



Figure 5. Excess surface entropy (S_E^σ) vs. surface mole fraction (x_1^σ) of benzene-dioxane on alumina.

therefore, reduces simply to $-R \sum x_i^\sigma \ln \rho_i^\sigma$ and, when plotted as a function of composition (x_i^σ), should yield a simple curve with one maximum. The curves obtained according to the S_E^σ values of two models are shown in Figure 5. Although the actual values from the two models are different, the natures of the curves are similar.

Conclusion

The adsorption of both benzene-cyclohexane and benzene-dioxane solutions decreases with a rise in the temperature range from 20 to 40 °C. This effect is more pronounced (~5 times) in benzene-dioxane than in benzene-cyclohexane. Benzene-cyclohexane solutions show multilayer adsorption.

The surface activity coefficients are independent of temperature in this temperature range.

Glossary

σ	surface phase
l	bulk liquid
i	component 1
j	component 2
n_0	total amount of liquid mixture
Δx_i^l	change in mole fraction due to adsorption
m	mass of the solid
a	thermodynamic activity
S	specific surface area of solid
A	surface area of adsorbate
S_E^σ	excess surface entropy
T	temperature
R	gas constant
E	Everett's model
SN	Schay and Nagy's model
ρ_i^σ	surface activity coefficient
x_i^σ	surface mole fraction
γ	surface tension

Registry No. Alumina, 1344-28-1; benzene, 71-43-2; cyclohexane, 110-82-7; dioxane, 123-91-1.

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Measurement of the Volumetric Properties of Sulfur Hexafluoride from 260 to 340 K at Pressures to 2.5 MPa

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The volumetric properties of sulfur hexafluoride are determined experimentally by a Burnett apparatus at 260, 285, 315, and 340 K at pressures to 2.5 MPa. Compression factors are calculated by fitting simultaneously the isothermal pressure series to the Burnett relation and the virial equation. The calculated second virial coefficients agree well with available data from the literature. Densities and compression factors at integral pressures are calculated from the experimental data. The purity of the gas is 99.7%.

Burnett Apparatus

The gas-phase PVT properties are obtained by the Burnett method (1) in which only temperature and pressure measurements are required. The apparatus is identical with the one

described by Fredenslund et al. (2) and Møllerup and Angelo (3). The Burnett cell consists of two thick-walled vessels of unspecified volume, here approximately 125 and 275 cm³, connected to each other by means of an expansion valve. The Burnett cell is located in a thermostat and the larger volume connected to the pressure measuring system by means of a diaphragm differential pressure indicator. The pressure is measured on either an oil dead-weight piston gauge or an air dead-weight piston gauge depending upon the pressure range.

The larger vessel in the Burnett apparatus is evacuated and filled with the gas or gas mixture to be studied, and the smaller vessel is evacuated and kept under high vacuum. After isothermal conditions are obtained, the temperature and pressure are measured. The valve on the vent line to the vacuum system is closed and the expansion valve is then opened, and the pressure is allowed to equalize between the two vessels. The temperature and pressure are measured again after isothermal

Table I. Experimental Pressures, Calculated Residuals, and Compression Factors of Sulfur Hexafluoride at 260, 285, 315, and 340 K

