

Registry No.  $\text{Me}_2\text{SO}$ , 67-68-5;  $\text{LiCl}$ , 7447-41-8;  $\text{H}_2\text{O}$ , 7732-18-5.

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**Supplementary Material Available:** Table entitled "Viscosity of Water- $\text{Me}_2\text{SO}$  Mixtures at 25 °C"; it reports viscosities of 38 different mole fractions and is compiled from four separate works (1 page). Ordering information is given on any current masthead page.

## Densities and Apparent Molar Volumes of Aqueous Magnesium Sulfate and Sodium Sulfate to 473 K and 100 bar

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Densities of aqueous solutions of  $\text{MgSO}_4$  and  $\text{Na}_2\text{SO}_4$  have been measured up to 473 K, 100 bar, and 1 *m* ionic strength. The derived apparent molar volumes along with the available literature data on these substances have been analyzed by using the ion-interaction (Pitzer) approach.  $\bar{V}_2^0$  for  $\text{MgSO}_4$  is obtained indirectly from values for  $\text{MgCl}_2$ ,  $\text{Na}_2\text{SO}_4$ , and  $\text{NaCl}$ . The term in  $\beta^{(2)}$  represents adequately the ion-association effect in  $\text{MgSO}_4$  up to 473 K. The present results yield the pressure effect on various thermodynamic properties.

### Introduction

Recently we (1) completed a study which yielded a comprehensive equation for the thermodynamic properties of aqueous magnesium sulfate at low pressure based on our own heat capacity measurements and heats of dilution and osmotic coefficients from the literature. In the present study we present the measurements of densities of aqueous magnesium sulfate and sodium sulfate to 473 K, 100 bar, and 1 *m* ionic strength. In addition to the intrinsic interest of the volumetric properties, they also yield the pressure dependence of the thermodynamic properties already determined for low pressure.

Sodium sulfate behaves as a normal strong electrolyte and the present results are readily fitted to the same form of ion-interaction (Pitzer) equations used earlier by Rogers and Pitzer (2) for the heat capacity of aqueous  $\text{Na}_2\text{SO}_4$ . Magnesium sulfate, however, shows significant electrostatic ion-pairing. Our study of the heat capacity (1) showed that it was possible to use the ion-interaction equations up to 473 K provided that a special binary interaction term was included. Also, one cannot unambiguously extrapolate the  $\text{MgSO}_4$  data to zero molality to

obtain  $\bar{V}_2^0$ ; rather this quantity for  $\text{Mg}^{2+} + \text{SO}_4^{2-}$  is obtained indirectly from the values of  $\bar{V}_2^0$  for  $\text{MgCl}_2$ ,  $\text{Na}_2\text{SO}_4$ , and  $\text{NaCl}$ . Values for  $\text{NaCl}$  are taken from the recent comprehensive studies (3, 4) from this laboratory. For  $\text{Na}_2\text{SO}_4$ , we have generated equations based on Pitzer's equations using our measurements and literature values (5-15). For  $\text{MgCl}_2$ , we have fitted the recent measurements of Saluja (15) along with the other available literature data to the same equations. The details of the  $\text{MgCl}_2$  calculations will be presented separately (16) along with the similar results for  $\text{CaCl}_2$  and  $\text{SrCl}_2$ , and their heat capacities.

With the infinitely dilute values fixed, our measurements along with literature data (5, 7-10, 12, 13) on aqueous  $\text{MgSO}_4$  were satisfactorily fitted by the ion-interaction equations. This type of equation has been very successful for mixed electrolytes. Thus it is valuable to have the required parameters for  $\text{Na}_2\text{SO}_4$  and  $\text{MgSO}_4$  since these ions are prominent in sea water and in other natural or industrial waters.

### Experimental Method

A dilatometer has been used to measure the changes in the volume of a solution as a function of temperature. The volume at room temperature is known from the literature. The construction details and operation of the dilatometer are described elsewhere (17).

The pressure vessel, which contains about 5.75  $\text{cm}^3$  of solution, was modified in design and rebuilt to stand higher pressures. The pressure vessel is connected through a stainless-steel capillary to a mercury reservoir made of a coil of stainless-steel tubing which in turn is connected to a glass capillary column. The entire system can be pressurized with nitrogen gas.

