0443 | 1.0402 |
| 1000 | 1.1835 | 1.0716 | 1.0514 | 1.0461 |
| 1100 | 1.1916 | 1.0865 | 1.0587 | 1.0521 |
| 1200 | 1.1979 | 1.1023 | 1.0665 | 1.0585 |
| 1300 | 1.2025 | 1.1180 | 1.0747 | 1.0652 |
| 1400 | 1.2063 | 1.1321 | 1.0831 | 1.0723 |
| 1500 | 1.2100 | 1.1440 | 1.0919 | 1.0795 |
|  |  |  |  |  |
| $a$ Relative to vacuum at 6328A. |  |  |  |  |

Table IV. Effect of Temperature and Pressure on the Compressibility Factor of Carbon Dioxide

|  | Compressibility factor ${ }^{a} z$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Pressure <br> psia | $100.2^{\circ} \mathrm{F}$ | $160.0^{\circ} \mathrm{F}$ | $219.8^{\circ} \mathrm{F}$ | $250.0^{\circ} \mathrm{F}$ |
| 100 | 0.971 | 0.980 | 0.984 | 0.987 |
| 200 | 0.940 | 0.959 | 0.969 | 0.974 |
| 300 | 0.908 | 0.937 | 0.953 | 0.961 |
| 400 | 0.874 | 0.913 | 0.938 | 0.949 |
| 500 | 0.839 | 0.890 | 0.922 | 0.936 |
| 600 | 0.799 | 0.865 | 0.906 | 0.923 |
| 700 | 0.755 | 0.840 | 0.890 | 0.910 |
| 800 | 0.708 | 0.814 | 0.874 | 0.896 |
| 900 | 0.657 | 0.788 | 0.857 | 0.884 |
| 1000 | 0.597 | 0.762 | 0.841 | 0.871 |
| 1050 | 0.560 | 0.748 | 0.834 | 0.864 |
| 1100 | 0.516 | 0.734 | 0.826 | 0.857 |
| 1150 | 0.465 | 0.720 | 0.818 | 0.850 |
| 1200 | 0.390 | 0.706 | 0.810 | 0.843 |
| 1250 | 0.305 | 0.691 | 0.801 | 0.837 |
| 1300 | 0.264 | 0.671 | 0.793 | 0.830 |
| 1350 | 0.254 | 0.662 | 0.785 | 0.823 |
| 1400 | 0.252 | 0.646 | 0.777 | 0.816 |
| 1500 | 0.256 | 0.616 | 0.764 | 0.803 |

$a$ Calcd using $V$ obtained from refractive index data where $\frac{n^{2}-1}{n^{2}+2} V=$
$6.600+1.25 / V-264 / V^{2}$. $6.600+1.25 / v-264 / V^{2}$.


Figure 3. Effect of temperature and pressure on the refractive index of ethane
volumes at each pressure and temperature, and from this the compressibility factors were calculated. Values are tabulated in Table IV for carbon dioxide and Table $V$ for ethane. Figures 4 and 5 show a comparison between the compressibility factors calculated from the refractive indices and those obtained by the direct volumetric measure-


Figure 4. Compressibility factors for carbon dioxide


Figure 5. Compressibility factors for ethane

Table V. Effect of Temperature and Pressure on the Compressibility Factor of Ethane

| Pressure, psia | Compressibility factor ${ }^{\text {a }} \boldsymbol{z}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $100.3^{\circ} \mathrm{F}$ | $160.4{ }^{\circ} \mathrm{F}$ | $210.6{ }^{\circ} \mathrm{F}$ | $249.6{ }^{\circ} \mathrm{F}$ |
| 100 | 0.953 | 0.966 | 0.976 | 0.979 |
| 200 | 0.902 | 0.932 | 0.951 | 0.957 |
| 300 | 0.846 | 0.896 | 0.926 | 0.936 |
| 400 | 0.788 | 0.859 | 0.901 | 0.914 |
| 500 | 0.721 | 0.819 | 0.874 | 0.892 |
| 600 | 0.639 | 0.776 | 0.848 | 0.871 |
| 700 | 0.521 | 0.729 | 0.820 | 0.849 |
| 750 | 0.433 | 0.705 | 0.807 | 0.838 |
| 800 | 0.277 | 0.682 | 0.793 | 0.827 |
| 850 | 0.246 | 0.958 | 0.778 | 0.816 |
| 900 | 0.245 | 0.633 | 0.765 | 0.805 |
| 1000 | 0.256 | 0.584 | 0.737 | 0.783 |
| 1100 | 0.272 | 0.534 | 0.711 | 0.762 |
| 1200 | . 0.288 | 0.491 | 0.685 | 0.741 |
| 1300 | 0.305 | 0.462 | 0.666 | 0.722 |
| 1400 | 0.322 | 0.446 | 0.639 | 0.704 |
| 1500 | 0.339 | 0.441 | 0.621 | 0.689 |
| a Calcd using $V$ obtained from refractive index data where $\frac{n^{2}-1}{n^{2}+2} \quad V=0.23 .2 / V-1670 / V^{2}$.$26+23$ |  |  |  |  |

ments on carbon dioxide and ethane by Reamer et al. (5, 6 ). Over most of the region, agreement between the two sets of data is very close.