P^* , MPa	$P^* - P$, MPa	Z	P^* , MPa	$P^* - P$, MPa	Z	P^* , MPa	$P^* - P$, MPa	Z
$T = 260$ K			$T = 315$ K			$T = 340$ K		
0.5466557	0.000004	0.89576	1.710358	-0.000553	0.81856	2.521419	0.000175	0.79229
0.4003751	0.000032	0.92598	1.287402	-0.000147	0.86920	1.913437	0.000545	0.84832
0.2902353	0.000028	0.94748	0.9512578	-0.000065	0.90619	1.422850	0.000362	0.89026
0.2088619	-0.000050	0.96275	0.6941875	0.000068	0.93296	1.043110	0.000215	0.92110
0.1496866	0.000016	0.97359	0.5020657	-0.000007	0.95221	0.7568559	-0.000125	0.94352
0.1068683	-0.000003	0.98128	0.3609449	-0.000025	0.96599	0.5455164	-0.000072	0.95970
0.07613516	0.0000011	0.98673	0.2584458	0.000021	0.97583	0.3912113	-0.000065	0.97130
0.05415350	0.0000049	0.99060	0.1844354	-0.000026	0.98283	0.2795657	-0.000058	0.97959
0.03842966	-0.0000381	0.99334	0.1316339	0.000243	0.98782	0.1992360	-0.000101	0.98550
0.02729816	-0.0000077	0.99528	1.421269	0.000067	0.85376	0.1418165	-0.000036	0.98971
0.6587463	-0.000048	0.87105	1.055902	0.000089	0.89496	0.1008070	-0.000013	0.99270
0.4867451	0.000016	0.90839	0.7733318	0.000081	0.92485	0.07161605	0.0000224	0.99482
0.3549307	0.000017	0.93497	0.5608474	0.000082	0.94639	0.05081747	0.0000096	0.99633
0.2565808	0.000060	0.95386	0.4038722	-0.000026	0.96183	0.03610165	0.0000607	0.99740
0.1842931	0.000004	0.96728	0.2894646	-0.000062	0.97286	1.841862	0.000230	0.85457
0.1318352	-0.000010	0.97680	0.2067826	-0.000063	0.98072	1.366705	0.000166	0.89489
0.09402423	-0.0000275	0.98356	0.1474082	-0.000020	0.98631	1.000393	0.000044	0.92448
0.06692514	-0.0000304	0.98835	1.359757	-0.000054	0.86091	0.7252507	-0.000068	0.94597
0.7796625	0.000021	0.84258	1.007713	0.000070	0.90017	0.5222763	-0.000096	0.96145
0.5822072	0.000037	0.88809	0.7367420	0.000055	0.92861	0.3744850	0.000059	0.97255
0.4273920	-0.000109	0.92052	0.5336335	0.000029	0.94909	0.2673715	-0.000110	0.98049
0.3104406	-0.000015	0.94360	0.3839925	-0.000019	0.96376	0.1906251	-0.000004	0.98614
0.2237983	0.000034	0.95999	0.2750979	-0.000010	0.97424	2.055568	-0.000399	0.83557
0.1604177	-0.000030	0.97163	0.1964623	-0.000001	0.98170	1.535674	-0.000031	0.88080
0.1146167	-0.000019	0.97989	0.1399417	-0.000045	0.98701	1.129498	0.000054	0.91418
0.08171551	0.0000163	0.98575	2.288660	0.000222	0.73872	0.8216163	-0.000004	0.93851
0.05816469	0.0000409	0.98990	1.778256	-0.000010	0.80998	0.5932180	0.000116	0.95609
0.04131815	0.0000179	0.99284	1.342709	0.000116	0.86290	0.4259017	0.000082	0.96871
0.02935869	0.0000380	0.99493	0.9942739	0.000086	0.90161	0.3045848	0.000037	0.97775
			0.7265668	0.000067	0.92966	0.2171083	-0.000115	0.98419
0.7947451	0.000055	0.88718	0.5260704	0.000021	0.94984	1.544284	-0.000258	0.88005
0.5835387	0.000048	0.91955	0.3784821	-0.000004	0.96430	1.136046	-0.000179	0.91364
0.4237752	0.000011	0.94274	0.2710513	-0.000054	0.97462	0.8266776	-0.000020	0.93812
0.3054696	0.000006	0.95931	0.1935612	-0.000021	0.98197	0.5968415	0.000000	0.95581
0.2190132	-0.000035	0.97111	0.1378516	-0.000072	0.98720	0.4286916	0.000150	0.96851
0.1565013	-0.000011	0.97950	0.09804344	-0.0000711	0.99092	0.3067239	0.000210	0.97760
0.1115259	-0.000020	0.98546	0.06972567	0.000065	0.99356	0.2188044	0.000169	0.98408
1.139260	0.000097	0.82902	0.04954069	0.0000375	0.99543			
0.8542892	-0.000004	0.87764	0.03521163	0.0000818	0.99676			
0.6293404	-0.000005	0.91270	0.02496167	0.0000415	0.99771			
0.4580812	-0.000023	0.93784						
0.3307199	-0.000013	0.95581						
0.2374314	0.000005	0.96862						
0.1697639	-0.000008	0.97773						
0.1210662	0.000005	0.98420						
1.538104	-0.000026	0.74888						
1.192117	0.000017	0.81933						
0.8972098	-0.000124	0.87062						
0.6626297	-0.000066	0.90765						
0.4831959	0.000004	0.93423						
0.3492649	0.000015	0.95323						
0.2509720	0.000051	0.96678						
0.1795262	0.000004	0.97643						
0.1280869	0.000023	0.98328						