Table I. Experimental Density of Aqueous Na<sub>2</sub>SO<sub>4</sub>

<i>T</i> /K	<i>P</i> /bar	<i>m</i> /mol·kg <sup>-1</sup>	$\rho$ /g·cm <sup>-3</sup>	$^*V$ /cm <sup>3</sup> ·mol <sup>-1</sup>	<i>T</i> /K	<i>P</i> /bar	<i>m</i> /mol·kg <sup>-1</sup>	$\rho$ /g·cm <sup>-3</sup>	$^*V$ /cm <sup>3</sup> ·mol <sup>-1</sup>
294.87	20.2	0.0578	1.006 08	14.49	293.99	20.5	0.2025	1.024 24	16.57
325.16	20.4	0.0578	0.995 03	18.25	324.39	20.7	0.2025	1.012 61	20.41
374.17	21.0	0.0578	0.965 53	17.42	374.14	21.0	0.2025	0.982 58	18.93
423.52	21.2	0.0578	0.924 87	6.17	423.72	21.6	0.2025	0.942 08	10.25
474.89	21.8	0.0578	0.870 76	-12.04	475.64	22.1	0.2025	0.888 30	-7.89
424.22	21.2	0.0578	0.924 27	4.92	424.10	21.5	0.2025	0.941 78	9.79
375.01	21.0	0.0578	0.964 97	16.58	374.60	21.0	0.2025	0.982 28	18.75
325.09	20.5	0.0578	0.995 07	18.08	324.05	20.8	0.2025	1.012 82	20.22
295.18	20.3	0.0578	1.006 04	14.08	294.33	20.4	0.2025	1.024 23	16.22
295.67	48.8	0.0578	1.007 17	14.96	294.12	48.4	0.2025	1.025 43	16.91
325.01	49.1	0.0578	0.996 34	18.33	324.17	48.8	0.2025	1.013 92	20.63
374.00	49.7	0.0578	0.967 05	16.62	374.08	49.4	0.2025	0.983 97	18.96
423.26	50.9	0.0578	0.926 81	6.20	423.44	50.2	0.2025	0.944 00	10.19
474.68	51.9	0.0578	0.873 35	-12.47	475.15	51.4	0.2025	0.891 12	-7.29
424.17	50.7	0.0578	0.926 01	4.73	424.23	50.4	0.2025	0.943 26	10.19
374.09	49.9	0.0578	0.967 03	15.97	374.29	50.0	0.2025	0.983 80	19.19
325.11	49.3	0.0578	0.996 31	18.28	324.42	49.1	0.2025	1.013 80	20.67
296.28	48.7	0.0578	1.007 02	14.83	294.35	48.4	0.2025	1.025 40	16.83
295.47	95.3	0.0578	1.009 28	15.42	294.61	96.0	0.2025	1.027 30	17.84
325.08	96.6	0.0578	0.998 33	18.99	324.34	96.9	0.2025	1.015 82	21.24
374.04	97.4	0.0578	0.969 28	15.71	373.99	98.4	0.2025	0.986 62	19.65
423.16	99.5	0.0578	0.929 67	4.89	423.19	100.7	0.2025	0.946 95	11.06
474.80	101.4	0.0578	0.877 03	-14.94	475.32	102.3	0.2025	0.894 62	-5.96
424.32	99.3	0.0578	0.928 66	3.42	424.11	99.9	0.2025	0.946 14	10.49
374.71	98.1	0.0578	0.968 81	16.27	374.12	98.8	0.2025	0.986 17	19.47
324.67	96.4	0.0578	0.998 54	18.33	325.13	96.6	0.2025	1.015 46	21.10
295.09	95.5	0.0578	1.009 39	15.43	294.52	95.9	0.2025	1.027 39	17.50
295.12	20.1	0.1085	1.012 39	15.25	295.12	20.4	0.3303	1.039 30	18.17
324.94	20.4	0.1085	1.001 23	19.15	325.27	20.8	0.3303	1.027 05	21.55
374.84	20.6	0.1085	0.971 10	17.44	374.45	21.2	0.3303	0.997 02	20.12
423.60	21.3	0.1085	0.930 94	8.35	423.81	21.6	0.3303	0.956 93	12.06
475.07	21.8	0.1085	0.877 10	-10.56	475.81	22.3	0.3303	0.903 77	-4.83
423.54	21.3	0.1085	0.931 04	7.93	424.67	21.8	0.3303	0.956 15	11.92
374.87	20.7	0.1085	0.971 10	17.17	374.89	21.3	0.3303	0.996 71	20.06
324.73	20.1	0.1085	1.001 38	18.51	325.18	20.8	0.3303	1.027 09	21.55
296.35	20.0	0.1085	1.012 07	15.40	294.99	20.3	0.3303	1.039 36	18.09
294.85	47.6	0.1085	1.013 66	15.63	295.08	47.7	0.3303	1.040 41	18.69
324.86	48.3	0.1085	1.002 45	19.53	325.33	48.2	0.3303	1.028 10	21.98
374.48	49.3	0.1085	0.972 65	18.16	374.46	49.2	0.3303	0.998 22	20.58
423.53	50.1	0.1085	0.932 65	8.48	423.69	50.3	0.3303	0.958 49	12.88
475.14	51.2	0.1085	0.879 29	-10.41	475.76	51.1	0.3303	0.905 90	-3.83
424.53	50.4	0.1085	0.931 75	8.06	424.02	49.5	0.3303	0.958 28	12.38
375.00	49.4	0.1085	0.972 30	17.90	375.33	48.7	0.3303	0.997 65	20.26
324.75	47.9	0.1085	1.002 52	19.21	324.61	47.9	0.3303	1.028 48	21.84
295.25	47.5	0.1085	1.013 56	15.62	294.86	47.1	0.3303	1.040 53	18.42
295.18	95.5	0.1085	1.015 67	16.34	296.26	95.6	0.3303	1.042 03	19.55
324.65	96.7	0.1085	1.004 63	19.66	324.65	96.5	0.3303	1.030 41	22.45
374.29	98.1	0.1085	0.975 14	17.40	374.55	98.0	0.3303	1.000 38	20.89
423.89	100.5	0.1085	0.935 19	7.94	423.97	100.0	0.3303	0.961 03	13.10
475.25	102.5	0.1085	0.883 13	-11.26	475.78	102.4	0.3303	0.909 59	-2.65
423.93	100.3	0.1085	0.935 18	7.47	425.02	100.0	0.3303	0.960 11	12.80
375.37	98.7	0.1085	0.974 37	17.58	375.72	98.7	0.3303	0.999 57	20.82
325.28	96.6	0.1085	1.004 36	19.46	325.42	97.0	0.3303	1.030 05	22.49
294.87	95.7	0.1085	1.015 76	16.32	296.28	95.6	0.3303	1.042 05	19.50

The mercury reservoir and most of the capillary connecting tube are thermostated with a water bath and a circulation pump at 30 °C. The solution cell is surrounded by a large aluminum block and is placed in a high-temperature bath. The aluminum block serves as a heat sink to reduce temperature fluctuations. Thus the temperature of the sample cell is stable to  $\pm 0.01$  K over a period of 30 min.

The density of the solution is determined by measuring the change in volume of the solution as the temperature is raised. Special care was taken to fill the sample cell with solution without introducing any air bubble. The cell was filled and emptied 6–7 times to eliminate any contamination by any old solution in the cell. The mercury and thereby also the solution was then pressurized with nitrogen gas. Any air bubbles in the solution could be detected by large displacement in the level of mercury column. The sample cell was then enclosed in the aluminum jacket and the solution was equilibrated at room temperature. The pressure of the system and the level of mercury were monitored overnight to check for any leakage.

The temperature of the solution was then raised in increments of 25 or 50 °C and the heights of mercury column were measured to  $\pm 0.01$  mm at each step with a Wilde cathetometer. The level of mercury menisci were noted at different temperatures until 200 °C. Then the temperature of the cell was decreased at the same intervals and the level of menisci noted until the initial temperature was reached. This sequence of measurements is followed in Tables I and II. The temperature of the solution cell was measured by using a 25- $\Omega$  platinum resistance thermometer using a G-2 Mueller Bridge. A bourdon-tube pressure gauge was used to determine the pressure of the system to  $\pm 0.1$  bar. The temperatures of the water bath and the room were measured with mercury-in-glass thermometers.

#### Experimental Results

The density of the solution can be related to the change in the height of the mercury column as follows. Let  $h$  be the