## Nomenclature

$B, C=$ second and third virial refractivity coefficients, respectively
$\rho=$ density, $g / \mathrm{cc}$
$M=$ molecular weight
$n=$ refractive index at $6328 \AA$ relative to vacuum
$R_{L L}=$ Lorentz-Lorenz molar refractivity, cc/g-mol
$R_{L L}{ }^{0}=$ zero density limit of the molar refractivity at $6328 \AA$
$V=$ molar volume, cc/g-mol

## Literature Cited

(1) Besserer, G. J., Robinson, D. B., Can. J. Chem. Eng., 49, 651 (1971).
(2) Buckingham, A. D., Pople, J. A., Discuss. Faraday Soc., 22, 17 (1956).
(3) Michels, A., Hamers, J., Physica, 4, 995 (1947).
(4) Phillips, P., Proc. Roy. Soc. (London). 97A, 225 (1920).
(5) Reamer, H. H., Sage, B. H., Lacey, W. N., Ind. Eng. Chem., 36 (1), 88 (1944).
(6) Reamer, H. H., Sage, B. H., Lacey, W. N., ibid., p 956.
(7) Schmidt, E., Traube, K., Proc. Intern. Res. Thermodyn. Transp. Prop., 2nd, Princeton, N.J., 1962, p 193.
(8) Sliwinski, P., Z. Phys. Chem., Neve Folge, 63, 263 (1969).
(9) Vukalovich, M. P., Kobelev, V. P., Timoshenko, N. I., Teploenergetika. 15 (6), 80 (1968).

# Vapor Pressure of $\alpha$-Samarium and $\alpha$-Ytterbium 

Alessandro Desideri, ${ }^{1}$ Vincenzo Piacente, and Sergio Nobili<br>Laboratorio di Chimica-Fisica ed Elettrochimica, Università di Roma, Rome, Italy

The torsion-effusion method was used to determine the vapor pressure of $\alpha$-samarium and $\alpha$-ytterbium in the temperature range $\mathbf{1 0 4 3 - 1 1 7 9 K}$ and $775-907 \mathrm{~K}$, respectively. The proposed heat of sublimation at 298 K , obtained as an estimated average between our secondand third-law values, was $49.56 \pm 0.50$ and $37.20 \pm$ $0.40 \mathrm{kcal} / \mathrm{mol}$ 'for $\alpha$-samarium and $\alpha$-ytterbium, respectively.

Few data on the vapor pressure of $\alpha$-samarium and $\alpha$ ytterbium are reported in the literature (3, 4, 6, 9, 10). A new series of tensimetric data, obtained by the Knudsen torsion-effusion method, is given in this paper. The basis of the method and the experimental apparatus have been described elsewhere (7, 8). Vapor pressures were calculated from the torsion-effusion data by means of the equation

$$
\begin{equation*}
P=2 K \beta /\left(a_{1} d_{1} f_{1}+a_{2} d_{2} f_{2}\right) \tag{1}
\end{equation*}
$$

where $P$ is the pressure; $\beta$ the deflection; $K$ the torsion constant of the suspension; $a_{1}, a_{2}, d_{1}$, and $d_{2}$ the orifice area and distances from the axis of rotation; and $f_{1}$ and $f_{2}$ the Freeman and Searcy (1) correction factor for orifice geometry.

The metals were vaporized from tantalum cells placed in the constant-temperature zone of the furnace. The samarium and ytterbium samples were from Koch-Light Lab certified to be at least $99 \%$ pure with respect to total impurities. The loading of the cell was carried out with samples ( $\sim 1$ gram) mechanically cleaned and reduced to small pieces $(\sim 2 \mathrm{~mm})$ to increase the evaporating surface. Loading in inert atmosphere was not necessary (4). An ambient vacuum less than $10^{-5} \mathrm{~mm} \mathrm{Hg}$ was maintained during the runs. The temperature of the operating cell was taken to be the temperature determined by a calibrated iron-constantan thermocouple placed inside an identical but empty cell located directly beneath the operating cell. This procedure was tested putting a second thermocouple in the operating cell and measuring the differences in emf values. All the temperature uncertainties were estimated to be $\pm 2^{\circ} \mathrm{C}$ at about 1200 K . The effusion cell was suspended on a $30-\mu$ diam

[^0]tungsten wire, 33.0 cm long, with a torsion constant of $0.358 \pm 0.004 \mathrm{dyn} \mathrm{cm} / \mathrm{rad}$. The details of the effusion cells are given as a footnote in Tablel.

The vapor pressures of $\alpha$-samarium and $\alpha$-ytterbium, determined in the temperature range $1043-1179 \mathrm{~K}$ and $775-907 \mathrm{~K}$, respectively, are reported in Table I, considering the vapor constituted by a monoatomic species only (10). The corresponding values of $\Delta H_{298}{ }^{\circ}$ (subl) of both elements, calculated using the free energy function, $\left(G_{T}{ }^{\circ}-H_{298}{ }^{\circ}\right) / T$, (fef), selected by Hultgren et al. (5), are reported in the same table. The average third-law $\Delta H_{298}{ }^{\circ}$ (subl) values for Sm and Yb are $49.60 \pm 0.20$ and $37.01+0.11 \mathrm{kcal} / \mathrm{mol}$, respectively, where the associated errors are standard deviations. The leastsquares lines through the pressure data over the experimental temperature range are given by the equations:
$\log P(\operatorname{atm})=$

$$
\begin{equation*}
5.33 \pm 0.19-\frac{(10.39 \mp 0.21) \times 10^{3}}{T}, \text { samarium } \tag{2}
\end{equation*}
$$

$\log P(a t m)=$

$$
\begin{equation*}
5.61 \pm 0.17-\frac{(7.97 \mp 0.14) \times 10^{3}}{T}, \text { ytterbium } \tag{3}
\end{equation*}
$$

The lines are drawn in Figure 1 and the relative slopes give the sublimation enthalpy values, $\Delta H_{298}{ }^{\circ}(\mathrm{Sm})=$


Figure 1. Vapor pressure of solid samarium and ytterbium


[^0]:    ${ }^{1}$ To whom correspondence should be addressed.