conditions are obtained, which generally takes 2–3 h depending on how pronounced the Joule–Thomson effect is. The expansion valve is then closed, and the smaller volume evacuated. These expansions are continued until a low pressure is reached.

Reduction of Burnett Data

The result of the experiment is a series of decreasing pressures $P_1, P_2, \dots, P_i, \dots, P_l$ at constant temperature. It was shown by Burnett that the pressures from an expansion series must satisfy the relation

$$Z_i/Z_1 = (P_i/P_1)N^{i-1} \quad (1)$$

where N is the apparatus constant defined as

$$N = (V_a + V_b)/V_a \quad (2)$$

Several schemes for least-squares reduction of Burnett data are possible. In this work we have applied the method of Dalton et al. (4), but replaced the Newton–Raphson optimization routine by a Levenberg–Marquardt routine.

It is assumed that the experimental data are subject to random errors only and that the virial equation is able to accurately represent the isothermal compression factors as a function of density. Furthermore, the apparatus constant must be independent of pressure; otherwise an explicit expression for the pressure dependence must be known. In this work we have neglected the pressure dependence of the apparatus constant as discussed elsewhere by Mollerup and Angelo (3).

Several isothermal expansion series can be reduced simultaneously. If we denote the number of expansion series by M , and the number of expansions in series m by Im , then series

Table II. Parameters Determined by a Least-Squares Fit of the Experimental Pressures to the Objective Function

	parameter	std error
T = 260 K		
N	1.41152972	1.2 × 10 ⁻⁴
B ₁ , L/mol	-3.72145275 × 10 ⁻¹	3.2 × 10 ⁻³
B ₂ , (L/mol) ²	1.02311881 × 10 ⁻²	5.3 × 10 ⁻³
T = 285 K		
N	1.41165616	6.3 × 10 ⁻⁵
B ₁ , L/mol	-3.05221560 × 10 ⁻¹	7.9 × 10 ⁻⁴
B ₂ , (L/mol) ²	1.78882606 × 10 ⁻²	7.1 × 10 ⁻⁴
T = 315 K		
N	1.41103060	9.4 × 10 ⁻⁵
B ₁ , L/mol	-2.40756177 × 10 ⁻¹	1.0 × 10 ⁻³
B ₂ , (L/mol) ²	1.67925246 × 10 ⁻²	7.0 × 10 ⁻⁴
T = 340 K		
N	1.41123823	9.8 × 10 ⁻⁵
B ₁ , L/mol	-2.03868896 × 10 ⁻¹	1.0 × 10 ⁻³
B ₂ , (L/mol) ²	1.71894819 × 10 ⁻²	7.4 × 10 ⁻⁴

m contains the pressures $P^*_{1m}, P^*_{2m}, \dots, P^*_{im}$.