Table II. Experimental Density of Aqueous MgSO<sub>4</sub>

T/K	P/bar	m/mol·kg <sup>-1</sup>	ρ/g·cm <sup>-3</sup>	φV/cm <sup>3</sup> ·mol <sup>-1</sup>
323.90	22.3	0.0723	0.99720	0.20
374.18	23.2	0.0723	0.96713	-1.45
422.84	24.1	0.0723	0.92695	-8.35
474.80	24.6	0.0723	0.87214	-21.01
423.58	24.0	0.0723	0.92635	-10.07
374.65	23.0	0.0723	0.96689	-3.01
324.48	22.3	0.0723	0.99699	-0.53
324.24	49.1	0.0723	0.99819	0.58
373.64	49.8	0.0723	0.96882	-2.24
423.17	50.6	0.0723	0.92830	-10.82
474.60	51.8	0.0723	0.87455	-22.53
423.08	50.7	0.0723	0.92837	-10.63
374.62	49.7	0.0723	0.96817	-3.04
324.17	49.0	0.0723	0.99822	0.52
324.26	98.3	0.0723	1.00020	2.24
374.49	100.0	0.0723	0.97043	-0.47
423.30	102.1	0.0723	0.93086	-7.77
475.57	104.1	0.0723	0.87714	-18.94
424.20	101.9	0.0723	0.93009	-8.92
374.94	100.0	0.0723	0.97018	-1.55
324.56	98.3	0.0723	1.00013	1.28
294.21	21.9	0.1693	1.01927	0.26
324.63	22.5	0.1693	1.00814	1.68
374.63	22.6	0.1693	0.97815	-1.79
423.58	23.3	0.1693	0.93776	-9.27
475.77	24.2	0.1693	0.88248	-20.39
424.24	24.0	0.1693	0.93712	-8.90
374.80	23.5	0.1693	0.97802	-1.47
324.91	23.1	0.1693	1.00799	2.00
294.23	23.1	0.1693	1.01924	0.70
293.84	47.6	0.1693	1.02043	0.90
324.76	47.5	0.1693	1.00917	1.81
374.41	48.2	0.1693	0.97949	-1.56
423.90	48.8	0.1693	0.93897	-9.56
475.55	50.0	0.1693	0.88477	-20.35
294.23	94.9	0.1693	1.02234	1.82
325.27	95.6	0.1693	1.01083	2.99
374.17	97.4	0.1693	0.98169	0.31
423.34	99.1	0.1693	0.94195	-6.84
475.38	101.5	0.1693	0.88850	-17.05
294.72	21.1	0.2503	1.02849	1.59
325.06	21.6	0.2503	1.01709	2.90
374.24	22.0	0.2503	0.98778	-1.28
295.01	47.8	0.2503	1.02954	2.06
324.20	48.1	0.2503	1.01863	3.10
374.16	49.2	0.2503	0.98887	-0.15
423.81	49.9	0.2503	0.94833	-7.57
475.10	51.0	0.2503	0.89479	-18.58
424.15	49.9	0.2503	0.94802	-7.65
375.00	48.9	0.2503	0.98832	-0.48
324.57	48.0	0.2503	1.01849	2.94
294.88	47.6	0.2503	1.02961	1.85
294.90	95.2	0.2503	1.03156	2.78
324.57	96.0	0.2503	1.02041	3.70
374.48	97.7	0.2503	0.99080	0.46
423.65	99.5	0.2503	0.95114	-6.73
475.10	101.5	0.2503	0.89849	-17.47
423.74	99.3	0.2503	0.95109	-6.92
375.02	97.9	0.2503	0.99043	0.37
324.74	96.3	0.2503	1.02033	3.80
295.03	95.4	0.2503	1.03156	2.71

experimental temperature,  $t_R$  the reference temperature, and  $t_B$  the water-bath temperature in degrees centigrade. At the experimental temperature, the density of the solution is given by

$$\rho_t = g_t/V_t \quad (1)$$

where  $g_t$  is the mass of the solution and  $V_t$  is the volume of the solution. The mass of the solution in the sample cell at any temperature is equal to the mass in the cell at the reference temperature,  $t_R$ , less the amount that has expanded from the cell. The mass of the solution expanded is given by

$$\Delta m = \Delta h A \rho_{t_B} \quad (2)$$

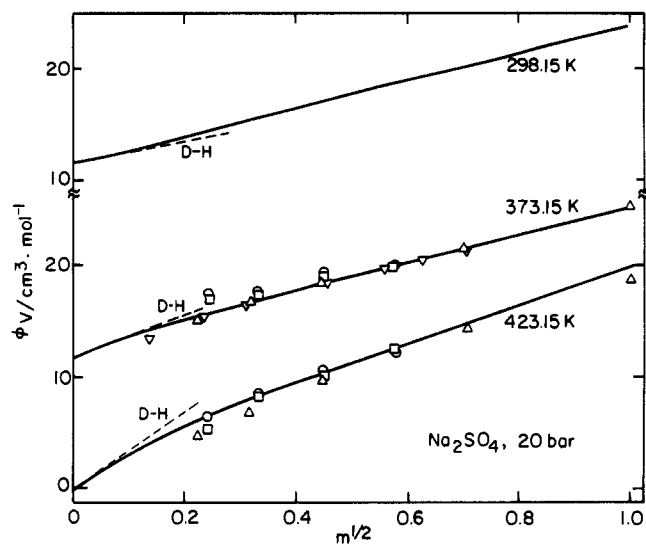


Figure 1. The apparent molar volume,  $\phi V$ , for Na<sub>2</sub>SO<sub>4</sub>: O, this study at 20 bar, rising temperature; □, this study at 20 bar, falling temperature; Δ, Ellis (7) at 20 bar; ▽, Sakuja (15) at 6 bar. Solid curves are calculated from the present fit (eq 19) at 20 bar; dashed lines are the Debye-Hückel slopes taken from Bradley and Pitzer (22).

Table III. Parameters for Eq 19 and 20 for Apparent Molar Volume of Aqueous Na<sub>2</sub>SO<sub>4</sub> and MgSO<sub>4</sub>

Na <sub>2</sub> SO <sub>4</sub> <sup>a</sup>	MgSO <sub>4</sub> <sup>b</sup>
$P_1 = -5.11686 \times 10^5$	$q_1 = -7.97433 \times 10^5$
$P_2 = 5.43000 \times 10^3$	$q_2 = 8.94735 \times 10^3$
$P_3 = -2.18218 \times 10^1$	$q_3 = -3.78219 \times 10^1$
$P_4 = 3.99915 \times 10^{-2}$	$q_4 = 7.15424 \times 10^{-2}$
$P_5 = -2.83513 \times 10^{-5}$	$q_5 = -5.14523 \times 10^{-5}$
$P_6 = 1.79274 \times 10^{-2}$	$q_6 = 7.42784$
$P_7 = 1.28610$	$q_7 = -6.12133 \times 10^{-2}$
$P_8 = -1.06978 \times 10^{-2}$	$q_8 = 1.77378 \times 10^{-4}$
$P_9 = 2.96022 \times 10^{-5}$	$q_9 = -1.82737 \times 10^{-7}$
$P_{10} = -2.70365 \times 10^{-8}$	$q_{10} = 2.06777 \times 10^3$
$P_{11} = 4.59690 \times 10^{-1}$	$q_{11} = -2.39043 \times 10^1$
$P_{12} = -1.31097 \times 10^{-3}$	$q_{12} = 1.03110 \times 10^{-1}$
	$q_{13} = -1.96675 \times 10^{-4}$
	$q_{14} = 1.40527 \times 10^{-7}$
	$q_{15} = 1.18481 \times 10^{-4}$

<sup>a</sup> Range: 0–200 °C; 0–1 m, to 20 bar; 0–0.33 m, to 100 bar.  
<sup>b</sup> Range: 0–200 °C; 0–0.25 m, to 100 bar.

where  $\Delta h$  is the total change in the height of the mercury column between temperatures  $t_R$  and  $t$ ,  $A$  is the cross-sectional area of the mercury column, and  $\rho_{t_B}$  is the density of the solution at the temperature of the water bath. Then the density of the solution is given by

$$\rho_t = (1/V_t)[\rho_{t_R} V_{t_R} - \Delta h A \rho_{t_B}] \quad (3)$$

We assume the temperature dependence of the volume of the cell as

$$V_t = V_0 e^{\alpha t} \quad (4)$$

where  $V_0$  is the volume of the cell at 0 °C and  $\alpha$  is the unknown expansivity of the cell. Substituting eq 4 into eq 3 we get

$$\rho_t = \rho_{t_R} e^{-\alpha(t-t_R)} - (A/V_0)\Delta h \rho_{t_B} e^{-\alpha t}$$

$$= \rho_{t_R}[1 - \alpha(t-t_R)] - \Delta h(A/V_0)\rho_{t_B}[1 - \alpha t]$$

or

$$\rho_t = \rho_{t_R} - \frac{A}{V_0}\Delta h \rho_{t_B} + \alpha \rho_{t_R}(t_R - t) + \alpha \frac{A}{V_0}t \Delta h \rho_{t_B} \quad (5)$$

Equation 5 was used to obtain the values of the density of the

Table IV. Apparent Molar Volumes ( $\text{cm}^3 \cdot \text{mol}^{-1}$ ) of Aqueous  $\text{Na}_2\text{SO}_4$ 

		A										
		$m/\text{mol} \cdot \text{kg}^{-1}$										
$t/^\circ\text{C}$	$P/\text{bar}$	0.00	0.01	0.05	0.10	0.20	0.30	0.40	0.50	0.60	0.80	1.00
0.0	1.0	2.12	3.17	4.94	6.42	8.61	10.32	11.78	13.07	14.26	16.42	18.39
10.0	1.0	6.72	7.79	9.46	10.82	12.77	14.27	15.53	16.64	17.64	19.45	21.07
20.0	1.0	9.99	11.08	12.68	13.93	15.71	17.04	18.15	19.11	19.98	21.52	22.90
25.0	1.0	11.22	12.32	13.89	15.11	16.81	18.08	19.12	20.03	20.85	22.29	23.57
30.0	1.0	12.21	13.33	14.88	16.06	17.70	18.91	19.91	20.77	21.55	22.91	24.11
40.0	1.0	13.59	14.76	16.28	17.41	18.95	20.09	21.01	21.81	22.53	23.78	24.88
50.0	1.0	14.30	15.54	17.06	18.16	19.65	20.73	21.62	22.39	23.08	24.28	25.34
60.0	1.0	14.48	15.80	17.35	18.44	19.90	20.97	21.86	22.62	23.31	24.52	25.59
70.0	1.0	14.22	15.65	17.24	18.34	19.81	20.90	21.79	22.58	23.28	24.54	25.67
80.0	1.0	13.59	15.14	16.80	17.93	19.44	20.56	21.49	22.31	23.06	24.40	25.61
90.0	1.0	12.62	14.32	16.08	17.26	18.83	20.00	20.99	21.86	22.66	24.10	25.41
100.0	1.0	11.34	13.21	15.09	16.34	17.99	19.23	20.28	21.22	22.08	23.64	25.07
110.0	1.5	9.74	11.81	13.85	15.18	16.94	18.27	19.39	20.39	21.32	23.01	24.56
120.0	2.0	7.80	10.12	12.35	13.78	15.67	17.08	18.28	19.36	20.35	22.16	23.82
130.0	2.7	5.50	8.09	10.55	12.11	14.13	15.65	16.93	18.08	19.13	21.06	22.83
140.0	3.6	2.77	5.70	8.42	10.13	12.32	13.94	15.30	16.52	17.63	19.66	21.50
150.0	4.8	-0.43	2.88	5.93	7.81	10.19	11.92	13.37	14.64	15.81	17.91	19.81
160.0	6.2	-4.18	-0.41	3.02	5.11	7.70	9.55	11.08	12.41	13.62	15.77	17.68
170.0	7.9	-8.56	-4.24	-0.35	1.98	4.82	6.80	8.41	9.79	11.03	13.19	15.08
180.0	10.0	-13.64	-8.68	-4.24	-1.62	1.51	3.64	5.33	6.76	8.01	10.16	11.98
190.0	12.6	-19.53	-13.80	-8.70	-5.73	-2.25	0.06	1.83	3.30	4.56	6.66	8.37
200.0	15.6	-26.32	-19.67	-13.77	-10.36	-6.46	-3.95	-2.07	-0.56	0.70	2.71	4.28

## B

		B							
		$m/\text{mol} \cdot \text{kg}^{-1}$							
$t/^\circ\text{C}$	$P/\text{bar}$	0.00	0.01	0.05	0.10	0.20	0.30	0.40	
0.0	100.0	3.89	4.94	6.69	8.17	10.35	12.06	13.51	
10.0	100.0	8.50	9.55	11.21	12.55	14.50	15.99	17.24	
20.0	100.0	11.77	12.84	14.42	15.66	17.43	18.75	19.85	
25.0	100.0	12.99	14.08	15.63	16.83	18.52	19.78	20.82	
30.0	100.0	13.98	15.09	16.61	17.78	19.40	20.61	21.60	
40.0	100.0	15.36	16.51	18.01	19.12	20.65	21.77	22.69	
50.0	100.0	16.07	17.29	18.78	19.86	21.32	22.40	23.28	
60.0	100.0	16.25	17.55	19.05	20.12	21.56	22.62	23.48	
70.0	100.0	16.00	17.38	18.93	20.01	21.45	22.51	23.39	
80.0	100.0	15.36	16.87	18.48	19.58	21.05	22.15	23.06	
90.0	100.0	14.40	16.04	17.74	18.88	20.41	21.55	22.51	
100.0	100.0	13.11	14.92	16.73	17.93	19.53	20.74	21.76	
110.0	100.0	11.51	13.50	15.46	16.74	18.43	19.71	20.80	
120.0	100.0	9.56	11.78	13.91	15.28	17.09	18.45	19.61	
130.0	100.0	7.24	9.72	12.06	13.54	15.48	16.93	18.16	
140.0	100.0	4.50	7.29	9.87	11.48	13.56	15.11	16.41	
150.0	100.0	1.27	4.42	7.30	9.07	11.31	12.95	14.33	
160.0	100.0	-2.50	1.07	4.29	6.24	8.67	10.42	11.86	
170.0	100.0	-6.91	-2.84	0.80	2.97	5.61	7.46	8.96	
180.0	100.0	-12.03	-7.37	-3.24	-0.81	2.08	4.05	5.61	
190.0	100.0	-17.96	-12.60	-7.87	-5.14	-1.95	0.16	1.78	
200.0	100.0	-24.80	-18.60	-13.15	-10.04	-6.49	-4.23	-2.53	

solution from the height of the mercury column.

Corrections in the mercury column due to the variations in water-bath temperature, room temperature, and hydrostatic head pressure were small and were applied.

Calibration runs were carried out using pure water in the solution cell. The calibration constants, i.e., the expansivity of the sample cell,  $\alpha$ , and the ratio of the cross-sectional area of the column to the volume of the cell at 0 °C,  $A/V_0$ , were evaluated as a function of pressure up to 100 bar by a least-squares regression of eq 5. Observed values of  $\Delta h$  from 25 to 200 °C were combined with the known values of the densities of water (18).

Observed values of  $\Delta h$  were used directly in eq 5 in combination with the known calibration constants to obtain densities of the solution. Measurements were made on solutions of magnesium sulfate and sodium sulfate up to 1  $m$  ionic strength at approximately 20, 50, and 100 bar pressures up to 200 °C. The experimental values of densities obtained for aqueous sodium sulfate and magnesium sulfate are given in Tables I and II.

The values of reference densities of sodium sulfate and magnesium sulfate solutions at reference temperature and at the water bath temperature are required in eq 5. These were obtained from literature (10). The values of the apparent molar volume,  $^{\phi}V$ , have been calculated from the density data as

$$^{\phi}V = \frac{1000 + mM_2}{m} \left( \frac{1}{\rho} \right) - \left( \frac{1000}{m} \right) \frac{1}{\rho_0} \quad (6)$$

where  $M_2$  is the molecular weight of the salt and  $\rho_0$  is the density of pure water.