The densities are calculated by minimizing the residual sum of squares

$$F = \sum_{m=1}^M \left(\sum_{i=1}^I w_i (P^*_i - P_i)^2 \right)_m \quad (3)$$

where P^*_i is the experimental pressure and P_i the calculated pressure that must satisfy the Burnett relation eq 1, that is

$$P_{im} = (Z_{im}/Z_{1m}) P_{1m} N^{1-im} \quad (4)$$

The compression factor $Z = P/(\rho RT)$ is calculated from the virial equation of state

$$Z_{im} = 1 + \sum_{j=1}^J B_j \rho^j_{im} \quad (5)$$

The parameters in the least-squares data reduction procedure are the virial coefficients B_j and if desired the apparatus constant N and the initial pressures P_{1m} . If the initial pressures P_{1m} are not parameters to be estimated, they are assumed to be exact and thus $P_{1m} = P^*_{1m}$.

The weight factor w_i is a function of the observed pressure

$$w_i = 1/\sum (a_j + b_j P_i)^2 \quad (6)$$

where a_j and b_j are constants that depend on the barometer, the differential pressure indicators, and the dead-weight gauge used.

Results and Discussion

The application of sulfur hexafluoride to high-voltage CGI problems requires accurate measurements of the volumetric properties of sulfur hexafluoride (5). Table I shows the experimental pressures of sulfur hexafluoride and the calculated residuals and compression factors. The compression factors are calculated from eq 5 at pressures that satisfy the Burnett relation eq 1. They are not calculated at the experimental pressures P^* .

The parameters and the standard errors of the parameters are shown in Table II. The first parameter is the apparatus constant of approximately 1.4; the other parameters are virial coefficients. Table III shows a comparison of calculated second virial coefficients and the smooth numbers from Dy-

Table III. Second Virial Coefficients

T, K	second virial coeff, cm ³ /mol	
	this work	ref 6
260	-372.1 ± 3	-380 ± 8
285	-305.2 ± 0.8	-310.8 ± 6
315	-240.8 ± 1	-244.9 ± 5
340	-203.9 ± 1	-203.4 ± 5

Table IV. Interpolated Densities and Compression Factors

P, MPa	ρ , mol/L	Z	P, MPa	ρ , mol/L	Z
T = 260 K			T = 315 K		
0.1	0.047 083	0.9825	0.1	0.038 539	0.9908
0.2	0.095 935	0.9644	0.2	0.077 814	0.9814
0.3	0.146 76	0.9456	0.3	0.117 86	0.9719
0.4	0.199 81	0.9261	0.4	0.158 73	0.9622
0.5	0.255 40	0.9056	0.5	0.200 45	0.9524
0.6	0.313 91	0.8842	0.6	0.243 08	0.9425
0.7	0.375 84	0.8616	0.7	0.286 66	0.9324
0.8	0.441 84	0.8376	0.8	0.331 27	0.9221
T = 285 K			0.9	0.376 95	0.9116
0.1	0.042 758	0.9870	1.0	0.423 78	0.9010
0.2	0.086 685	0.9737	1.5	0.678 25	0.8444
0.3	0.131 87	0.9601	2.0	0.978 37	0.7805
0.4	0.178 42	0.9461	T = 340 K		
0.5	0.226 45	0.9318	0.1	0.035 633	0.9928
0.6	0.276 10	0.9171	0.2	0.071 794	0.9855
0.7	0.327 52	0.9020	0.3	0.108 50	0.9781
0.8	0.380 91	0.8863	0.4	0.145 78	0.9707
0.9	0.436 47	0.8702	0.5	0.183 64	0.9631
1.0	0.494 48	0.8535	0.6	0.222 12	0.9556
1.5	0.835 79	0.7574	0.7	0.261 23	0.9479
			0.8	0.301 00	0.9402
			0.9	0.341 46	0.9324
			1.0	0.382 63	0.9245
			1.5	0.600 39	0.8838
			2.0	0.841 67	0.8406
			2.5	1.113 4	0.7943

mond and Smith (6); the agreement is generally very good.

Table IV shows calculated densities and compression factors at integral pressures.

The experimental error of the measured compression factors is estimated to be 0.1% or less and to be mainly due to the uncertainty in the temperature and pressure. The absolute error in the temperature is less than 0.03 K at the triple point of water.

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