The uncertainty in the experimentally measured density increases with temperature, being  $\pm 4 \times 10^{-5} \text{ g} \cdot \text{cm}^{-3}$  at 50 °C and  $\pm 1 \times 10^{-4} \text{ g} \cdot \text{cm}^{-3}$  at 200 °C. This results in an uncertainty in the apparent molar volume of 0.1  $\text{cm}^3 \cdot \text{mol}^{-1}$  for the measurement at 0.33  $\text{mol} \cdot \text{kg}^{-1}$  to 0.7  $\text{cm}^3 \cdot \text{mol}^{-1}$  for the dilute solutions below 0.06  $\text{mol} \cdot \text{kg}^{-1}$  at 50 °C. The uncertainty increases with temperature and the corresponding values for the same solutions at 200 °C are 0.4 and 2.2  $\text{cm}^3 \cdot \text{mol}^{-1}$ .

Table V. Apparent Molar Volumes ( $\text{cm}^3 \cdot \text{mol}^{-1}$ ) of Aqueous  $\text{MgSO}_4$ 

		A							
$t/^\circ\text{C}$	$P/\text{bar}$	$m/\text{mol}\cdot\text{kg}^{-1}$							
		0.00	0.01	0.02	0.05	0.10	0.15	0.20	0.25
0.0	1.0	-13.85	-10.26	-9.11	-7.39	-5.79	-4.64	-3.72	-2.94
10.0	1.0	-10.41	-7.18	-6.14	-4.57	-3.11	-2.08	-1.25	-0.55
20.0	1.0	-8.44	-5.26	-4.25	-2.77	-1.42	-0.48	0.27	0.89
25.0	1.0	-7.87	-4.63	-3.62	-2.16	-0.86	0.04	0.76	1.35
30.0	1.0	-7.54	-4.19	-3.16	-1.71	-0.44	0.42	1.10	1.67
40.0	1.0	-7.43	-3.74	-2.65	-1.19	0.02	0.82	1.45	1.97
50.0	1.0	-7.87	-3.75	-2.57	-1.06	0.11	0.86	1.44	1.92
60.0	1.0	-8.72	-4.09	-2.81	-1.23	-0.09	0.63	1.17	1.61
70.0	1.0	-9.85	-4.69	-3.29	-1.64	-0.50	0.18	0.69	1.10
80.0	1.0	-11.22	-5.50	-3.98	-2.23	-1.10	-0.44	0.04	0.42
90.0	1.0	-12.81	-6.51	-4.85	-3.00	-1.86	-1.22	-0.77	-0.41
100.0	1.0	-14.64	-7.71	-5.91	-3.95	-2.78	-2.16	-1.73	-1.39
110.0	1.5	-16.75	-9.12	-7.16	-5.06	-3.87	-3.26	-2.85	-2.53
120.0	2.0	-19.21	-10.77	-8.62	-6.36	-5.12	-4.53	-4.14	-3.84
130.0	2.7	-22.13	-12.70	-10.32	-7.85	-6.56	-5.97	-5.59	-5.32
140.0	3.6	-25.62	-14.95	-12.28	-9.55	-8.18	-7.58	-7.22	-6.96
150.0	4.8	-29.83	-17.56	-14.51	-11.45	-9.97	-9.36	-9.01	-8.76
160.0	6.2	-34.89	-20.58	-17.05	-13.54	-11.92	-11.28	-10.93	-10.69
170.0	7.9	-40.98	-24.04	-19.88	-15.81	-13.99	-13.30	-12.93	-12.69
180.0	10.0	-48.27	-27.98	-23.01	-18.22	-16.12	-15.35	-14.94	-14.68
190.0	12.6	-56.96	-32.41	-26.42	-20.69	-18.23	-17.33	-16.86	-16.56
200.0	15.6	-67.25	-37.36	-30.06	-23.14	-20.19	-19.11	-18.53	-18.15

B

		B							
$t/^\circ\text{C}$	$P/\text{bar}$	$m/\text{mol}\cdot\text{kg}^{-1}$							
		0.00	0.01	0.02	0.05	0.10	0.15	0.20	0.25
0.0	100.0	-13.85	-9.00	-7.55	-5.63	-3.99	-2.84	-1.92	-1.14
10.0	100.0	-10.41	-5.87	-4.53	-2.75	-1.26	-0.22	0.60	1.29
20.0	100.0	-8.44	-3.91	-2.59	-0.89	0.49	1.43	2.17	2.79
25.0	100.0	-7.87	-3.26	-1.94	-0.26	1.08	1.98	2.68	3.27
30.0	100.0	-7.54	-2.80	-1.45	0.22	1.52	2.38	3.05	3.61
40.0	100.0	-7.43	-2.31	-0.89	0.80	2.04	2.83	3.45	3.96
50.0	100.0	-7.87	-2.28	-0.76	0.98	2.17	2.92	3.48	3.95
60.0	100.0	-8.72	-2.58	-0.95	0.86	2.03	2.73	3.25	3.68
70.0	100.0	-9.85	-3.14	-1.39	0.50	1.66	2.32	2.81	3.21
80.0	100.0	-11.22	-3.92	-2.03	-0.05	1.10	1.73	2.19	2.56
90.0	100.0	-12.81	-4.89	-2.87	-0.79	0.37	0.98	1.40	1.74
100.0	100.0	-14.64	-6.06	-3.90	-1.70	-0.53	0.05	0.45	0.76
110.0	100.0	-16.75	-7.45	-5.13	-2.80	-1.61	-1.05	-0.68	-0.40
120.0	100.0	-19.21	-9.09	-6.58	-4.09	-2.88	-2.34	-2.00	-1.74
130.0	100.0	-22.13	-11.02	-8.27	-5.60	4.35	-3.82	-3.51	-3.28
140.0	100.0	-25.62	-13.28	-10.24	-7.33	-6.02	-5.51	-5.22	-5.02
150.0	100.0	-29.83	-15.91	-12.51	-9.28	-7.90	-7.39	-7.12	-6.94
160.0	100.0	-34.89	-18.97	-15.10	-11.46	-9.97	-9.45	-9.20	-9.04
170.0	100.0	-40.98	-22.49	-18.01	-13.85	-12.20	-11.66	-11.41	-11.27
180.0	100.0	-48.27	-26.51	-21.25	-16.42	-14.54	-13.95	-13.70	-13.56
190.0	100.0	-56.96	-31.05	-24.81	-19.11	-16.92	-16.25	-15.97	-15.82
200.0	100.0	-67.25	-36.14	-28.64	-21.84	-19.24	-18.43	-18.08	-17.89

### Review of Equations

We have the dual objectives of representation in simple equations of the present measurements and the expression of the resulting pressure dependency of various thermodynamic properties. Both of these objectives are met by use of the same form of equations used for the other thermodynamic properties (1). These are the virial-series-plus-Debye-Hückel (Pitzer) equations as developed by Pitzer (19, 20) and Pitzer and Mayorga (21). The general set of equations for various properties are given in several papers (1, 2, 4, 20) and will not be duplicated here. The excess Gibbs energy is related to the total Gibbs energy in the molality system by the equation

$$G = n_1 G^\circ_1 + n_2 \bar{G}^\circ_2 + G^{\text{EX}} + n_2 \nu RT (\ln m - 1) \quad (7)$$

where  $n_1$  and  $n_2$  are the numbers of moles of water and salt of standard molar Gibbs energies  $G^\circ_1$  and  $\bar{G}^\circ_2$ . The composition dependence of  $G^{\text{EX}}$  is then given by the following equation

$$1000G^{\text{EX}}/n_1 M_w RT = -A_\phi(4I/b) \ln(1 + bI^{1/2}) + 2\nu_M \nu_X [m^2 B_{\text{MX}}(I) + m^3 \nu_M Z_M C_{\text{MX}}] \quad (8a)$$

$$B_{\text{MX}}(I) = \beta_{\text{MX}}^{(0)} + \beta_{\text{MX}}^{(1)} g(\alpha_1 I^{1/2}) + \beta_{\text{MX}}^{(2)} g(\alpha_2 I^{1/2}) \quad (8b)$$

$$g(x) = 2[1 - (1 + x) \exp(-x)]/x^2 \quad (8c)$$

where the electrolyte MX contains  $\nu_M$  and  $\nu_X$  ions of charge  $Z_M$  and  $Z_X$ , and  $\nu = \nu_M + \nu_X$ . In some papers the quantity  $C^\phi = 2|Z_M Z_X|^{1/2} C$  is used instead of  $C$ .  $I$  is the ionic strength

$$I = \frac{1}{2} \sum_i m_i Z_i^2 \quad (9)$$

and  $A_\phi$  is the Debye-Hückel slope for the osmotic coefficient for which values are given by Bradley and Pitzer (22):

$$A_\phi = (1/3)(2\pi N_0 d_w / 1000)^{1/2} (e^2 / DK T)^{3/2} \quad (10)$$

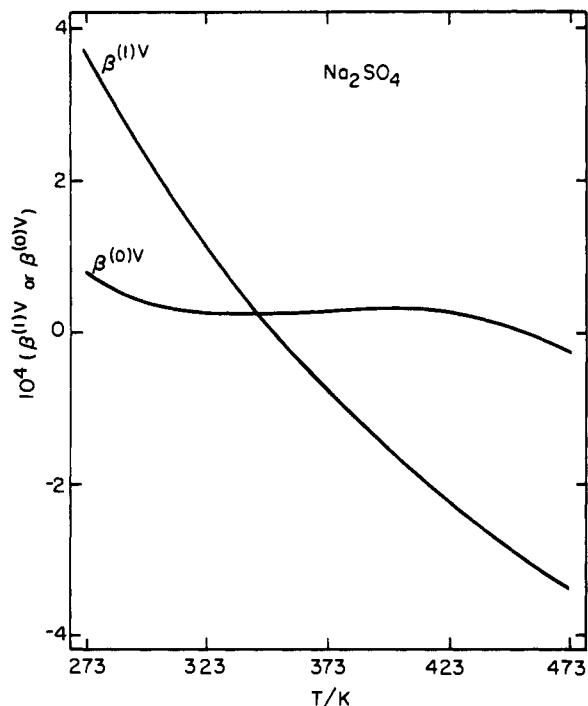


Figure 2. Temperature dependence of the fitting parameters  $\beta^{(0)V}$  and  $\beta^{(1)V}$  for  $\text{Na}_2\text{SO}_4$ .

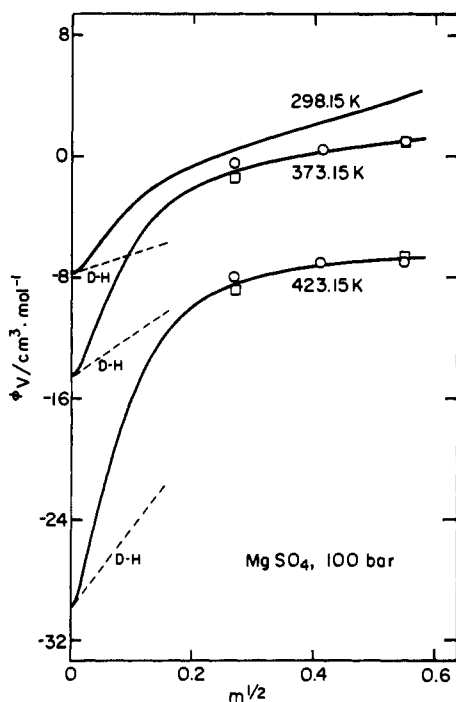


Figure 3. The apparent molar volume,  $\phi_V$ , for  $\text{MgSO}_4$  at 100 bar: O, rising temperature;  $\square$ , falling temperature. Solid curves are calculated from the fit (eq 20); dashed lines are the Debye-Hückel slopes.

Here  $d_w$  is the density and  $D$  the dielectric constant of pure water.

The parameter  $b$  has a fixed value of  $1.2 \text{ kg}^{1/2} \cdot \text{mol}^{-1/2}$  and  $\alpha_1$  and  $\alpha_2$  have values of 1.4 and  $12.0 \text{ kg}^{1/2} \cdot \text{mol}^{-1/2}$  for a 2:2 electrolyte. For a 2:1 or a 1:1 electrolyte  $\alpha_1 = 2.0 \text{ kg}^{1/2} \cdot \text{mol}^{-1/2}$ , and the term in  $\beta^{(2)}$  is omitted. The only quantities fitted for a particular electrolyte are those with the MX subscript in eq 8.

The total volume of the solution,  $V$ , is given by the pressure derivative of the total Gibbs energy of the solution

$$V = (\partial G / \partial P)_T \quad (11)$$

$$V = n_1 V_1^0 + n_2 \bar{V}_2^0 + (\partial G^{\text{EX}} / \partial P)_T \quad (12)$$

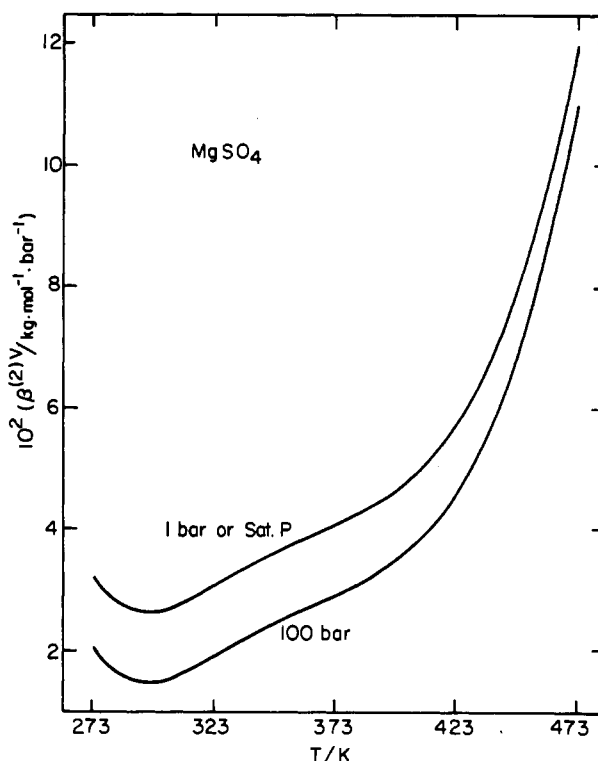


Figure 4. Temperature dependence of the fitting parameter  $\beta^{(2)V}$  for  $\text{MgSO}_4$ .

The apparent molar volume is defined by

$$\phi_V = (V - n_1 V_1^0) / n_2 \quad (13)$$

By using eq 12 in the above equation

$$\phi_V = \bar{V}_2^0 + (1/n_2) (\partial G^{\text{EX}} / \partial P)_T \quad (14)$$

$\bar{V}_2^0$  is the partial molar volume of the solute at infinite dilution. Substitution of eq 8 into the above equation yields the parametric form of the equation for the apparent molar volume

$$\phi_V = \bar{V}_2^0 + \nu |Z_M Z_X| A_\nu h(I) + 2\nu_M \nu_X RT [m B_{\text{MX}}^V(I) + m^2 (\nu_M Z_M) C_{\text{MX}}^V] \quad (15a)$$

where

$$h(I) = \ln(1 + bI^{1/2}) / 2b \quad (15b)$$

$$B_{\text{MX}}^V(I) = (\partial B_{\text{MX}}(I) / \partial P)_{T,I} \quad (16a)$$

$$C_{\text{MX}}^V = (\partial C_{\text{MX}} / \partial P)_T \quad (16b)$$

Also

$$A_\nu = -4RT(\partial A_\phi / \partial P)_T \quad (17)$$

There is some theoretical basis for making the parameter  $\alpha_2$  proportional to the Debye-Hückel slope  $A_\phi$  rather than a constant. This is discussed in detail elsewhere (1, 21). The alternate basis yields an improved fit to some measurements at very low molality and improved agreement with treatments based on an ion-pairing reaction equilibrium, but there is no difference above molality  $0.1 \text{ mol} \cdot \text{kg}^{-1}$  for a 2-2 electrolyte. Thus, for most practical purposes the assumption  $\alpha_2 = 12 \text{ kg}^{1/2} \cdot \text{mol}^{-1/2}$  is adequate and avoids many complications. The resulting  $\beta^{(2)V}$  from this paper gives the correct pressure dependency for  $\beta^{(2)}$  on the constant  $\alpha_2$  basis in ref 1.

The equations on the basis  $\alpha_2$  proportional to  $A_\phi$  are complicated by the pressure and temperature dependencies of  $A_\phi$ . While this can all be handled without real difficulty, the advantages are so limited that it does not seem worthwhile to treat

the present data on that basis. The present  $\beta^{(2)V}$  values will give the correct pressure dependency of the quantity  $\beta^{(2)*}$ , defined for use with variable  $\alpha_2$ , through the equation

$$(\partial\beta^{(2)*}/\partial P)_T = \partial/\partial P [\beta^{(2)}A_\phi^2(30.65/12 \text{ kg}^{1/2}\cdot\text{mol}^{-1/2})^2] = [\beta^{(2)V}A_\phi^2 - \beta^{(2)}A_\phi A_V/2RT](30.65/12 \text{ kg}^{1/2}\cdot\text{mol}^{-1/2})^2 \quad (18)$$

The second term in the brackets will be very small and can, probably, be neglected in comparison with the first term.

### Sodium Sulfate

Our results for the apparent molar volumes of sodium sulfate extend up to 0.33 *m*, to 473 K, and to 100 bar. We used our results along with the literature values up to 1 *m* and 100 bar in order to get an equation for the apparent molar volumes up to 473 K. These data were fitted to eq 15. The fitting parameters  $\bar{V}_2^0$ ,  $\beta^{(0)V}$ , and  $\beta^{(1)V}$  in eq 15 were assumed to be temperature- and pressure-dependent and the triple interaction parameter  $C_{MX}^V$  was assumed zero since our data extend only to 1 mol·kg<sup>-1</sup>. The parameter  $\beta^{(2)V}$  was also zero in this case—it is needed only for 2:2 electrolytes where there is a significant tendency toward ion pairing. Several different functional forms of the three fitting parameters were tried and the following parameters were found to be necessary to fit all the data within uncertainties.

$$\bar{V}_2^0 = p_1/T + p_2 + p_3T + p_4T^2 + p_5T^3 + p_6P \quad (19a)$$

$$\beta^{(0)V} = p_7/T + p_8 + p_9T + p_{10}T^2 \quad (19b)$$

$$\beta^{(1)V} = p_{11}/T + p_{12} \quad (19c)$$

Here the  $p_i$ 's are the fitting parameters, *T* is the temperature in kelvin, and *P* is the pressure in bar.

The uncertainties in the data were given as quoted by the authors or as judged by us. In general, the uncertainties in the data were increased with increasing temperature. Also a few values from the literature at very low molalities showed larger deviations and therefore were assigned larger uncertainties. The data of Fabuss et al. (6) showed substantial deviations from the fit and from other literature values and were not used in the final fit. The overall standard deviation of the fit was 0.28 cm<sup>3</sup>·mol<sup>-1</sup>.

### Magnesium Sulfate

Our results for magnesium sulfate extend up to 0.25 *m*, to 473 K, and to 100 bar. Again we used our results along with the literature values up to 0.25 *m* and 100 bar in order to get an equation for the apparent molar volumes up to 473 K. These data were fitted to eq 15.

Because of the rapid change of ion association at low molalities in MgSO<sub>4</sub>, it is not possible to obtain an accurate value of  $\bar{V}_2^0$  by extrapolation. Therefore this quantity was obtained indirectly as follows

$$\bar{V}_2^0(\text{MgSO}_4) = \bar{V}_2^0(\text{MgCl}_2) + \bar{V}_2^0(\text{Na}_2\text{SO}_4) - 2\bar{V}_2^0(\text{NaCl})$$

where the equation for NaCl is available from literature (3, 4), for Na<sub>2</sub>SO<sub>4</sub> is described in this paper, and for MgCl<sub>2</sub> has been developed recently from the results of Saluja (15) and others in literature and is being prepared for publication separately (16).  $\bar{V}_2^0$  for MgSO<sub>4</sub> was assumed pressure-independent because the pressure effect for the sum of MgCl<sub>2</sub> and Na<sub>2</sub>SO<sub>4</sub> is roughly the same as twice that for NaCl and therefore cancels out. The values of  $\bar{V}_2^0$  for MgSO<sub>4</sub> thus obtained from 273 to 473 K were fitted to a temperature-dependent equation. The parameters  $\beta^{(0)V}$  and  $C_{MX}^V$  were not needed in our concentration range and therefore the fitting parameters of eq 15 were

only  $\beta^{(1)V}$  and  $\beta^{(2)V}$ . The following equations were found necessary to fit the array of data.

$$\bar{V}_2^0 = q_1/T + q_2 + q_3T + q_4T^2 + q_5T^3 \quad (20a)$$

$$\beta^{(1)V} = q_6/T + q_7 + q_8T + q_9T^2 \quad (20b)$$

$$\beta^{(2)V} = q_{10}/T + q_{11} + q_{12}T + q_{13}T^2 + q_{14}T^3 + q_{15}P \quad (20c)$$

where  $q_i$ 's are the fitting parameters.

The pattern of uncertainties for MgSO<sub>4</sub> was the same as that for Na<sub>2</sub>SO<sub>4</sub> and the data of Fabuss et al. (6) were ignored in this case also. The overall standard deviation of the fit was 0.29 cm<sup>3</sup>·mol<sup>-1</sup>.

### Results

Table III gives the parameters for the eq 19 and 20 for the apparent molar volume of Na<sub>2</sub>SO<sub>4</sub> and MgSO<sub>4</sub>.

Sections A and B of Table IV give the values of the apparent molar volume of Na<sub>2</sub>SO<sub>4</sub> at rounded molalities, temperatures, and pressures. Table IVa gives the values at 1 atm or saturation pressure up to 1 mol·kg<sup>-1</sup>. Table IVb gives the values at 100 bar up to 0.4 mol·kg<sup>-1</sup>. Our equation for Na<sub>2</sub>SO<sub>4</sub> is valid only up to 0.33 mol·kg<sup>-1</sup> to 100 bar, and up to 1 mol·kg<sup>-1</sup> to 20 bar.

Sections A and B of Table V give the values of the apparent molar volume of MgSO<sub>4</sub> at rounded molalities, temperatures, and pressures.

### Discussion

Figure 1 shows experimental values at 373 and 423 K for Na<sub>2</sub>SO<sub>4</sub> and the curves from the equations adopted as well as the limiting Debye–Hückel slopes. The curve for 298 K is also shown for comparison. Even at 423 K the extrapolation to zero molality is straightforward while the curves at lower temperatures are nearly linear. Figure 2 shows the two ion-interaction parameters as a function of temperature. There is no indication of ion pairing since that would yield a positive  $\beta^{(1)V}$  whereas the actual trend of  $\beta^{(1)V}$  is to negative values at higher temperature.

Figure 3 shows the experimental points, calculated curves, and Debye–Hückel limiting slopes for MgSO<sub>4</sub> at 373 and 423 K with the 298 K curve for comparison. Here the large positive departures from the Debye–Hückel slopes are a clear indication of ion-pairing. This effect is also apparent in Figure 4 which shows a large, positive  $\beta^{(2)V}$  which also increases rapidly above 423 K. The  $\beta^{(1)V}$  parameter for MgSO<sub>4</sub> is relatively small and shows a negative slope with *T*.

The large, positive value of  $\beta^{(2)V}$  indicates that ion-pairing decreases with increase in pressure, which is the expected effect.

Registry No. Na<sub>2</sub>SO<sub>4</sub>, 7757-82-6; MgSO<sub>4</sub>, 7487-88-9.

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## Ultrasonic Velocities in, and Adiabatic Compressibilities for, Binary Liquid Mixtures of 1,2-Dichloroethane with Benzene, Toluene, *p*-Xylene, Quinoline, and Cyclohexane

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Measurements of ultrasonic velocities ( $u$ ) and adiabatic compressibilities ( $k_s$ ) have been made for mixtures of 1,2-dichloroethane ( $\text{CH}_2\text{ClCH}_2\text{Cl}$ ) with benzene, toluene, and quinoline at 303.15 and 313.15 K, for mixtures of  $\text{CH}_2\text{ClCH}_2\text{Cl}$  with *p*-xylene at 303.15 and 318.15 K, and for mixtures of  $\text{CH}_2\text{ClCH}_2\text{Cl}$  with cyclohexane at 308.15 K. The values of the quantity  $\Delta k_s$ , which refers to the deviations of the experimental values of  $k_s$  for mixtures from the mole fraction mixture law values, have been found to be positive for  $\text{CH}_2\text{ClCH}_2\text{Cl}$ -benzene,  $\text{CH}_2\text{ClCH}_2\text{Cl}$ -toluene,  $\text{CH}_2\text{ClCH}_2\text{Cl}$ -*p*-xylene, and  $\text{CH}_2\text{ClCH}_2\text{Cl}$ -cyclohexane, and negative for  $\text{CH}_2\text{ClCH}_2\text{Cl}$ -quinoline.

### Introduction

Quite recently, Nath and Singh (1) made measurements of excess volumes for binary liquid mixtures of 1,2-dichloroethane ( $\text{CH}_2\text{ClCH}_2\text{Cl}$ ) with benzene, toluene, *p*-xylene, quinoline, and cyclohexane at different temperatures, and the results obtained have been discussed from the viewpoint of the existence of specific interaction between the components of the various mixtures. It has been indicated (1) that there exists specific interaction leading to the formation of molecular complexes of  $\text{CH}_2\text{ClCH}_2\text{Cl}$  with the aromatic hydrocarbons and quinoline in the liquid state. Since the adiabatic compressibilities as obtained from ultrasonic velocities in binary liquid mixtures are known (2-5) to shed light on the existence of specific interaction between the components, it was thought worthwhile to get further information concerning the formation of adducts between the components of the binary mixtures of  $\text{CH}_2\text{ClCH}_2\text{Cl}$  with the aromatic hydrocarbons, and quinoline, from measurements of ultrasonic velocities. Hence, in the present program, we have made measurements of ultrasonic velocities in, and adiabatic compressibilities for, the binary liquid mixtures of  $\text{CH}_2\text{ClCH}_2\text{Cl}$  with benzene, toluene, *p*-xylene, quinoline, and cyclohexane and the results obtained have been interpreted in this paper.

Table I. Values of Ultrasonic Velocities in, and Adiabatic Compressibilities<sup>a</sup> for, Various Pure Liquids at 303.15 K

liquid	$u$ , m s <sup>-1</sup>			$10^{12}k_s$ , Pa <sup>-1</sup>		
	this work	lit.	ref	this work	lit.	ref
benzene	1278	1278	8	705	705	8
toluene	1284	1284.5	8	707	706.6	8
<i>p</i> -xylene	1289	1288	3	706	707	3
1,2-dichloroethane	1175	1175 <sup>b</sup>	5	585	585	5

<sup>a</sup> Values of  $k_s$  calculated from eq 1, by using densities obtained from the data reported in ref 7. <sup>b</sup> Value obtained by interpolation.

### Experimental Section

**Materials.** The methods of purifying the various components and checking their purity have been described earlier (1).

**Method.** The ultrasonic velocities,  $u$ , in pure liquids and their binary mixtures were measured with a single-crystal interferometer (supplied by Mittal Enterprises, New Delhi) at a frequency of 2 MHz s<sup>-1</sup>. The accuracy in the values of  $u$  is of the order of  $\pm 1.0$  m s<sup>-1</sup>. The adiabatic compressibilities,  $k_s$ , were calculated from the relation (3-6)

$$k_s = u^{-2}\rho^{-1} \quad (1)$$

where  $\rho$  refers to the density. The densities used to calculate  $k_s$  for pure liquids at various temperatures were obtained from the data reported by Timmermans (7), whereas the densities used to calculate  $k_s$  for mixtures were estimated from the densities of pure liquids, and the data on excess volumes for the various mixtures, as reported earlier (1). The accuracy in the values of  $10^{12}k_s$  is of the order of  $\pm 1.0$  Pa<sup>-1</sup>.

The working of the interferometer was tested by making measurements for pure liquids benzene, toluene, *p*-xylene, and 1,2-dichloroethane at 303.15 K. Table I shows that present values of  $u$  and  $k_s$  for these liquids are in good agreement with the data available in literature (3, 5, 8).

### Results and Discussion

The experimental values of ultrasonic velocities in, and adi-